

# Non-equilibrium phases of matter-light systems

Jonathan Keeling



University of  
St Andrews  

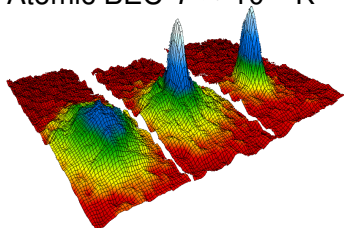
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1413-2013

Potsdam, October 2015

# Coherent states of matter and light

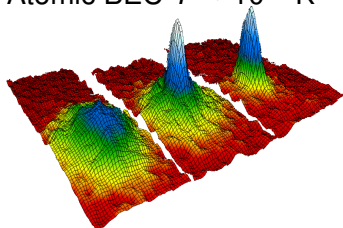
Atomic BEC  $T \sim 10^{-7}\text{K}$



[Anderson *et al.* Science '95]

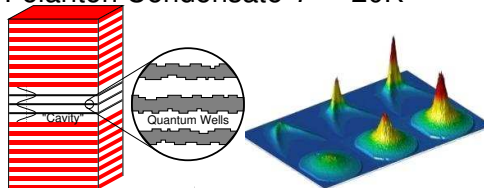
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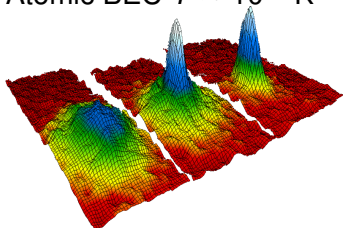
Polariton Condensate  $T \sim 20\text{K}$



[Kasprzak *et al.* Nature, '06]

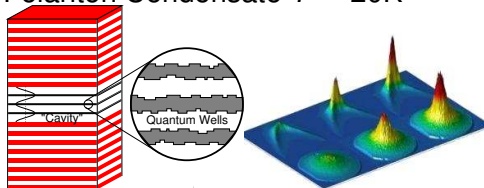
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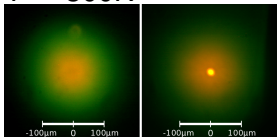
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## Photon Condensate

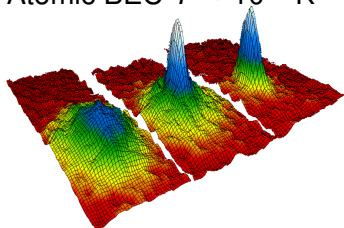
$T \sim 300\text{K}$



[Klaers *et al.* Nature, '10]

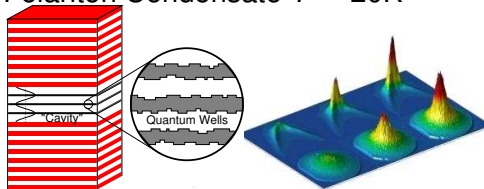
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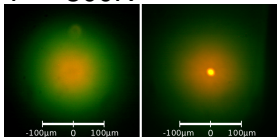
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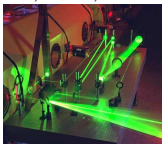
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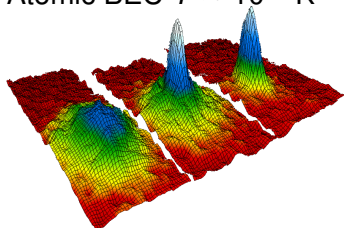
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Laser  
 $T \sim ?, < 0, \infty$



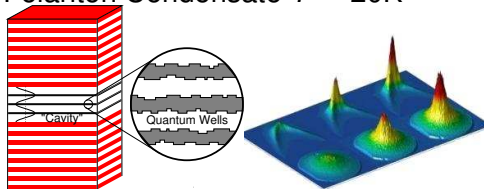
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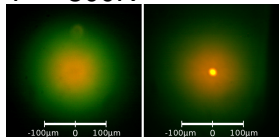
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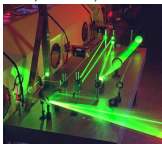
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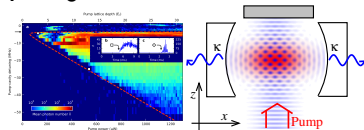


[Klaers *et al.* Nature, '10]

Laser  
 $T \sim ?, < 0, \infty$



Superradiance transition  
 $T \sim 0$



[Baumann *et al.* Nature, '10]

# Driven systems

Open quantum system

$$\partial_t \rho = -i[\hat{H}, \rho] + \sum_i \kappa_i \mathcal{L}[\hat{X}_i], \quad \mathcal{L}[\hat{X}_i] = 2\hat{X}_i \rho \hat{X}_i^\dagger - \hat{X}_i^\dagger \hat{X}_i \rho - \rho \hat{X}_i^\dagger \hat{X}_i$$

Need **drive** to balance **loss**

External coherent drive:

$$H \rightarrow H + V \cos(\Omega t)$$

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1 External **coherent** drive:

$$\hat{H} \rightarrow \hat{H} + \hat{V} \cos(\Omega t)$$

- $\tilde{\rho} = e^{-i\Omega N t} \rho e^{i\Omega N t} - \Omega \hat{N}$
- Neglect fast  $e^{2i\Omega t}$  terms — fast
- Rotating frame — breaks detailed balance with bath.



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Need **drive** to balance **loss**

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$$\tilde{\hat{H}} = \begin{pmatrix} h_0 & v_{01} \cos(\Omega t) & 0 & \dots \\ v_{01}^\dagger \cos(\Omega t) & h_1 & v_{12} \cos(\Omega t) & \dots \\ 0 & v_{12}^\dagger \cos(\Omega t) & h_2 & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

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- ▶  $\tilde{\hat{H}} = e^{-i\Omega \hat{N}t} \hat{H} e^{i\Omega \hat{N}t} - \Omega \hat{N}$
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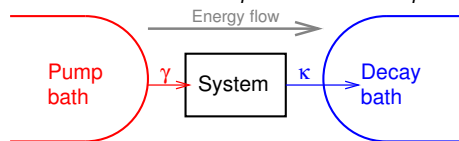
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# Non-equilibrium steady state

## 2 External **incoherent** drive:

$$\partial_t \rho = -i[\hat{H}, \rho] + \sum_i \kappa_i \mathcal{L}[\hat{X}_i] + \sum_i \gamma_i \mathcal{L}[\hat{X}_i^\dagger]$$



• Energy flow through system

• Not thermodynamics — attractors of dynamics

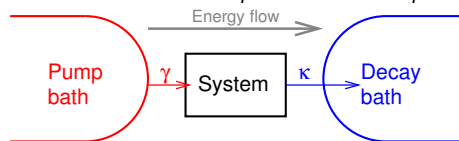
• Stationary points — extrema of energy?

• Nontrivial attractors

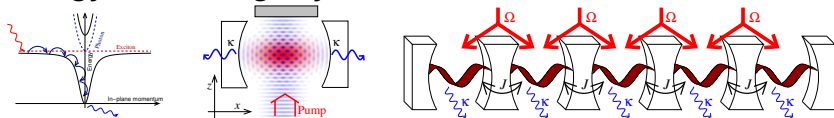
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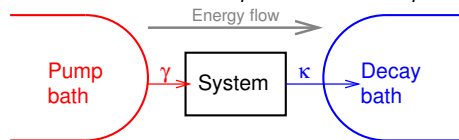
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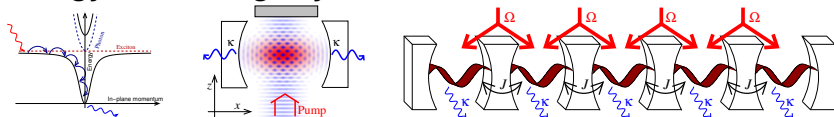
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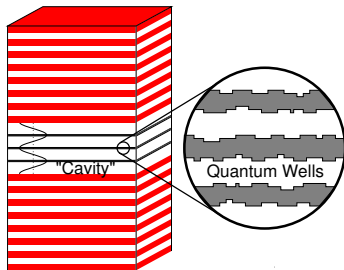
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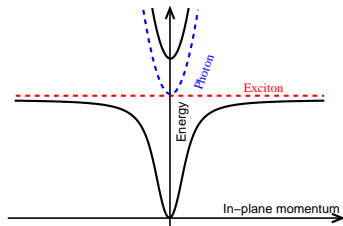
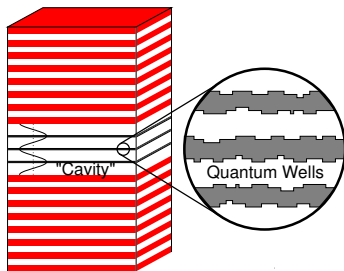
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# Microcavity polaritons

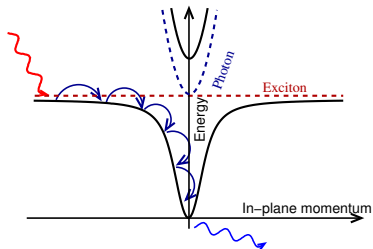
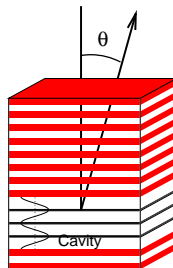
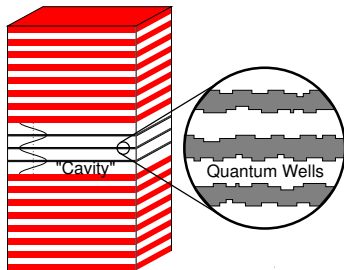


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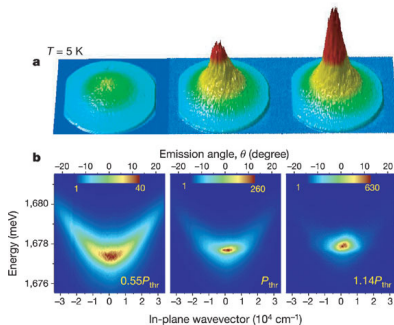
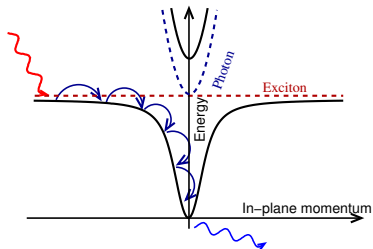
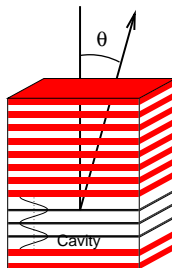
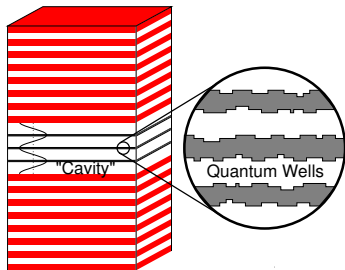




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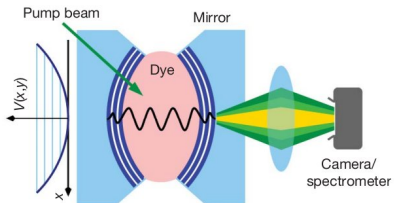


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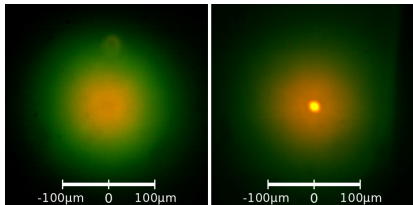
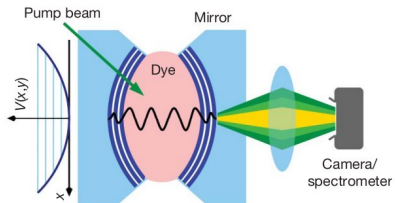
[Kasprzak, *et al.* Nature, '06]

# Photon Condensates



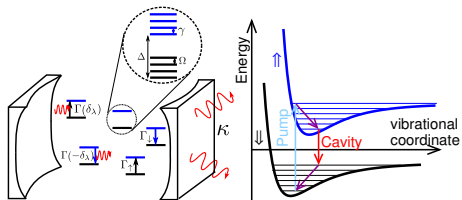
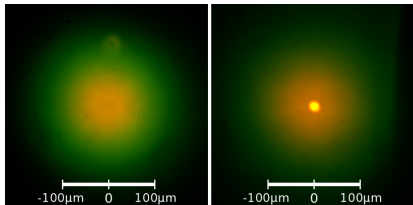
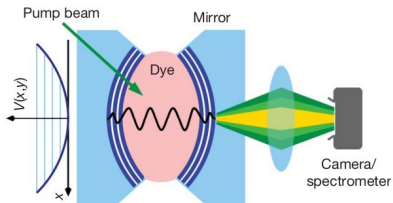
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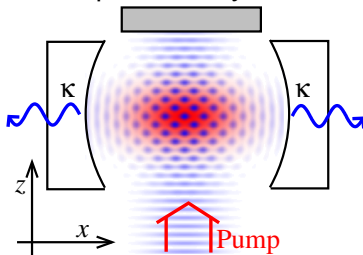
# Atoms and light

- **Cold atoms, optical lattice:** static light, dynamic matter
- **Quantum optics:** (Static) atoms, quantum dynamics of light.
- Coupling atomic motion to optical cavity

[Experiments: MIT, ETH, CQT Singapore, Stanford, Hamburg]

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# Overview

## 1 Non-equilibrium phases of matter-light systems

- Physical systems

## 2 Cold atoms in optical cavities – Dicke model

- Superradiance and the Dicke model
- Cold atoms & open Dicke model
- Dynamics and attractors
- Beyond Dicke: Chaotic dynamics

## 3 Beyond mean-field theory

- Coupled cavity arrays
- Multimode cavities



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# Coupling many atoms to light

**Old question:** *What happens to radiation when many atoms interact “collectively” with light.*

**Superradiance** — dynamical and steady state.

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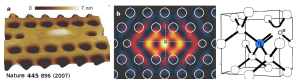
**Superradiance** — dynamical and steady state.

**New relevance**

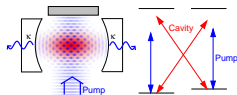
- Superconducting qubits



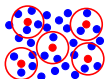
- Quantum dots & NV centres



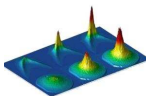
- Ultra-cold atoms



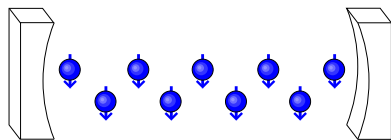
- Rydberg atoms/polaritons



- Microcavity Polaritons



# Dicke model and Dicke-Hepp-Lieb transition

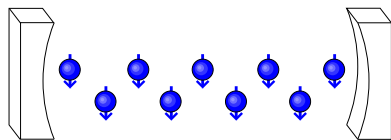


$$H = \omega \hat{a}^\dagger \hat{a} + \sum_i \omega_0 S_i^z + g(\hat{a} + \hat{a}^\dagger)(S_i^+ + S_i^-)$$

- Coherent state:  $|\hat{a}\rangle \rightarrow e^{i\alpha \hat{a}^\dagger + \eta S_i^-} |\Omega\rangle$
- Small  $g$ , min at  $\alpha, \eta = 0$

[Hepp, Lieb, Ann. Phys. '73]

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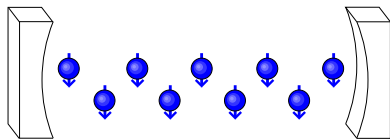


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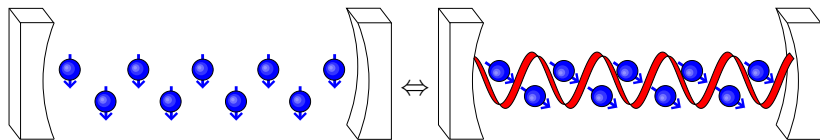
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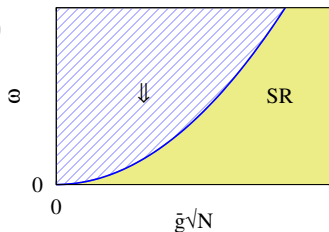
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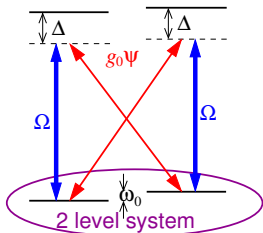
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Non-zero cavity field if:  $4Ng^2 > \omega\omega_0$



[Hepp, Lieb, Ann. Phys. '73]

# Raman scheme, decoupling $g, \omega_0$



$$H = \omega_0 S^z + g(\hat{a} + \hat{a}^\dagger)(S^- + S^+) + \omega \hat{a}^\dagger \hat{a}$$

- 2 Level system,  $|\downarrow\rangle, |\uparrow\rangle$

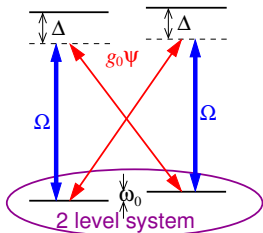
- Coupling  $g = \frac{g_0 \Omega}{2\Delta}$

- Rotating frame of pump,  $\omega = \omega_{\text{cavity}} - \omega_{\text{pump}}$

[Dimer *et al.* PRA '07 ]



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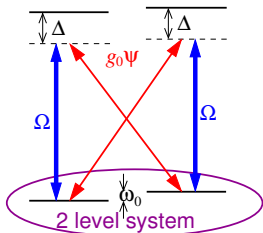
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[Dimer *et al.* PRA '07 ]

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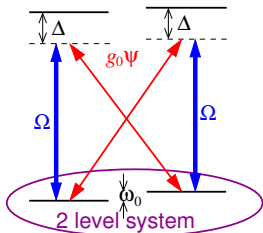


$$H = \omega_0 S^z + g(\hat{a} + \hat{a}^\dagger)(S^- + S^+) + \omega \hat{a}^\dagger \hat{a}$$

- 2 Level system,  $|\downarrow\rangle, |\uparrow\rangle$
- Coupling  $g = \frac{g_0 \Omega}{2\Delta}$
- Rotating frame of pump,  $\omega = \omega_{\text{cavity}} - \omega_{\text{pump}}$

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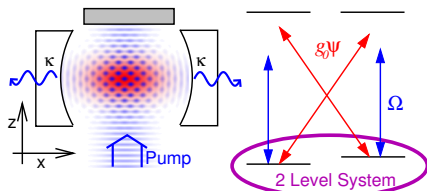
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## Open system Dicke transition

- $4Ng^2 > \frac{(\omega^2 + \kappa^2)}{\omega} \omega_0$
- Need  $H = \dots + g(S^+ a^\dagger + \text{H.c.}) + \dots$

[Dimer *et al.* PRA '07 ]

# Mapping transverse pumping to Dicke model

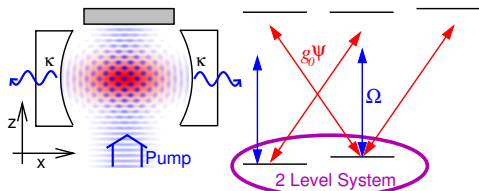


Reduced basis:  $\Psi(x, z) \propto \begin{cases} 1 & \Downarrow \\ \cos(qz) \cos(qz) & \Uparrow \end{cases}$

$$H = \omega \hat{a}^\dagger \hat{a} + \omega_0 S^z + g(\hat{a} + \hat{a}^\dagger)(S^- + S^+)$$

[Baumann *et al* Nature '10]

# Mapping transverse pumping to Dicke model



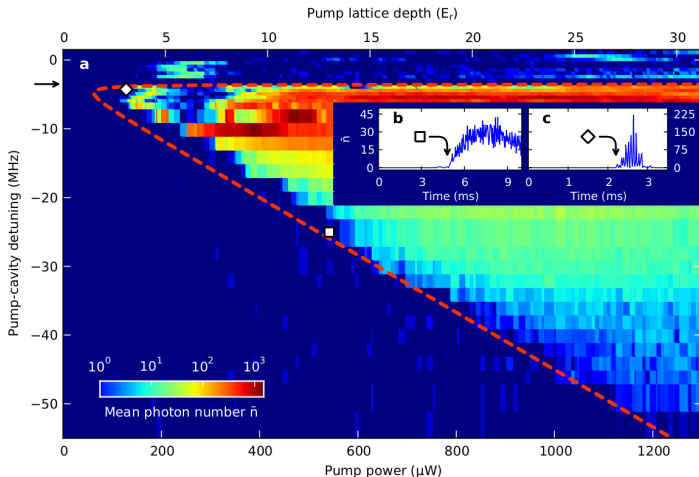
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$$H = \omega \hat{a}^\dagger \hat{a} + \omega_0 S^z + g(\hat{a} + \hat{a}^\dagger)(S^- + S^+) + U S_z \hat{a}^\dagger \hat{a}.$$

“Feedback” due to extra states  $U = -\frac{g_0^2}{4\Delta}$

[Baumann *et al* Nature '10]

# Experimental phase diagram



- Pump power  $g \propto \sqrt{\text{Power}}$
- NB  $\omega_{\text{Pump}} - \omega_{\text{cavity}} = -\omega$

[Baumann *et al* Nature '10]

# Cold atoms in optical cavities – Dicke model

## 1 Non-equilibrium phases of matter-light systems

- Physical systems

## 2 Cold atoms in optical cavities – Dicke model

- Superradiance and the Dicke model
- Cold atoms & open Dicke model
- **Dynamics and attractors**
- Beyond Dicke: Chaotic dynamics

## 3 Beyond mean-field theory

- Coupled cavity arrays
- Multimode cavities

# Classical dynamics of the extended Dicke model

Open dynamical system:

$$H = \omega \hat{a}^\dagger \hat{a} + \omega_0 S^z + g(\hat{a} + \hat{a}^\dagger)(S^- + S^+) + U S_z \hat{a}^\dagger \hat{a}.$$
$$\partial_t \rho = -i[H, \rho] + \kappa \mathcal{L}[\hat{a}]$$



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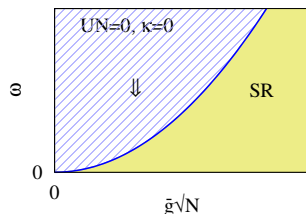
$$H = \omega \hat{a}^\dagger \hat{a} + \omega_0 \mathbf{S}^z + g(\hat{a} + \hat{a}^\dagger)(\mathbf{S}^- + \mathbf{S}^+) + U S_z \hat{a}^\dagger \hat{a}.$$
$$\partial_t \rho = -i[H, \rho] + \kappa \mathcal{L}[\hat{a}]$$

Classical EOM  
( $|\mathbf{S}| = N/2 \gg 1$ )

$$\dot{\mathbf{S}}^- = -i(\omega_0 + U|\alpha|^2)\mathbf{S}^- + 2ig(\alpha + \alpha^*)\mathbf{S}^z$$
$$\dot{\mathbf{S}}^z = ig(\alpha + \alpha^*)(\mathbf{S}^- - \mathbf{S}^+)$$
$$\dot{\alpha} = -[\kappa + i(\omega + US^z)]\alpha - ig(\mathbf{S}^- + \mathbf{S}^+)$$

# Steady state phase diagram

$$0 = i(\omega_0 + U|\alpha|^2)S^- + 2ig(\alpha + \alpha^*)S^z$$
$$0 = ig(\alpha + \alpha^*)(S^- - S^+)$$
$$0 = -[\kappa + i(\omega + US^z)]\alpha - ig(S^- + S^+)$$



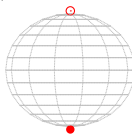
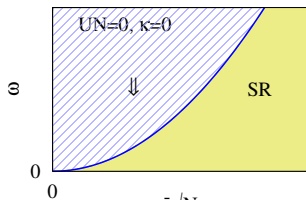
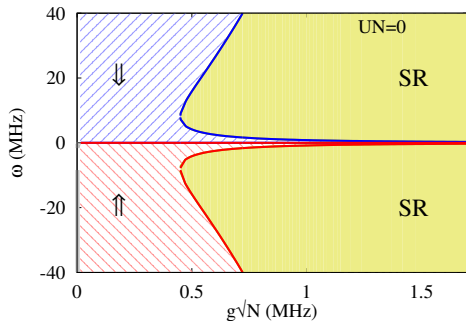
See also Domokos and Ritsch PRL '02, Domokos *et al.* PRL '10

# Steady state phase diagram

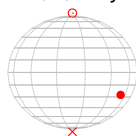
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$g\sqrt{N}$   
SR(A):  $S_y = 0$



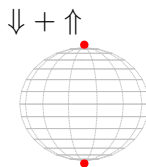
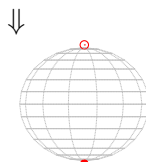
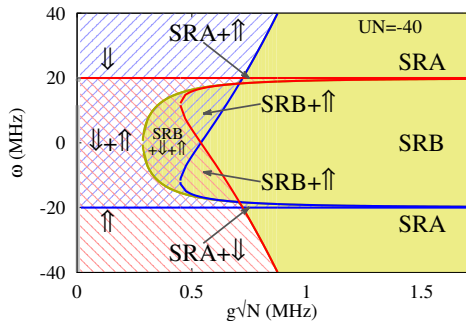
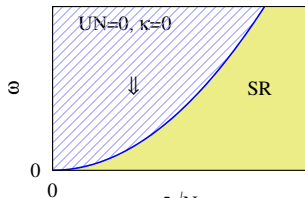
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# Steady state phase diagram

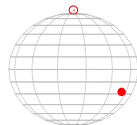
$$0 = i(\omega_0 + U|\alpha|^2)S^- + 2ig(\alpha + \alpha^*)S^z$$

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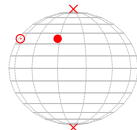
$$0 = -[\kappa + i(\omega + US^z)]\alpha - ig(S^- + S^+)$$



$\text{SR(A): } S_y = 0$



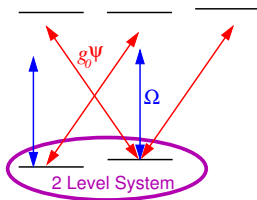
$\text{SR(B): } \alpha' = 0$



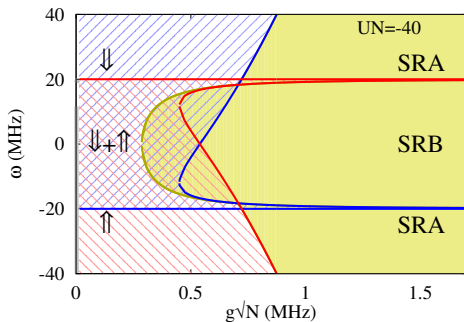
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# Regions without fixed points

Changing  $U$ :

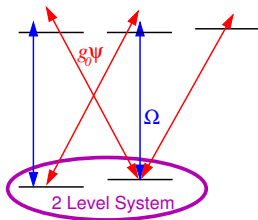


$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$

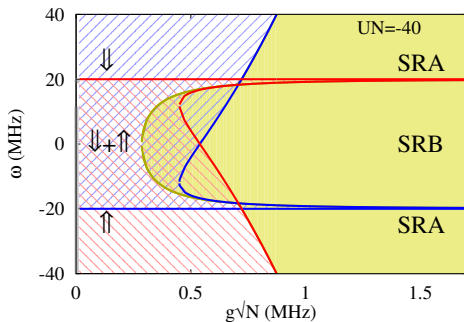


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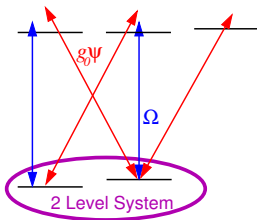


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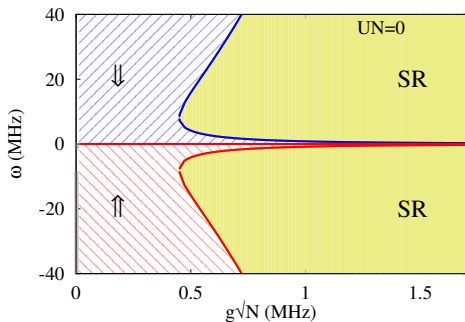


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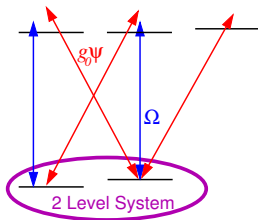


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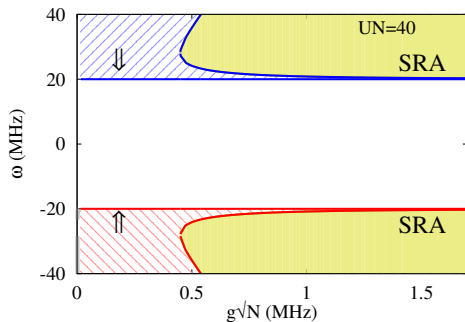


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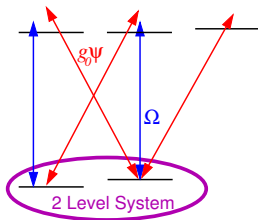
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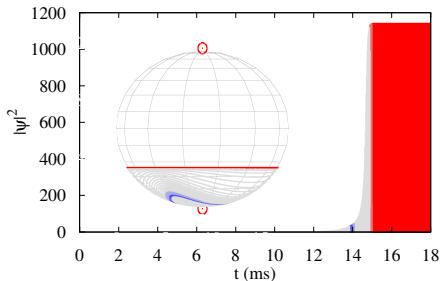
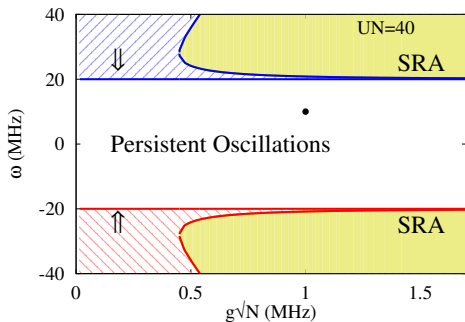


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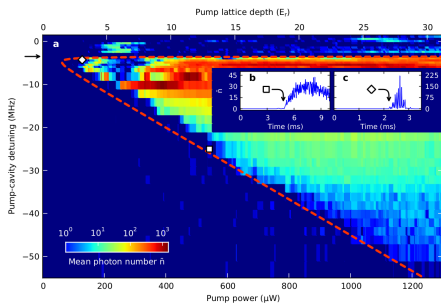
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# Bosons beyond Dicke — single mode

So far  $\Psi(\mathbf{r}) = \chi_0 + \chi_1 2 \cos(qx) \cos(qz) \rightarrow \mathbf{S} = \chi^\dagger \vec{\sigma} \chi$ .

Generally  $\Psi(\mathbf{r}) = \sum_{\mathbf{n}} \chi_{\mathbf{n}} e^{i\mathbf{q}\mathbf{n}\cdot\mathbf{r}}$ .

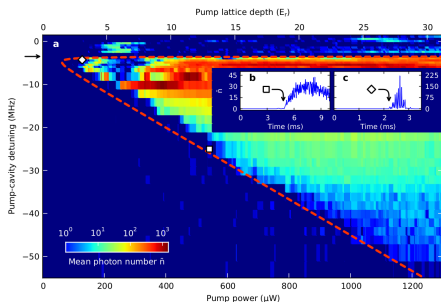


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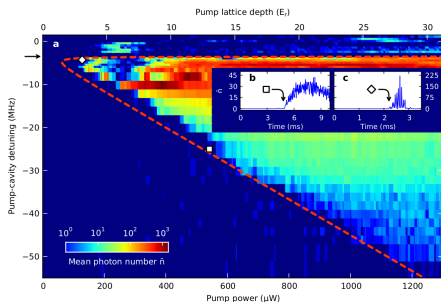
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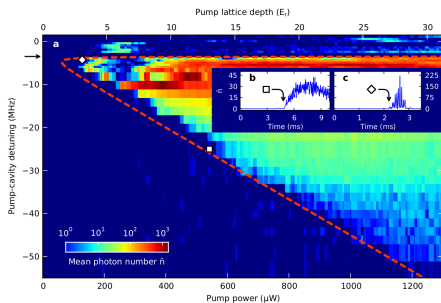
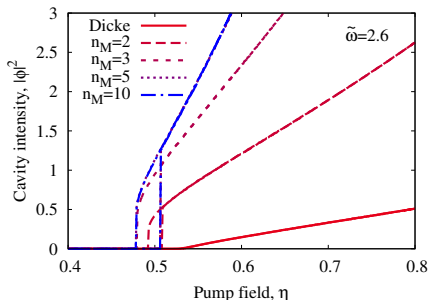
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Truncate  $|\mathbf{n}| < n_M$  — Hysteresis at intermediate  $\omega$



# Bosons beyond Dicke — single mode

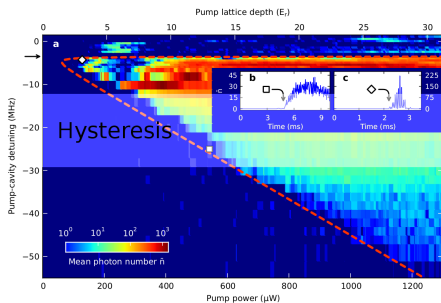
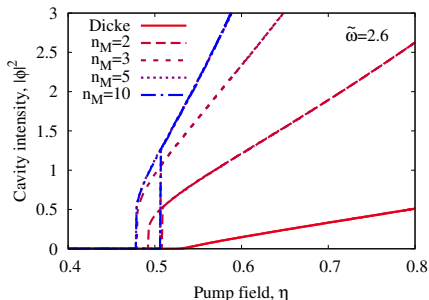
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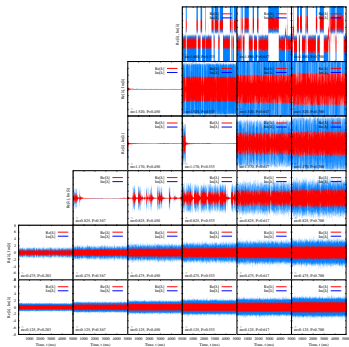
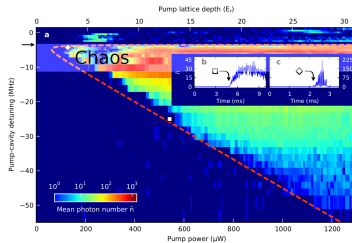
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# Bosons beyond Dicke — chaos

Near resonance: irregular dynamics

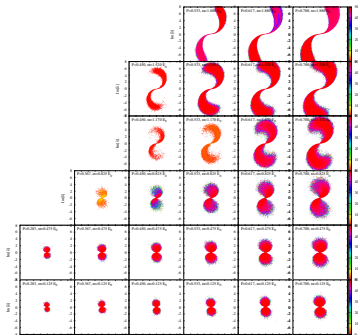
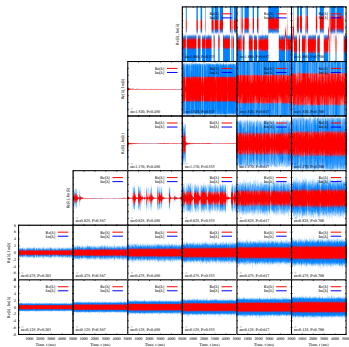
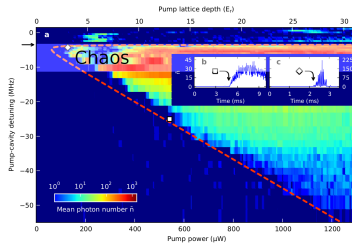
(NB  $\omega_{\text{Pump}} - \omega_{\text{cavity}} = -\omega$ )



# Bosons beyond Dicke — chaos

Near resonance: irregular dynamics

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# Beyond mean-field theory

## 1 Non-equilibrium phases of matter-light systems

- Physical systems

## 2 Cold atoms in optical cavities – Dicke model

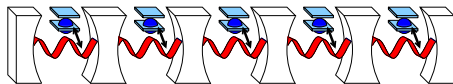
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## 3 Beyond mean-field theory

- Coupled cavity arrays
- Multimode cavities

# Coupled cavity arrays

- Control photon dispersion — lattice

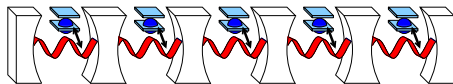


[Hartmann *et al.* Nat. Phys. '06; Greentree *et al.* *ibid* 06; Angelakis *et al.* PRA '07]

• X-Hubbard Model,  $\hat{H} = \sum_I \hat{H}_{X,site} - J \sum_{\langle i,j \rangle} \hat{a}_i^\dagger \hat{a}_j$   
[X=Bose, Jaynes-Cummings, Rabi, ...]

# Coupled cavity arrays

- Control photon dispersion — lattice



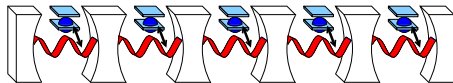
[Hartmann *et al.* Nat. Phys. '06; Greentree *et al.* *ibid* 06; Angelakis *et al.* PRA '07]

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# Coupled cavity arrays

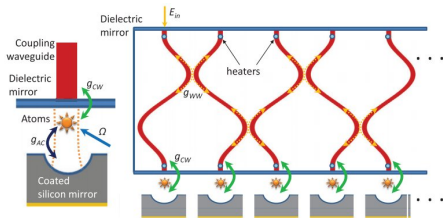
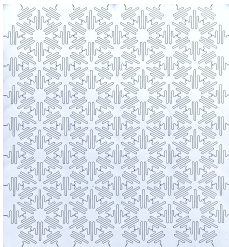
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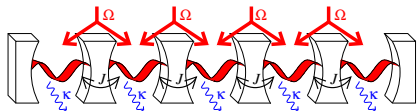
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[Lepert *et al.* NJP '11; APL '13]

[Underwood *et al.* PRA '12; Nat. Phys '12]

## e.g. Parametric pumping $\rightarrow$ Ising model

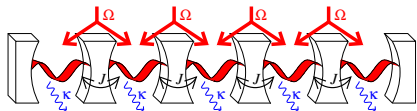


[Bardyn & Imamoglu, PRL '12]

$$\hat{H} = -J \sum \left[ \hat{\tau}_i^+ \hat{\tau}_{i+1}^- + \hat{\tau}_{i+1}^+ \hat{\tau}_i^- + g \hat{\tau}_i^z \right. \\ \left. + \Delta \left( \hat{\tau}_i^+ \hat{\tau}_{i+1}^+ + \hat{\tau}_{i+1}^- \hat{\tau}_i^- \right) \right]$$

[Joshi, Nissen, Keeling. PRA '13]

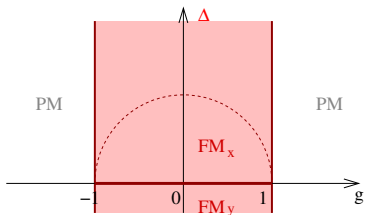
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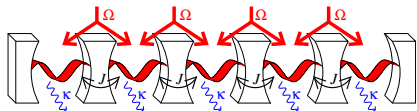
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Ground state:



[Joshi, Nissen, Keeling. PRA '13]

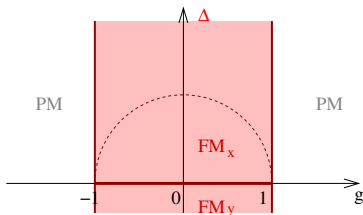
# e.g. Parametric pumping $\rightarrow$ Ising model



[Bardyn & Imamoglu, PRL '12]

$$\hat{H} = -J \sum \left[ \hat{\tau}_i^+ \hat{\tau}_{i+1}^- + \hat{\tau}_{i+1}^+ \hat{\tau}_i^- + g \hat{\tau}_i^z \right. \\ \left. + \Delta \left( \hat{\tau}_i^+ \hat{\tau}_{i+1}^+ + \hat{\tau}_{i+1}^- \hat{\tau}_i^- \right) \right]$$

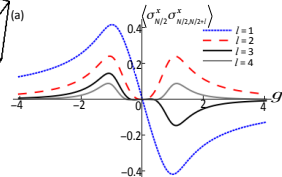
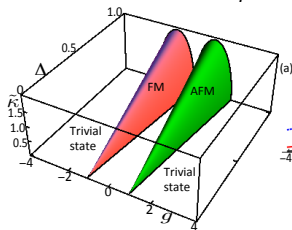
Ground state:



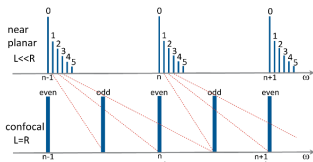
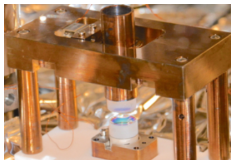
[Joshi, Nissen, Keeling. PRA '13]

Open system:

$$\partial_t \rho = -i[\hat{H}, \rho] + \sum_i \kappa \mathcal{L}[\hat{\tau}_i^-]$$



# Multimode cavity



[Kollár *et al.* NJP '15]

- Degenerate limit, on-axis pump

$$i\partial_t \psi = -\frac{\nabla^2}{2m} \psi + \int dr' U(r-r') [|\psi(r')|^2 + |\psi(-r')|^2] \psi(r).$$

with  $U(s) = U_0 \left( \delta(s) + \frac{\sin[s^2/\ell^2]}{\pi \ell^2} \right)$ .

[Staffini, Lev, Keeling *et al.* in prep.]

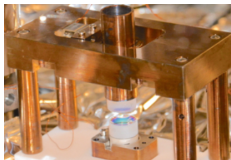
- Degenerate limit, transverse pump:

$$i\partial_t \psi_k = \left[ \Delta + \lambda(|k| - q)^2 \right] \psi_k + U_{\text{contact}} \sum_{k', q} \psi_{k'+q}^* \psi_{k'} \psi_{k-q}$$

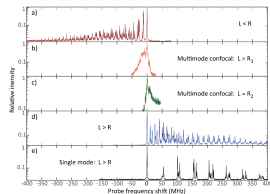
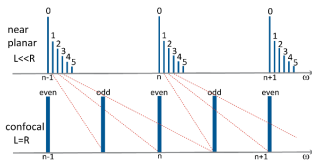
[Gopalakrishnan, Lev, Goldbart. N. Phys '09]



# Multimode cavity



[Kollár *et al.* NJP '15]



- Degenerate limit, on-axis pump

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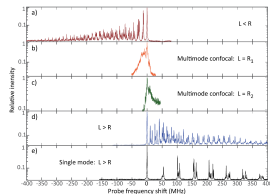
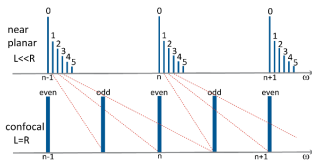
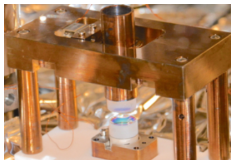
[Staffini, Lev, Keeling *et al.* in prep.]

- Degenerate limit, transverse pump:

$$i\partial_t \psi_{\mathbf{k}} = \left[ \Delta + \lambda(|\mathbf{k}| - q)^2 \right] \psi_{\mathbf{k}} + U_{\text{contact}} \sum_{\mathbf{k}', q} \psi_{\mathbf{k}'+q}^* \psi_{\mathbf{k}} \psi_{\mathbf{k}-q}$$

[Gopalakrishnan, Lev, Goldbart. N. Phys '09]

# Multimode cavity



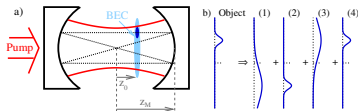
[Kollár *et al.* NJP '15]

- Degenerate limit, on-axis pump

$$i\partial_t \Psi = -\frac{\nabla^2}{2m} \Psi + \int d\mathbf{r}' U(\mathbf{r} - \mathbf{r}') \left[ |\Psi(\mathbf{r}')|^2 + |\Psi(-\mathbf{r}')|^2 \right] \Psi(\mathbf{r}),$$

with  $U(\mathbf{s}) = U_0 \left( \delta(\mathbf{s}) + \frac{\sin[s^2/\ell^2]}{\pi\ell^2} \right)$ ,

[Staffini, Lev, Keeling *et al.* in prep.]

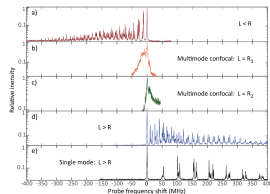
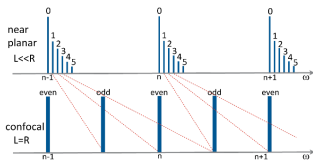
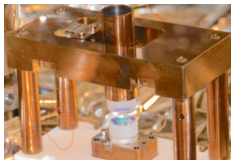


• Degenerate limit, transverse pump:

$$i\partial_t \Psi_{\mathbf{k}} = \left[ \Delta + \lambda(|\mathbf{k}| - q)^2 \right] \Psi_{\mathbf{k}} + U_{\text{contact}} \sum_{\mathbf{k}' < \mathbf{q}} \Psi_{\mathbf{k}+\mathbf{q}} \Psi_{\mathbf{k}'} \Psi_{-\mathbf{q}}$$

[Gopalakrishnan, Lev, Goldbart, N. Phys '19]

# Multimode cavity



[Kollár *et al.* NJP '15]

- Degenerate limit, on-axis pump

$$i\partial_t \Psi = -\frac{\nabla^2}{2m} \Psi + \int d\mathbf{r}' U(\mathbf{r} - \mathbf{r}') \left[ |\Psi(\mathbf{r}')|^2 + |\Psi(-\mathbf{r}')|^2 \right] \Psi(\mathbf{r}),$$

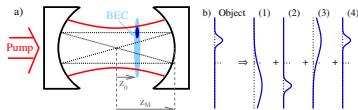
with  $U(\mathbf{s}) = U_0 \left( \delta(\mathbf{s}) + \frac{\sin[s^2/\ell^2]}{\pi\ell^2} \right)$ ,

[Staffini, Lev, Keeling *et al.* in prep.]

- Degenerate limit, transverse pump:

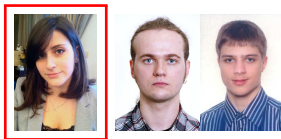
$$i\partial_t \Psi_{\mathbf{k}} = \left[ \Delta + \lambda(|\mathbf{k}| - q)^2 \right] \Psi_{\mathbf{k}} + U_{\text{contact}} \sum_{\mathbf{k}', \mathbf{q}} \Psi_{\mathbf{k}'+\mathbf{q}}^* \Psi_{\mathbf{k}'} \Psi_{\mathbf{k}-\mathbf{q}}$$

[Gopalakrishnan, Lev, Goldbart. N. Phys '09]



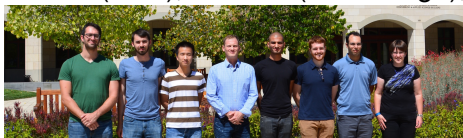
# Acknowledgements

GROUP:



COLLABORATORS: Bhaseen (KCL), Simons (Cambridge),

Lev group (Stanford)



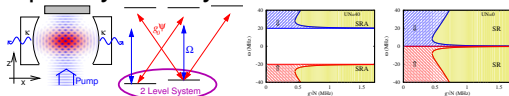
FUNDING:



The Leverhulme Trust

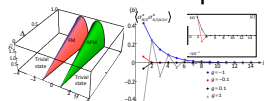
# Summary

- Collective behaviour in driven-dissipative systems
- Open system dynamics of Dicke model



JK *et al.* PRL '10, Bhaseen *et al.* PRA '12

- Beyond Dicke model: hysteresis, chaos  
[Staffini *et al.* in preparation]
- CCA — non-equilibrium transverse field Ising model



Joshi *et al.* PRA '13, Schiro *et al.* arXiv

- Multimode physics  $\rightarrow$  beyond mean-field theory  
[Staffini, Lev *et al.* in preparation]





# “Textbook” Laser: Maxwell Bloch equations

$$H = \omega_0 \hat{a}^\dagger \hat{a} + \sum_{\alpha} \epsilon_{\alpha} \sigma_{\alpha}^z + g_{\alpha, \mathbf{k}} \left( \hat{a} \sigma_{\alpha}^{+} + \hat{a}^{\dagger} \sigma_{\alpha}^{-} \right)$$

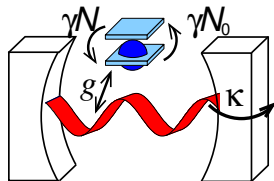
Maxwell-Bloch eqns:

$$\alpha = \langle \hat{a} \rangle, P = -i \langle \sigma^{-} \rangle, N = 2 \langle \sigma^z \rangle$$

$$\partial_t \alpha = -i \omega_0 \alpha - \kappa \alpha + \sum_{\alpha} g_{\alpha} P_{\alpha}$$

$$\partial_t P_{\alpha} = -2i \epsilon_{\alpha} P_{\alpha} - 2\gamma P_{\alpha} + g_{\alpha} \alpha N_{\alpha}$$

$$\partial_t N_{\alpha} = 2\gamma(N_0 - N_{\alpha}) - 2g_{\alpha}(\alpha^* P_{\alpha} + P_{\alpha}^* \alpha)$$





# “Textbook” Laser: Maxwell Bloch equations

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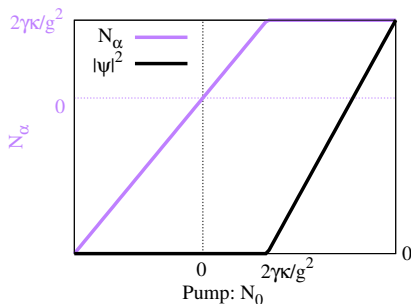
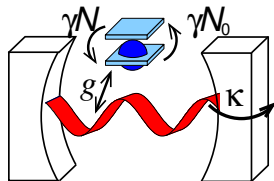
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$$\partial_t N_{\alpha} = 2\gamma(N_0 - N_{\alpha}) - 2g_{\alpha}(\alpha^* P_{\alpha} + P_{\alpha}^* \alpha)$$



$$|\alpha|^2 > 0 \text{ if } N_0 g^2 > 2\gamma\kappa$$