

Spatial dynamics, thermalization and breakdown of thermalization in photon condensates

Jonathan Keeling



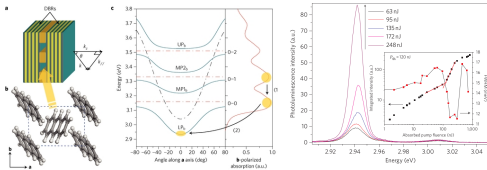
University of
St Andrews
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1413



Condensates of Light, January 2016

Motivation: polariton condensates

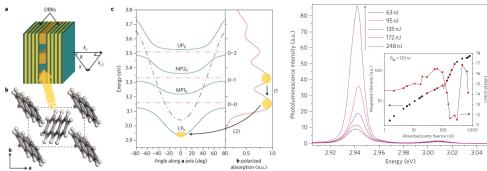
- Anthracene Polariton Lasing
 $T \sim 300\text{K}$



[Kena Cohen and Forrest, Nat. Photon '10]

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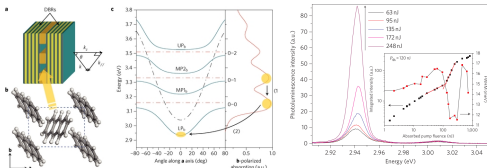


- Q1. Vibrational replicas?
- Q2. Relevance of disorder?
- Q3. Lasing vs condensation?

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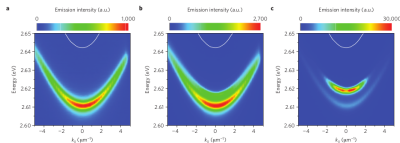
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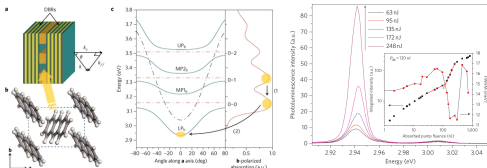
- Polariton condensates, other materials, e.g. polymers:



[Plumhoff *et al.* Nat. Materials '14, Daskalakis *et al.* *ibid* '14]

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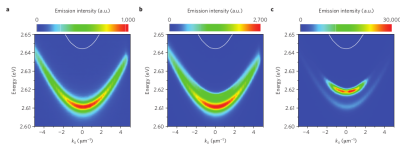
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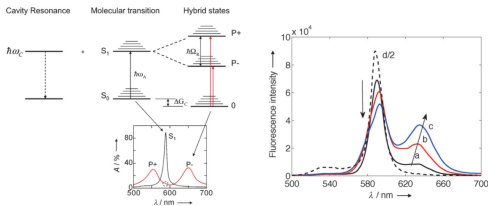


- Q1. Frenkel to Wannier crossover?
- Q2. Optimal vibrational properties?
- Q3. Nonlinearities?

[Plumhoff *et al.* Nat. Materials '14, Daskalakis *et al.* *ibid* '14]

Motivation: vacuum-state strong coupling

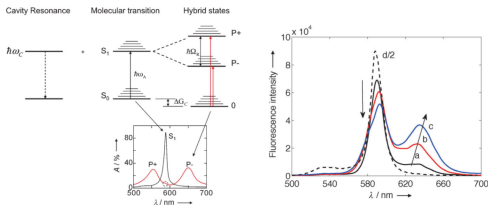
- Linear response (no pump, no condensate): effects of matter-light coupling alone.



[Canaguier-Durand *et al.* Angew. Chem. '13;
Baumberg group]

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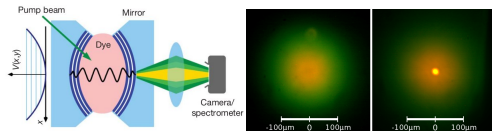


[Canaguier-Durand *et al.* Angew. Chem. '13;
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- Q1. Can **ultra-strong** coupling to light change:
- ▶ charge distribution?
 - ▶ vibrational configuration?
 - ▶ molecular orientation?
 - ▶ crystal structure?
- Q2. Are changes collective (\sqrt{N} factor) or not?

Motivation: photon condensates

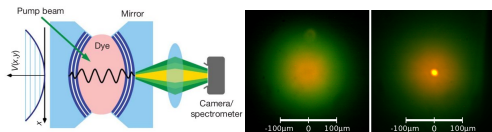
- Photon Condensate $T \sim 300\text{K}$



[Klaers *et al.* Nature, '10, Marelic *et al.* '15]

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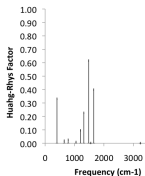
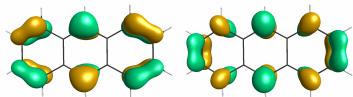


- Q1. Relation to dye laser?
- Q2. Relation to polaritons?
- Q3. Thermalisation breakdown?

[Klaers *et al.* Nature, '10, Marelic *et al.* '15]

Toy models

1 Full molecular spectra electronic structure & Raman spectrum



2 Simplified archetypal model: Dicke-Holstein

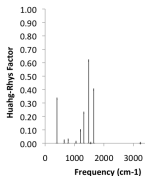
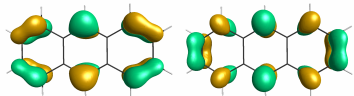
• Each molecule: two DoF

→ Electronic state: 2LS

→ Vibrational state: harmonic oscillator

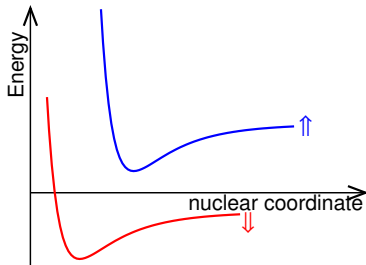
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- 1 Full molecular spectra electronic structure & Raman spectrum



- 2 Focus on low-energy effective theory

- Two-level system, HOMO/LUMO
- Single DoF PES



See also [Galego, Garcia-Vidal, Feist. PRX '15]

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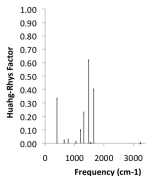
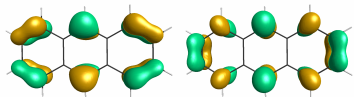
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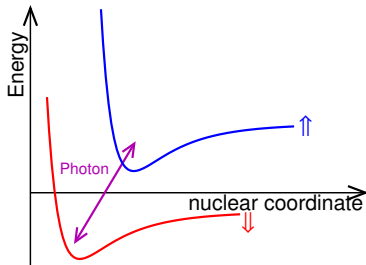
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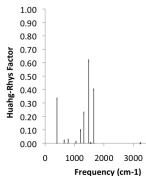
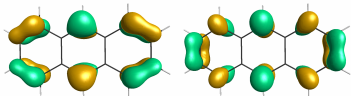
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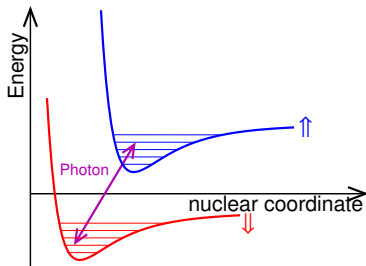
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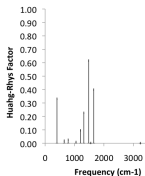
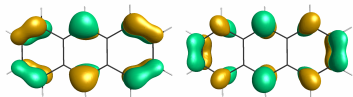
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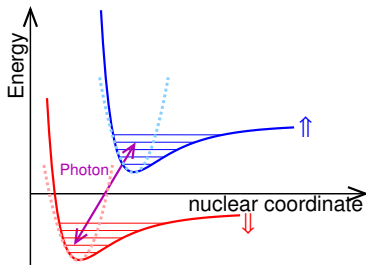


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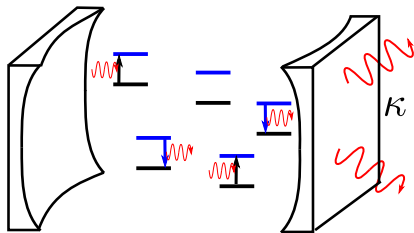


See also [Galego, Garcia-Vidal, Feist. PRX '15]

Dicke Holstein Model

- Dicke model: 2LS \leftrightarrow photons

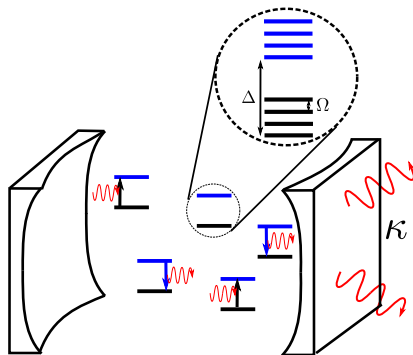
- Molecular vibrational mode
 - Phonon frequency Ω
 - Huang-Rhys parameter S — coupling strength



$$H_{\text{sys}} = \omega \psi^\dagger \psi + \sum_{\alpha} \left[\frac{\epsilon}{2} \sigma_{\alpha}^z + g (\psi + \psi^\dagger) (\sigma_{\alpha}^+ + \sigma_{\alpha}^-) \right]$$

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Modelling photon BEC

- 1 Modelling photon BEC
 - Uniform pumping results
- 2 Modelling steady-state spatial profile
 - Spatial profile vs spot size
 - Threshold vs spot size
- 3 Modelling spatial oscillations
 - Toy problem; validating model
 - Oscillation results

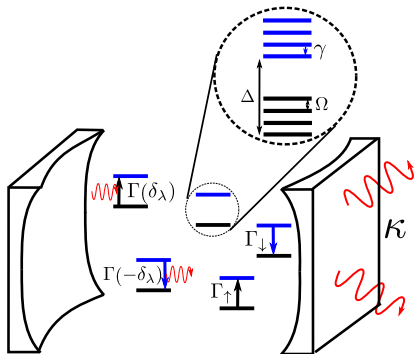
Photon: Microscopic Model

$$H_{\text{sys}} = \sum_m \omega_m \psi_m^\dagger \psi_m + \sum_\alpha \left[\frac{\epsilon}{2} \sigma_\alpha^z + g (\psi_m \sigma_\alpha^+ + \text{H.c.}) \right] \\ + \sum_\alpha \Omega \left\{ b_\alpha^\dagger b_\alpha + \sqrt{S} \sigma_\alpha^z (b_\alpha^\dagger + b_\alpha) \right\}$$

- **2D** harmonic oscillator

$$\omega_m = \omega_{\text{cutoff}} + m\omega_{H.O.}$$

- Incoherent processes: excitation, decay, loss, vibrational thermalisation.
- Weak coupling, perturbative in g

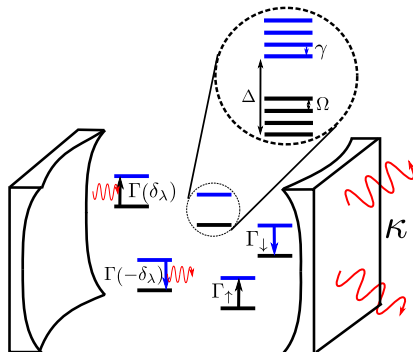


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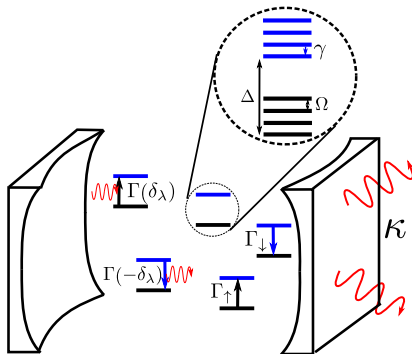
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Microscopic model – all orders in S

- Polaron transform (exact), $H = \sum_m \omega_m \psi_m^\dagger \psi_m + \sum_\alpha h_\alpha$,

$$h_\alpha = \frac{\epsilon}{2} \sigma_\alpha^z + g (\psi_m \sigma_\alpha^+ D_\alpha + \text{H.c.}) + \Omega b_\alpha^\dagger b_\alpha, \quad D_\alpha = e^{2\sqrt{S}(b_\alpha^\dagger - b_\alpha)}$$

- Master equation

$$\dot{\rho} = -i[H_0, \rho] + \sum_m \frac{\kappa}{2} \mathcal{L}[\psi_m] + \sum_\alpha \left[\frac{\Gamma_1}{2} \mathcal{L}[\sigma_\alpha^+] + \frac{\Gamma_1}{2} \mathcal{L}[\sigma_\alpha^-] \right] + \sum_{m,\alpha} \left[\frac{\Gamma(\delta_m - \omega_m - \epsilon)}{2} \mathcal{L}[\sigma_\alpha^+ \psi_m] + \frac{\Gamma(-\delta_m - \epsilon - \omega_m)}{2} \mathcal{L}[\sigma_\alpha^- \psi_m^\dagger] \right]$$

- Correlation function:

$$f(t) = 2g^2 \Re \left[\int dt e^{-|t| - (\Gamma_+ + \Gamma_-)t/2} \langle \sigma_\alpha^+(t) \sigma_\alpha^-(0) \rangle \right]$$

[Marthaler et al PRL '11; Kirton & JK PRL '13]

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$$G(\delta) = 2\sigma^2 \pi \int dt e^{-i\delta t - (\Gamma_\uparrow + \Gamma_\downarrow)|t|^2/2} \langle \sigma_\alpha^-(t) \sigma_\alpha^-(0) \rangle$$

[Marthaler et al PRL '11, Kirtou & JK PRL '13]

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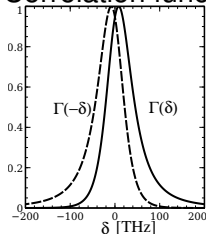
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[Marthaler et al PRL '11, Kirton & JK PRL '13]

Steady state populations and equilibrium

- Rate equation:

$$\partial_t n_m = -\kappa n_m + \Gamma(-\delta_m)(n_m + 1)N_{\uparrow} - \Gamma(\delta_m)n_m N_{\downarrow}$$

- Steady state distribution:

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Steady state populations vs loss

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- Bose-Einstein distribution without losses

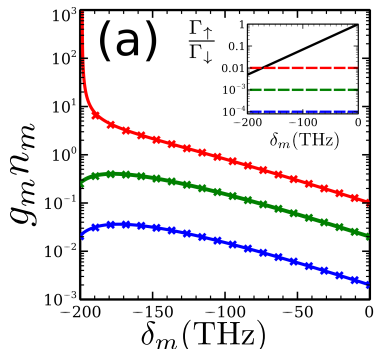
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Low loss: Thermal

[Kirton & JK PRL '13]

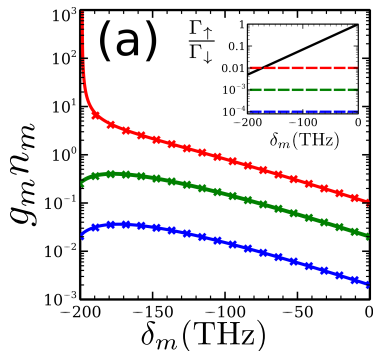
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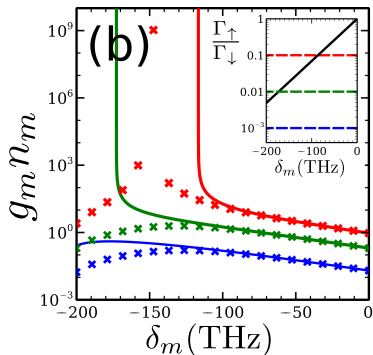
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Low loss: Thermal



High loss \rightarrow Laser

[Kirton & JK PRL '13]

Chemical potential?

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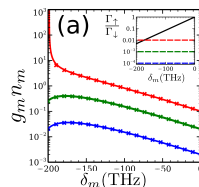
- $\kappa \ll N\Gamma(\delta)$, Kennard-Stepanov

$$\frac{n_m}{n_m + 1} = e^{-\beta\delta_m + \beta\mu}, \quad e^{\beta\mu} \equiv \frac{N_\uparrow}{N_\downarrow} = \frac{\Gamma_\uparrow + \sum_m \Gamma(\delta_m) n_m}{\Gamma_\downarrow + \sum_m \Gamma(-\delta_m) (n_m + 1)}$$

- Below threshold,

$$\mu = k_B T \ln[\Gamma_\uparrow / \Gamma_\downarrow]$$

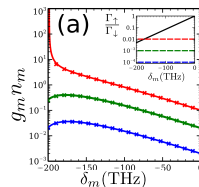
- At/above threshold, $\mu \rightarrow \delta_0$



Chemical potential?

- Steady state distribution:

$$\frac{n_m}{n_m + 1} = \frac{\Gamma(-\delta_m) N_{\uparrow}}{\kappa + \Gamma(\delta_m) N_{\downarrow}}$$



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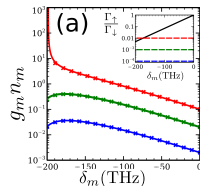
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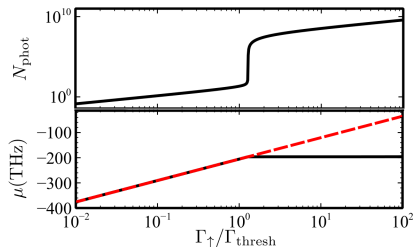
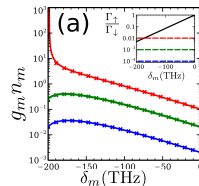
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[Kirton & JK, PRA '15]

Modelling steady-state spatial profile

- 1 Modelling photon BEC
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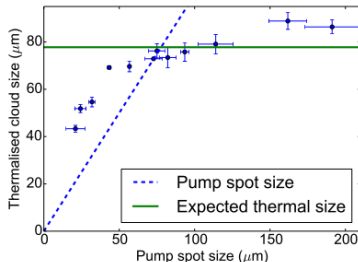
Spatially varying pump intensity

- Consider effects of pump profile, $\Gamma_{\uparrow}(\mathbf{r}) = \frac{\Gamma_{\uparrow} \exp(-r^2/2\sigma_p^2)}{(2\pi\sigma_p^2)^{d/2}}$

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Modelling spatial profile.

- Varying excited density – differential coupling to modes

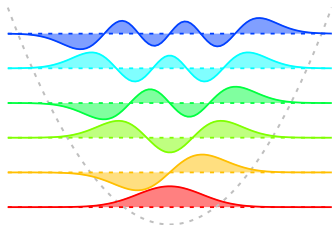
$$\partial_t n_m = -\kappa n_m + \Gamma(-\delta_m) O_m (n_m + 1) - \Gamma(\delta_m) (\rho_m - O_m) n_m$$

$$O_m = \int d\mathbf{r} \rho_{\uparrow}(\mathbf{r}) |\psi_{2m}(\mathbf{r})|^2, \quad \rho_{\uparrow} + \rho_{\downarrow} = \rho_m$$

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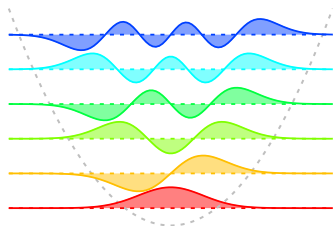
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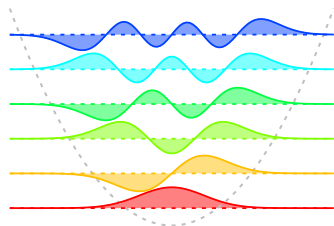
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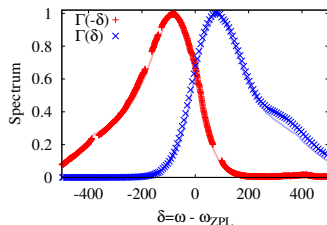
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- Use exact R6G spectrum



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Spatially varying pump: below threshold

- Far below threshold:

- ▶ If $\kappa \ll \rho_m \Gamma(\delta_m)$,
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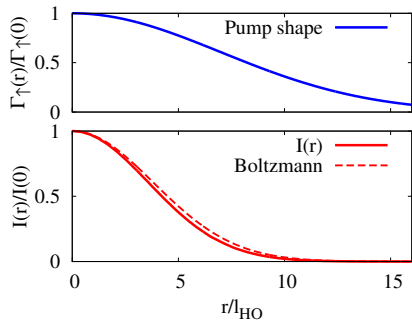
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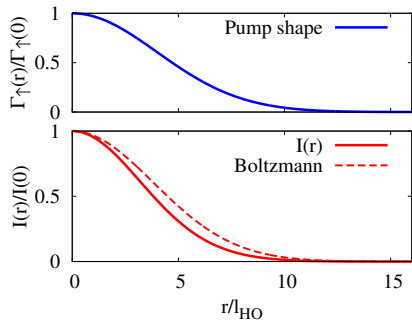
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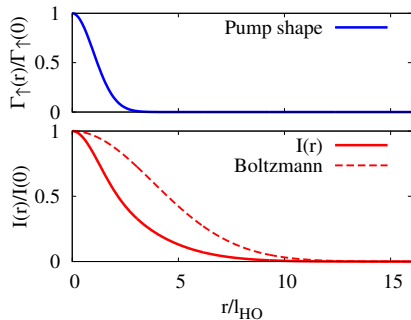


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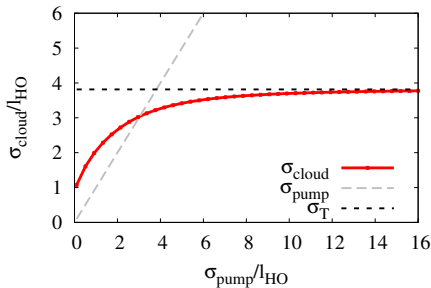
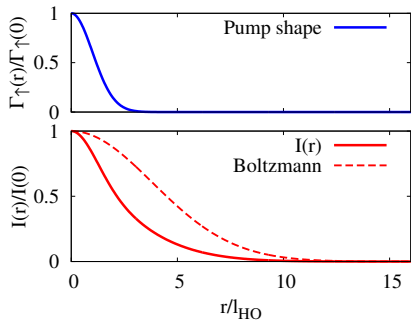


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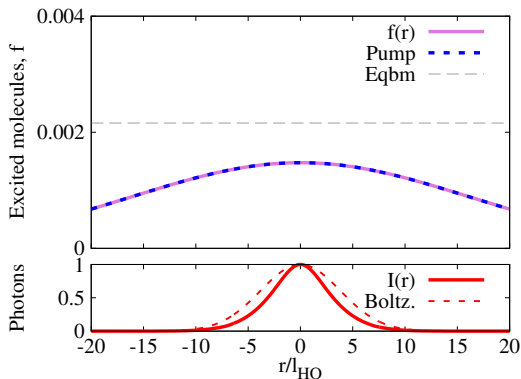
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Near threshold behaviour

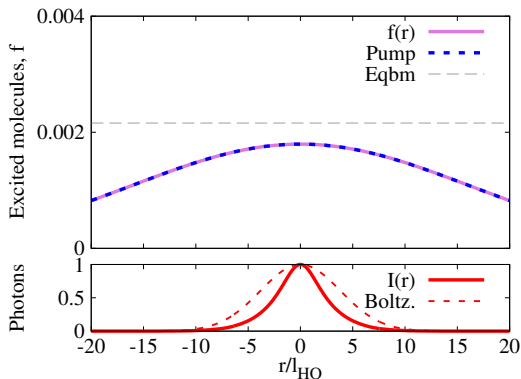


- Large spot, $\sigma_p \gg l_{HO}$

- "Gain saturation" at centre

- Saturation of $f(r) = 1/(1 + e^{-\beta r})$ — spatial equilibration

Near threshold behaviour

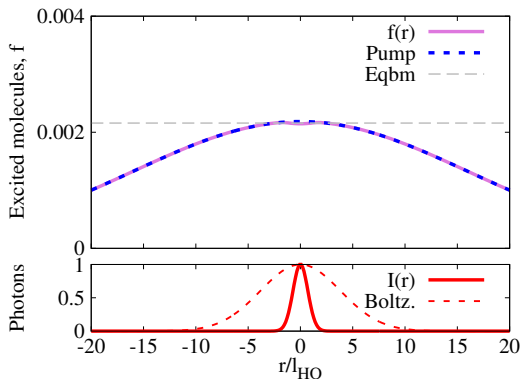


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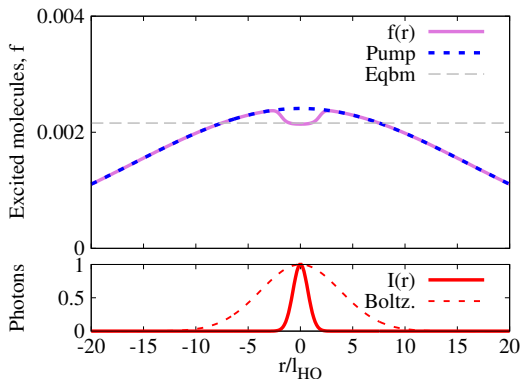
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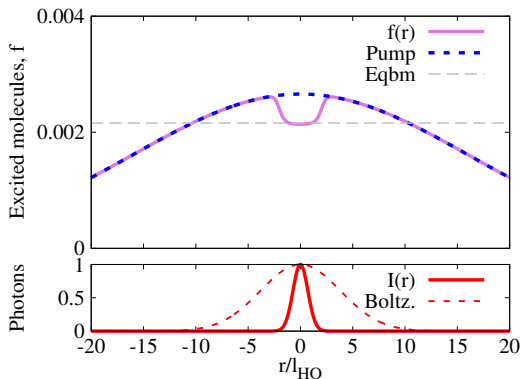
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- Lasing threshold, dependence on spot size.

- ▶ Equilibrium: $\mu = \delta_c$

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- ↳ Dependence on ω_c — experimental spectrum

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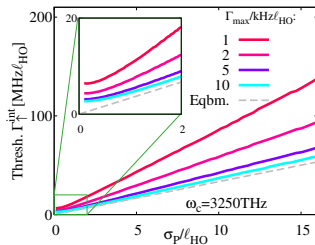
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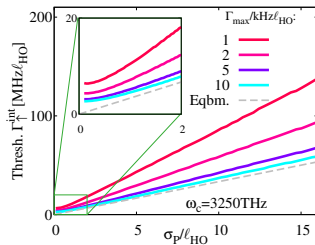


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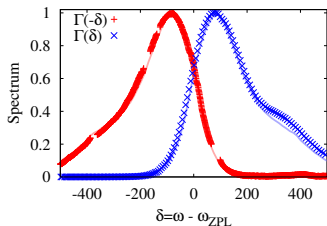
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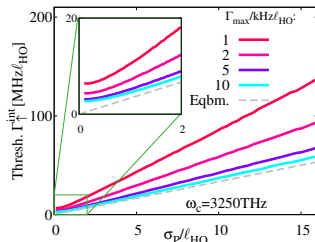
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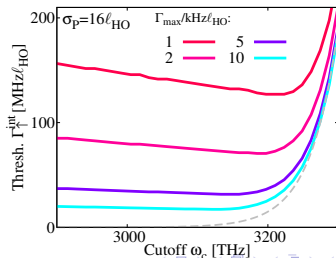
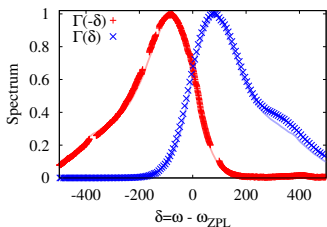
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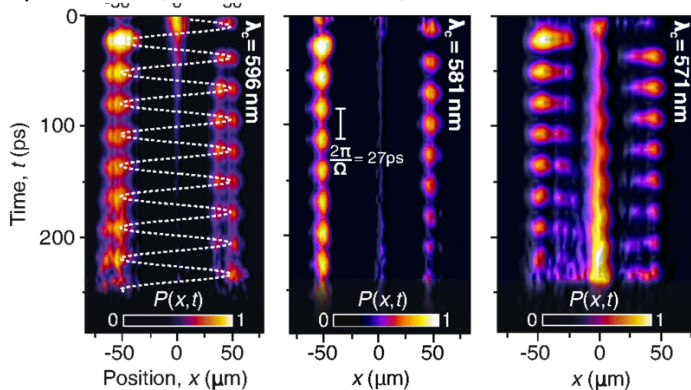


Threshold vs spot size

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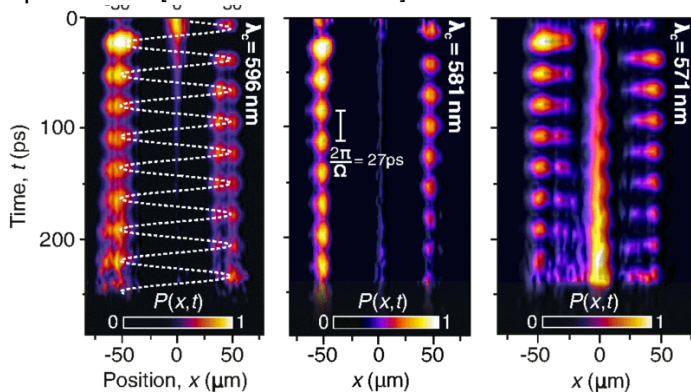
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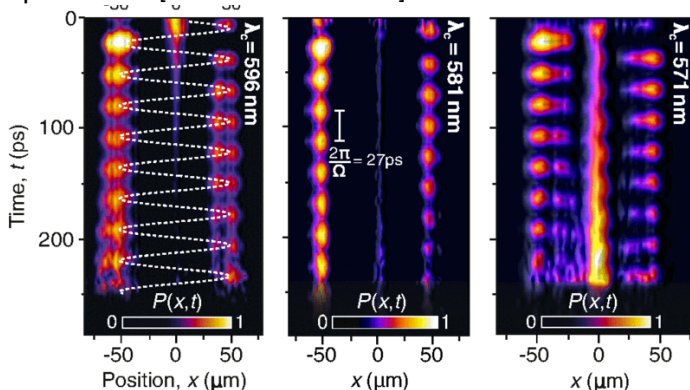


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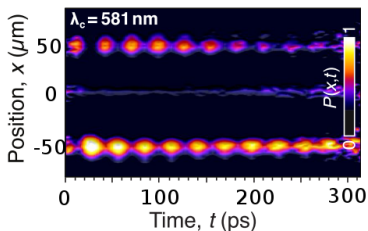
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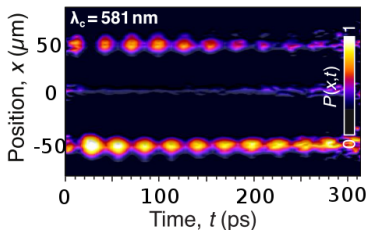
Limit of rate equations



$$\partial_t n_m = -\kappa n_m + \Gamma(-\delta_m)(n_m + 1)N_{\uparrow} - \Gamma(\delta_m)n_m N_{\downarrow}$$

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- Need $I(x) = \sum_{m,m'} n_{m,m'} \psi_m(x) \psi_{m'}(x)$
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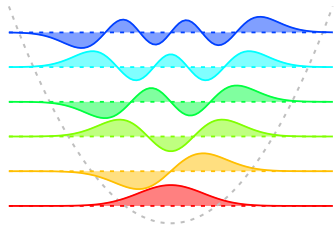


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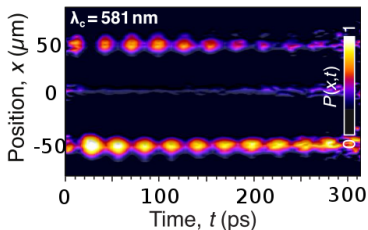
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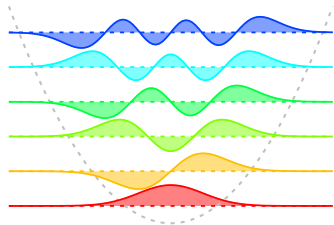
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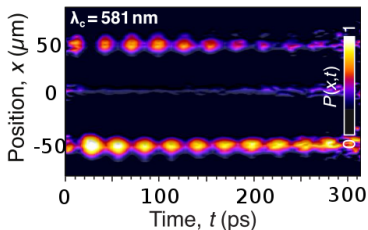
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Limit of rate equations



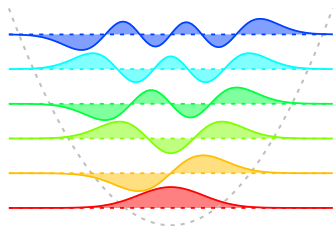
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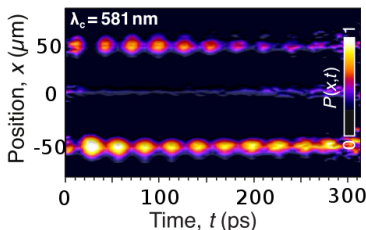
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Emission must create coherence between non-degenerate modes.

Limit of rate equations

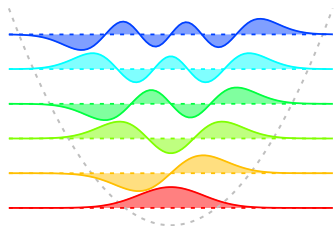


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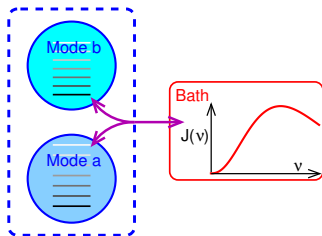
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Emission must create coherence between non-degenerate modes.

Toy problem: two bosonic modes

- Basic problem: Emission from thermal bath



$$H = \omega_a \hat{\psi}_a^\dagger \hat{\psi}_a + \omega_b \hat{\psi}_b^\dagger \hat{\psi}_b + H_{\text{Bath}} \\ + (\varphi_a^* \hat{\psi}_a^\dagger + \varphi_b^* \hat{\psi}_b^\dagger) \sum_i g_i \hat{c}_i + \text{H.c.}$$

Toy problem: naïve solutions

- Two “expected” behaviours:
 - ▶ At resonance: “weak lasing” — coupling to bath dominates

$$\frac{d}{dt}\rho = \Gamma^\downarrow \mathcal{L}[\varphi_a \hat{\psi}_a + \varphi_b \hat{\psi}_b] + \Gamma^\uparrow \mathcal{L}[\varphi_a^* \hat{\psi}_a^\dagger + \varphi_b^* \hat{\psi}_b^\dagger]$$

• Far from resonance: pointer states are eigenstates

$$\frac{d}{dt}\rho = \sum_{i=a,b} \Gamma^\downarrow \mathcal{L}[\hat{\psi}_i] + \Gamma^\uparrow \mathcal{L}[\hat{\psi}_i^\dagger]$$

- Explicit derivation → Redfield theory

$$\begin{aligned} \partial_t \rho = & -i[H, \rho] + \sum_j \mathcal{L}_j^\downarrow \left(2\hat{\psi}_j \rho \hat{\psi}_j^\dagger - [\rho, \hat{\psi}_j^\dagger \hat{\psi}_j]_+ \right) \\ & + \sum_j \mathcal{L}_j^\uparrow \left(2\hat{\psi}_j^\dagger \rho \hat{\psi}_j - [\rho, \hat{\psi}_j \hat{\psi}_j^\dagger]_+ \right). \end{aligned}$$

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- Two “expected” behaviours:
 - ▶ At resonance: “weak lasing” — coupling to bath dominates

$$\frac{d}{dt}\rho = \Gamma^\downarrow \mathcal{L}[\varphi_a \hat{\psi}_a + \varphi_b \hat{\psi}_b] + \Gamma^\uparrow \mathcal{L}[\varphi_a^* \hat{\psi}_a^\dagger + \varphi_b^* \hat{\psi}_b^\dagger]$$

- ▶ Far from resonance: pointer states are eigenstates

$$\frac{d}{dt}\rho = \sum_{i=a,b} \Gamma_i^\downarrow \mathcal{L}[\hat{\psi}_i] + \Gamma_i^\uparrow \mathcal{L}[\hat{\psi}_i^\dagger]$$

- Explicit derivation → Redfield theory

$$\begin{aligned} \partial_t \rho = & -i[\hat{H}, \rho] + \sum_{ij} L_{ij}^\downarrow \left(2\hat{\psi}_j \rho \hat{\psi}_i^\dagger - [\rho, \hat{\psi}_i^\dagger \hat{\psi}_j]_+ \right) \\ & + \sum_{ij} L_{ij}^\uparrow \left(2\hat{\psi}_j^\dagger \rho \hat{\psi}_i - [\rho, \hat{\psi}_i \hat{\psi}_j^\dagger]_+ \right). \end{aligned}$$

Toy problem: exact solution

- Solve via Laplace transform. Find $F_{ij}(t) = \langle \hat{\psi}_i^\dagger(t) \hat{\psi}_j(t) \rangle$

- Steady state:

- Time evolution —

$$F_{ab}(t) \sim \exp(-\alpha \Delta^2 t)$$

- Always some coherence

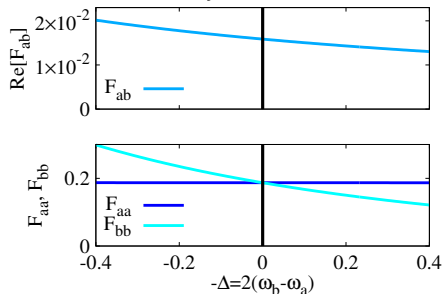
- (individual always wrong)

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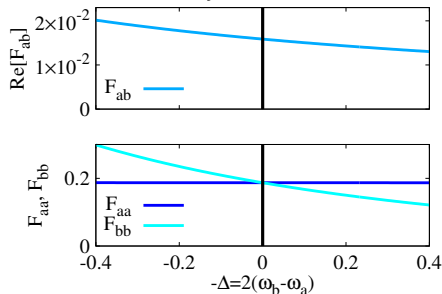


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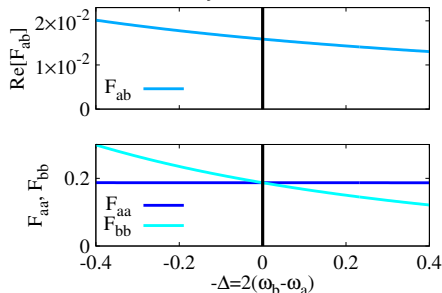
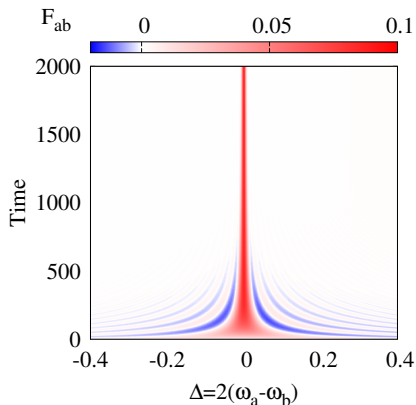
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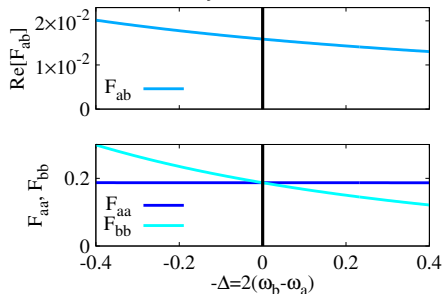
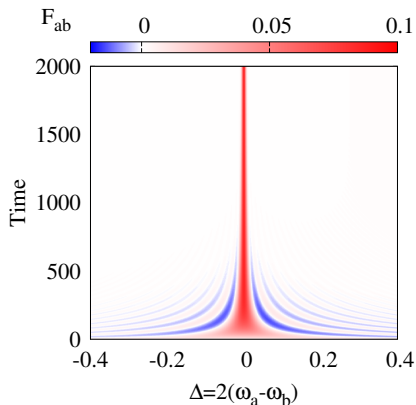
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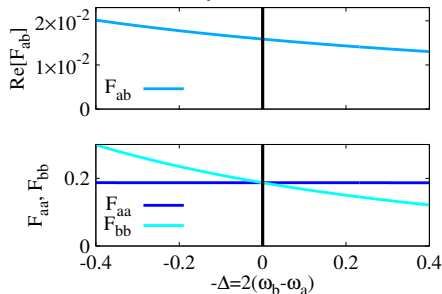
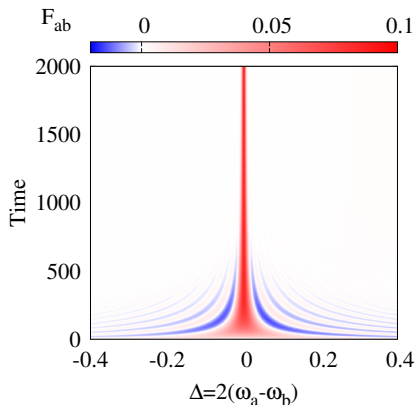


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Toy problem: Redfield theory

Unsecularised Redfield theory:

$$\partial_t \rho = -i[\hat{H}, \rho] + \sum_{ij} \varphi_i^* \varphi_j \left[\mathcal{K}_{ij}^\downarrow \left(2\hat{\psi}_j \rho \hat{\psi}_i^\dagger - [\rho, \hat{\psi}_i^\dagger \hat{\psi}_j]_+ \right) + \mathcal{K}_{ij}^\uparrow \left(2\hat{\psi}_j^\dagger \rho \hat{\psi}_i - [\rho, \hat{\psi}_i \hat{\psi}_j^\dagger]_+ \right) \right].$$

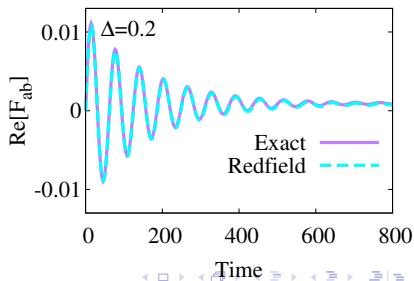
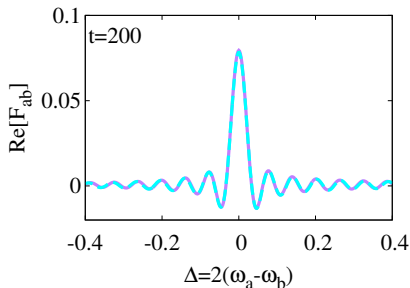
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Toy problem: Secularisation

- Secularisation (in eigenbasis of \hat{H}): $L_{ij}^{\uparrow,\downarrow} \rightarrow L_{ii}^{\uparrow,\downarrow} \delta_{ij} \rightarrow F_{ab} = 0$

- Secularisation often invoked to cure negative eigenvalues of $L_{ij}^{\uparrow,\downarrow}$.

- Check stability: consider $f = (F_{aa}, F_{bb}, \Re[F_{ab}], \Im[F_{ab}])$

$$\partial_t f = -Mf + f_0$$

- Eigenvalues of M exist in closed form:

- Unstable (negative only if $dJ(\nu)/d\nu \gg 1$
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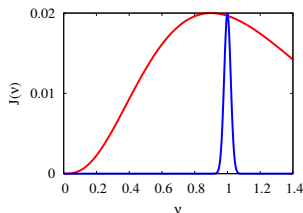
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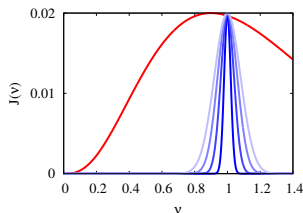


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Beyond Redfield: Schrödinger picture Bloch Redfield

- Is BR the best (time-local) theory we can find?

- Hints it is not:

- Eigenvalues of M vs exact sol'n near $\Delta = 0$.

- Sum rule [Salmilehto *et al.* PRA '12; Hell *et al.* PRB '14]:

- *"For \hat{X} s.t. $[\hat{X}, H_{\text{system-bath}}] = 0$, then $\partial_t \langle \hat{X} \rangle$ should match closed system."*

- Alternate approach:

- BR assumes $\hat{\rho}(t)$ is "slow" in interaction picture

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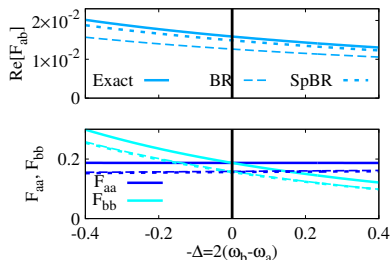
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Modelling spatial oscillations

- 1 Modelling photon BEC
 - Uniform pumping results
- 2 Modelling steady-state spatial profile
 - Spatial profile vs spot size
 - Threshold vs spot size
- 3 Modelling spatial oscillations
 - Toy problem; validating model
 - Oscillation results

Modelling

- Following toy model, use Redfield theory:

$$\partial_t \rho = -i \left[\sum_m \omega_m \hat{a}_m^\dagger \hat{a}_m, \rho \right] + \sum_{m,m',i} \psi_m^*(r_i) \psi_{m'}(r_i) \left(K(\delta_{m'}) [\hat{a}_{m'} \hat{\sigma}_i^+ \hat{\rho}, \hat{a}_m^\dagger \hat{\sigma}_i^-] \right. \\ \left. + K(-\delta_m) [\hat{a}_m^\dagger \hat{\sigma}_i^- \hat{\rho}, \hat{a}_{m'} \hat{\sigma}_i^+] \right) + \text{H.c.} + (\text{pumping, decay} \dots),$$

- $K(\delta)$ analytic continuation of $\Gamma(\delta)$.

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• Semiclassical equations for $n_{m,m'} = \langle \hat{a}_m^\dagger \hat{a}_{m'} \rangle$ and $f(r)$.

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Modelling

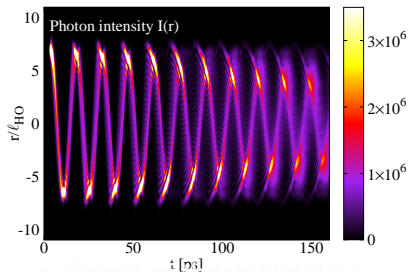
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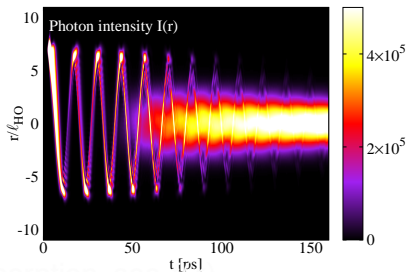
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Dynamics from model

Longer cavity



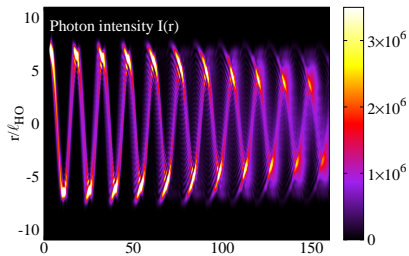
Shorter cavity



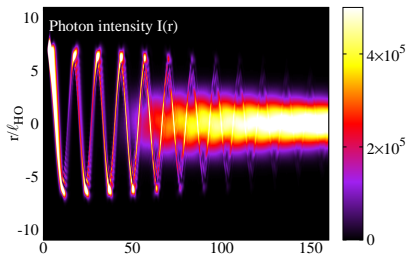
• Origin of thermalisation — reabsorption, see Fig. 10

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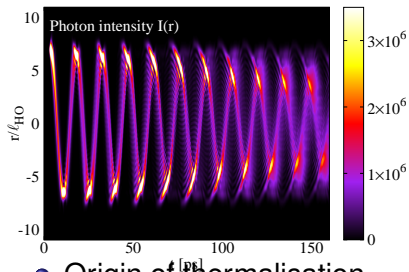
Shorter cavity



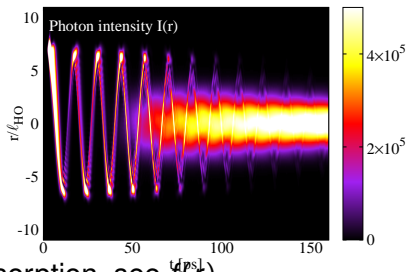
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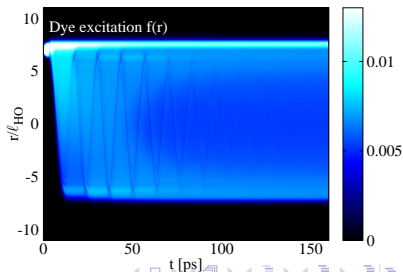
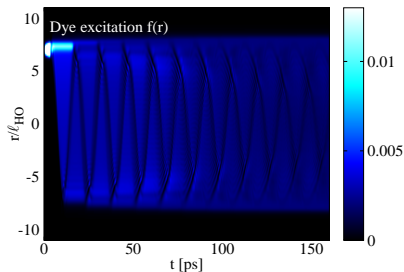
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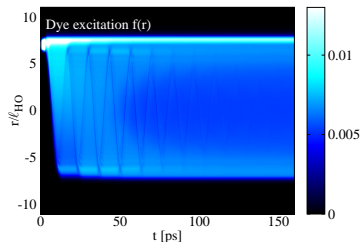
● Origin of thermalisation — reabsorption, see $f(r)$



Thermalisation at late times

- Reabsorption “fills-in” excited molecules

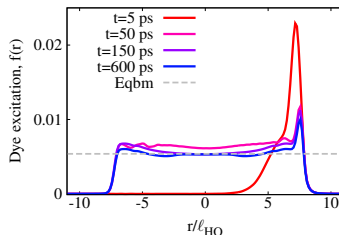
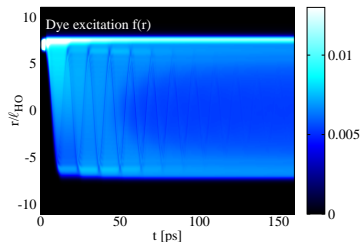
● Reach thermal equilibrium, $f = [e^{-\beta h\nu} + 1]^{-1}$



● Photon occupation thermalises later

Thermalisation at late times

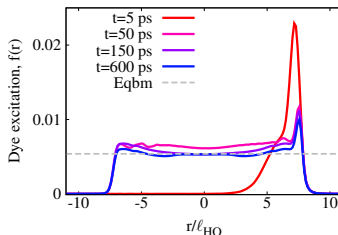
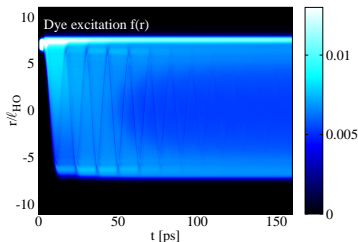
- Reabsorption “fills-in” excited molecules
- Reach thermal equilibrium, $f = [e^{-\beta\delta_0} + 1]^{-1}$



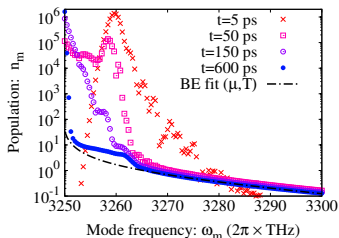
● Photon occupation thermalises later

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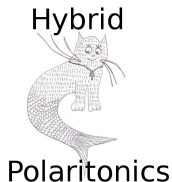


Acknowledgements

GROUP:



FUNDING:



ICSCE 8

Edinburgh, 25th–29th April, 2016.



Plenary speakers: Ataç İmamoğlu, Peter Zoller.

Invited speakers: Ehud Altman, Mete Atatüre, Natasha Berloff, Charles Bardin, Jacqueline Bloch, Iacopo Carusotto, Cristiano Ciuti, Michele Devoret[†], Thomas Ebbesen, Thierry Giamarchi, Jan Klärs, Dmitry Krizhanovskii, Xiaoqin (Elaine) Li, Peter Littlewood, Allan MacDonald, Francesca Marchetti, Keith Nelson, Pavlos Lagoudakis, Vivien Zapf.

([†] To be confirmed)

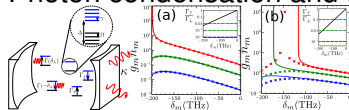
Early-bird registration & abstract deadline: 31st January 2016.

Final registration deadline: 31st March 2016.

<http://www.st-andrews.ac.uk/~icsce8>

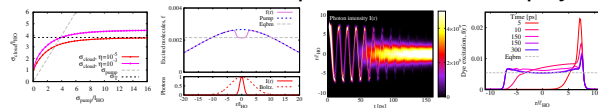
Summary

- Photon condensation and thermalisation



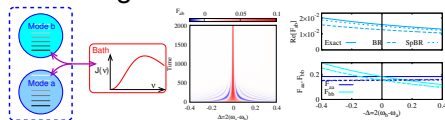
[Kirton & JK, PRL '13, PRA '15]

- Photon condensation, pattern formation physics



[JK & Kirton, PRA '16]

- Modelling incoherent emission into non-degenerate modes



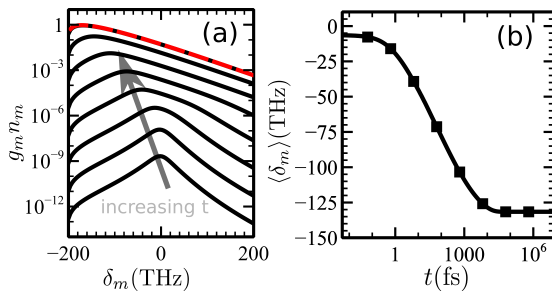
[Eastham, Kirton, Cammack, Lovett, JK arXiv:1508.04744]

Extra Slides

- 4 Approach to steady state
- 5 Threshold vs temperature
- 6 Beyond semiclassics
- 7 Toy problem
- 8 More oscillations
- 9 Polariton spectral weight

Time evolution

- Initial state: excited molecules
 - Initial emission, follows gain peak
 - Thermalisation by repeated absorption

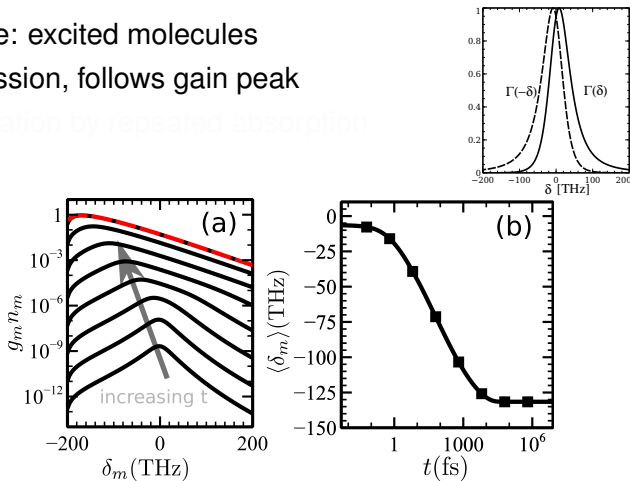


[Kirton & JK PRA '15]

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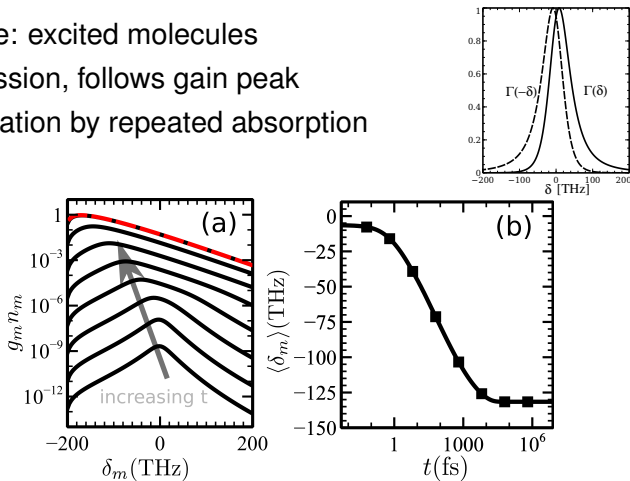
◦ Thermalisation by repeated absorption



[Kirton & JK PRA '15]

Time evolution

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[Kirton & JK PRA '15]

Threshold condition

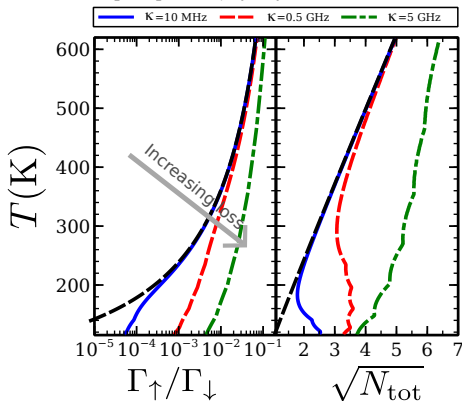
Use: $\max[n_m] = 1/(\beta\epsilon) \rightarrow k_B T_c = \sqrt{6/\pi^2\epsilon}\sqrt{N}$.

- Pump rate (Laser)
- Critical density (condensate)

- Thermal at low κ /high temperature
- High loss, κ competes with $\Gamma(\pm\delta_0)$
- Low temperature, $\Gamma(\pm\delta_0)$ shrinks

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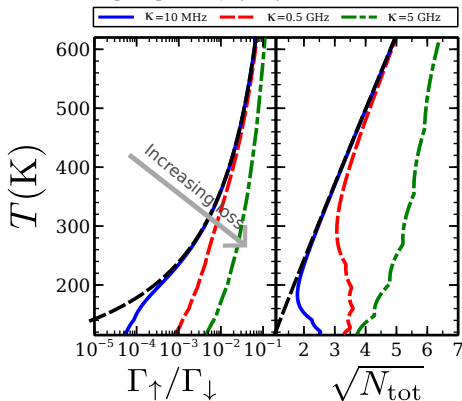
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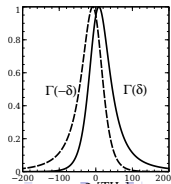


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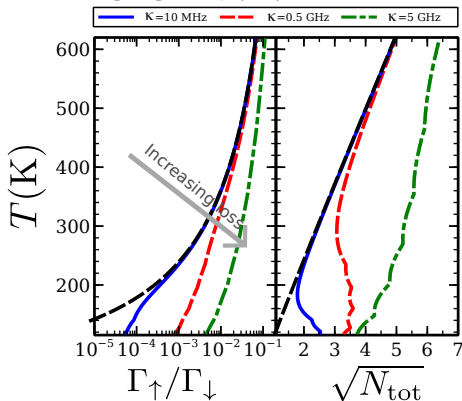
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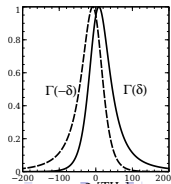
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Quantum model, linewidth

Full Master equation:

$$\dot{\rho} = -i[H_0, \rho] - \frac{\kappa}{2} \mathcal{L}[\psi] - \sum_{\alpha} \left[\frac{\Gamma_{\uparrow}}{2} \mathcal{L}[\sigma_{\alpha}^{+}] + \frac{\Gamma_{\downarrow}}{2} \mathcal{L}[\sigma_{\alpha}^{-}] \right] \\ - \sum_{\alpha} \left[\frac{\Gamma(\delta = \omega - \epsilon)}{2} \mathcal{L}[\sigma_{\alpha}^{+} \psi] + \frac{\Gamma(-\delta = \epsilon - \omega)}{2} \mathcal{L}[\sigma_{\alpha}^{-} \psi^{\dagger}] \right]$$

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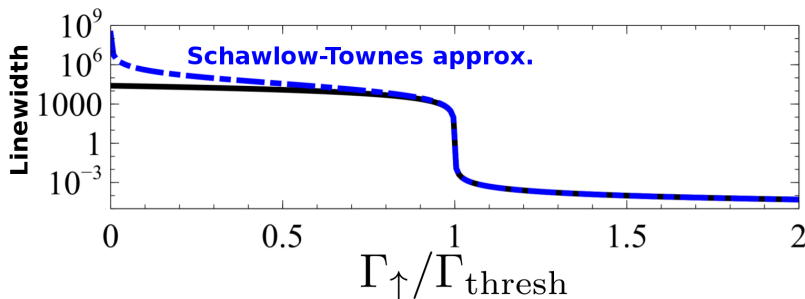
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Beyond Redfield: Schrödinger picture Bloch Redfield

- Is BR the best (time-local) theory we can find?

- Hints it is not:

- Eigenvalues of M vs exact sol'n near $\Delta = 0$.

- Sum rule [Salmilehto *et al.* PRA '12; Hell *et al.* PRB '14]:

- *"For \hat{X} s.t. $[\hat{X}, H_{\text{system-bath}}] = 0$, then $\partial_t \langle \hat{X} \rangle$ should match closed system."*

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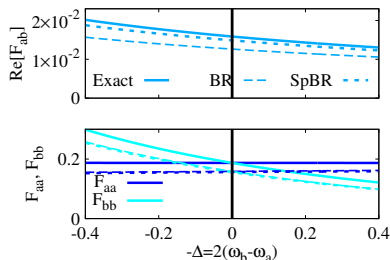
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Thermalisation of spectrum:

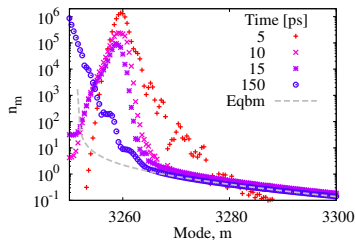
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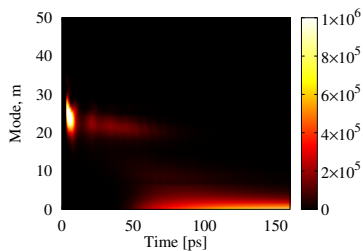
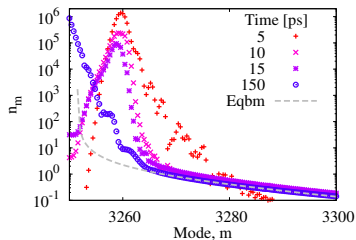
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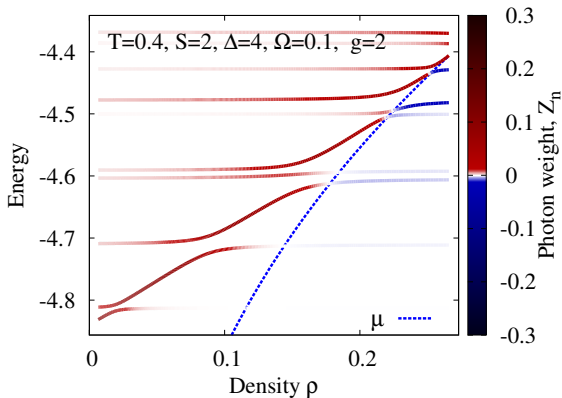


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Polariton spectrum: photon weight

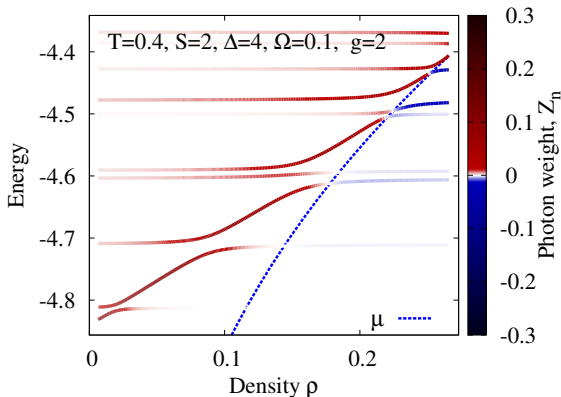


• What is nature of polariton mode?

$$G^R(t) = -i \langle \psi^\dagger(t) \psi(0) \rangle, \quad G^R(\omega) = \sum_n \frac{Z_n}{\omega - \omega_n}$$

[Cwik *et al.* EPL '14]

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