

Ground and excited states of vibrationally dressed polaritons

Jonathan Keeling



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St Andrews

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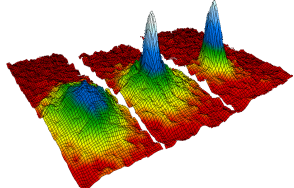


SCOM, October 2016



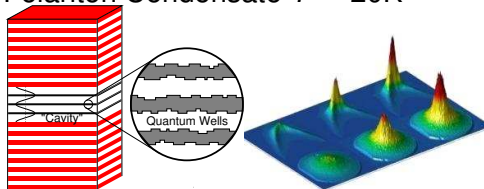
Condensation, Lasing, Superradiance

Atomic BEC $T \sim 10^{-7}K$



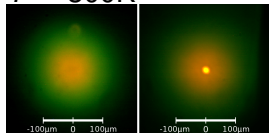
[Anderson *et al.* Science '95]

Polariton Condensate $T \sim 20K$



[Kasprzak *et al.* Nature, '06]

Photon Condensate
 $T \sim 300K$

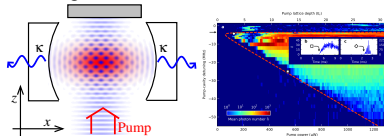


[Klaers *et al.* Nature, '10]

Laser
 $T \sim ?, < 0, \infty$



Superradiance transition
 $T \sim 0$



[Baumann *et al.* Nature '10]

Paradigms & Models

- Weakly interacting dilute Bose gas

$$H = \int d^d r \hat{\psi}^\dagger (-\mu - \nabla^2) \hat{\psi} + U \hat{\psi}^\dagger \hat{\psi}^\dagger \hat{\psi} \hat{\psi}$$

- ▶ Single field — assumes strong coupling
- ▶ Continuum model, hard to include molecular physics

- Laser rate equations

- ▶ Emission, absorption — assumes weak coupling, lasing

- Complex Gross-Pitaevskii/Ginzburg Landau equations

$$i\partial_t \psi = \left(-\nabla^2 \psi + V(r) + U|\psi|^2 \right) \psi + i(P(\psi, n, r) - \kappa) \psi$$

- ▶ Applies to laser, condensate — fluids of light
- ▶ Continuum theory

- Microscopic model ...

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Simple models, N body physics

- Model capable of lasing & condensation
 - ▶ Tavis-Cummings / Dicke model

$$H = \omega \hat{a}^\dagger \hat{a} + \sum_{i=1}^N \left[\omega_X \sigma_i^+ \sigma_i^- + g (\sigma_i^+ (\hat{a} + \hat{a}^\dagger) + \text{H.c.}) \right]$$

- ▶ Weak pumping \rightarrow Superradiance/BEC transition
- ▶ High temperature: Maxwell-Bloch laser

- Including molecular physics

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Szymanska et al. PRL 06; Keeling et al. book chapter 1010.3338

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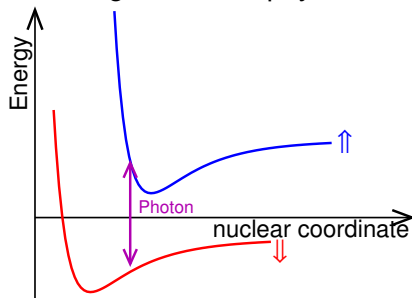
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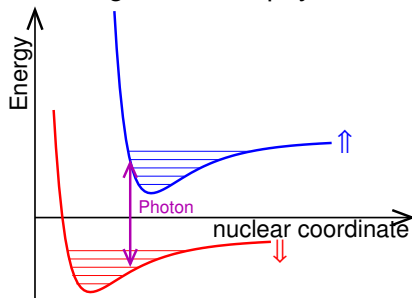
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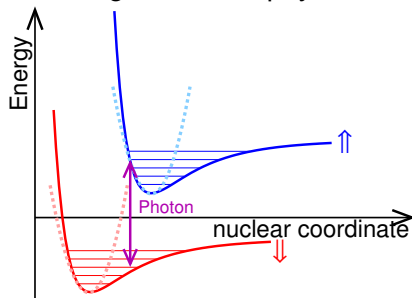
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Holstein-Tavis-Cummings & Holstein-Dicke model

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- Few emitters (molecules/quantum dots)

Wilson-Rae & Imamoglu PRB 2002; McCutcheon & Nazir PRB 2011; Roy & Hughes PRB 2011; Bera *et al.* PRB 2014; Pollock *et al.* NJP 2013; Hornecker *et al.* arXiv:1609.09754; ...

- Weak coupling

Kirton & JK, PRL 2013, PRA 2015; PRA 2016 ...

- Full model

Cwik *et al.* EPL 105 2014; Spano, J. Chem. Phys 2015; Galego *et al.* PRX 2015; Cwik *et al.* PRA 2016; Herrera & Spano PRL 2016; Wu *et al.* arXiv:1608.08019; Zeb *et al.* arXiv:1608.08929; Herrera & Spano arXiv:1610.04252

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Introduction and models

- 1 Introduction and models
 - Holstein-Dicke model
 - Weak coupling: photon BEC
- 2 Strong coupling: polariton states
 - Exact solutions
 - Scaling with N
- 3 Ultrastrong coupling: ground state changes

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HTC weak coupling

- Polaron transform (exact), $\mathcal{D}_i = e^{\lambda_0(\hat{b}_i^\dagger - \hat{b}_i)}$

$$H = \sum_m \omega_m \hat{a}_m^\dagger \hat{a}_m + \sum_i \left[\omega_X \sigma_i^+ \sigma_i^- + g (\hat{a}_m \sigma_i^+ \mathcal{D}_i + \text{H.c.}) + \omega_V \hat{b}_i^\dagger \hat{b}_i \right],$$

- Perturbation in g & dissipative processes

$$\begin{aligned} \dot{\rho} = & -i[H_0, \rho] + \sum_m \frac{\kappa}{2} \mathcal{L}[\hat{a}_m] + \sum_j \left[\frac{\Gamma_1}{2} \mathcal{L}[\sigma_j^+] + \frac{\Gamma_2}{2} \mathcal{L}[\sigma_j^-] \right] \\ & + \sum_{m,j} \left[\frac{\Gamma(\delta_m = \omega_m - \omega_X)}{2} \mathcal{L}[\sigma_j^+ \hat{a}_m] + \frac{\Gamma(-\delta_m = \omega_X - \omega_m)}{2} \mathcal{L}[\sigma_j^- \hat{a}_m^\dagger] \right] \end{aligned}$$

- Correlation function:

$$G(t) = 2g^2 \text{Re} \left[\int dt e^{-|t| - (\Gamma_1 + \Gamma_2)t/2} \langle \mathcal{D}_i^\dagger(t) \mathcal{D}_i(0) \rangle \right]$$

[Marthaler et al PRL '11, Kirton & JK PRL '13]

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$$C(t) = 2g^2 \text{Re} \left[\int dt e^{-iEt - (\Gamma_\uparrow + \Gamma_\downarrow)t/2} \langle \sigma_i^+(0) \sigma_i^-(0) \rangle \right]$$

[Marthaler et al PRL 111, Kirton & JK PRL 113]

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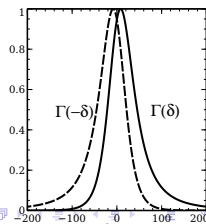
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$$\Gamma(\delta) = 2g^2 \text{Re} \left[\int dt e^{-i\delta t - (\Gamma_\uparrow + \Gamma_\downarrow)t/2} \langle D_i^\dagger(t) D_i(0) \rangle \right]$$

[Marthaler et al PRL '11, Kirton & JK PRL '13]



HTC weak-coupling and photon BEC

- Rate equation:

$$\partial_t n_m = -\kappa n_m + \Gamma(-\delta_m)(n_m + 1) \sum_i \langle \sigma_i^+ \sigma_i^- \rangle - \Gamma(\delta_m) n_m \sum_i \langle \sigma_i^- \sigma_i^+ \rangle$$

- Steady state distribution: $\frac{n_m}{n_m + 1} = \frac{\Gamma(-\delta_m) \sum_i \langle \sigma_i^+ \sigma_i^- \rangle}{\Gamma(\delta_m) \sum_i \langle \sigma_i^- \sigma_i^+ \rangle}$
- Kennard-Stepanov, $\Gamma(\delta) = \Gamma(-\delta) e^{\beta \hbar \delta}$

[Kirton & JK, PRL 2013, PRA 2015],

Expt: Weitz (Bonn), Nyman (Imperial), van Oosten (Utrecht)

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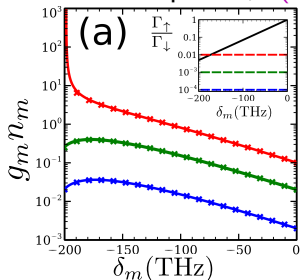
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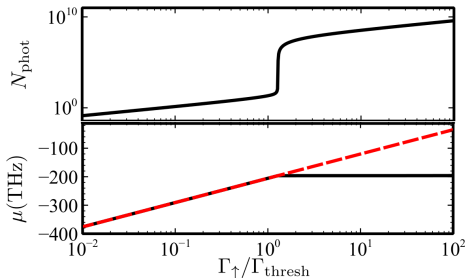
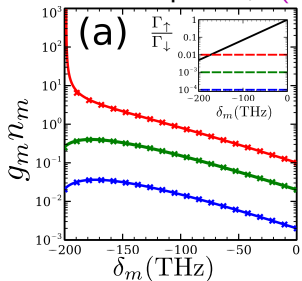
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Strong coupling: polariton states

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One excitation subspace, questions

$$H = \omega \hat{a}^\dagger \hat{a} + \sum_{i=1}^N \left[\omega_X \sigma_i^+ \sigma_i^- + g (\sigma_i^+ \hat{a} + \text{H.c.}) \right. \\ \left. + \omega_V \left(\hat{b}_i^\dagger \hat{b}_i - \lambda_0 \sigma_i^+ \sigma_i^- (\hat{b}_i^\dagger + \hat{b}_i) \right) \right]$$

- Rotating wave approximation — Holstein Tavis Cummings

- Restrict, $\hat{a}^\dagger \hat{a} + \sum_i \sigma_i^+ \sigma_i^- = 1$.

- Questions:

- Competition of $g\sqrt{N}$ vs ω_V , $\omega_V \lambda_0^2$

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 - ▶ Scaling with N

Exact solution, $N = 2$

Vibrational Wigner function:

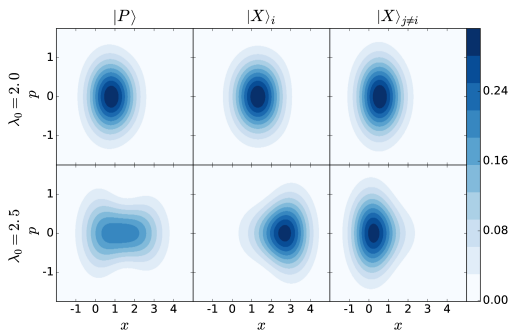
$$W(x, p) = \int dy \langle x + y/2 | \rho | x - y/2 \rangle_i e^{iyp}, \quad \left(\frac{\hat{b}_i + \hat{b}_i^\dagger}{\sqrt{2}} \right) |x\rangle_i = x|x\rangle_i$$

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Conditioned on Photon $|P\rangle$ /Exciton at i , $|X\rangle_i$ /Other site $|X\rangle_{j \neq i}$



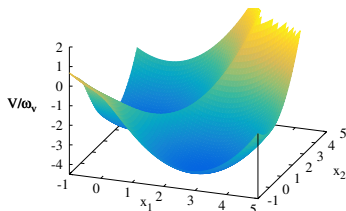
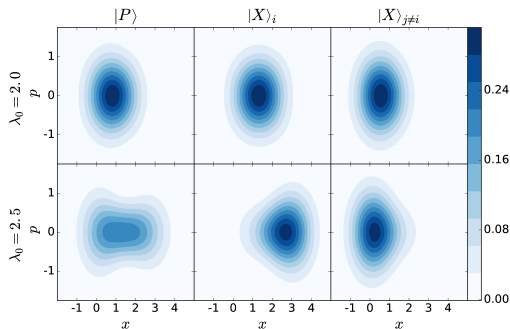
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Exact solution, $N = 2$

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$$N = 2, \omega = \omega_X, \omega_R \equiv g/\sqrt{N} = 1$$

Exact solution, larger N

- Brute force approach, N sites, $\hat{b}^\dagger \hat{b} < M$, $D_{\text{Hilbert}} = M^N$

• Permutation symmetry. $D_{\text{Hilbert}} \sim N^M$, typical $M \sim 5 - 6$

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- Distinct behaviour vs λ_0
- Exact energy and state vs ω_R, λ_0 for validation

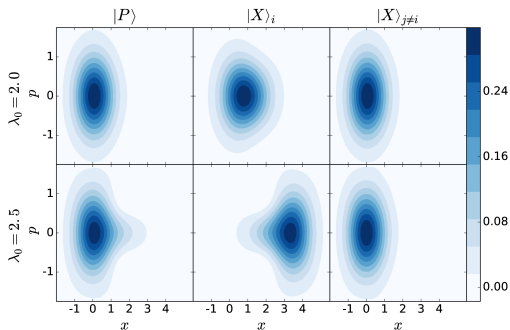
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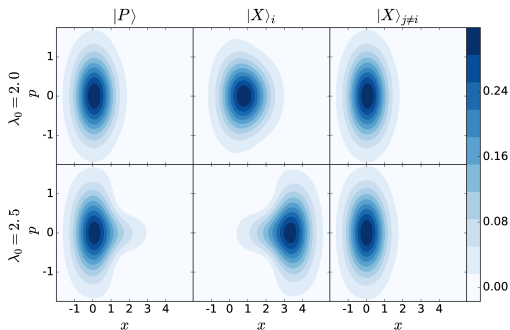
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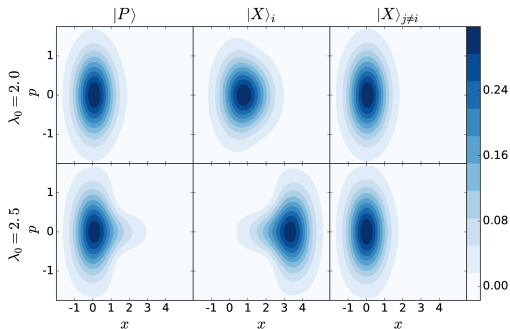
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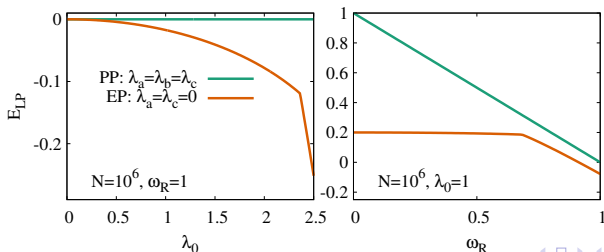
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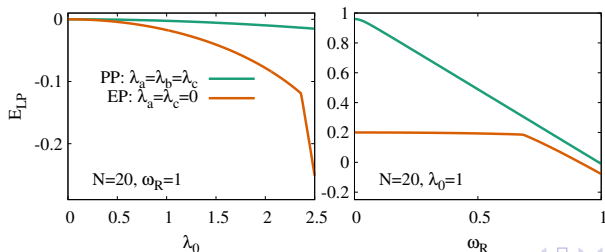
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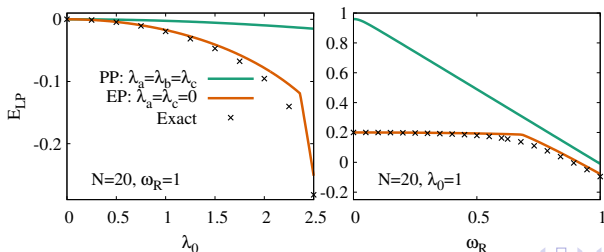
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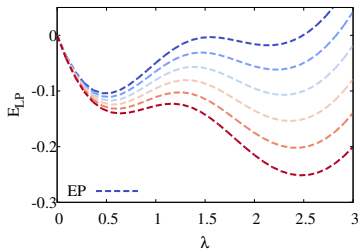
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Polaron crossover

- Crossover near $\omega_R \simeq \omega_V \lambda_0^2$

[Silbey and Harris, J. Chem. Phys. 1984]

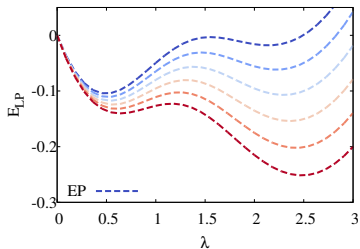


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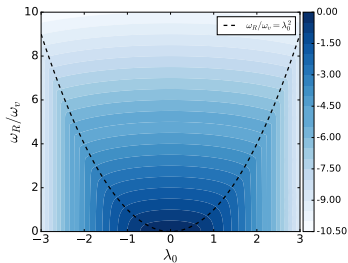
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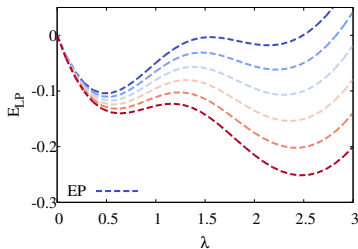


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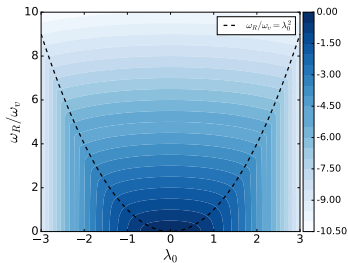
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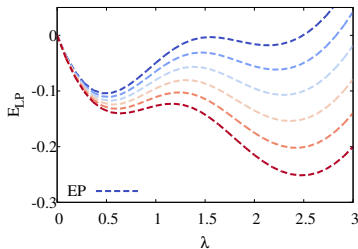
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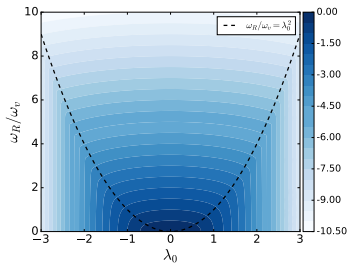
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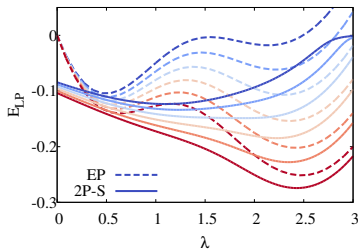


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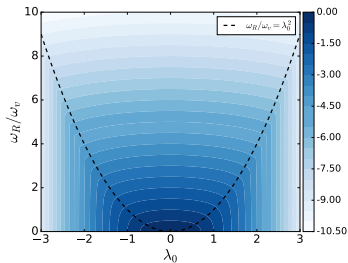
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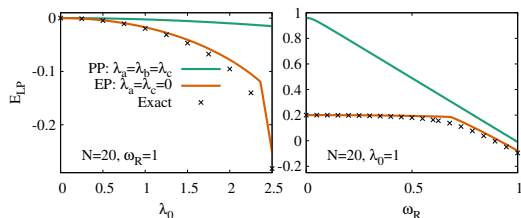


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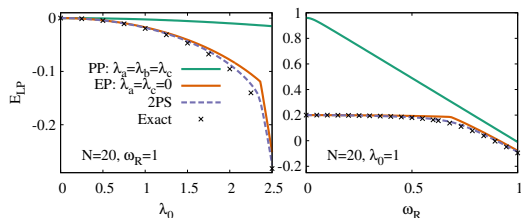
- Accurate energy & wavefunction



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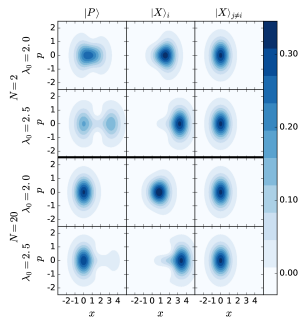
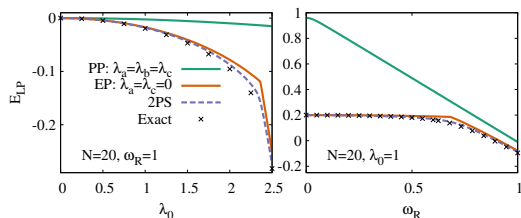
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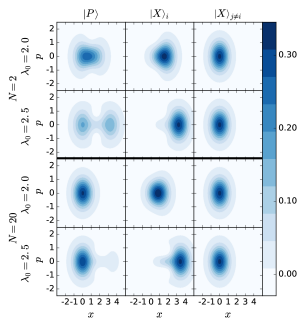
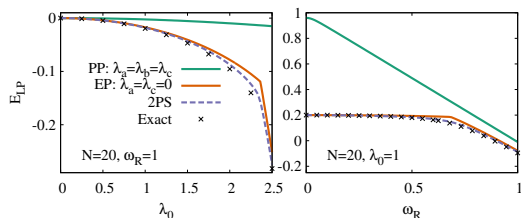
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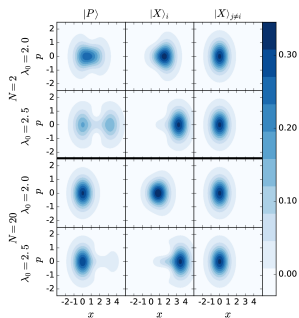
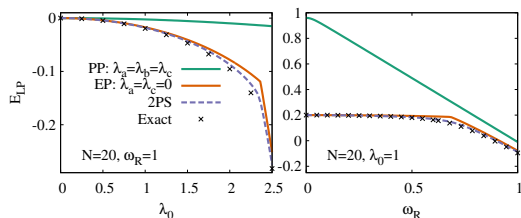
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Ultrastrong coupling: ground state changes

- 1 Introduction and models
 - Holstein-Dicke model
 - Weak coupling: photon BEC
- 2 Strong coupling: polariton states
 - Exact solutions
 - Scaling with N
- 3 Ultrastrong coupling: ground state changes

Ground state molecular reconfiguration

- Dicke model: beyond rotating wave approximation

$$H = \sum_K \omega_k \hat{a}_k^\dagger \hat{a}_k + \sum_{i=1}^N \left[\omega_X \sigma_i^+ \sigma_i^- + \sum_k g_{\mathbf{k}} \left(\sigma_i^+ (\hat{a}_k + \hat{a}_k^\dagger) + \text{H.c.} \right) + \dots \right]$$

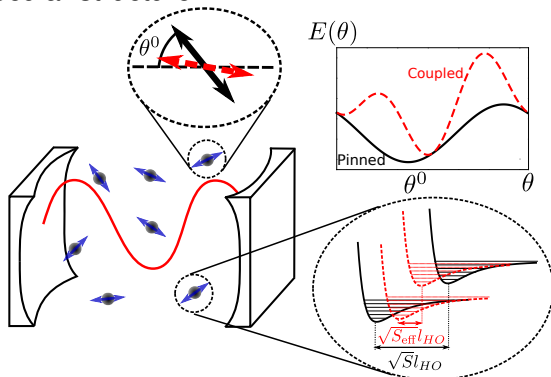
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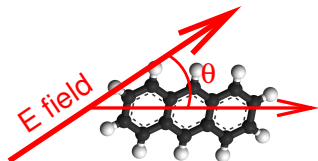
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Rotational reorientation

- Rotational degrees of freedom



- Effective Hamiltonian

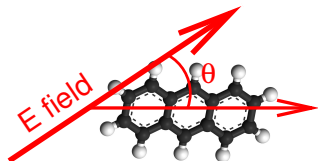
$$H = \dots + \sum_{i,k} \left[\dots + g_{i,k} \cos(\theta_i) (\hat{a}_k^\dagger + \hat{a}_{-k}) \sigma_i^z + E_0(\theta_i) \right]$$

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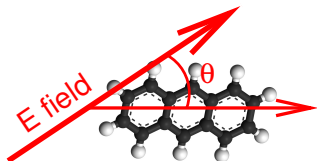
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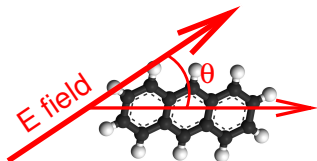
- Schrieffer-Wolff, $\delta H = \sum_{i,k} g_{i,\mathbf{k}} (\hat{a}_k^\dagger \sigma_i^+ + \text{H.c.})$:

$$H_{\text{eff}} = \dots + \sum_i \left[-K_0 \cos^2(\theta_i) + E_0(\theta_i) \right], \quad K_0 = \sum_k \frac{g_k^2}{\omega_k + \omega_X}$$

• No \sqrt{N} enhancement $\rightarrow K_0$ small, independent of density

Rotational reorientation

- Rotational degrees of freedom



- Effective Hamiltonian

$$H = \dots + \sum_{i,k} \left[\dots + g_{i,\mathbf{k}} \cos(\theta_i) (\hat{a}_{\mathbf{k}}^\dagger + \hat{a}_{-\mathbf{k}}) \sigma_i^x + E_0(\theta_i) \right]$$

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- ▶ No \sqrt{N} enhancement — K_0 small, independent of density

Vibrational reconfiguration

- Schrieffer-Wolff – mixes vibrational states

$$\delta H = - \sum_{i,k} \frac{g_k^2}{2(\omega_X + \omega_k)} \left\{ 1 - \frac{\omega_V \lambda_0 (b_i + b_i^\dagger)}{\omega_X + \omega_k} + \mathcal{O} \left[\left(\frac{\omega_V}{\omega_X} \right)^2, \frac{g\sqrt{N}}{\omega_X} \right] \right\}$$

- Reduced vibrational offset

$$\lambda_0 \rightarrow \lambda_0 (1 - K_V), \quad K_V = \sum_k \frac{g_k^2}{(\omega_k + \omega_X)^2}$$

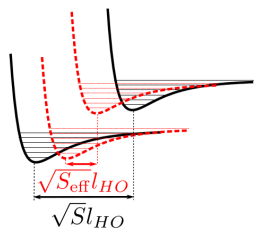
Vibrational reconfiguration

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- Reduced vibrational offset

$$\lambda_0 \rightarrow \lambda_0(1 - K_1), \quad K_1 = \sum_k \frac{g_k^2}{(\omega_k + \omega_X)^2}$$



- Increased effective coupling: $g_{\text{eff}}^2 = g^2 \exp(-S)$
- Again, $K_1 \ll 1$, independent of density.

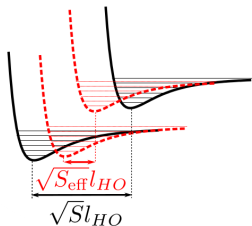
Vibrational reconfiguration

- Schrieffer-Wolff – mixes vibrational states

$$\delta H = - \sum_{i,k} \frac{g_k^2}{2(\omega_X + \omega_k)} \left\{ 1 - \frac{\omega_V \lambda_0 (b_i + b_i^\dagger)}{\omega_X + \omega_k} + \mathcal{O} \left[\left(\frac{\omega_V}{\omega_X} \right)^2, \frac{g_V \sqrt{N}}{\omega_X} \right] \right\}$$

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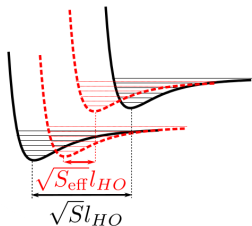
Vibrational reconfiguration

- Schrieffer-Wolff – mixes vibrational states

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- Reduced vibrational offset

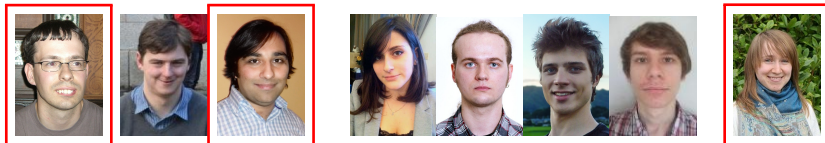
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- ▶ Increased effective coupling: $g_{\text{eff}}^2 = g^2 \exp(-S)$
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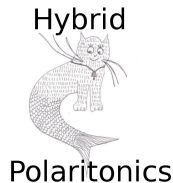
Acknowledgements

GROUP:



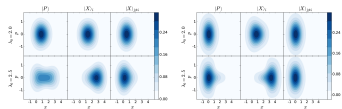
COLLABORATORS: S. De Liberato (Southampton)

FUNDING:

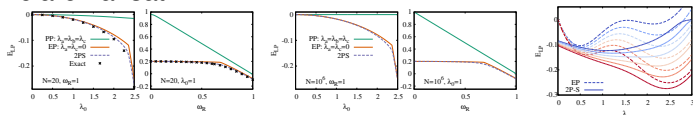


Summary

- Holstein-Dicke and Holstein-Tavis-Cummings models
- Single polariton state
 - ▶ Exact solution

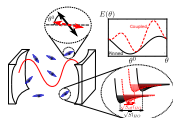


- ▶ Polaron ansatz



[Zeb, Kirton, JK, arXiv:1608.08929]

- Ground state configuration



[Cwik, Kirton, De Liberato, JK, PRA 2016]