

Collective behaviour and driven-dissipative systems

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1413-2013

CUNY, April 2015

Acknowledgements

GROUP (&ALUMNI):



COLLABORATORS: Fazio (Pisa & CQT), Schiro (CNRS), Tureci (Princeton), Eastham (TCD), Lovett (St Andrews).

FUNDING:



The Leverhulme Trust

Collective behaviour and driven-dissipative systems

1 Nonequilibrium quantum matter

2 Collective behaviour in driven–dissipative systems

- Transverse field Ising
- Rabi-Hubbard model

3 Collective dissipation

- Coupled qubit-cavity systems
- Bath induced coherence

Nonequilibrium quantum matter

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2 Collective behaviour in driven–dissipative systems

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Driven systems

Open quantum system

$$\partial_t \rho = -i[\hat{H}, \rho] + \sum_i \kappa_i \mathcal{L}[X_i], \quad \mathcal{L}[X_i] = 2X_i \rho X_i^\dagger - X_i^\dagger X_i \rho - \rho X_i^\dagger X_i$$

Need **drive** to balance loss

$$\hat{H} \rightarrow \hat{H} + \tilde{V} \cos(\Omega t)$$

Driven systems

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- ① External **coherent** drive:

$$\hat{H} \rightarrow \hat{H} + \hat{V} \cos(\Omega t)$$

- $\hat{H} = e^{-i\hat{H}_0 t} H_0 e^{i\hat{H}_0 t} - g \hat{N}$
- Neglect fast $e^{i\Omega t}$ terms — fast
- Rotating frame — breaks detailed balance with bath

Driven systems

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- ① External **coherent** drive:

$$\tilde{\hat{H}} = \begin{pmatrix} h_0 & v_{01} \cos(\Omega t) & 0 & \dots \\ v_{01}^\dagger \cos(\Omega t) & h_1 & v_{12} \cos(\Omega t) & \dots \\ 0 & v_{12}^\dagger \cos(\Omega t) & h_2 & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

$$\sim \tilde{H} = e^{-i\Omega t} H e^{i\Omega t} - g \tilde{N}$$

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Rotating frame — break detailed balance with bath

Driven systems

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Need **drive** to balance loss

- ① External **coherent** drive:

$$\tilde{\hat{H}} \simeq \begin{pmatrix} h_0 & v_{01} & 0 & \dots \\ v_{01}^\dagger & h_1 - \Omega & v_{12} & \dots \\ 0 & v_{12}^\dagger & h_2 - 2\Omega & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

- $\tilde{\hat{H}} = e^{-i\Omega \hat{N}t} \hat{H} e^{i\Omega \hat{N}t} - \Omega \hat{N}$
- Neglect fast $e^{2i\Omega t}$ terms — fast

Rotating frame — leads to steady-state solution

Driven systems

Open quantum system

$$\partial_t \rho = -i[\hat{H}, \rho] + \sum_i \kappa_i \mathcal{L}[X_i], \quad \mathcal{L}[X_i] = 2X_i \rho X_i^\dagger - X_i^\dagger X_i \rho - \rho X_i^\dagger X_i$$

Need **drive** to balance loss

- ① External **coherent** drive:

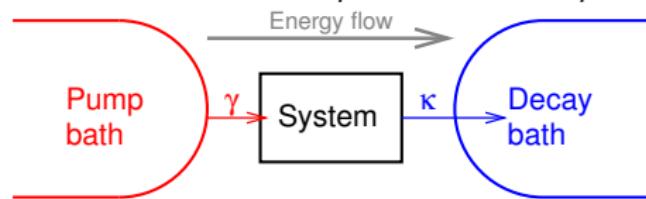
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Non-equilibrium steady state

- ② External **incoherent** drive:

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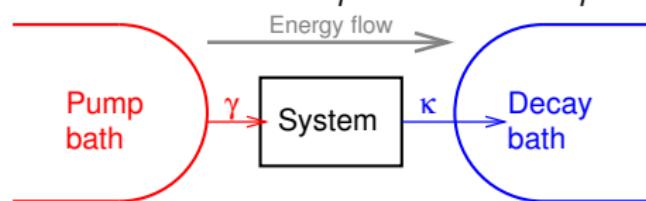
• Energy flow through system

- Not thermodynamics — attractors of dynamics
 - Stationary points — extrema of energy?
 - Non-trivial attractors

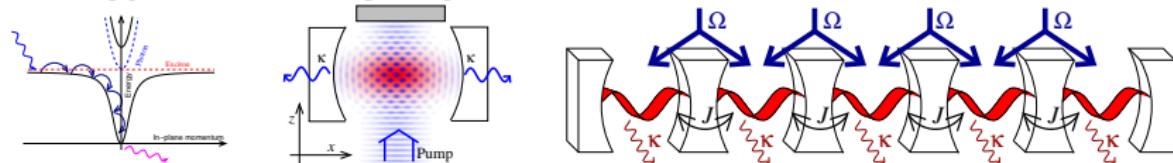
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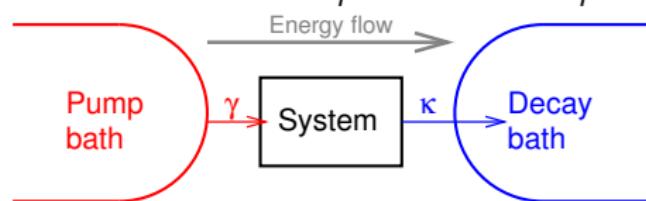


- Non-thermodynamics — attractors of dynamics
 - Stationary points — extrema of energy?
 - Non-adiabatic effects?

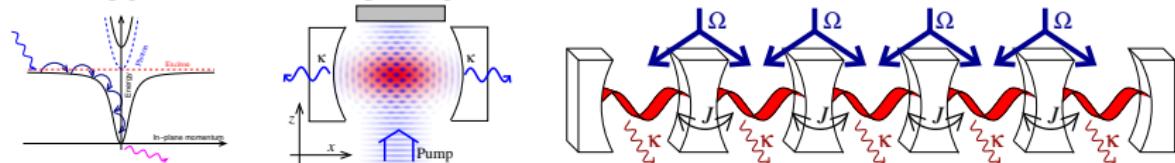
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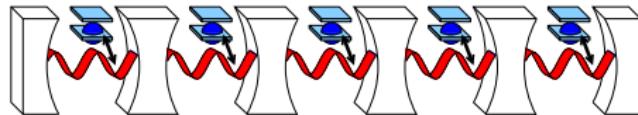
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Coupled cavity arrays

- Control photon dispersion — lattice



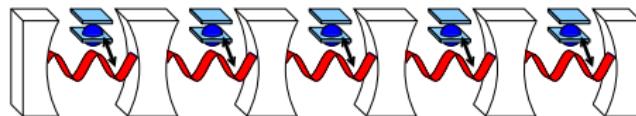
[Hartmann *et al.* Nat. Phys. '06; Greentree *et al.* *ibid* 06; Angelakis *et al.* PRA '07]

• Coupling between two coupled cavity arrays

[X-Bose, Jaynes-Cummings, Rabi, ...]

Coupled cavity arrays

- Control photon dispersion — lattice

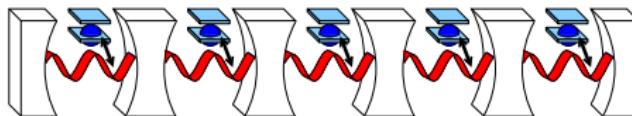


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- X-Hubbard Model, $H = \sum_i H_{X,site} - J \sum_{\langle ij \rangle} \psi_i^\dagger \psi_j$
[X=Bose, Jaynes-Cummings, Rabi, ...]

Coupled cavity arrays

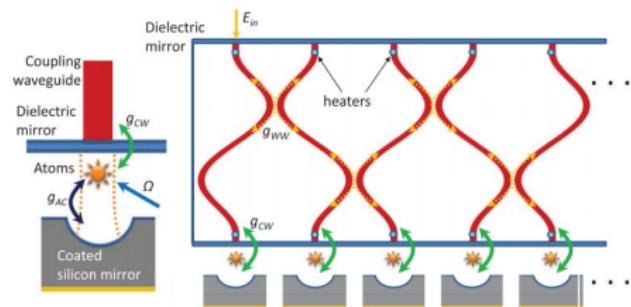
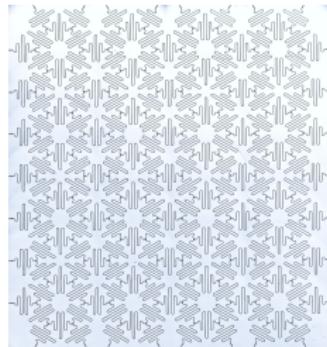
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[Lepert *et al.* NJP '11; APL '13]

[Underwood *et al.* PRA '12; Nat. Phys '12]

Collective behaviour in driven–dissipative systems

1 Nonequilibrium quantum matter

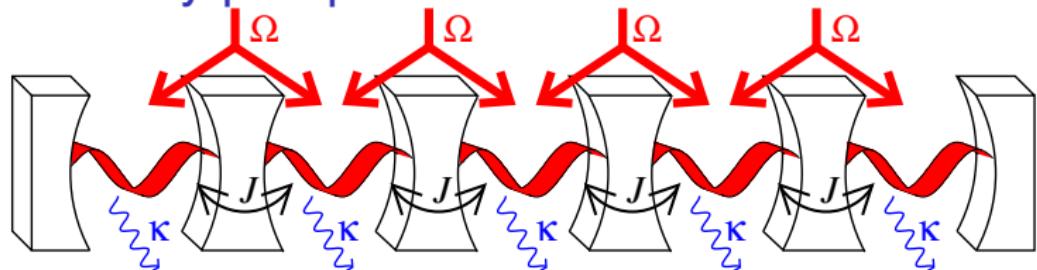
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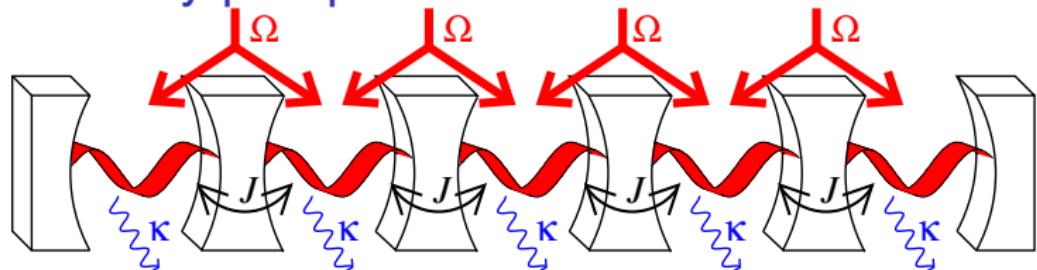
Parametrically pumped BHM



$$H = -\frac{J}{z} \sum_{\langle ij \rangle} \psi_i^\dagger \psi_j + \sum_i \left[\omega_c \psi_i^\dagger \psi_i + U \psi_i^\dagger \psi_i^\dagger \psi_i \psi_i - \Omega (\psi_i^\dagger \psi_{i+1}^\dagger e^{-2i\omega_p t} + \text{H.c.}) \right]$$

[Bardyn & Immamoglu, PRL '12]

Parametrically pumped BHM



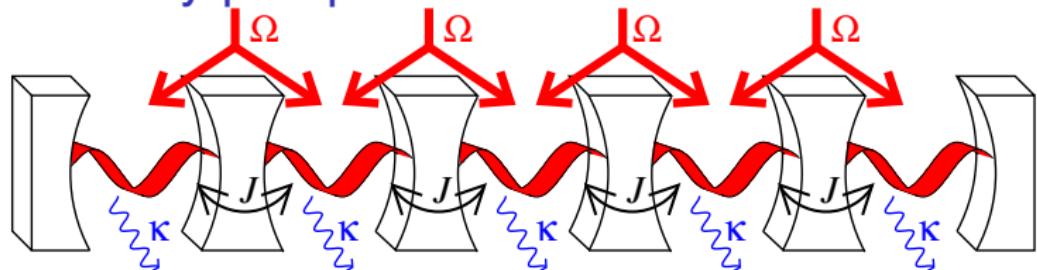
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Rotating frame, blockade approximation, rescale:

$$H = -J \sum \left[\tau_i^+ \tau_{i+1}^- + \tau_{i+1}^+ \tau_i^- + g \tau_i^z + \Delta \left(\tau_i^+ \tau_{i+1}^+ + \tau_{i+1}^- \tau_i^- \right) \right]$$

[Bardyn & Immamoglu, PRL '12]

Parametrically pumped BHM



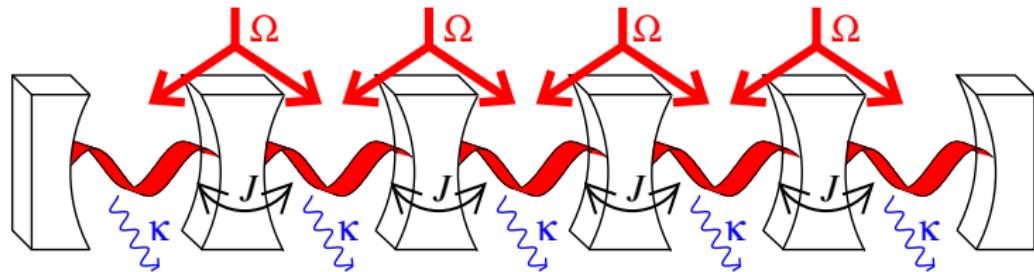
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[Bardyn & Immamoglu, PRL '12]

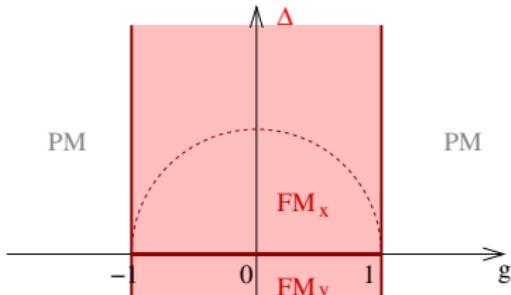
Parametric pumping – equilibrium



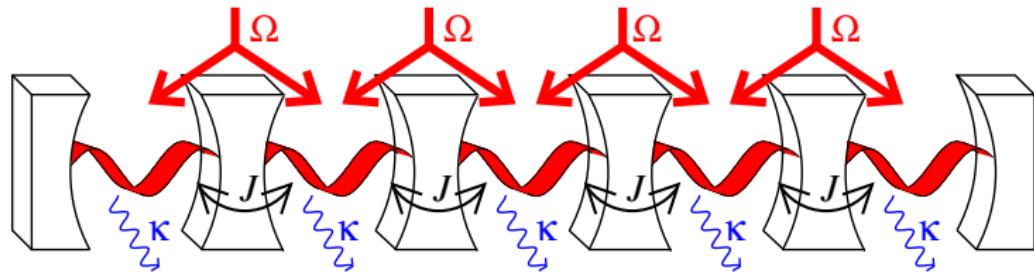
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- Equilibrium – transverse field Ising model
 - ▶ g – transverse field, $g_{\text{crit}} = 1$.

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Parametric pumping – equilibrium



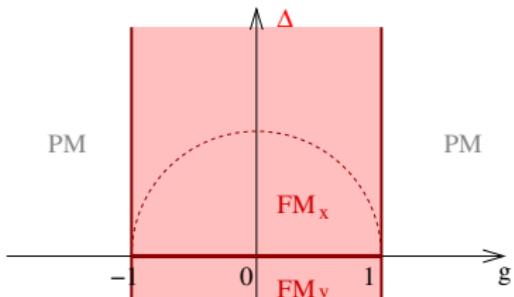
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- Equilibrium – transverse field Ising model

- ▶ g – transverse field, $g_{\text{crit}} = 1$.
- ▶ Δ – anisotropy.

$\Delta = 0$: XY, $|\Delta| > 0$: Ising (X,Y).

[Bardyn & Immamoglu, PRL '12]



Parametric pumping – open system

$$H = -J \sum \left[\tau_i^+ \tau_{i+1}^- + \tau_{i+1}^+ \tau_i^- + g \tau_i^z + \Delta (\tau_i^+ \tau_{i+1}^+ + \tau_{i+1}^- \tau_i^-) \right]$$

$$\partial_t \rho = -i[H, \rho] + \sum_i \kappa \mathcal{L}[\tau_i^-]$$

- Mean-field EOM: $\partial_t \langle \tau_i^\alpha \rangle = F_\alpha(\langle \tau_{i-1}^\beta \rangle, \langle \tau_i^\beta \rangle, \langle \tau_{i+1}^\beta \rangle)$

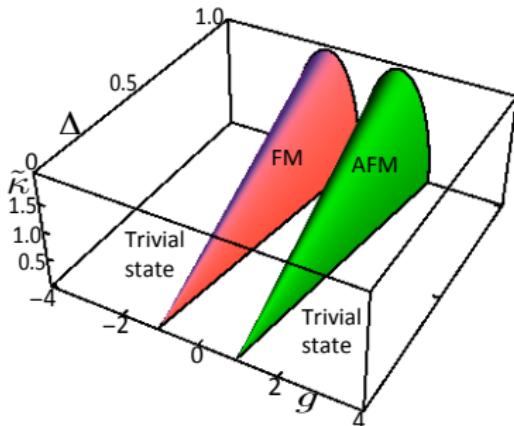
⇒ Dynamical instabilities, linear stability

Parametric pumping – open system

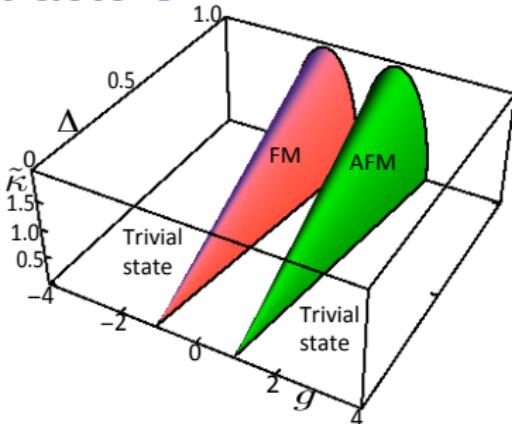
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- Dynamical attractors, linear stability:



Why AFM/FM attractors



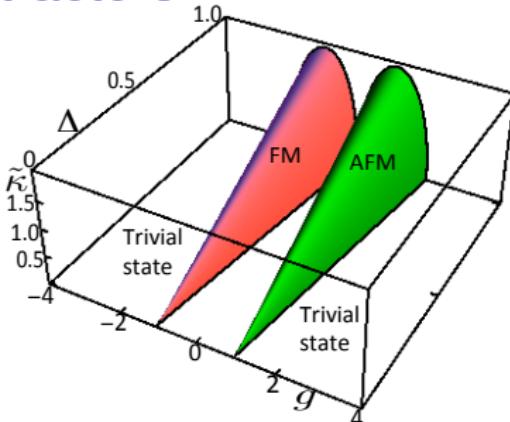
• Linear stability, fluctuation $\sim \exp(-i\omega_t + ik)$ Linear stability

$$\omega_k = -i\kappa + 2J/\sigma^2 + 2g\cos k + (1 - \delta^2)\cos^2 k$$

- $g < -1$, Dissipation matches ground state
 - Most unstable mode, $k=0$
- $g > +1$, Dissipation matches max energy
 - Most unstable mode, $k=\pi$

[Joshi, Nissen, Keeling, PRA '13]

Why AFM/FM attractors



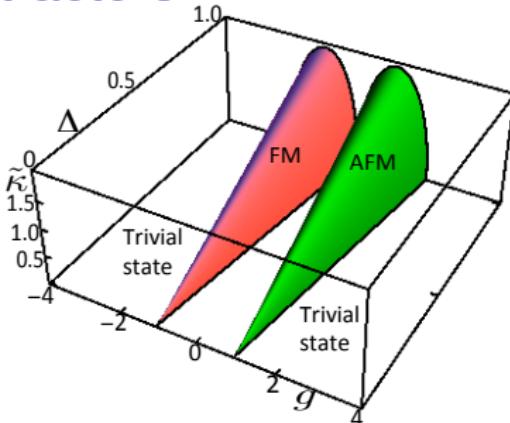
- Linear stability, fluctuation $\sim \exp(-i\nu_k t + ikr_i)$ Linear stability

$$\nu_k = -i\kappa \pm 2J\sqrt{g^2 + 2g \cos k + (1 - \Delta^2) \cos^2 k}$$

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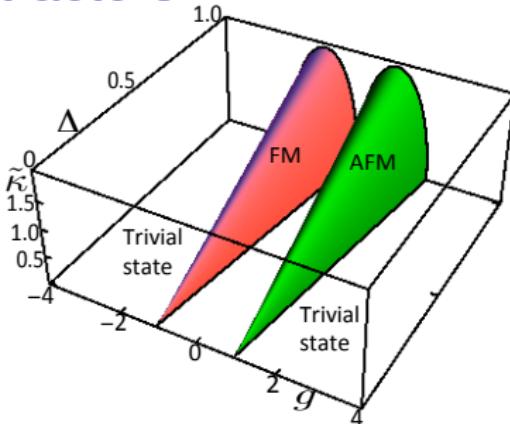


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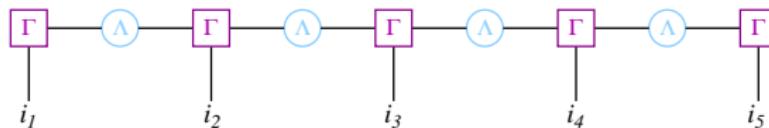
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[Joshi, Nissen, Keeling, PRA '13]

Beyond mean-field

- Matrix-product-operator representation of

$$\rho = \sum_{\{i_1, i_2, \dots, i_N\}} \left(\sum_{\{\alpha_j\}} \Gamma_{1, \alpha_1}^{[1] i_1} \Lambda_{\alpha_1}^{[1]} \Gamma_{\alpha_1, \alpha_2}^{[2] i_2} \dots \Gamma_{\alpha_{N-2}, \alpha_{N-1}}^{[N-1] i_{N-1}} \Lambda_{\alpha_{N-1}}^{[N-1]} \Gamma_{\alpha_{N-1}, 1}^{[N] i_N} \right) \bigotimes_{j=1}^N \tau_j^{i_j}$$



Vidal, White, Schollwöck, et al. Density matrices: [Zwolak & Vidal, PRL '04]

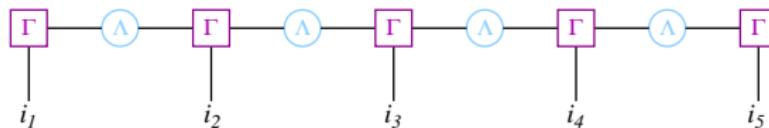
→ broken symmetry → correlations

$$\langle \sigma_x \rangle = 1, \langle \sigma_z \rangle = 0.001$$

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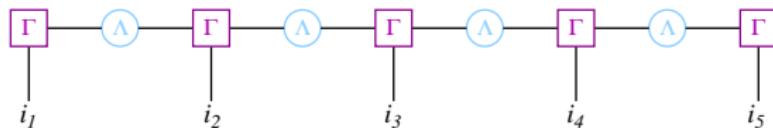
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$$\Delta = 1, \kappa = 0.5J:$$

Beyond mean-field

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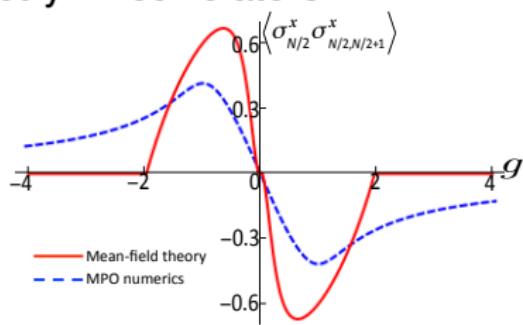
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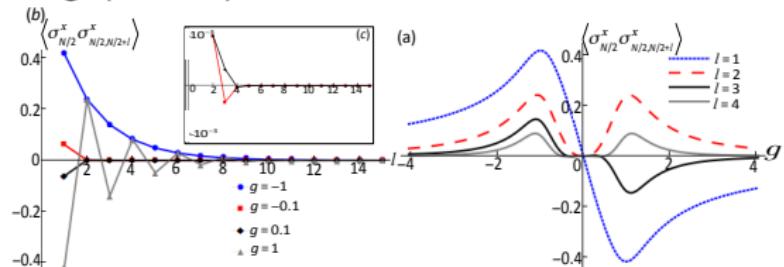
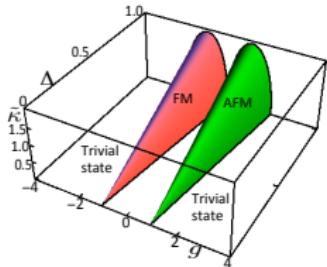
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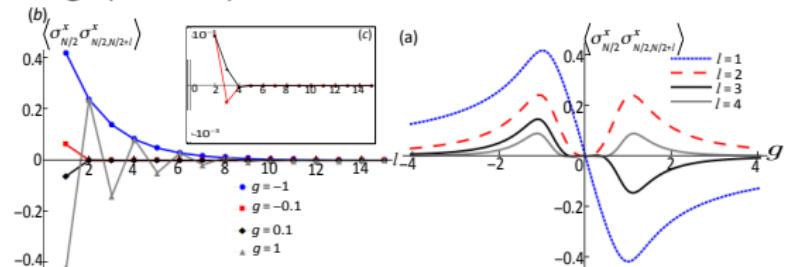
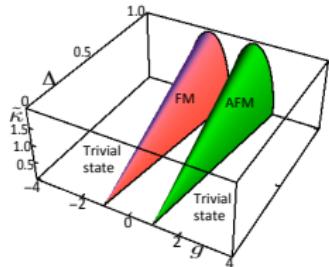
Correlations

- AFM vs FM from sign of g ($\Delta = 1$)



Correlations

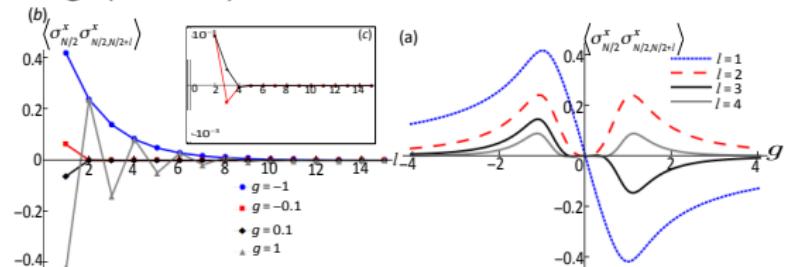
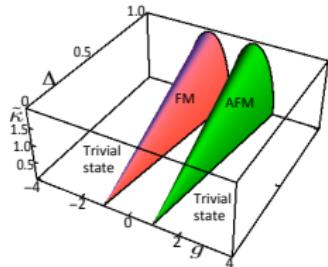
- AFM vs FM from sign of g ($\Delta = 1$)



- $\Delta \rightarrow 0$, Analytic spin-wave,
$$|\langle \tau_i^- \tau_{i+I}^\pm \rangle| \propto \exp(-\xi_c I)$$

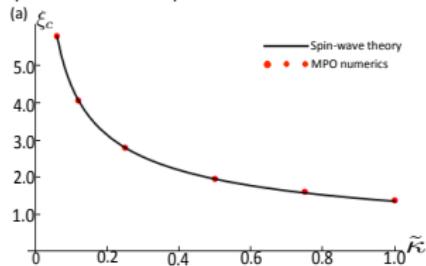
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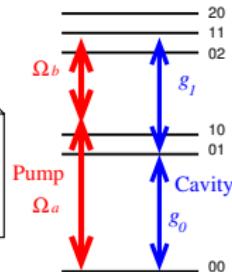
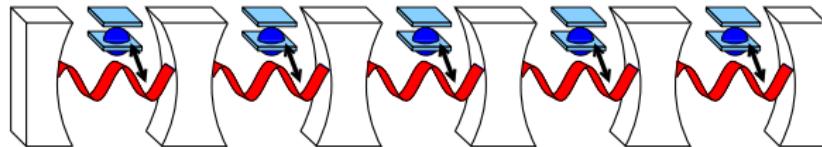


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Rabi Hubbard model



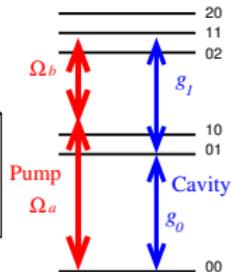
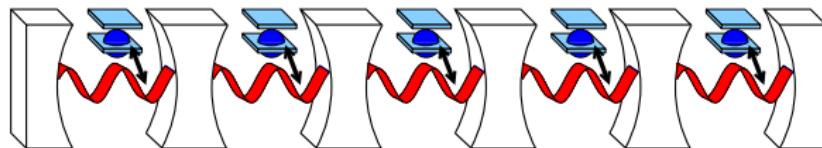
$$H = -J \sum_{\langle ij \rangle} a_i^\dagger a_j + \sum_i h_i^{\text{Rabi}}$$

$$h_i^{\text{Rabi}} = \omega a^\dagger a + \frac{\omega_0}{2} \sigma^z + [a^\dagger (g \sigma^- + g' \sigma^+) + \text{H.c.}]$$

• $\omega = \omega_{\text{cavity}} - \omega_{\text{pump}}$

• g, g' separately tunable

Rabi Hubbard model



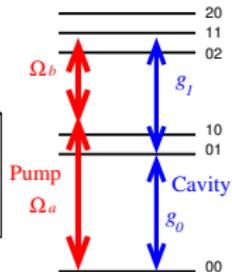
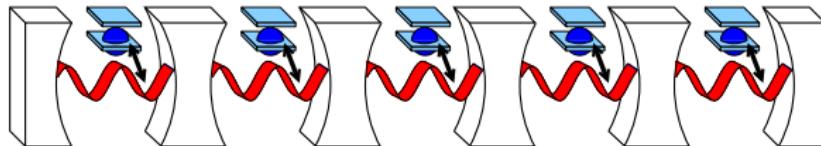
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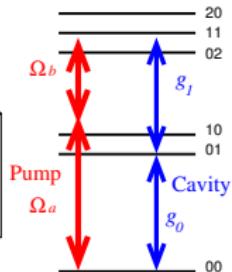
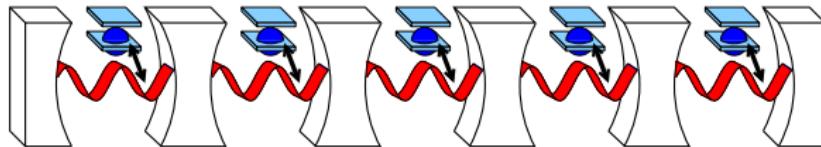


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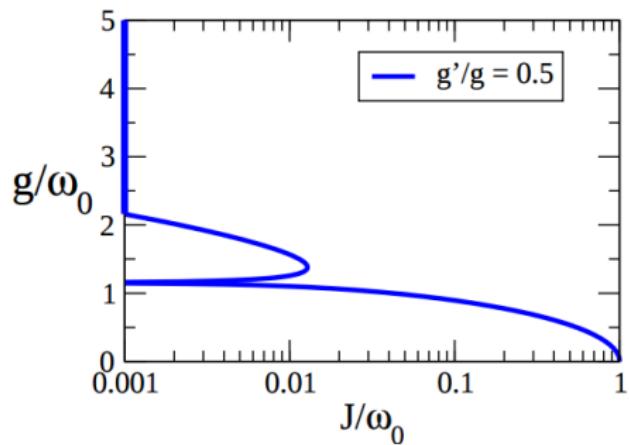
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Rabi Hubbard model – equilibrium



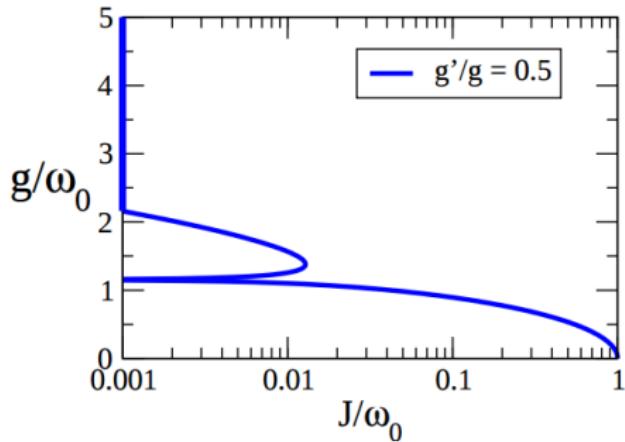
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– Parity Mott lobes

– $g = g'$, never degenerate — never superfluid

[Schiró *et al.* PRL '12]

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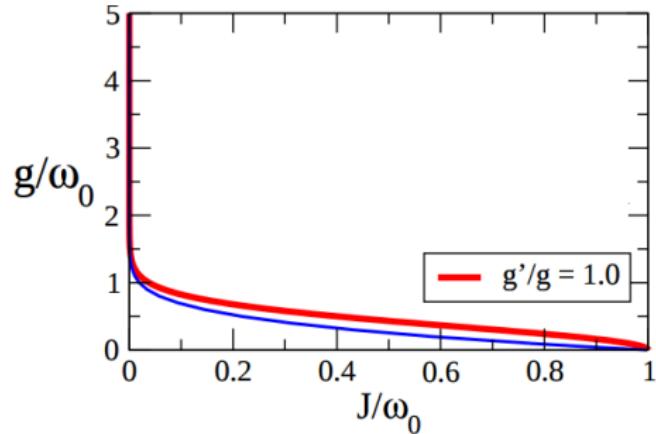
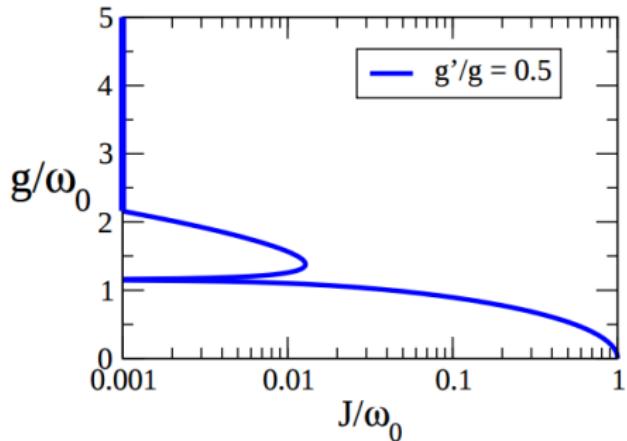


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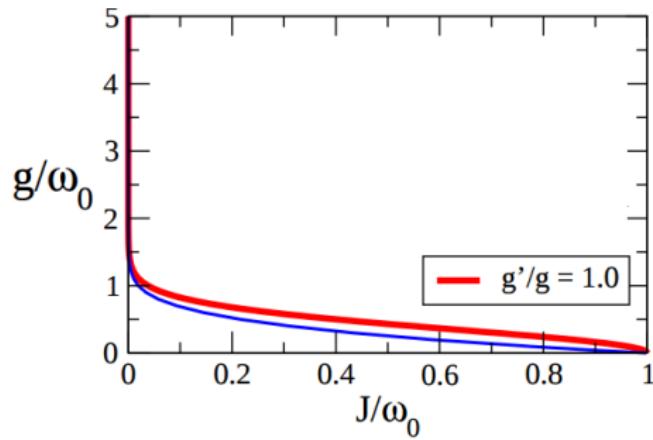
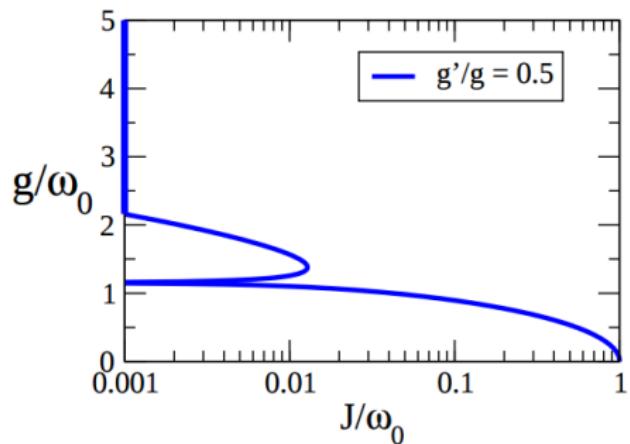


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Driven-dissipative system — linear stability

Mean field theory — still large Hilbert space.

Form of steady-state solution $\psi = \sum_k \delta p_k e^{ikx - i\omega t} + \text{H.c.}$

$\omega = \nu_k$ Eigenvalues of $M = M_0 + k_x M_1$, $k_x = -2J \cos(k)$

Unstable if $\Im[\nu_k] > 0$

- Given J , $|k_x| < 2J$
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Follow [Boité *et al.*, PRA 2014]

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Follow [Boité *et al.*, PRA 2014]

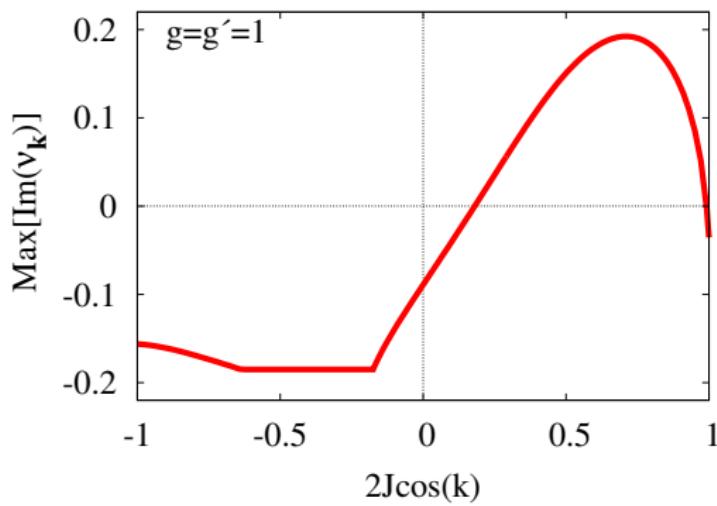
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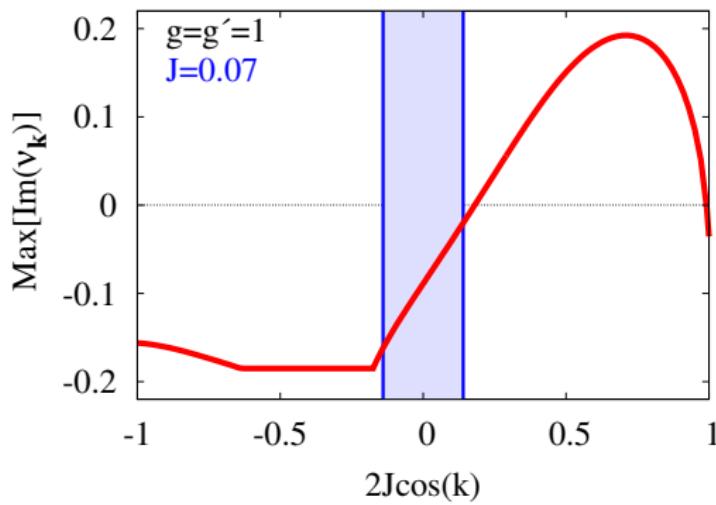
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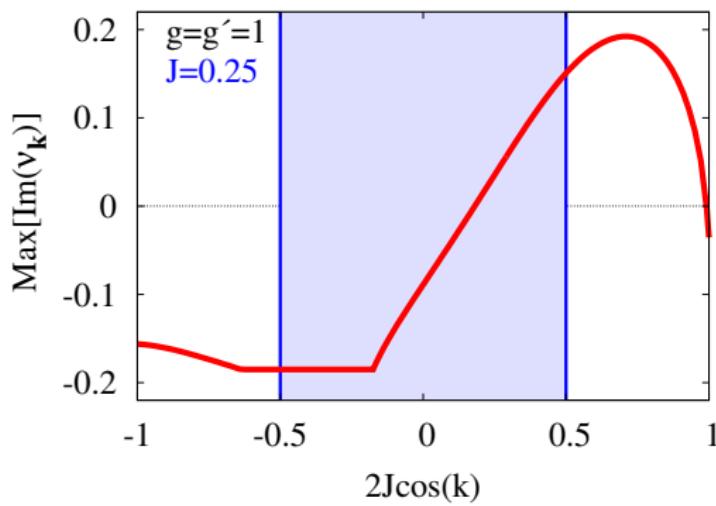
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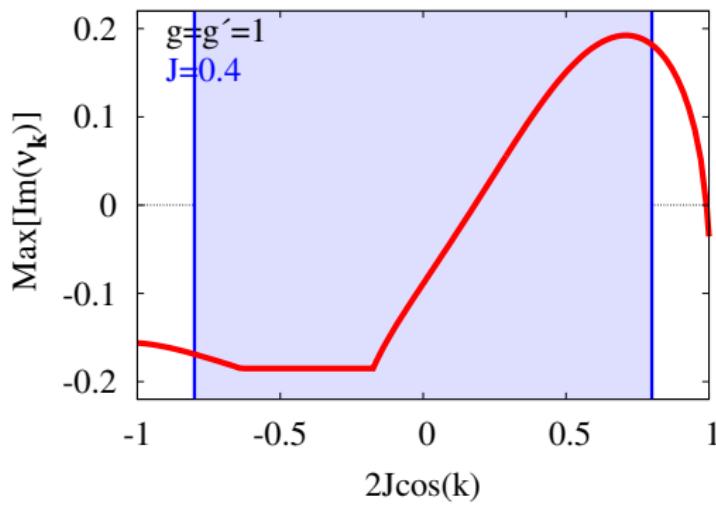
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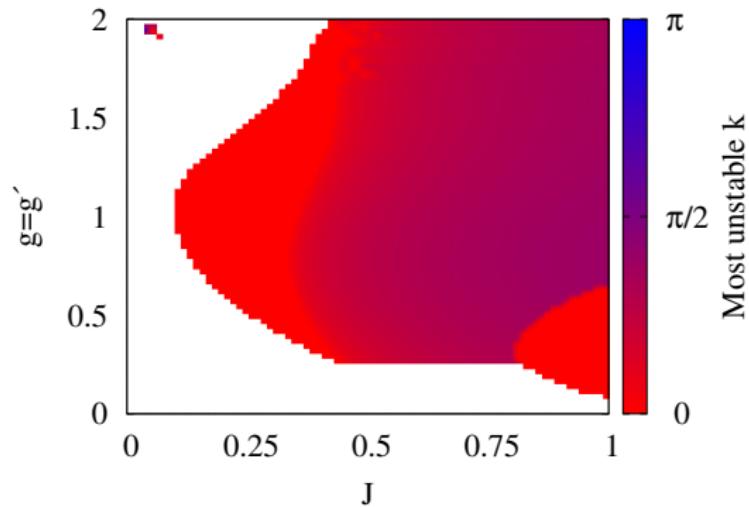
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Rabi-Hubbard model — linear stability

Stability phase diagram:

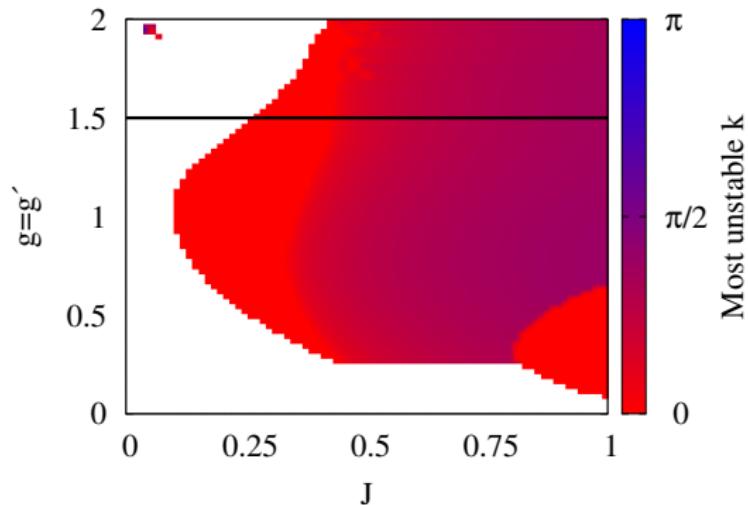


[Schiró *et al.* arXiv:1503.04456]

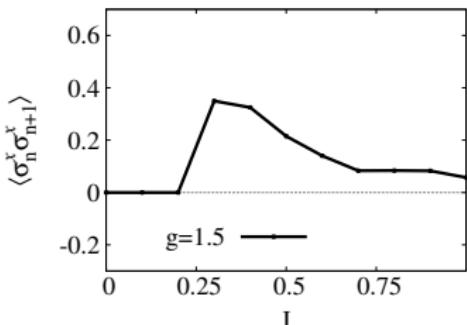
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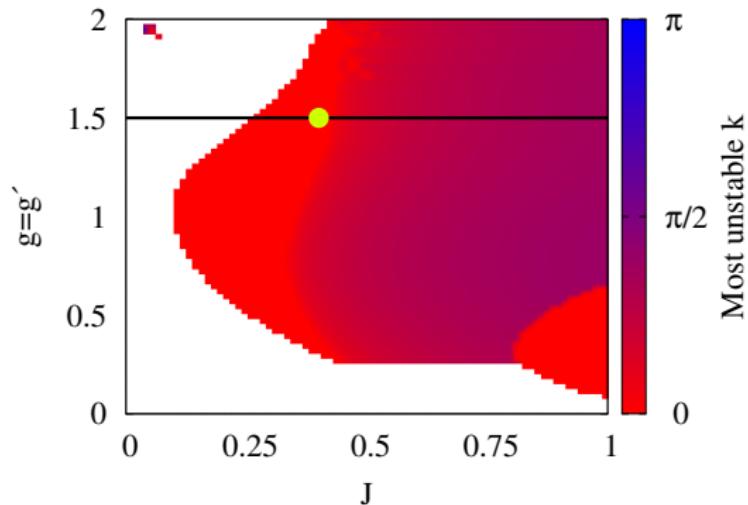
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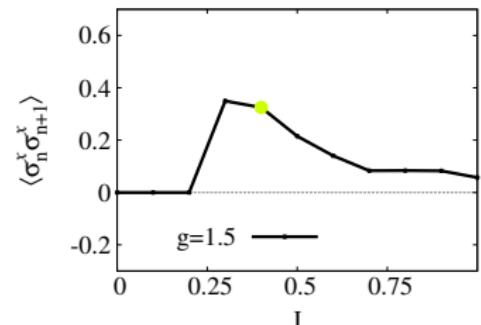
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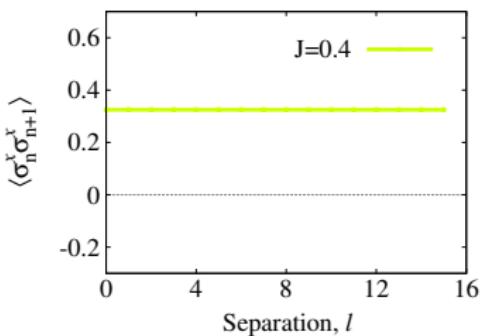
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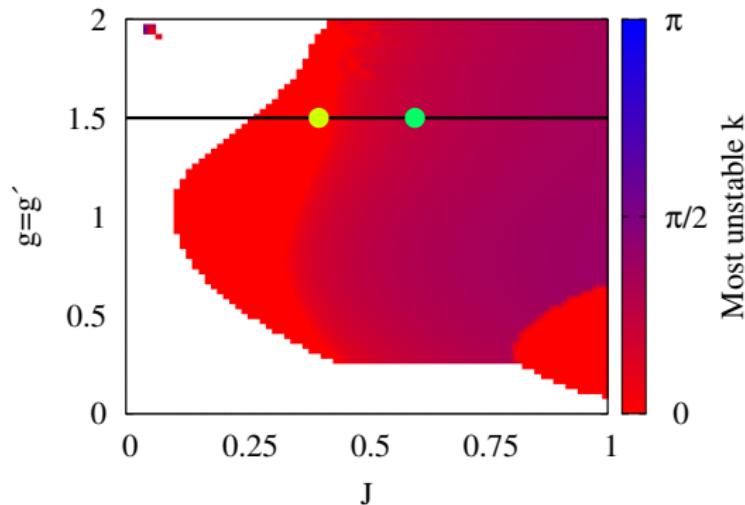
... vs $|i - j| = \uparrow$



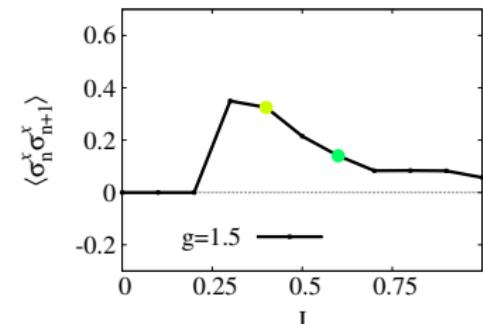
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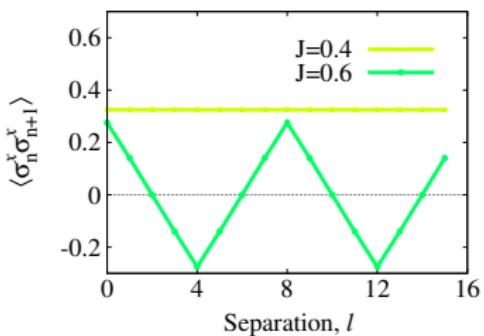
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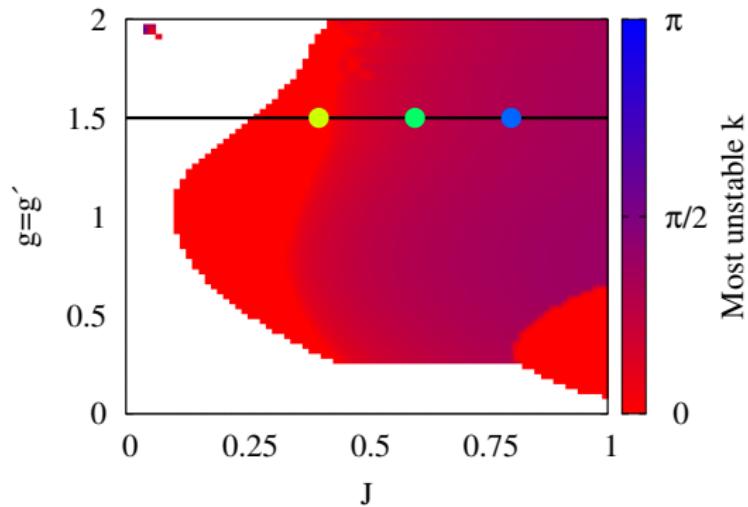
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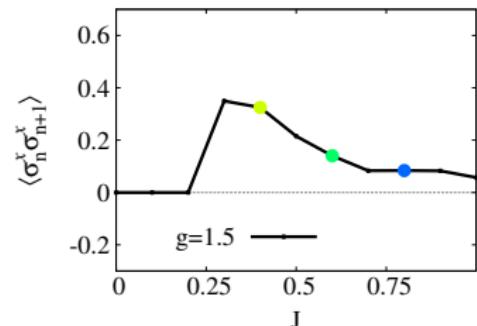
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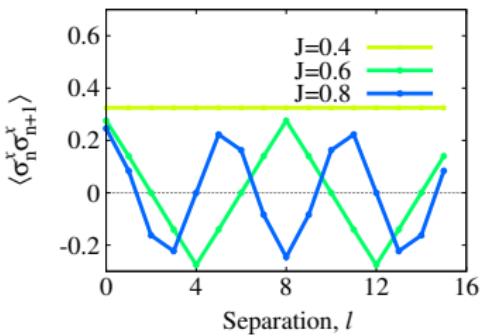
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Linear stability – limit cycles

- If $\nu_k = \pm\nu'_k + i\nu''_k$ at instability \rightarrow Limit Cycle

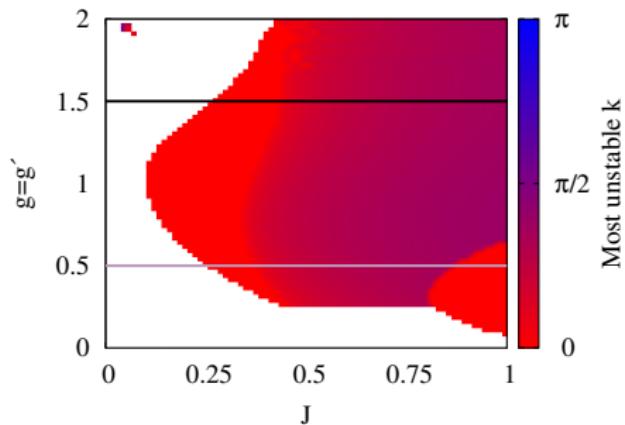
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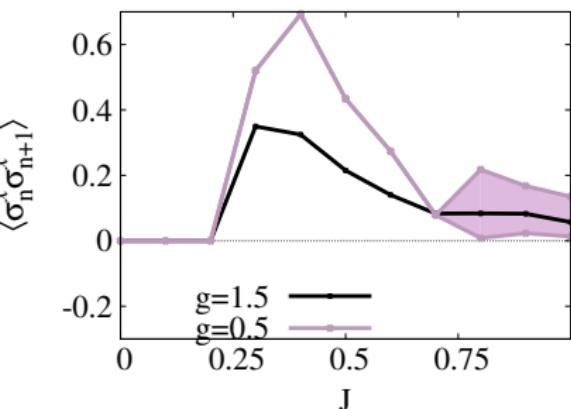
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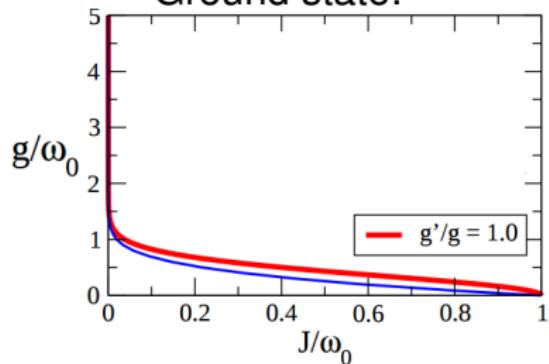
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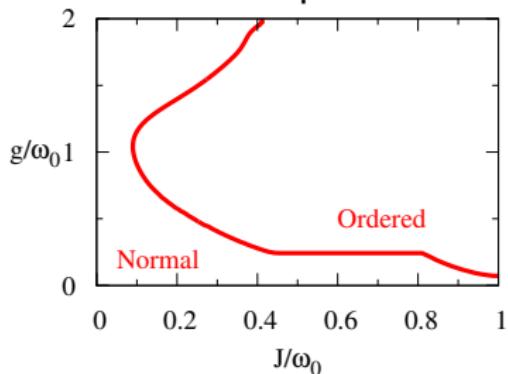
Phase-boundary Effective model

- Compare phase boundaries

Ground state:



Driven dissipative:



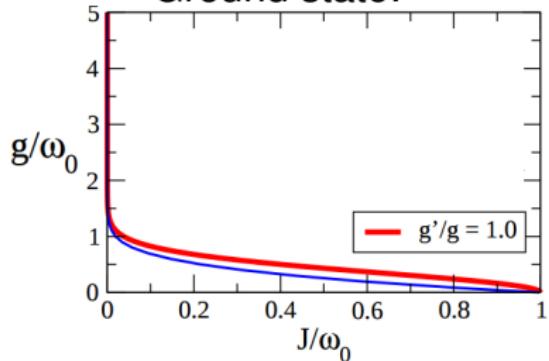
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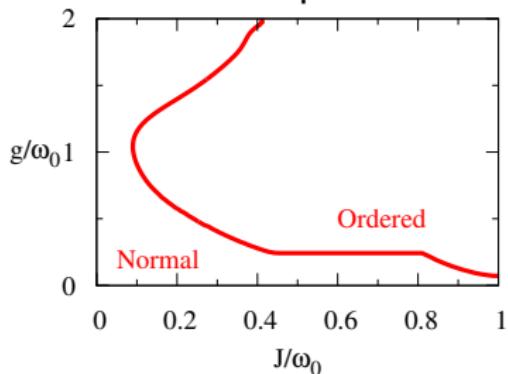
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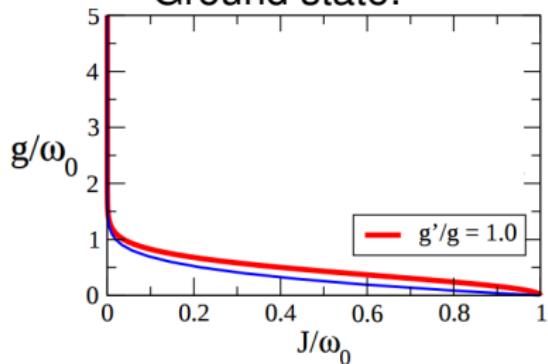


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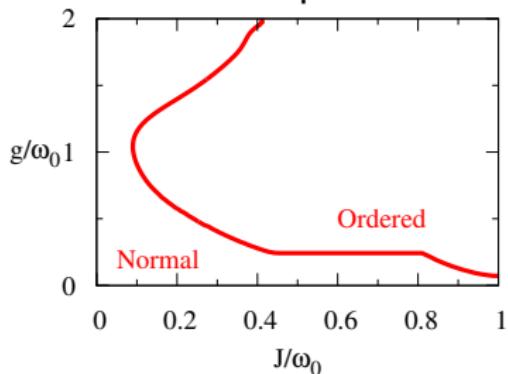
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• Level populations:

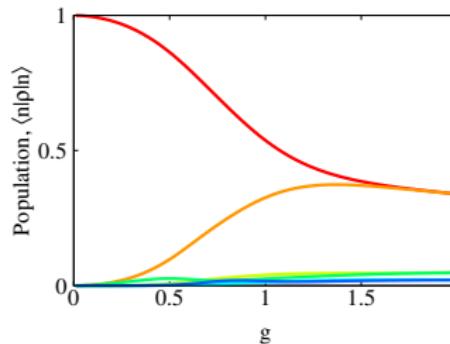
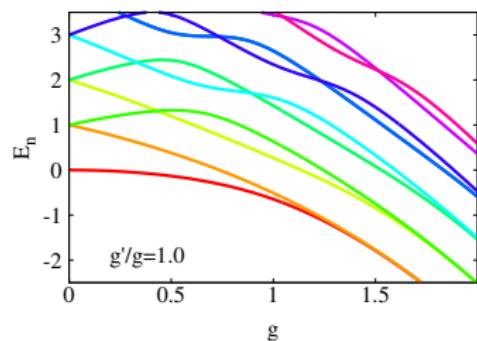
• If $\Delta \sim \omega_0 e^{-2g^2/\omega^2} \ll 1$ $J_{\text{eff}} = \frac{Rg^2}{\omega^2} + \frac{\omega^2}{16g^2}$

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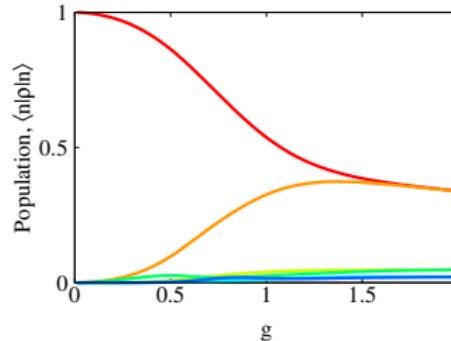
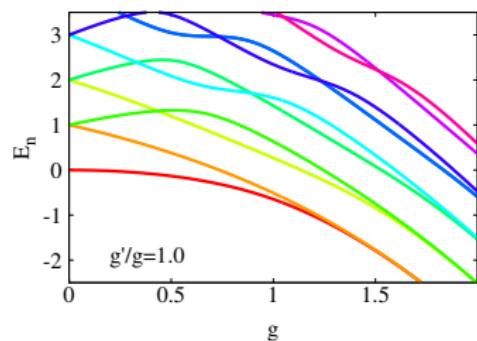


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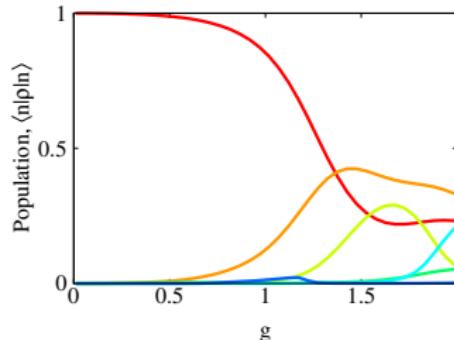
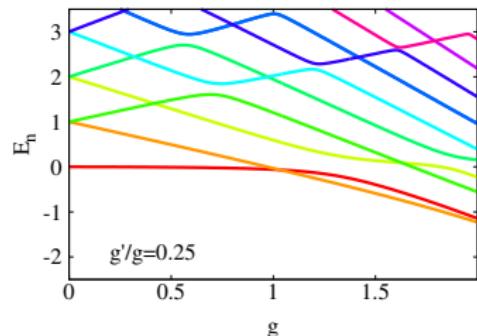


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$g' \neq g$, Level crossings

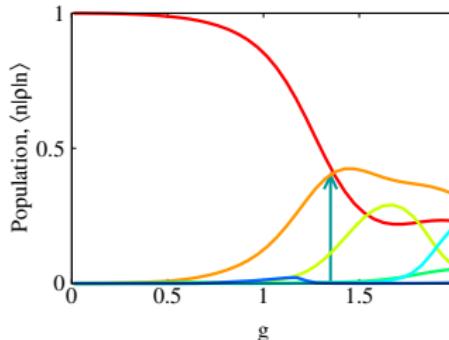
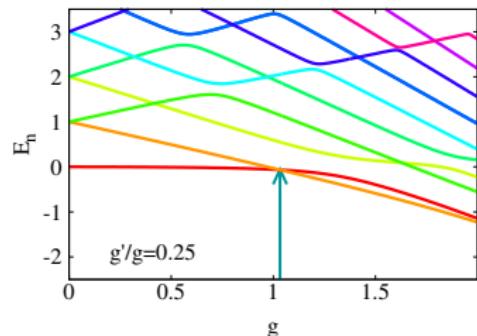
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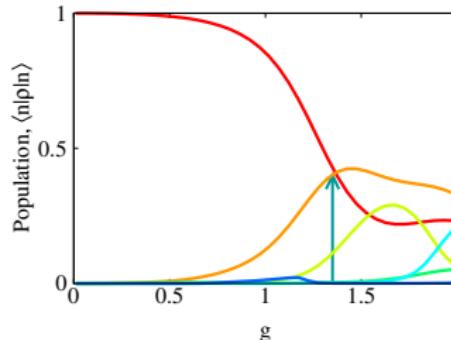
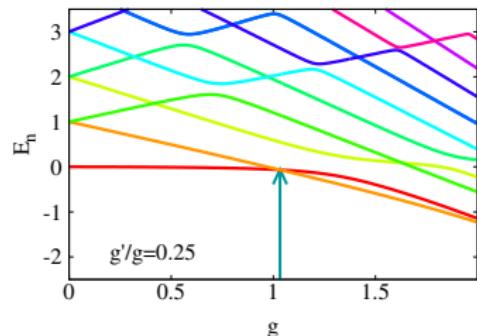
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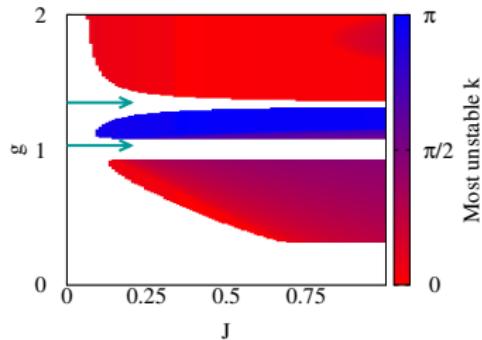
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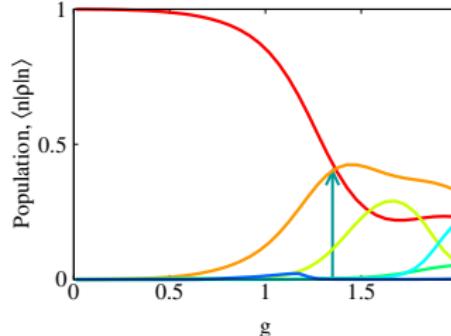
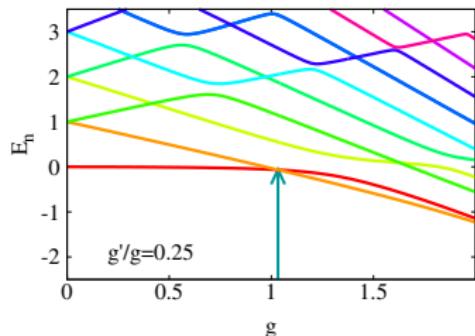


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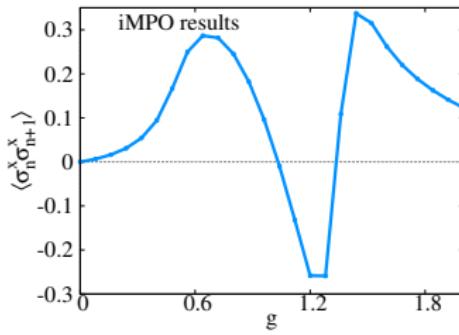
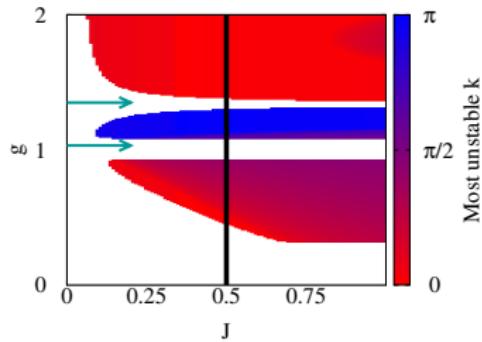


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Collective dissipation

1 Nonequilibrium quantum matter

2 Collective behaviour in driven–dissipative systems

- Transverse field Ising
- Rabi-Hubbard model

3 Collective dissipation

- Coupled qubit-cavity systems
- Bath induced coherence

Collective dephasing

- Real environment is not Markovian
 - ▶ [Carmichael & Walls JPA '73] Requirements for correct equilibrium
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- Cannot assume fixed λ_i/γ_i
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Dicke model linewidth:

$$H = \omega\psi^\dagger\psi + \sum_{i=1}^N \frac{\epsilon_i}{2}\sigma_i^z + g(\sigma_i^+\psi + \text{h.c.})$$
$$+ \sum_i \sigma_i^z \sum_q \gamma_q (b_q^\dagger + b_q) + \sum_q \beta_q b_{iq}^\dagger b_q.$$

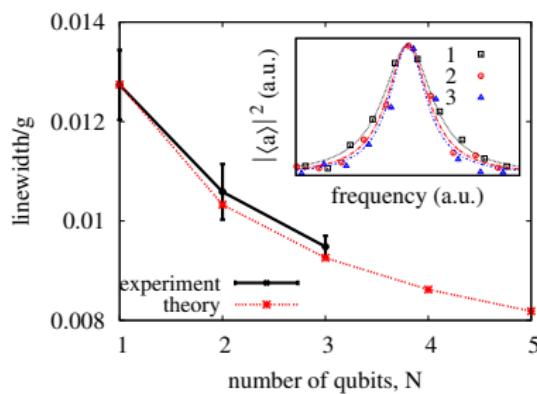
[Nissen, Fink *et al.* PRL '13]

Collective dephasing

- Real environment is not Markovian
 - ▶ [Carmichael & Walls JPA '73] Requirements for correct equilibrium
 - ▶ [Ciuti & Carusotto PRA '09] Dicke SR and emission
- Cannot assume fixed κ, γ
- Phase transition \rightarrow soft modes

Dicke model linewidth:

$$H = \omega\psi^\dagger\psi + \sum_{i=1}^N \frac{\epsilon_i}{2}\sigma_i^z + g(\sigma_i^+\psi + \text{h.c.}) + \sum_i \sigma_i^z \sum_q \gamma_q (b_q^\dagger + b_q) + \sum_q \beta_q b_{iq}^\dagger b_q.$$



[Nissen, Fink et al. PRL '13]

Collective dephasing & weak lasing

- Weak lasing, [Aleiner, Altshuler, Rubo PRB '12] — dissipation selects collective modes.

Key problem:

$$\dot{H} = \omega_a a^\dagger a + \omega_b b^\dagger b + (a^\dagger + b^\dagger) \sum_i g_i c_i + H.c. + H_{\text{int}}$$

Standard picture:

$$\dot{\rho} = -i[H_0, \rho] + \begin{cases} -\gamma_a L[a^\dagger + b^\dagger] + \gamma_b L[a + b] + \dots & \text{degenerate} \\ \gamma_{a,b} L[a^\dagger] + \gamma_{b,a} L[b^\dagger] + \dots & \text{secularised} \end{cases}$$

Exactly solvable problem – which is correct? Consider $\langle a^\dagger b \rangle$

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Bath induced coherence

- Steady state:

- ▶ If $\omega_a = \omega_b$, then $\langle a^\dagger b \rangle = \langle a^\dagger a \rangle = \langle b^\dagger b \rangle$
- ▶ If $\omega_a \neq \omega_b$ then $\langle a^\dagger b \rangle \ll \langle a^\dagger a \rangle, \langle b^\dagger b \rangle$

- ▶ $\langle a^\dagger b \rangle$ strong if $\omega_a \approx \omega_b$

- ▶ Residual coherence – non-flat DoS

- ▶ Requires non-secular master eqn.

- Approaching steady state:

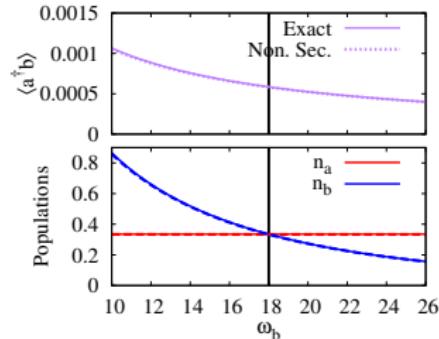
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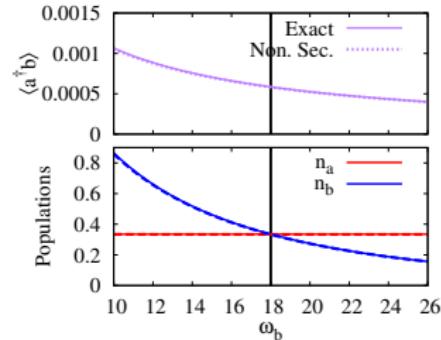
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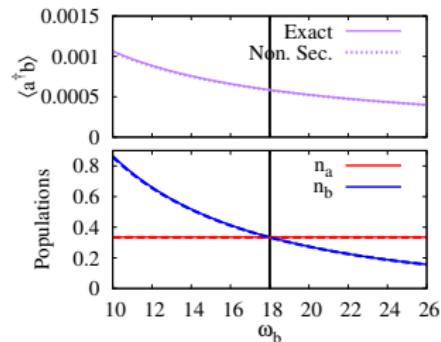
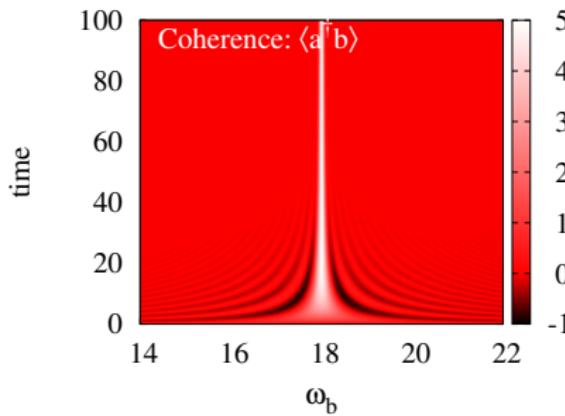
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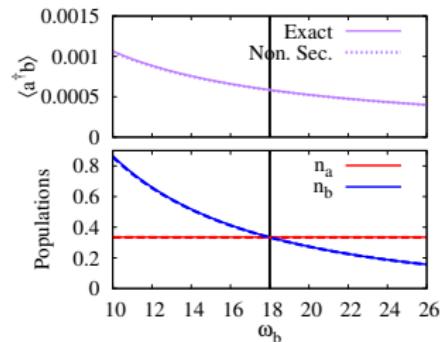
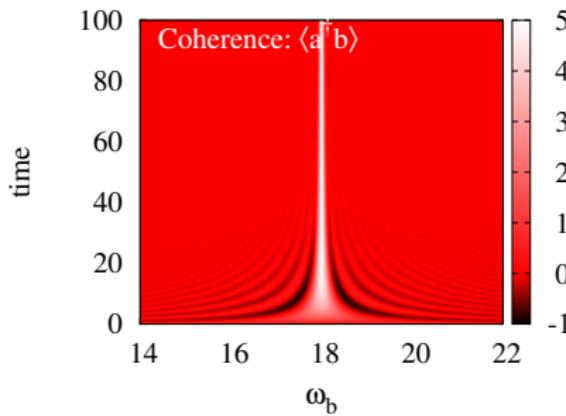


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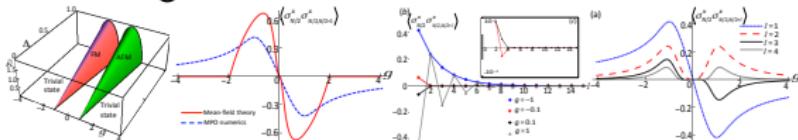


$$\langle a^\dagger b \rangle \simeq \exp \left[-C(\omega_a - \omega_b)^2 t \right]$$

$$+ \langle a^\dagger b \rangle_{t \rightarrow \infty}$$

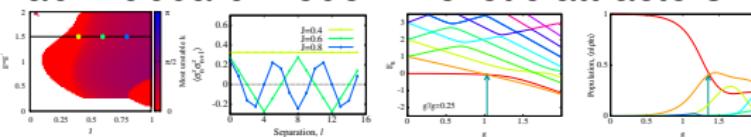
Summary

- Parametric pumping — non-equilibrium “phases” of transverse field Ising model



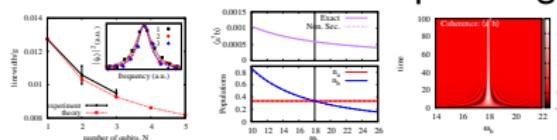
Joshi *et al.* PRA '13

- Rabi Hubbard model — exotic attractors.



Schiró *et al.* arXiv:1503.04456

- Collective effects in dephasing



Nissen *et al.* PRL '13

