

Photon and polariton condensation with organic molecules

Jonathan Keeling

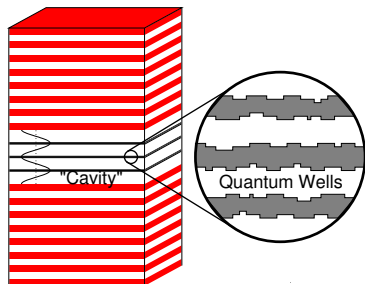


University of
St Andrews

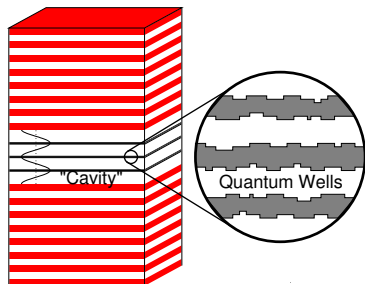
1413-2013

UK-NL meeting, Amsterdam, September 2015

Microcavity polaritons

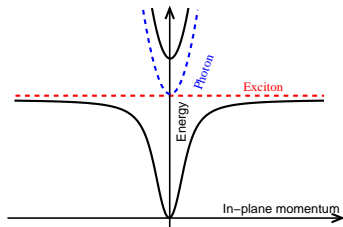


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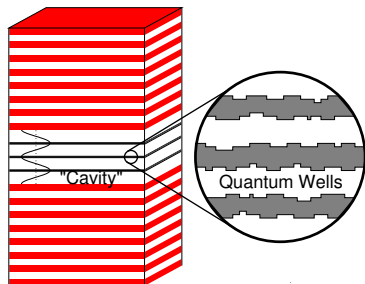


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2/2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$

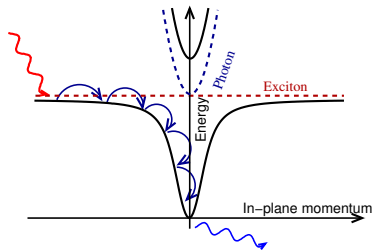


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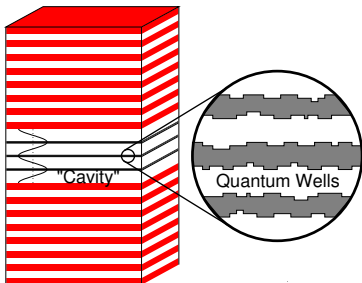


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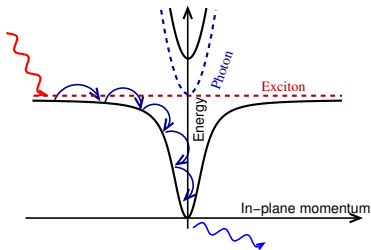
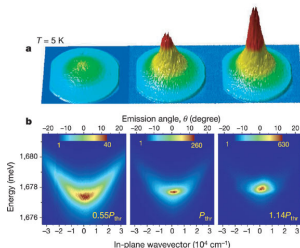


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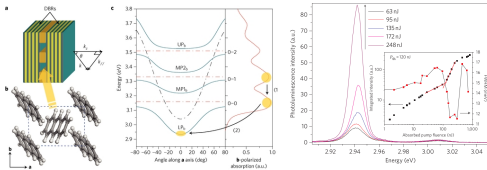
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Motivation: polariton condensates

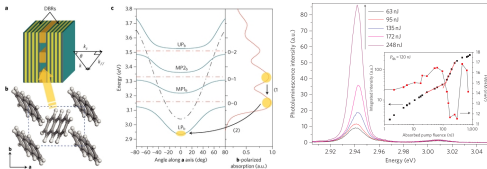
- Anthracene Polariton Lasing
 $T \sim 300\text{K}$



[Kena Cohen and Forrest, Nat. Photon '10]

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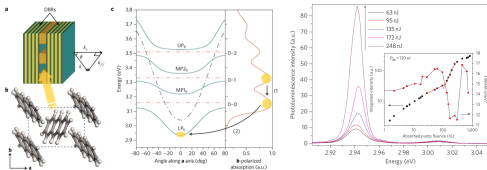


- Q1. Vibrational replicas?
- Q2. Relevance of disorder?
- Q3. Lasing vs condensation?

[Kena Cohen and Forrest, Nat. Photon '10]

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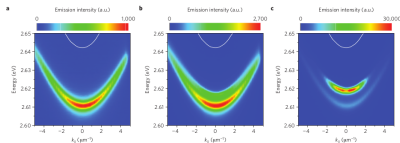
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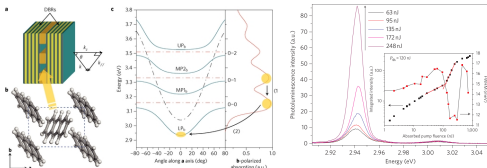
- Polariton condensates, other materials, e.g. polymers:



[Plumhoff *et al.* Nat. Materials '14, Daskalakis *et al.* *ibid* '14]

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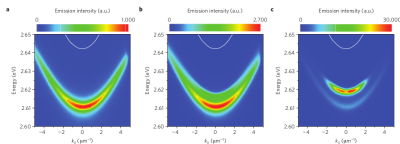
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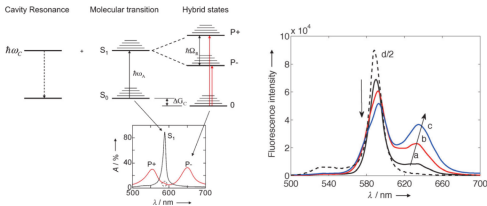


- Q1. Frenkel to Wannier crossover?
- Q2. Optimal vibrational properties?
- Q3. Nonlinearities?

[Plumhoff *et al.* Nat. Materials '14, Daskalakis *et al.* *ibid* '14]

Motivation: vacuum-state strong coupling

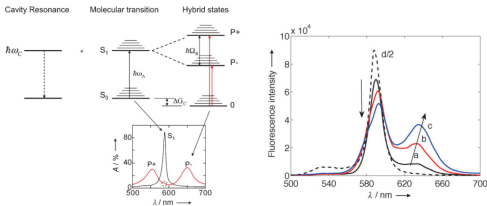
- Linear response (no pump, no condensate): effects of matter-light coupling alone.



[Canaguier-Durand *et al.* Angew. Chem. '13]

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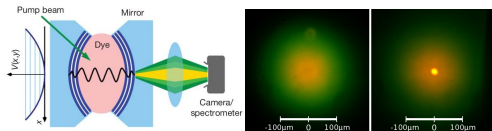
Q1. Can **ultra-strong** coupling to light change:

- ▶ charge distribution?
- ▶ vibrational configuration?
- ▶ molecular orientation?
- ▶ crystal structure?

Q2. Are changes collective (\sqrt{N} factor) or not?

Motivation: photon condensates

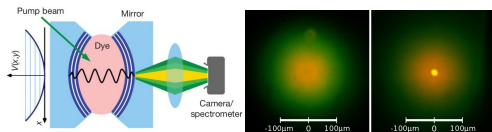
- Photon Condensate $T \sim 300\text{K}$



[Klaers *et al.* Nature, '10]

Motivation: photon condensates

● Photon Condensate $T \sim 300\text{K}$



[Klaers *et al.* Nature, '10]

- Q1. Relation to dye laser?
- Q2. Relation to polaritons?
- Q3. Thermalisation breakdown?

- 1 Modelling photon BEC
 - Steady state
- 2 Spatial profile
 - Steady state
 - Spatial oscillations
- 3 (Toward) strong coupling

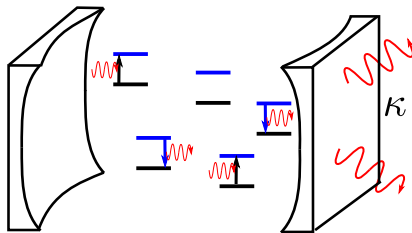
Dicke Holstein Model

- Dicke model: 2LS \leftrightarrow photons

- Molecular vibrational mode

- Phonon frequency Ω

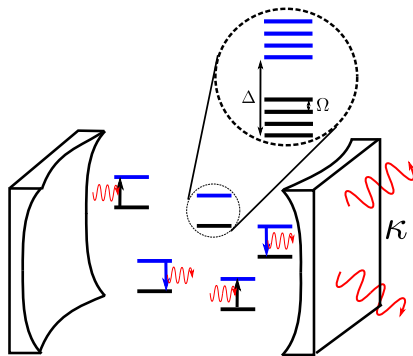
- Huang-Rhys parameter S —
coupling strength



$$H_{\text{sys}} = \omega \psi^\dagger \psi + \sum_{\alpha} \left[\frac{\epsilon}{2} \sigma_{\alpha}^z + g (\psi + \psi^\dagger) (\sigma_{\alpha}^+ + \sigma_{\alpha}^-) \right]$$

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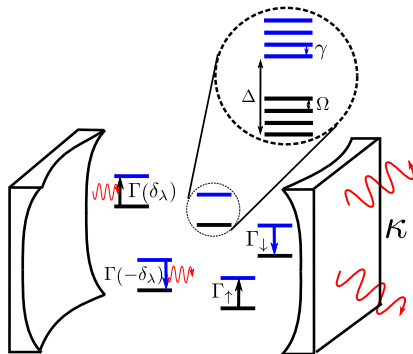
Photon: Microscopic Model

$$H_{\text{sys}} = \sum_m \omega_m \psi_m^\dagger \psi_m + \sum_\alpha \left[\frac{\epsilon}{2} \sigma_\alpha^z + g (\psi_m \sigma_\alpha^+ + \text{H.c.}) \right] \\ + \sum_\alpha \Omega \left\{ b_\alpha^\dagger b_\alpha + \sqrt{S} \sigma_\alpha^z (b_\alpha^\dagger + b_\alpha) \right\}$$

- **2D** harmonic oscillator

$$\omega_m = \omega_{\text{cutoff}} + m\omega_{H.O.}$$

- Incoherent processes: excitation, decay, loss, vibrational thermalisation.
- Weak coupling, perturbative in g

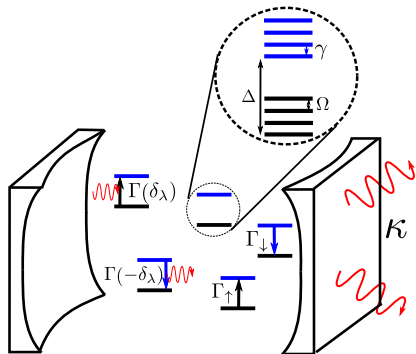


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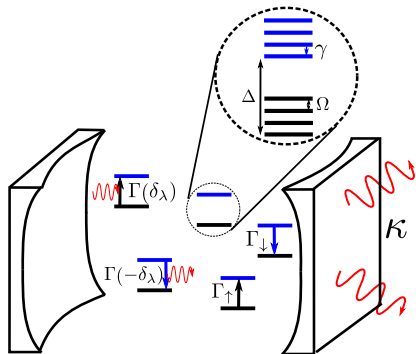
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Microscopic model – all orders in S

- Polaron transform (exact), $H = \sum_m \omega_m \psi_m^\dagger \psi_m + \sum_\alpha h_\alpha$,

$$h_\alpha = \frac{\epsilon}{2} \sigma_\alpha^z + g (\psi_m \sigma_\alpha^+ D_\alpha + \text{H.c.}) + \Omega b_\alpha^\dagger b_\alpha, \quad D_\alpha = e^{2\sqrt{S}(b_\alpha^\dagger - b_\alpha)}$$

- Master equation

$$\dot{\rho} = -i[H_0, \rho] + \sum_m \frac{\kappa}{2} \mathcal{L}[\psi_m] + \sum_\alpha \left[\frac{\Gamma_1}{2} \mathcal{L}[\sigma_\alpha^+] + \frac{\Gamma_1}{2} \mathcal{L}[\sigma_\alpha^-] \right] + \sum_{m,\alpha} \left[\frac{\Gamma(\delta_m - \omega_m - \epsilon)}{2} \mathcal{L}[\sigma_\alpha^+ \psi_m] + \frac{\Gamma(-\delta_m - \epsilon - \omega_m)}{2} \mathcal{L}[\sigma_\alpha^- \psi_m^\dagger] \right]$$

- Correlation function:

$$f(\delta) = 2g^2 \pi \left[\int dt e^{-i\delta t - (\Gamma_+ + \Gamma_-) |t|/2} \langle \sigma_\alpha^+(t) \sigma_\alpha^-(0) \rangle \right]$$

[Marthaler et al PRL '11, Kirtou & JK PRL '13]

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$$G(\delta) = 2\sigma^2 \pi \int dt e^{-i\delta t - (\Gamma_\uparrow + \Gamma_\downarrow)t/2} \langle \sigma_\alpha^-(t) \sigma_\alpha^-(0) \rangle$$

[Marthaler et al PRL '11, Kirtou & JK PRL '13]

Microscopic model – all orders in S

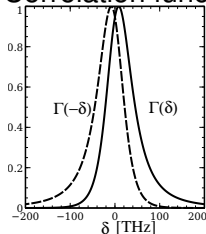
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$$\Gamma(\delta) = 2g^2 \Re \left[\int dt e^{-i\delta t - (\Gamma_\uparrow + \Gamma_\downarrow)t/2} \langle D_\alpha^\dagger(t) D_\alpha(0) \rangle \right]$$

[Marthaler et al PRL '11, Kirton & JK PRL '13]

Steady state populations and equilibrium

- Rate equation:

$$\partial_t n_m = -\kappa n_m + \Gamma(-\delta_m)(n_m + 1)N_{\uparrow} - \Gamma(\delta_m)n_m N_{\downarrow}$$

- Steady state distribution:

$$\frac{n_m}{n_m + 1} = \frac{\Gamma(-\delta_m)N_{\uparrow}}{\Gamma(\delta_m)N_{\downarrow}}$$

- Microscopic conditions for equilibrium:

- Emission/absorption rate:

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Steady state populations vs loss

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- Bose-Einstein distribution without losses

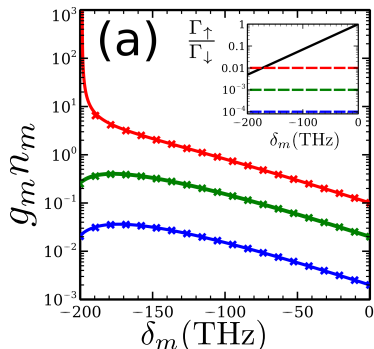
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Low loss: Thermal

[Kirton & JK PRL '13]

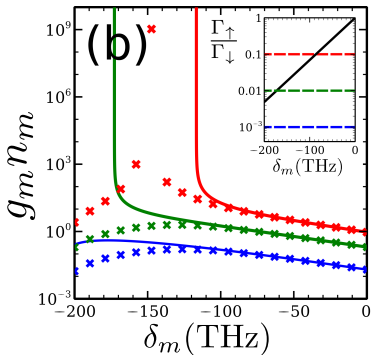
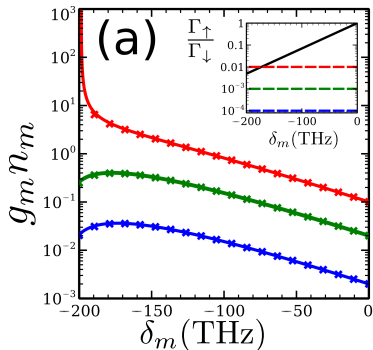
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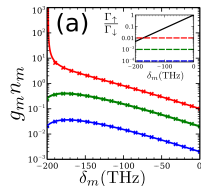
- $\kappa \ll N\Gamma(\delta)$, Kennard-Stepanov

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- Below threshold,

$$\mu = k_B T \ln[\Gamma_{\uparrow}/\Gamma_{\downarrow}]$$

- Above threshold, $\mu \rightarrow \delta_0$



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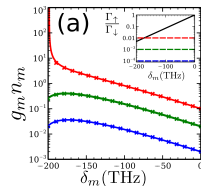
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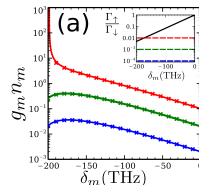
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- At/above threshold, $\mu \rightarrow \delta_0$



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$$\frac{n_m}{n_m + 1} = \frac{\Gamma(-\delta_m)N_{\uparrow}}{\kappa + \Gamma(\delta_m)N_{\downarrow}}$$

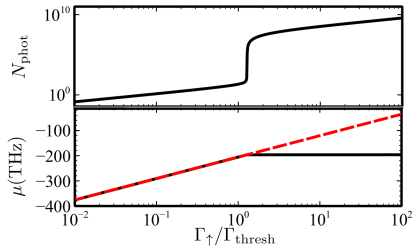
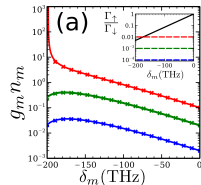
- $\kappa \ll N\Gamma(\delta)$, Kennard-Stepanov

$$\frac{n_m}{n_m + 1} = e^{-\beta\delta_m + \beta\mu}, \quad e^{\beta\mu} \equiv \frac{N_{\uparrow}}{N_{\downarrow}} = \frac{\Gamma_{\uparrow} + \sum_m \Gamma(\delta_m)n_m}{\Gamma_{\downarrow} + \sum_m \Gamma(-\delta_m)(n_m + 1)}$$

- Below threshold,

$$\mu = k_B T \ln[\Gamma_{\uparrow}/\Gamma_{\downarrow}]$$

- At/above threshold, $\mu \rightarrow \delta_0$



Spatial profile

1 Modelling photon BEC

- Steady state

2 Spatial profile

- Steady state
- Spatial oscillations

3 (Toward) strong coupling

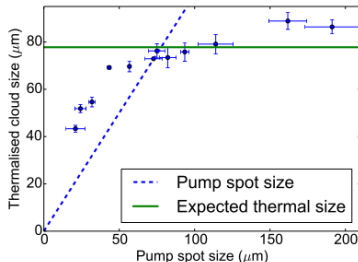
Spatially varying pump intensity

- Consider effects of pump profile, $\Gamma_{\uparrow}(\mathbf{r}) = \frac{\Gamma_{\uparrow} \exp(-r^2/2\sigma_p^2)}{(2\pi\sigma_p^2)^{d/2}}$

• Experiments: [Marelic & Nyman, PRA 15]

Spatially varying pump intensity

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- Experiments: [Marelic & Nyman, PRA '15]



Modelling spatial profile.

- Varying excited density – differential coupling to modes

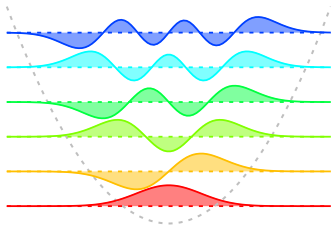
$$\partial_t n_m = -\kappa n_m + \Gamma(-\delta_m) O_m (n_m + 1) - \Gamma(\delta_m) (\rho_m - O_m) n_m$$

$$O_m = \int d\mathbf{r} \rho_{\uparrow}(\mathbf{r}) |\psi_{m1}(\mathbf{r})|^2, \quad \rho_{\uparrow} + \rho_{\downarrow} = \rho_m$$

Modelling spatial profile.

- Gauss-Hermite modes

$$I(\mathbf{r}) = \sum_m n_m |\psi_m(\mathbf{r})|^2$$



- Varying excited density – differential coupling to modes

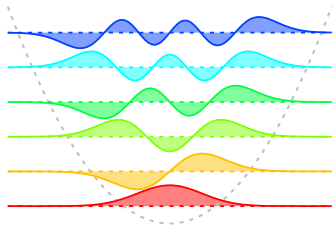
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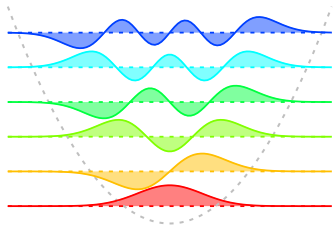
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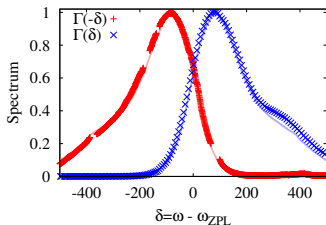
Modelling spatial profile.

- Gauss-Hermite modes

$$I(\mathbf{r}) = \sum_m n_m |\psi_m(\mathbf{r})|^2$$



- Use exact R6G spectrum



- Varying excited density – differential coupling to modes

$$\partial_t n_m = -\kappa n_m + \Gamma(-\delta_m) O_m (n_m + 1) - \Gamma(\delta_m) (\rho_m - O_m) n_m$$

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$$\partial_t \rho_{\uparrow}(\mathbf{r}) = -\tilde{\Gamma}_{\downarrow}(\mathbf{r}) \rho_{\uparrow}(\mathbf{r}) + \tilde{\Gamma}_{\uparrow}(\mathbf{r}) \rho_{\downarrow}(\mathbf{r})$$

Spatially varying pump: below threshold

- Far below threshold:

- ▶ If $\kappa \ll \rho_m \Gamma(\delta_m)$,
$$\frac{n_m}{n_m + 1} \simeq e^{-\beta \delta_m} \times \int d\mathbf{r} \rho_{\uparrow}(\mathbf{r}) |\psi_m(\mathbf{r})|^2$$

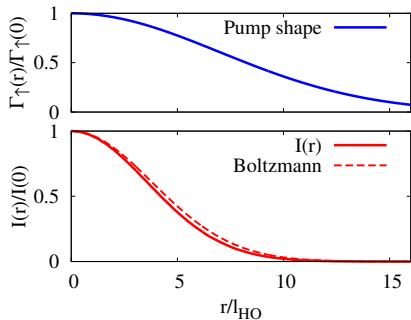
● Resulting profile, $I(\mathbf{r}) = \sum_m n_m |\psi_m(\mathbf{r})|^2$

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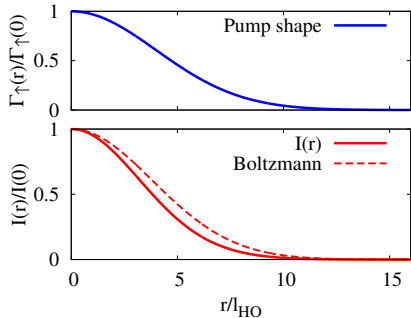


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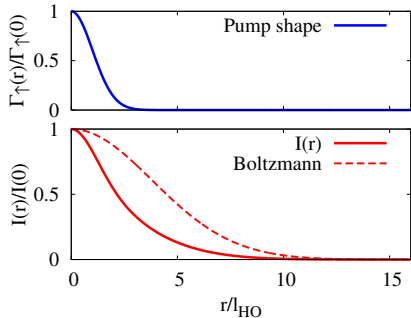


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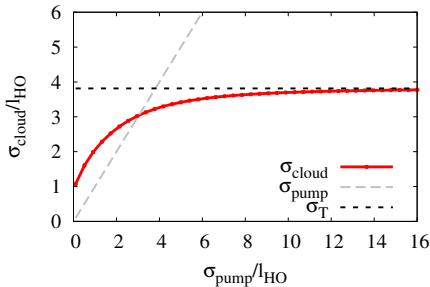
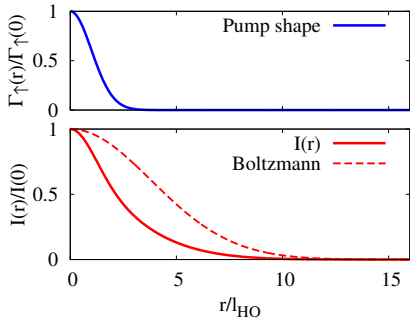
Spatially varying pump: below threshold

- Far below threshold:

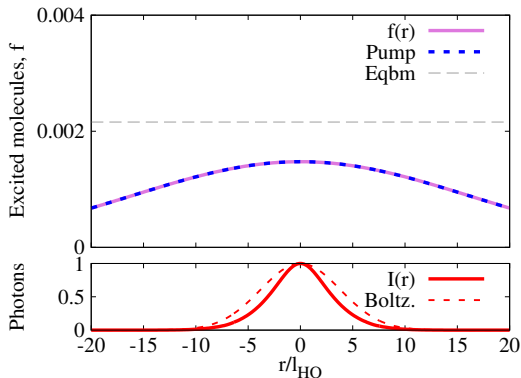
▶ If $\kappa \ll \rho_m \Gamma(\delta_m)$,

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- Resulting profile, $I(\mathbf{r}) = \sum_m n_m |\psi_m(\mathbf{r})|^2$



Near threshold behaviour

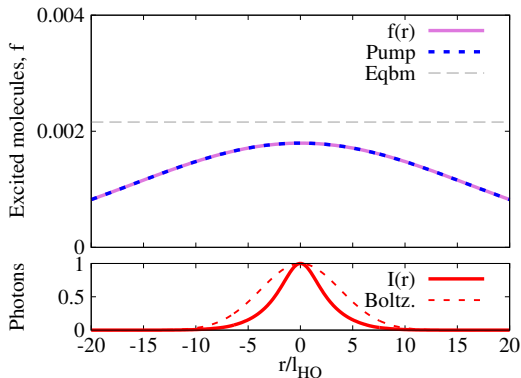


- Large spot, $\sigma_p \gg l_{HO}$

- "Gain saturation" at centre

- Saturation of $f(r) = 1/(1 + e^{-\beta r})$ — spatial equilibration

Near threshold behaviour

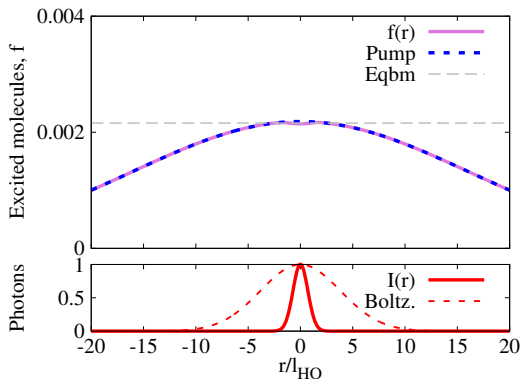


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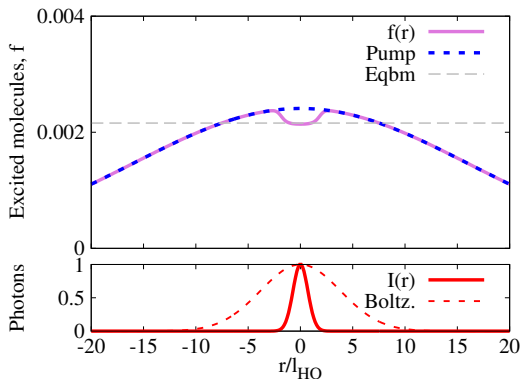
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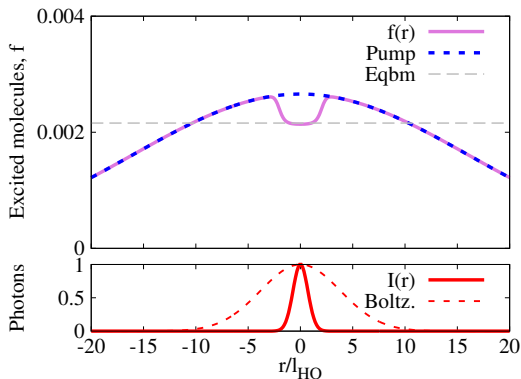
- Large spot, $\sigma_p \gg l_{HO}$
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Spatial oscillations

1 Modelling photon BEC

- Steady state

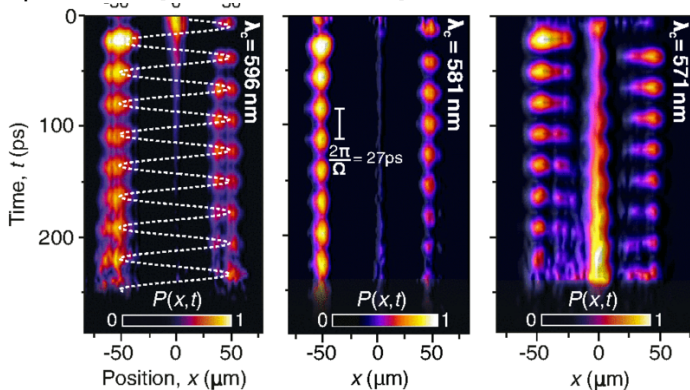
2 Spatial profile

- Steady state
- **Spatial oscillations**

3 (Toward) strong coupling

Off centre pumping; oscillations

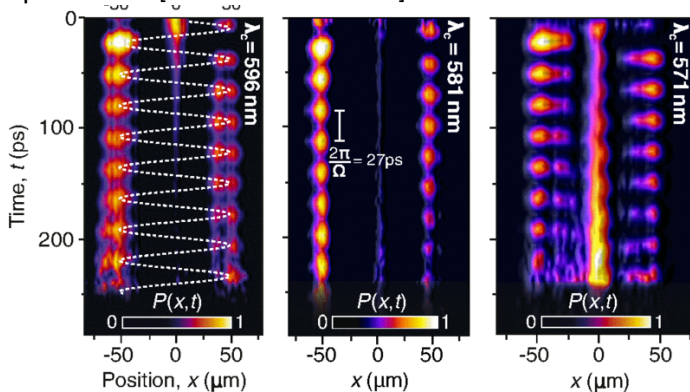
- Experiments [Schmitt *et al.* PRA '15]



- Oscillations in space – beating of normal modes
- Thermalisation depends on cutoff

Off centre pumping; oscillations

- Experiments [Schmitt *et al.* PRA '15]

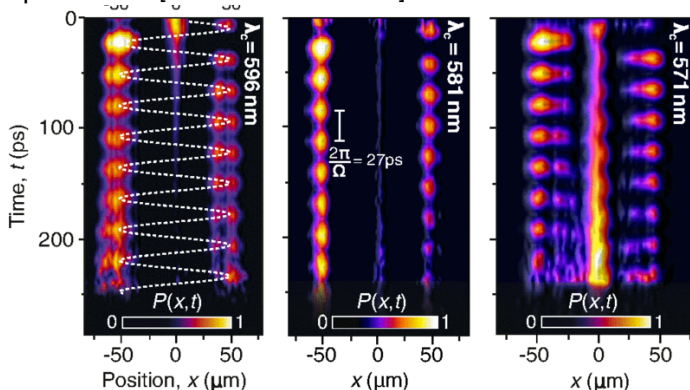


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Off centre pumping; oscillations

- Experiments [Schmitt *et al.* PRA '15]



- Oscillations in space – beating of normal modes
- Thermalisation depends on cutoff

Modelling

- Full master equation required

$$\partial_t \rho = -i \left[\sum_m \omega_m \hat{a}_m^\dagger \hat{a}_m, \rho \right] + \sum_{m,m',i} \psi_m^*(r_i) \psi_{m'}(r_i) \left(K(\delta_{m'}) [\hat{a}_{m'} \hat{\sigma}_i^+ \hat{\rho}, \hat{a}_m^\dagger \hat{\sigma}_i^-] \right. \\ \left. + K(-\delta_m) [\hat{a}_m^\dagger \hat{\sigma}_i^- \hat{\rho}, \hat{a}_{m'} \hat{\sigma}_i^+] \right) + \text{H.c.} + (\text{pumping, decay} \dots),$$

- Not secular approximation

Modelling

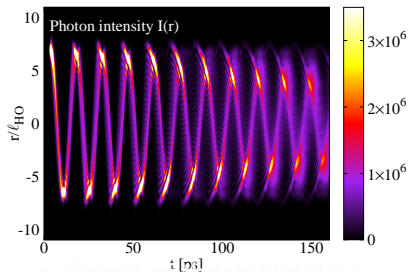
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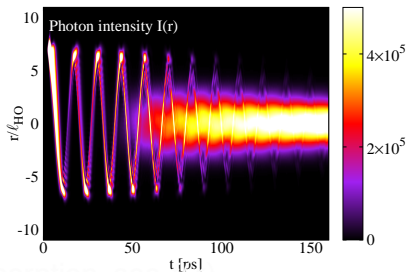
- Not secular approximation
 - ▶ **Must** have emission into m, m' superposition
 - ▶ **Must** have $K = K(\delta_m)$ (Kennard-Stepanov)

Dynamics from model

Longer cavity



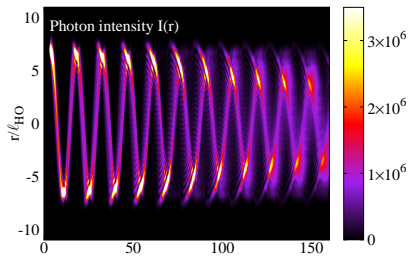
Shorter cavity



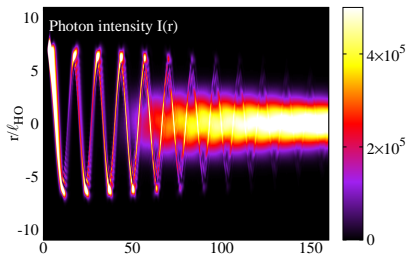
Origin of thermalisation — reabsorption, see Fig. 1

Dynamics from model

Longer cavity



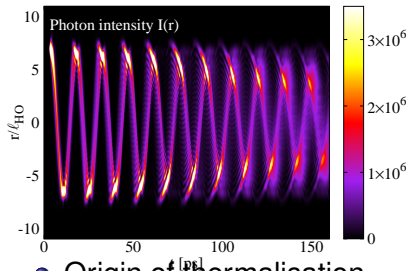
Shorter cavity



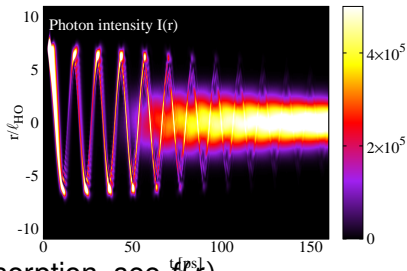
- Origin of thermalisation — reabsorption, see $f(r)$

Dynamics from model

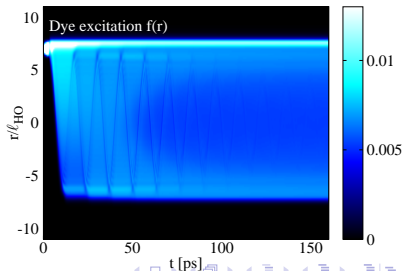
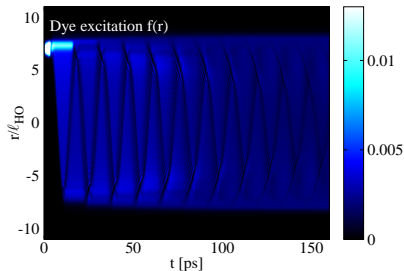
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Shorter cavity



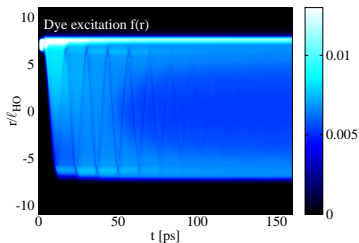
● Origin of thermalisation — reabsorption, see $f(r)$



Thermalisation at late times

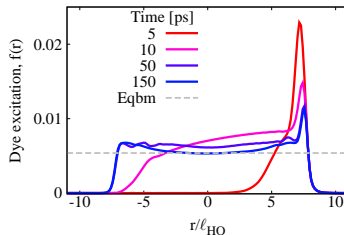
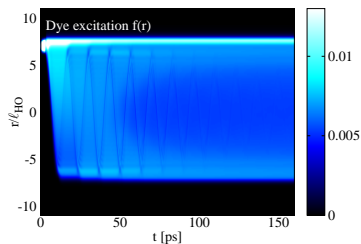
- Reabsorption “fills-in” excited molecules

● Reach thermal equilibrium, $f = [e^{-\beta\epsilon} + 1]^{-1}$



Thermalisation at late times

- Reabsorption “fills-in” excited molecules
- Reach thermal equilibrium, $f = [e^{-\beta\delta_0} + 1]^{-1}$



(Toward) strong coupling

1 Modelling photon BEC

- Steady state

2 Spatial profile

- Steady state
- Spatial oscillations

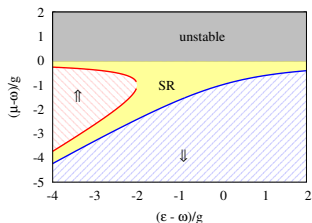
3 (Toward) strong coupling

Strong coupling phase diagram — mean field

- Mean field — single photon mode

$$H = \omega \psi^\dagger \psi + \sum_{\alpha} \left[\frac{\epsilon}{2} \sigma_{\alpha}^z + g \left(\psi \sigma_{\alpha}^+ + \psi^\dagger \sigma_{\alpha}^- \right) + \Omega \left\{ b_{\alpha}^\dagger b_{\alpha} + \sqrt{S} \left(b_{\alpha}^\dagger + b_{\alpha} \right) \sigma_{\alpha}^z \right\} \right]$$

- Dicke phase diagram vs μ

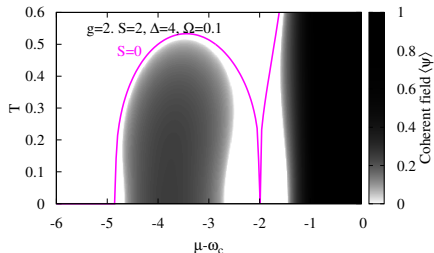


- S reduces g_{eff}
- Reentrant behaviour — Min μ at $k_B T \sim 0.1\Omega$

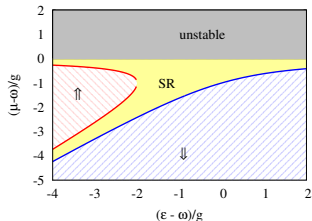
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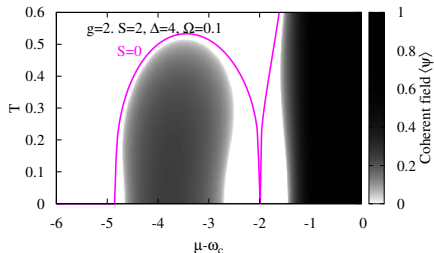
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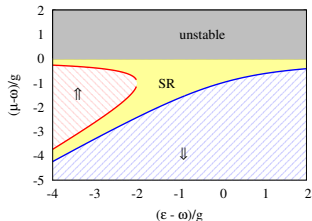
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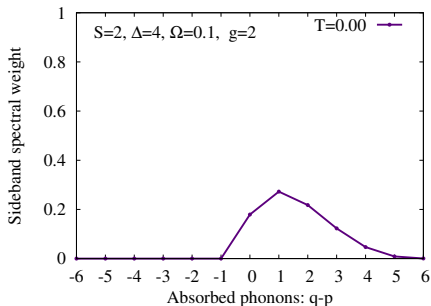


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Polariton spectrum: what condensed

- Vibrational composition of condensate

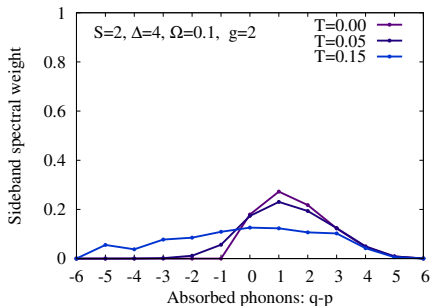


[Cwik *et al.* EPL '14]

• Optimal $T \sim 2\Omega$

Polariton spectrum: what condensed

- Vibrational composition of condensate

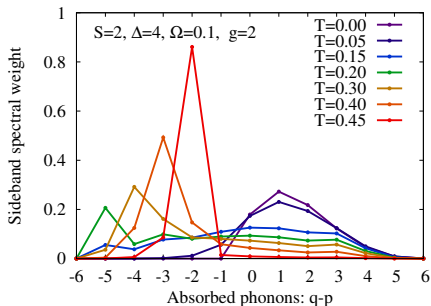


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[Cwik *et al.* EPL '14]

Polariton spectrum: what condensed

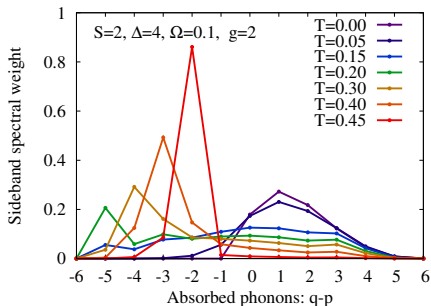
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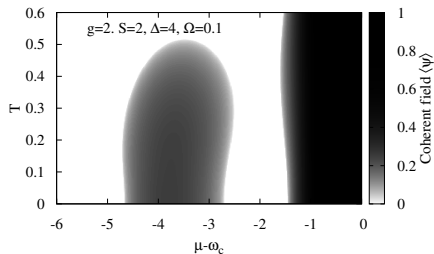
[Cwik *et al.* EPL '14]

Polariton spectrum: what condensed

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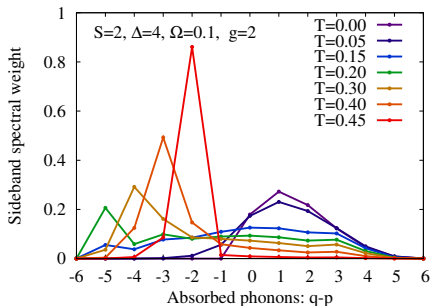
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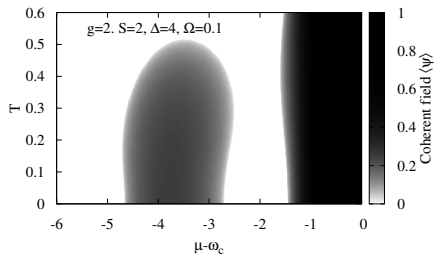
[Cwik *et al.* EPL '14]

Polariton spectrum: what condensed

- Vibrational composition of condensate



- Optimal $T \sim 2\Omega$



[Cwik *et al.* EPL '14]

Acknowledgements

GROUP:



FUNDING:



The Leverhulme Trust

CONDENSATES OF LIGHT

13TH-15TH JANUARY 2016, CHICHELEY HALL, BUCKINGHAMSHIRE, UK

The scope of the workshop covers all physical realisations of condensates of light, including:

- Bose-Einstein condensation of light in dye-filled microcavities
- Condensates of semiconductor exciton-polaritons
- Organic polariton condensation
- Classical condensation phenomena in optics
- Superfluid light

Keynote speakers:

- Jason Fleischer (Princeton, USA)
- Elisabeth Giacobino (LKB, Paris, France)
- Stephane Kena-Cohen (Polytechnique, Montreal, Canada)
- Henk Stoof (Utrecht, Netherlands)
- Martin Weitz (Bonn, Germany)

Invited speakers:

- Alberto Amo (Marcoussis, France)
- Iacopo Carusotto (Trento, Italy)
- Natalia Berloff (Cambridge, UK/Moscow)
- Baruch Fischer (Technion, Haifa, Israel)
- Jonathan Keeling (St Andrews, UK)
- Pavlos Lagoudakis (Southampton, UK)
- Rainer Mahrth (IBM Zurich, Switzerland)
- Gian-Luca Oppo (Strathclyde, UK)
- Antonio Picozzi (Bourgogne, France)
- Daniele Sanvitto (Lecce, Italy)
- David Snoke (Pittsburgh, USA)
- Marzena Szymanska (UCL, UK)
- Jacob Taylor (Maryland, USA)

Important Dates:

- 31st October 2015: abstract submission deadline
- 30th November 2015: registration deadline

Organising committee:

- Robert Nyman (Imperial College, London)
- Peter Kirton (St Andrews)
- Jonathan Keeling (St Andrews)
- Martin Weitz (Bonn)

<http://condensates-of-light.org>

ICSCE 8

Edinburgh, 25th–29th April, 2016.



Plenary speakers: Ataç İmamoğlu, Peter Zoller.

Invited speakers: Ehud Altman, Mete Atatüre, Natasha Berloff, Charles Bardin[†], Jacqueline Bloch, Iacopo Carusotto[†], Cristiano Ciuti, Michele Devoret[†], Thomas Ebbesen, Thierry Giamarchi, Jan Klärs, Dmitry Krizhanovskii, Xiaoqin (Elaine) Li, Peter Littlewood, Allan MacDonald, Francesca Marchetti, Keith Nelson, Pavlos Lagoudakis, Vivien Zapf.

([†] To be confirmed)

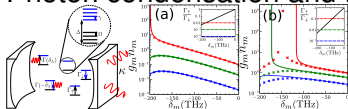
Early-bird registration & abstract deadline: 31st January 2016.

Final registration deadline: 31st March 2016.

<http://www.st-andrews.ac.uk/~icsce8>

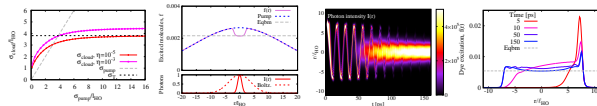
Summary

- Photon condensation and thermalisation



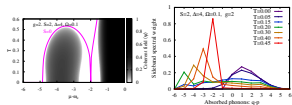
[Kirton & JK, PRL '13, PRA '15]

- Photon condensation, pattern formation physics



[JK & Kirton, arXiv:1506:00280]

- Reentrance, phonon assisted transition, 1st order at $S \gg 1$



[Cwik *et al.* EPL '14]

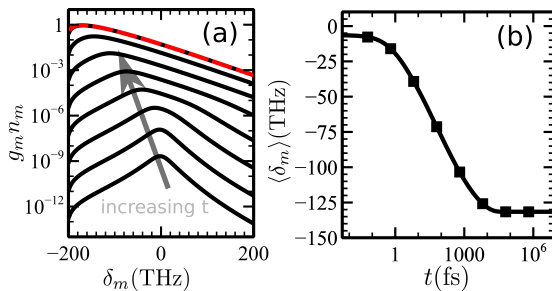
PDRA position available on organic polaritons

Extra Slides

- 4 Approach to steady state
- 5 Threshold vs temperature
- 6 Threshold vs spot size
- 7 Beyond semiclassics
- 8 More oscillations
- 9 Toy problem – two bosonic modes
- 10 Polariton spectral weight

Time evolution

- Initial state: excited molecules
 - Initial emission, follows gain peak
 - Thermalisation by repeated absorption

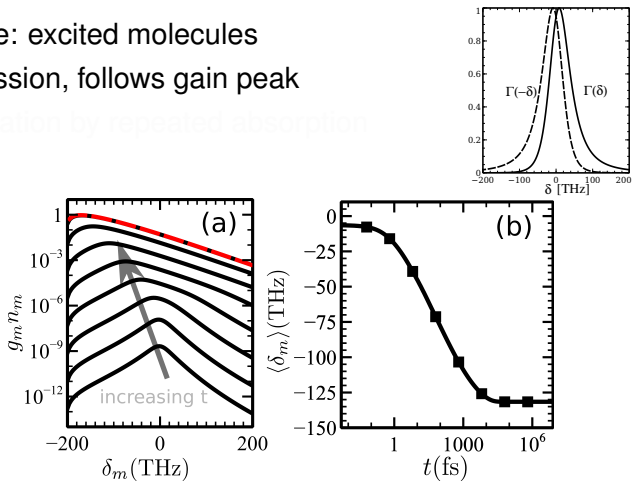


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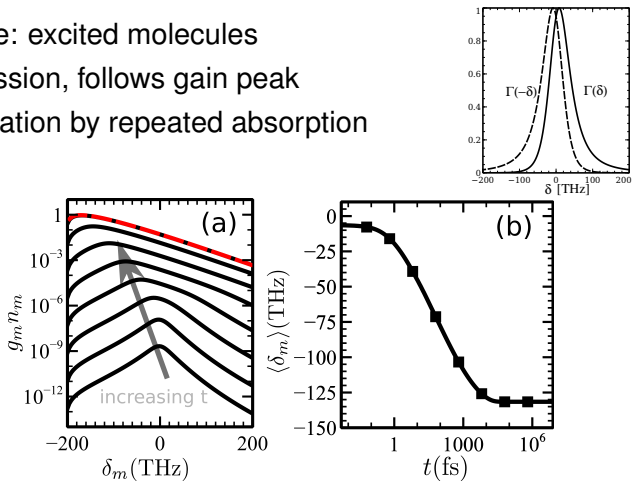
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Threshold condition

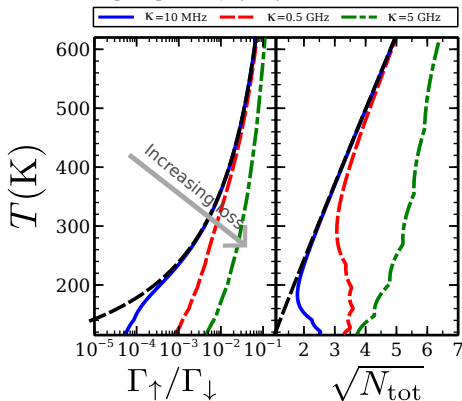
$$\text{Use: } \max[n_m] = 1/(\beta\epsilon) \quad \rightarrow \quad k_B T_c = \sqrt{6/\pi^2\epsilon}\sqrt{N}.$$

- Pump rate (Laser)
- Critical density (condensate)

- Thermal at low κ /high temperature
- High loss, κ competes with $\Gamma(\pm\delta_0)$
- Low temperature, $\Gamma(\pm\delta_0)$ shrinks

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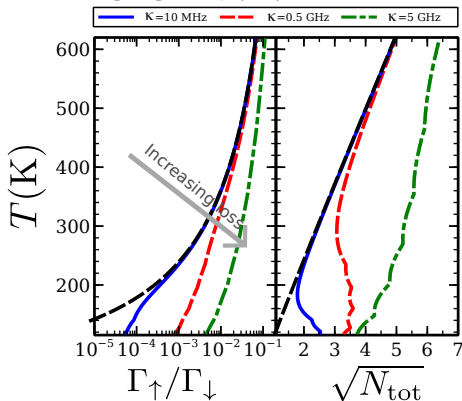
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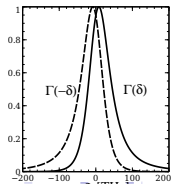


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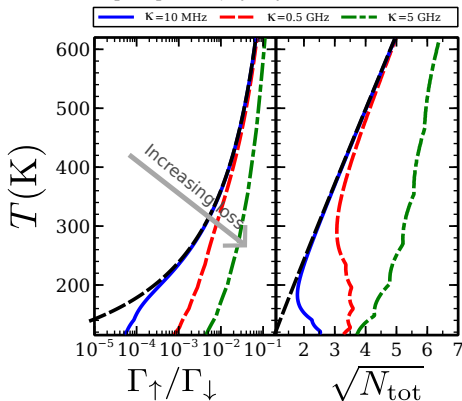
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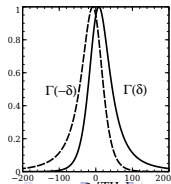
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- Lasing threshold, dependence on spot size.

- ▶ Equilibrium: $\mu = \delta_c$

- Gives $\Gamma_+(r=0) = \Gamma_+ e^{2\delta_c}$

- Dependence on ω_c — experimental spectrum

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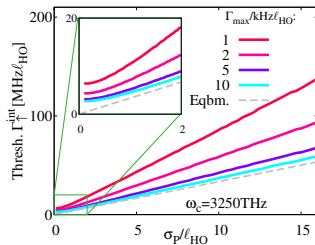
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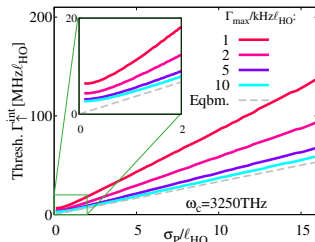


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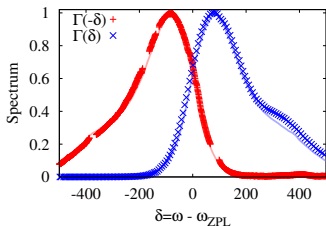
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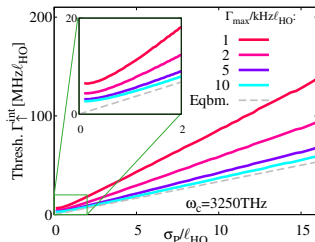
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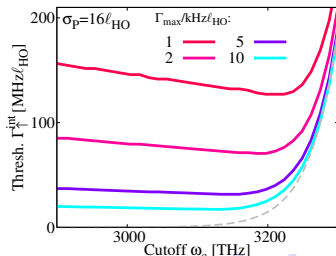
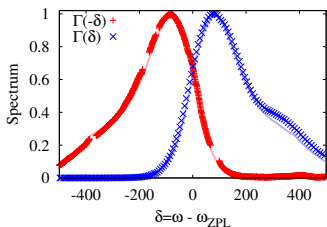
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Quantum model, linewidth

Full Master equation:

$$\dot{\rho} = -i[H_0, \rho] - \frac{\kappa}{2} \mathcal{L}[\psi] - \sum_{\alpha} \left[\frac{\Gamma_{\uparrow}}{2} \mathcal{L}[\sigma_{\alpha}^{+}] + \frac{\Gamma_{\downarrow}}{2} \mathcal{L}[\sigma_{\alpha}^{-}] \right] \\ - \sum_{\alpha} \left[\frac{\Gamma(\delta = \omega - \epsilon)}{2} \mathcal{L}[\sigma_{\alpha}^{+} \psi] + \frac{\Gamma(-\delta = \epsilon - \omega)}{2} \mathcal{L}[\sigma_{\alpha}^{-} \psi^{\dagger}] \right]$$

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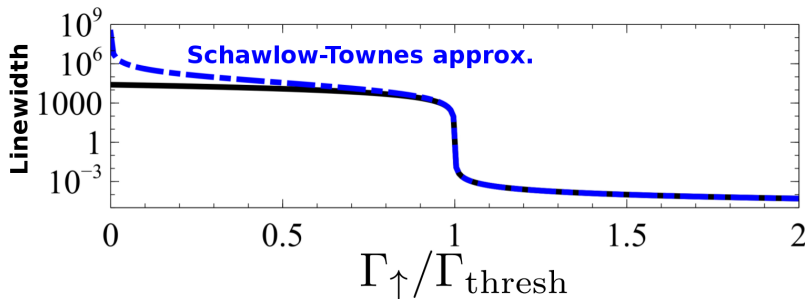
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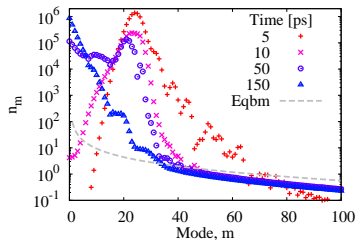
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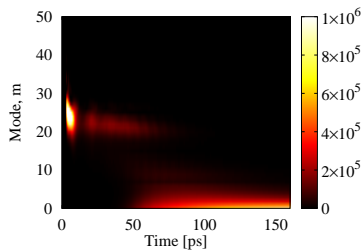
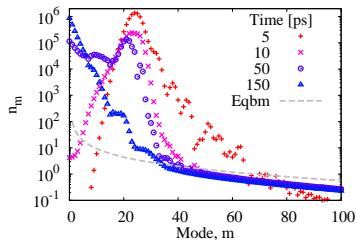
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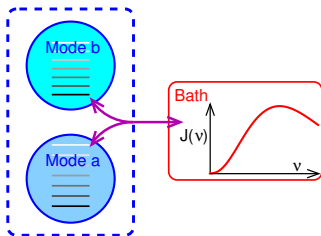
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Toy problem: two bosonic modes

- Basic problem: Emission from thermal bath



$$H = \omega_a \hat{\psi}_a^\dagger \hat{\psi}_a + \omega_b \hat{\psi}_b^\dagger \hat{\psi}_b + H_{\text{Bath}} \\ + (\varphi_a^* \hat{\psi}_a^\dagger + \varphi_b^* \hat{\psi}_b^\dagger) \sum_i g_i \hat{c}_i + \text{H.c.}$$

Toy problem: naïve solutions

Two “expected” behaviours:

- At resonance: “weak lasing” — coupling to bath dominates

$$\frac{d}{dt}\rho = \Gamma^\downarrow \mathcal{L}[\varphi_a \hat{\psi}_a + \varphi_b \hat{\psi}_b] + \Gamma^\uparrow \mathcal{L}[\varphi_a^* \hat{\psi}_a^\dagger + \varphi_b^* \hat{\psi}_b^\dagger]$$

- Far from resonance: pointer states are eigenstates

$$\frac{d}{dt}\rho = \sum_{i=a,b} \Gamma_i^\downarrow \mathcal{L}[\hat{\psi}_i] + \Gamma_i^\uparrow \mathcal{L}[\hat{\psi}_i^\dagger]$$

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- How does crossover work?
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Toy problem: exact solution

- Solve via Laplace transform. Find $F_{ij}(t) = \langle \hat{\psi}_i^\dagger(t) \hat{\psi}_j(t) \rangle$

- Steady state:

- Time evolution —

$$F_{ab}(t) \sim \exp(-\alpha \Delta^2 t)$$

- Always some coherence

- (individual always wrong)

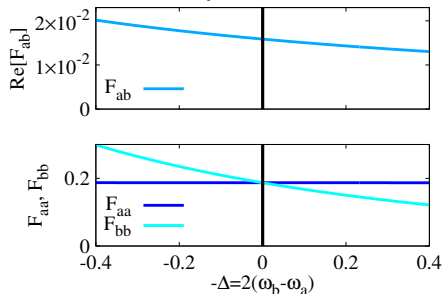
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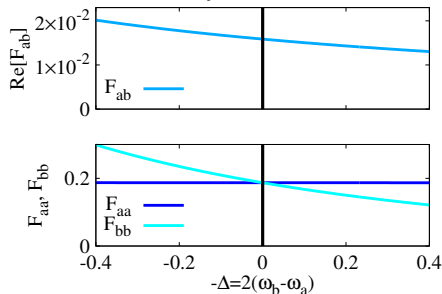


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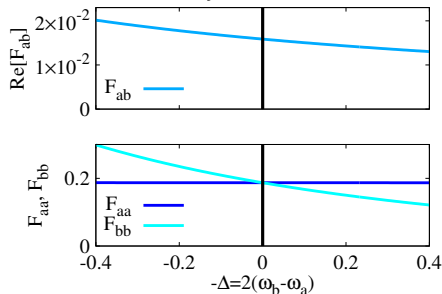
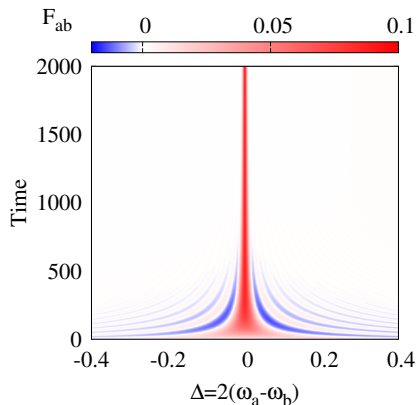
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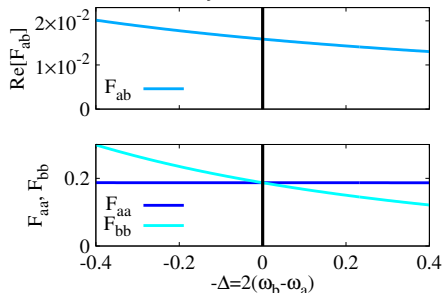
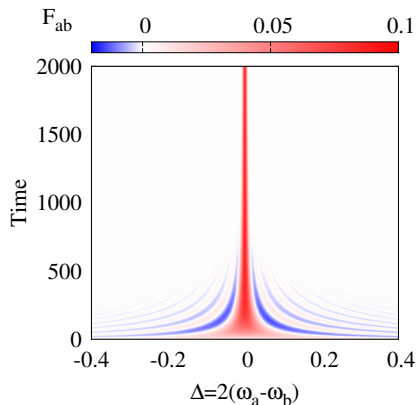
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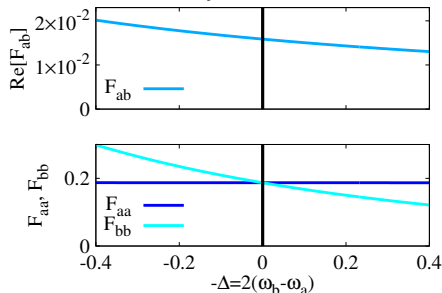
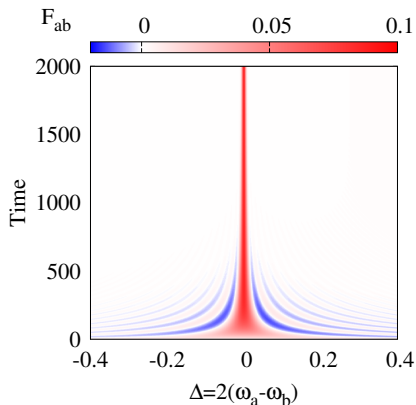


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Toy problem: Bloch-Redfield theory

Unsecularised Bloch-Redfield theory:

$$\begin{aligned} \partial_t \rho = & -i[\hat{H}, \rho] + \sum_{ij} L_{ij}^{\downarrow} \varphi_i^* \varphi_j \left(2\hat{\psi}_j \rho \hat{\psi}_i^{\dagger} - [\rho, \hat{\psi}_i^{\dagger} \hat{\psi}_j]_+ \right) \\ & + \sum_{ij} L_{ij}^{\uparrow} \varphi_i^* \varphi_j \left(2\hat{\psi}_j^{\dagger} \rho \hat{\psi}_i - [\rho, \hat{\psi}_i \hat{\psi}_j^{\dagger}]_+ \right). \end{aligned}$$

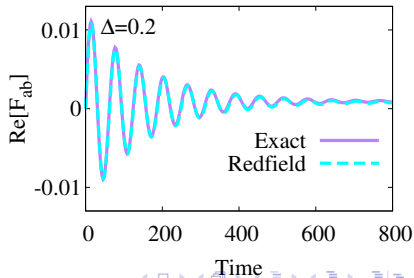
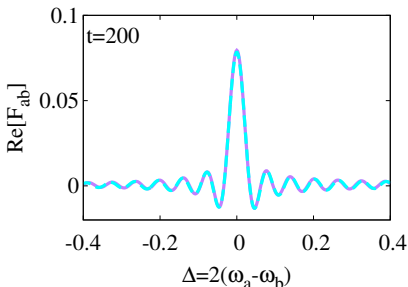
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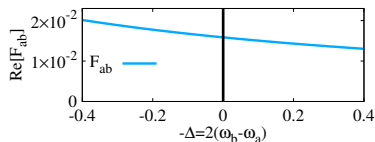
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Toy problem: Secularisation

- Secularisation (in eigenbasis of \hat{H}): $L_{ij}^{\uparrow,\downarrow} \rightarrow L_{ij}^{\uparrow,\downarrow} \delta_{ij}$



- Leads to $F_{ab}(t \rightarrow \infty) = 0$. Exact:

- Secularisation often invoked to cure negative eigenvalues of $L_{ij}^{\uparrow,\downarrow}$

- Check stability: consider $f = (F_{aa}, F_{bb}, \Re[F_{ab}], \Im[F_{ab}])$

$$\partial_t f = -Mf + f_0$$

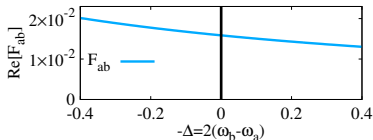
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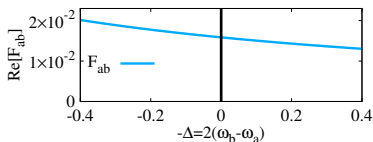
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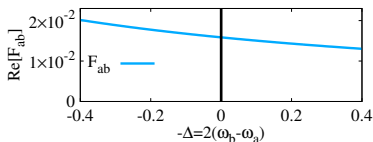
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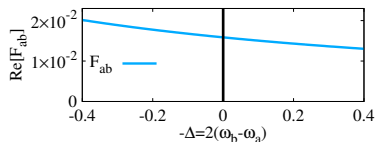
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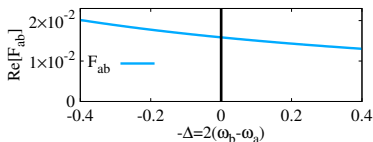
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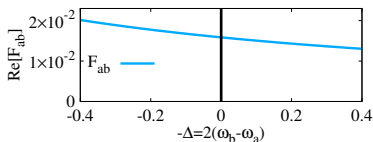
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- Secularisation (in eigenbasis of \hat{H}): $L_{ij}^{\uparrow,\downarrow} \rightarrow L_{ii}^{\uparrow,\downarrow} \delta_{ij}$

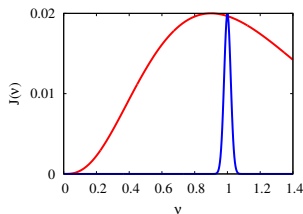
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 - Non-positivity of density matrix,
 - Unstable (unbounded growth).
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$$\partial_t \mathbf{f} = -\mathbf{M} \mathbf{f} + \mathbf{f}_0$$

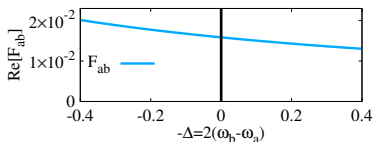
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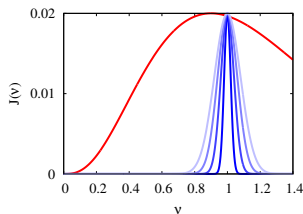
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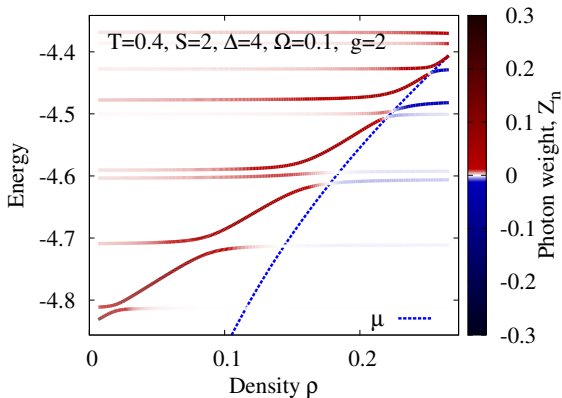
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Polariton spectrum: photon weight

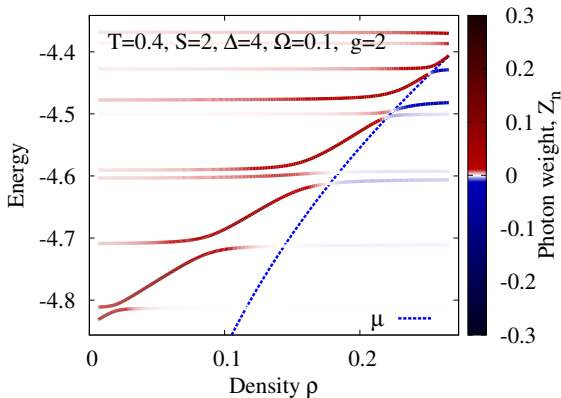


• What is nature of polariton mode?

$$G^R(t) = -i \langle \psi^\dagger(t) \psi(0) \rangle, \quad G^R(\omega) = \sum_n \frac{Z_n}{\omega - \omega_n}$$

[Cwik *et al.* EPL '14]

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