

From weak to strong matter-light coupling with organic materials

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St Andrews

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Paradigms & Models

- Weakly interacting dilute Bose gas

$$H = \int d^d r \hat{\psi}^\dagger (-\mu - \nabla^2) \hat{\psi} + U \hat{\psi}^\dagger \hat{\psi}^\dagger \hat{\psi} \hat{\psi}$$

- ▶ Single field — assumes strong coupling
- ▶ Continuum model, hard to include molecular physics

- Laser rate equations

- ▶ Emission, absorption — assumes weak coupling, lasing

- Complex Gross-Pitaevskii/Ginzburg Landau equations

$$i\partial_t \psi = \left(-\nabla^2 \psi + V(r) + U|\psi|^2 \right) \psi + i(P(\psi, n, r) - \kappa) \psi$$

- ▶ Applies to laser, condensate — fluids of light
- ▶ Continuum theory

- Microscopic model ...

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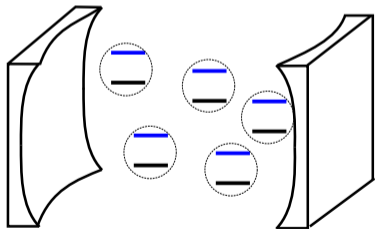
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Holstein-Tavis-Cummings model

Model capable of lasing & condensation

$$H = \omega \hat{a}^\dagger \hat{a} + \sum_{i=1}^N \left[\omega_X \sigma_i^+ \sigma_i^- + g (\sigma_i^+ \hat{a} + \text{H.c.}) \right]$$

• Including molecular physics

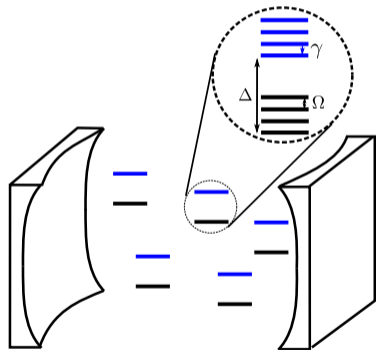


Cwik *et al.* EPL 105 '14; Spano, J. Chem. Phys '15; Galego *et al.* PRX '15; Cwik *et al.* PRA '16; Herrera & Spano PRL '16; Wu *et al.* PRB '16; Zeb *et al.* arXiv:1608.08929; Herrera & Spano arXiv:1610.04252;

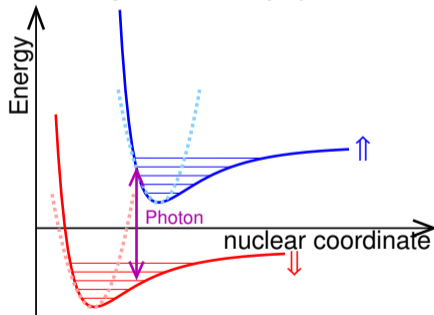
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Introduction and models

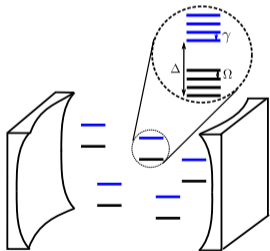
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 - Holstein-Dicke model
- 2 Weak coupling
 - Photon BEC
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 - Exact polariton states, scaling with N
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Photon: Microscopic Model

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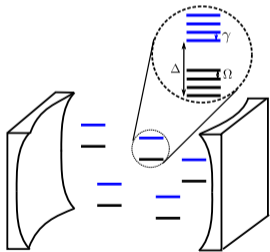


- **2D** harmonic oscillator $\omega_m = \omega_{\text{cutoff}} + m\omega_{H.O.}$

- Incoherent processes: excitation, decay, loss, vibrational thermalisation.
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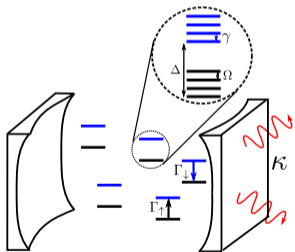


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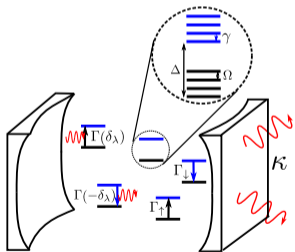
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Steady state populations and equilibrium

Rate equation: $\partial_t n_m = -\kappa n_m + \Gamma(-\delta_m)(n_m + 1)N_{\uparrow} - \Gamma(\delta_m)n_m N_{\downarrow}$

• Steady state distribution:

$$\frac{n_m}{n_{m+1}} = \frac{\Gamma(-\delta_m)N_{\uparrow}}{\Gamma(\delta_m)N_{\downarrow}}$$

• Microscopic conditions for equilibrium:

• Emission/absorption rate:

$$\Gamma(\delta) \simeq 2g^2 \text{Re} \left[\int dt e^{-i\delta t} \langle D_{\downarrow}^{\dagger}(t) D_{\downarrow}(0) \rangle \right] \quad D_{\downarrow} = \exp \left(2\lambda_0 (\delta_{\downarrow} - \delta_{\downarrow}^{\dagger}) \right)$$

[Kirton & JK PRL '13]

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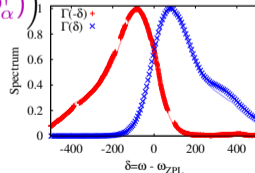
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$$\langle D_{\alpha}(0) D_{\alpha}(t) \rangle = \langle D_{\alpha}(-t - i\eta) D_{\alpha}(0) \rangle \leftrightarrow \Gamma(+\delta) = \Gamma(-\delta) e^{\beta \hbar \delta}$$

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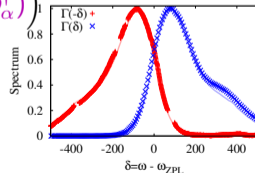
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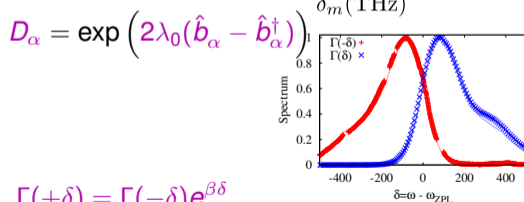
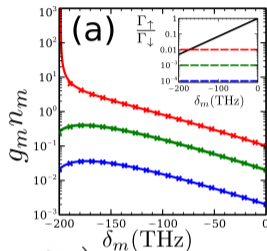
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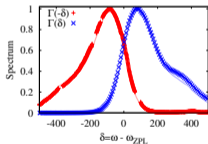
Strong coupling

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Strong coupling: One excitation subspace

$$H = \omega \hat{a}^\dagger \hat{a} + \sum_{i=1}^N \left[\omega_X \sigma_i^+ \sigma_i^- + g (\sigma_i^+ \hat{a} + \text{H.c.}) + \omega_V \left(\hat{b}_i^\dagger \hat{b}_i - \lambda_0 \sigma_i^+ \sigma_i^- (\hat{b}_i^\dagger + \hat{b}_i) \right) \right]$$

Strong coupling: fate of spectrum

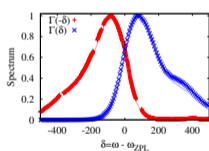


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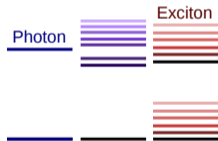
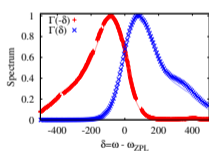


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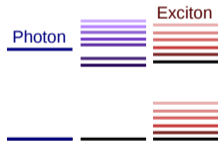
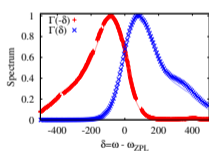


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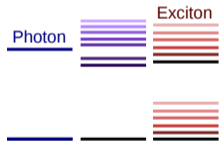
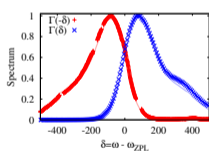
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Exact solution, $N = 20$

- Vibrational Wigner function:

$$W(x, p) = \int dy \langle x + y/2 | \rho | x - y/2 \rangle_i e^{iyp}, \quad \left(\frac{\hat{b}_i + \hat{b}_i^\dagger}{\sqrt{2}} \right) |x\rangle_i = x|x\rangle_i$$

- Conditioned on Photon $|P\rangle$ /Exciton at i $|X\rangle_i$ /Other site $|X\rangle_{j \neq i}$

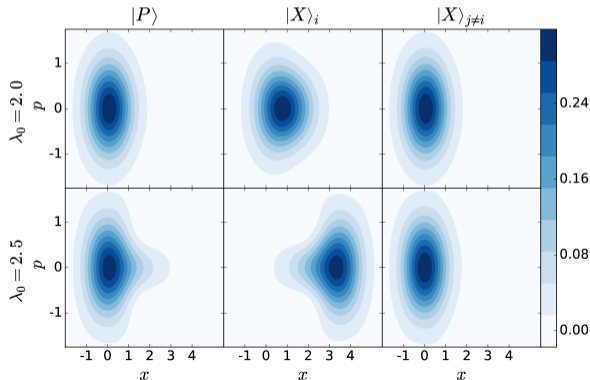
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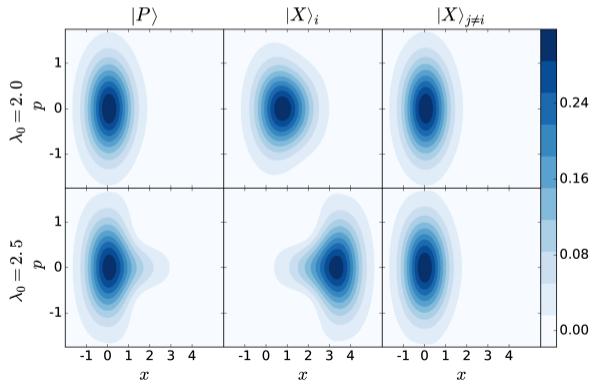
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Extending to arbitrary N , polaron ansatz

- Polaron transform, $\mathcal{D}_i(\lambda) = \exp(\lambda(\hat{b}_i^\dagger - \hat{b}_i))$

• N site polaron ansatz

$$|\Psi\rangle = \left[\alpha |P\rangle \prod_I \mathcal{D}_I(\lambda_a) + \frac{\beta}{\sqrt{N}} \sum_I |\mathcal{X}\rangle_I \mathcal{D}_I(\lambda_b) \prod_{I \neq I'} \mathcal{D}_{I'}(\lambda_c) \right] |0\rangle_V$$

[Wu *et al.* PRB '16, Zeb *et al.* arXiv:1608.08929]

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→ Allows distinct Wigner functions $|P\rangle, |X\rangle_i, |X\rangle_{j \neq i}$

→ Polaron energy: $E_{LP} = \frac{\tilde{\omega}_X + \tilde{\omega}_P}{2} - \sqrt{\left(\frac{\tilde{\omega}_X + \tilde{\omega}_P}{2}\right)^2 + \tilde{\omega}_B^2}$

$$\tilde{\omega}_X = \omega_X + \omega_v(\lambda_a^2 - 2\lambda_a\lambda_b + (N-1)\lambda_b^2), \quad \tilde{\omega}_P = \omega + \omega_v N \lambda_a^2$$

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Extending to arbitrary N , polaron ansatz

- Polaron transform, $\mathcal{D}_i(\lambda) = \exp(\lambda(\hat{b}_i^\dagger - \hat{b}_i))$
- N site polaron ansatz

$$|\Psi\rangle = \left[\alpha |P\rangle \prod_j \mathcal{D}_j(\lambda_a) + \frac{\beta}{\sqrt{N}} \sum_i |X\rangle_i \mathcal{D}_i(\lambda_b) \prod_{j \neq i} \mathcal{D}_j(\lambda_c) \right] |0\rangle_V$$

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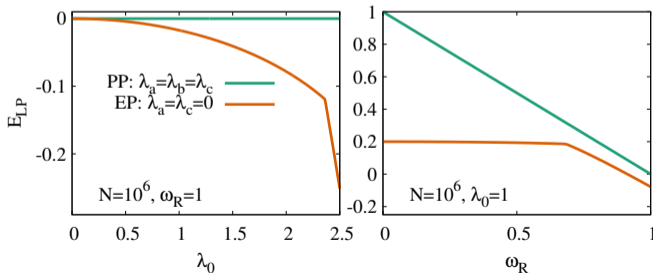
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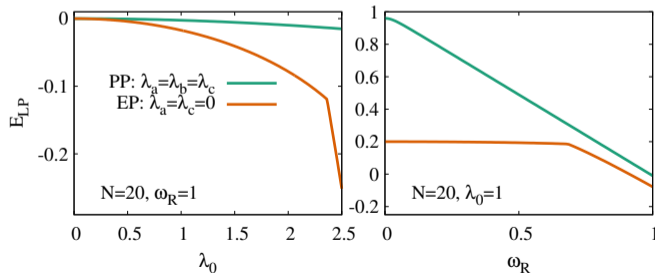
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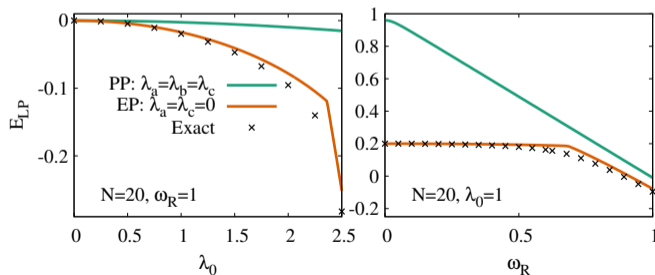
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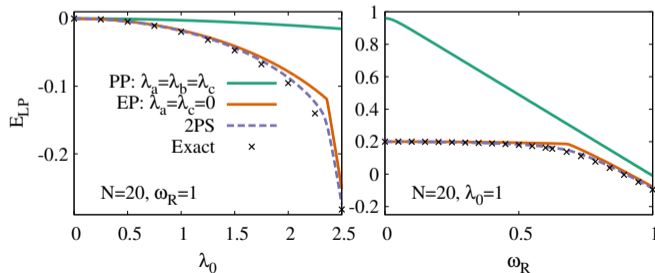
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- Minimisation: Multipolaron ansatz: bimodal Wigner



Strong coupling

- 1 Introduction and models
 - Holstein-Dicke model
- 2 Weak coupling
 - Photon BEC
- 3 Strong coupling**
 - Exact polariton states, scaling with N
 - Spectrum
- 4 Ultra strong coupling
 - Vibrational reconfiguration
 - Vibrations and disorder

Tavis-Cummings-Holstein spectrum

- Direct calculation

$$D^R(t) = -i \langle 0 | \left[\hat{a}(t), \hat{a}^\dagger(0) \right] | 0 \rangle \theta(t)$$

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- Mean-field Green's function

- Multiple excitation $\sim 1/N$.

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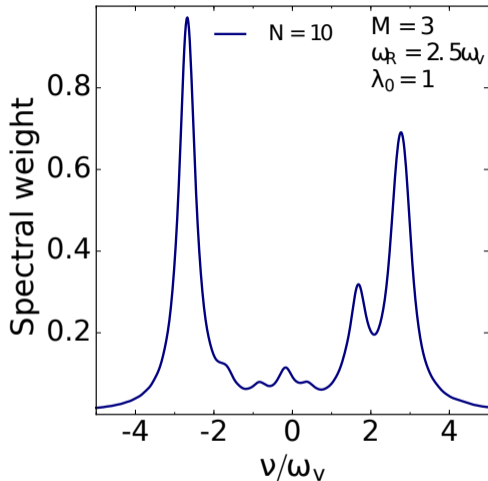
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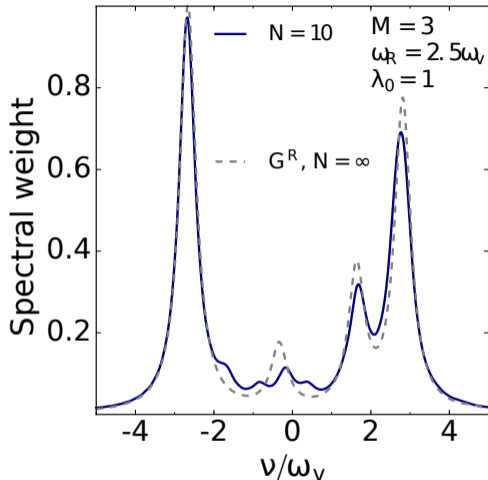
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(Classical expression)



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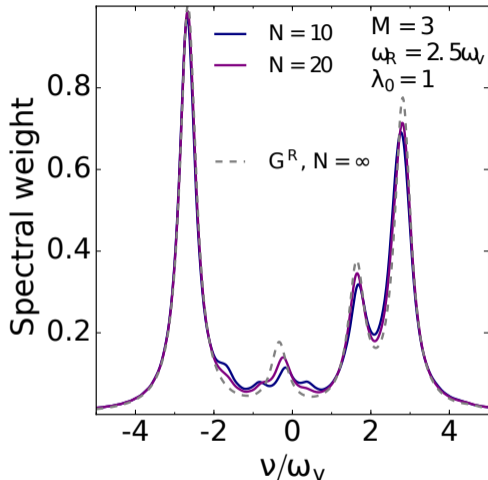
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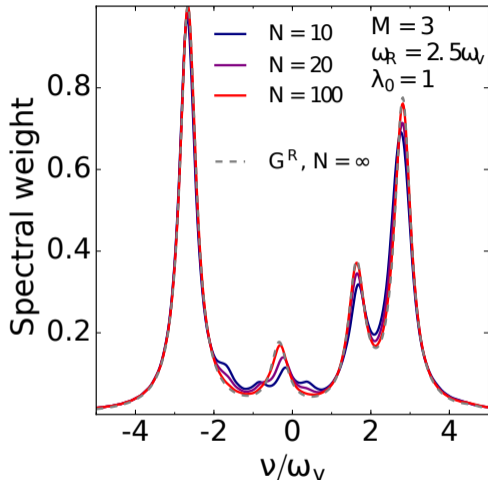
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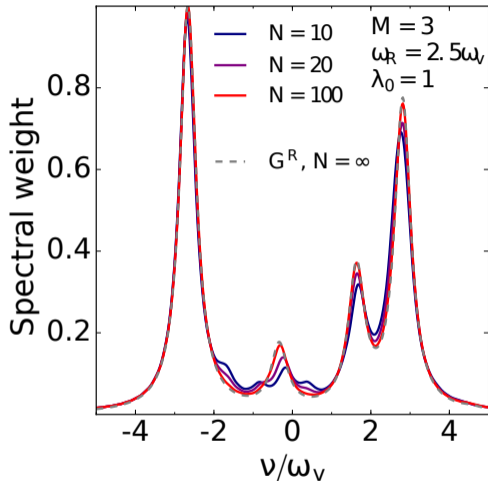
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Ultra strong coupling: vibrational reconfiguration

- 1 Introduction and models
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- 4 **Ultra strong coupling**
 - **Vibrational reconfiguration**
 - **Vibrations and disorder**

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Vibrational reconfiguration

- Many photon modes, **beyond RWA perturbatively**

$$H = \sum_k \omega_k \hat{a}_k^\dagger \hat{a}_k + \sum_{i=1}^N \left[\omega_X \sigma_i^+ \sigma_i^- + \sum_k g_{\mathbf{k}} \left(\sigma_i^+ (\hat{a}_k + \hat{a}_k^\dagger) + \text{H.c.} \right) \right]$$

- Reduced vibrational offset

$$\lambda_0 \rightarrow \lambda_0 (1 - K_1), \quad K_1 = \sum_k \frac{g_k^2}{(\omega_k + \omega_X)^2}$$

[Cwik *et al.* PRA '16]

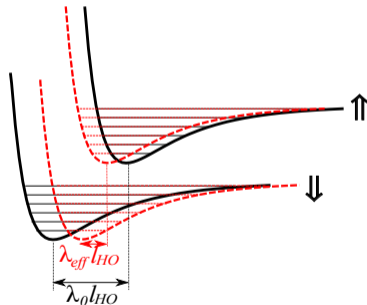
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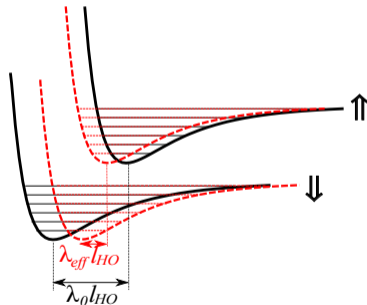
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• But, no collective effect: $\delta H \approx K_1 N$



[Cwik *et al.* PRA '16]

Vibrational reconfiguration

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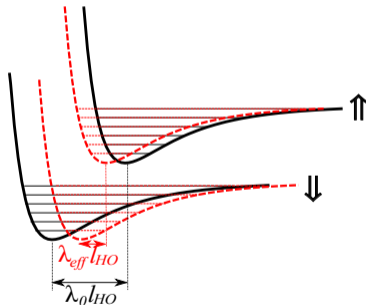
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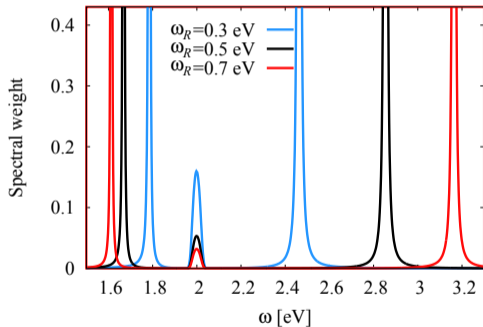
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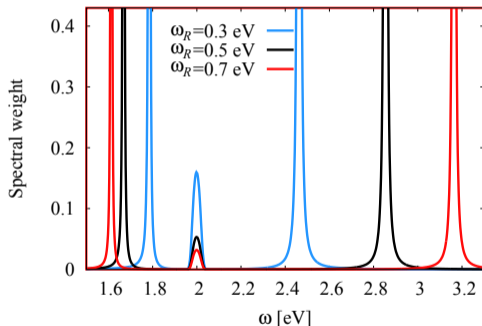
Bumps in the middle of the spectrum

- Origin of bumps in middle of spectrum: Disorder



Bumps in the middle of the spectrum

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- Central peak:

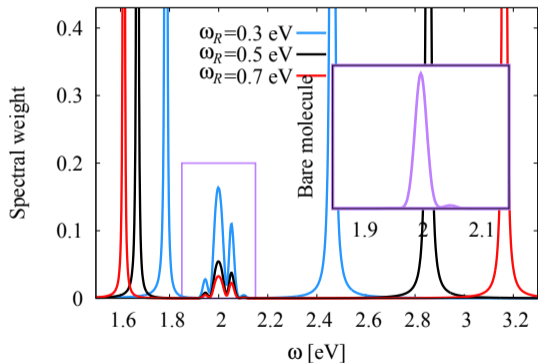
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Gaussian $\rho(x)$, variance σ_x [Houdré *et al.*, PRA '96]

Disorder + Vibrations + Strong coupling

- Disordered spectrum + vibrations,
 $\lambda_0^2 = 0.02 \ll 1, \sigma_x = 0.01\text{eV}$

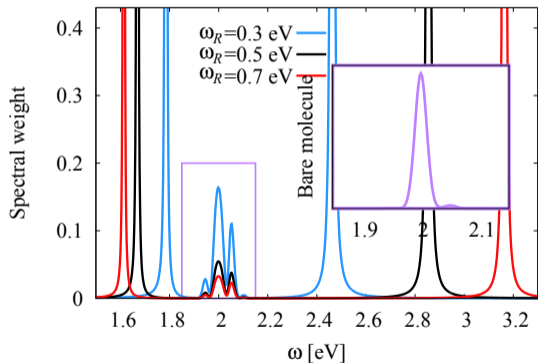
- Stronger disorder,
 $\lambda_0^2 = 0.5, \sigma = 0.025\text{eV}$



[Cwik *et al.* PRA '16]

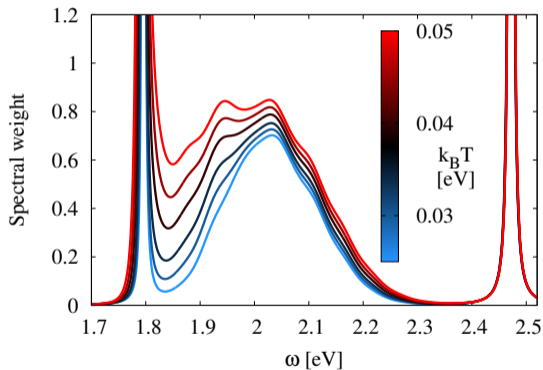
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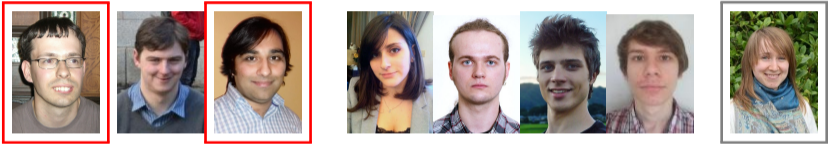
[Cwik *et al.* PRA '16]

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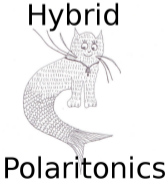


Acknowledgements

GROUP:

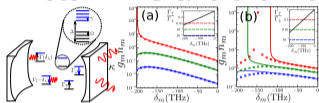


FUNDING:



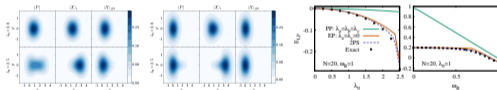
Conclusions

- Matter-light coupling & organic molecules: Holstein-Tavis-Cummings model
- Photon BEC and thermalisation [Kirton & JK, PRL '13, PRA '15]

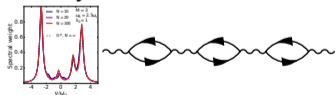


- Strong coupling, one excitation

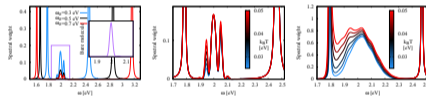
▶ Exact solution & Polaron Ansatz



▶ Validity of mean-field Green's functions



- Vibrations + disorder + USC



[Cwik *et al.* PRA '16]

[Zeb, Kirton, JK, arXiv:1608.08929]

Conclusions

- 5 Photon BEC Chemical potential
- 6 Spatial profile
- 7 Spectrum

Chemical potential?

- Steady state, thermalised:

$$\frac{n_m}{n_m + 1} = \frac{\Gamma(-\delta_m)N_{\uparrow}}{\kappa + \Gamma(\delta_m)N_{\downarrow}} \simeq e^{-\beta\delta_m + \beta\mu},$$

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Below threshold,

$$\mu = k_B T \ln[\Gamma_{\uparrow}/\Gamma_{\downarrow}]$$

At/above threshold, $\mu \rightarrow \delta_0$

[Kirton & JK, PRA '15]

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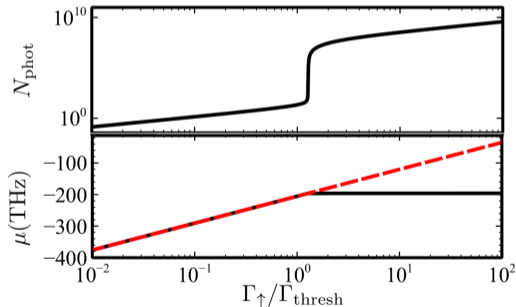
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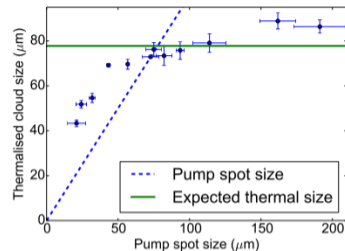


[Kirton & JK, PRA '15]

Spatial profile, pump-spot dependence

- Dependence on pump profile, $\Gamma_{\uparrow}(\mathbf{r}) = \frac{\Gamma_{\uparrow} \exp(-r^2/2\sigma_p^2)}{(2\pi\sigma_p^2)^{d/2}}$

[Marelic & Nyman, PRA '15]



• Use Gauss-Hermite modes $\hat{A}\hat{y} I(\mathbf{r}) = \sum_m n_m |\psi_m(\mathbf{r})|^2$

• Varying excited density – differential coupling to modes

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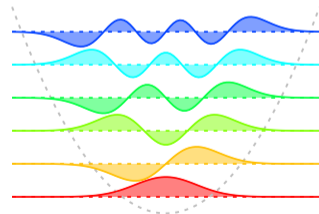
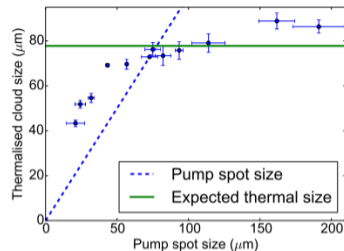
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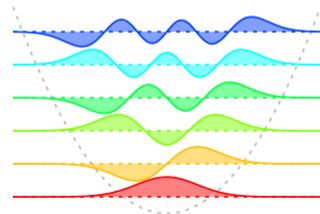
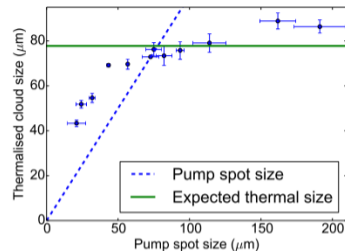
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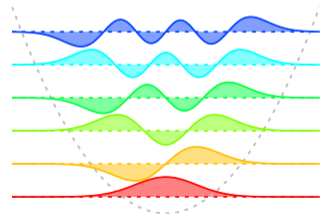
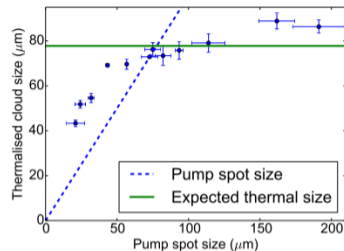
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Spatially varying pump: below threshold

- Far below threshold:

- ▶ If $\kappa \ll \rho_M \Gamma(\delta_m)$,
$$\frac{n_m}{n_m + 1} \simeq e^{-\beta \delta_m} \times \int d\mathbf{r} \frac{1}{\rho_M} \rho_{\uparrow}(\mathbf{r}) |\psi_m(\mathbf{r})|^2$$

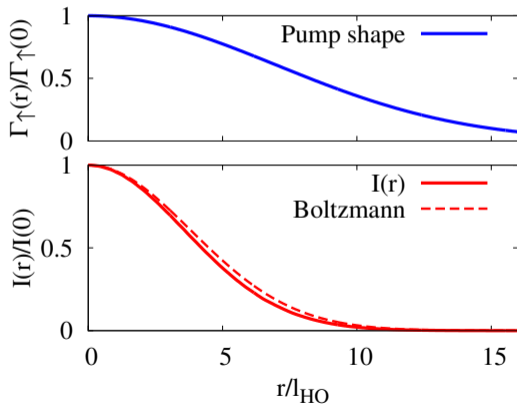
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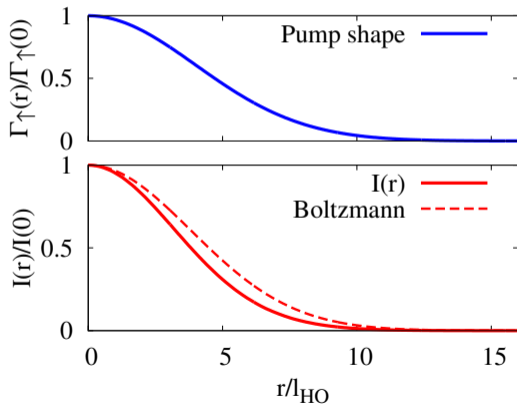


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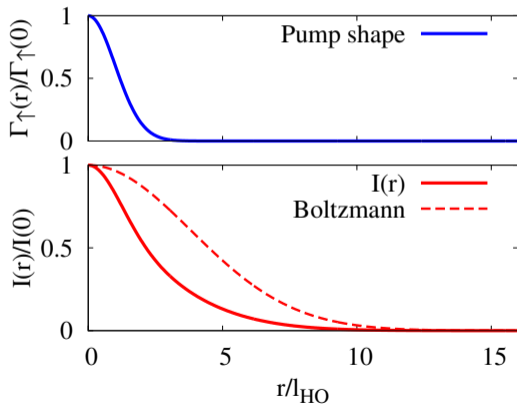


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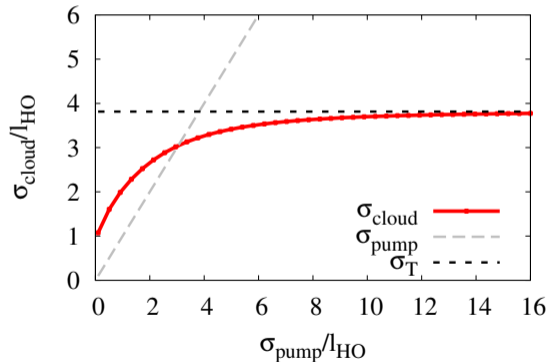
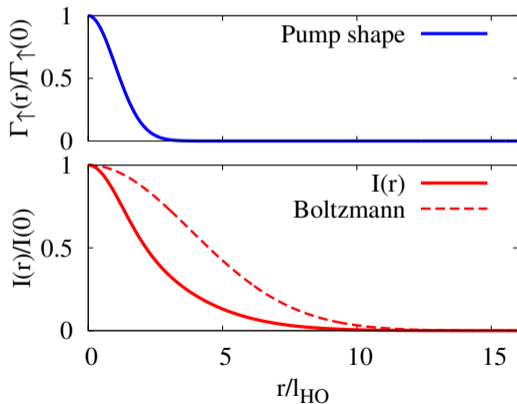


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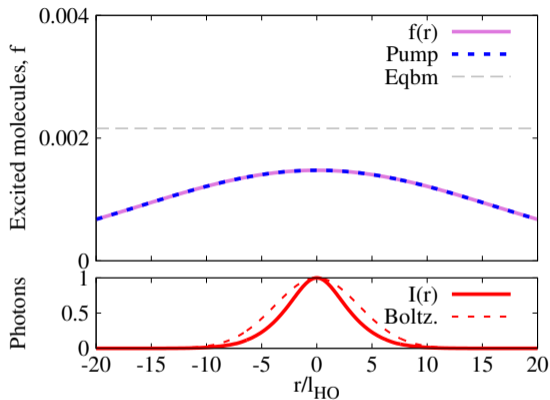
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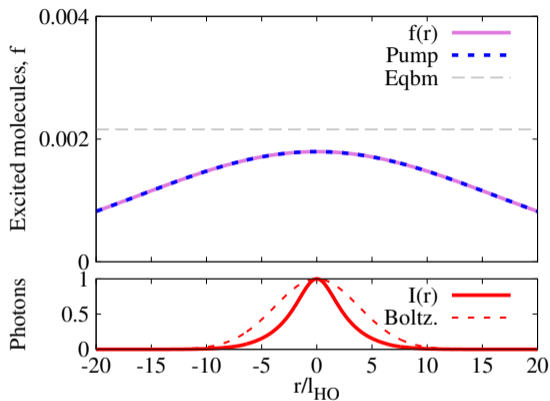


Near threshold behaviour



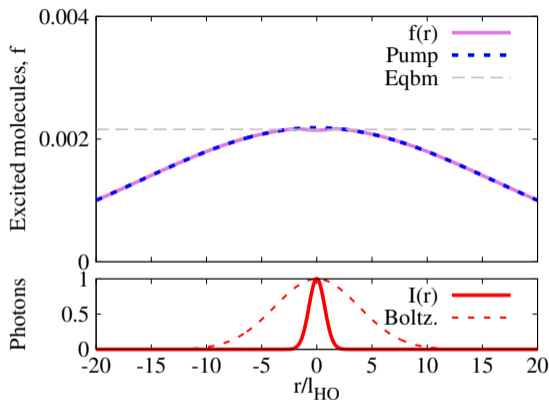
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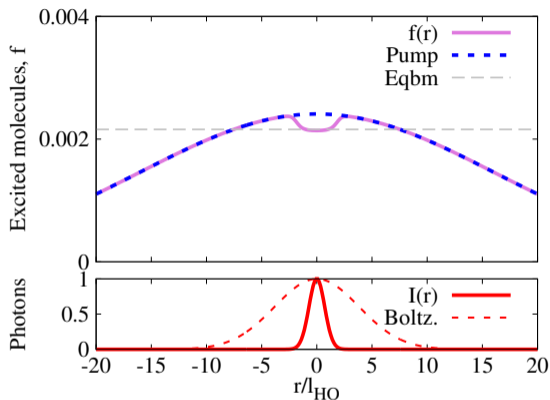
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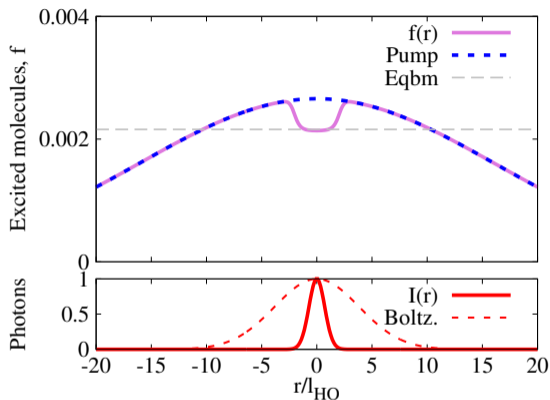
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Calculating spectra: Input-Output formalism

- Observable features: absorption spectrum, $A(\nu) = 1 - T(\nu) - R(\nu)$

- Scattering matrix gives:

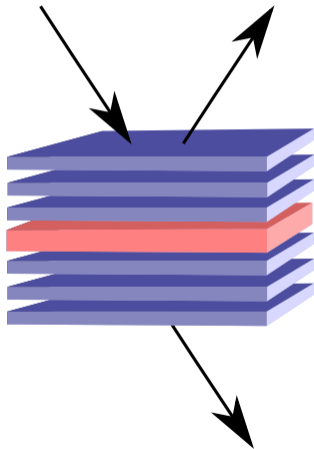
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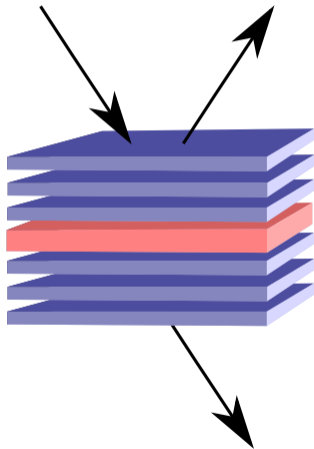
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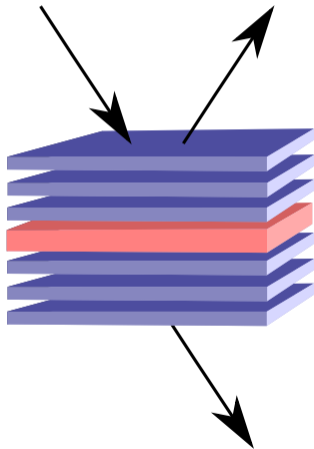
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Tavis-Cummings-Holstein vs Coupled Oscillators

- Diagrammatic approach — equivalent to coupled oscillator model:

$$H = \omega_P \hat{a}^\dagger \hat{a} + \sum_i \left[\frac{\omega_R}{\sqrt{N}} \left(\hat{a} \sum_n f_n(\lambda_0) \sigma_i^{n0} + \text{H.c.} \right) + \omega_n \sigma_i^{nn} \right]$$

$$\omega_n = \omega_X + n\omega_V, \quad f_n(\lambda_0) = \langle n | D(\lambda_0) | 0 \rangle$$

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$$\chi(\nu) = - \sum_n \frac{\omega_R^2 |f_n(\lambda_0)|^2}{\nu + i\gamma/2 - \omega_n}$$

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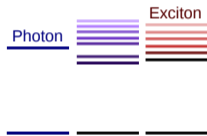
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