

Photon and polariton condensates with organic molecules

Jonathan Keeling

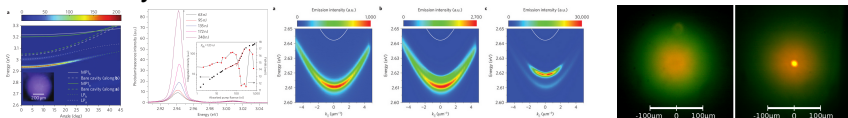


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St Andrews
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Matter-Light coupling with organic molecules

• What & why?



[Kena Cohen and Forrest, Nat. Photon '10; Plumhoff *et al.* Nat. Materials '14, Daskalakis *et al.* *ibid* '14] [Klaers *et al.* Nature '10]

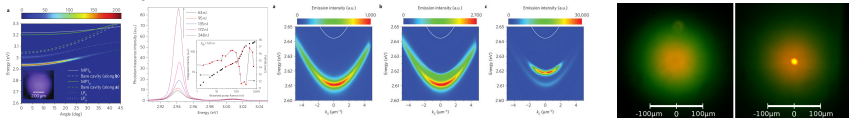
- Wide variety of systems:
 - ▶ polymers, fluorenes, J-aggregates, molecular crystals.
 - ▶ Often large polariton splitting, $g\sqrt{N} \sim 0.1 \text{ eV} \leftrightarrow 1000K$

• Theory questions/challenges

- ▶ Ultrastrong coupling
- ▶ Vibrational modes
- ▶ (Partial) thermalisation

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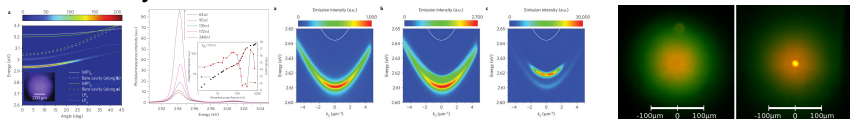
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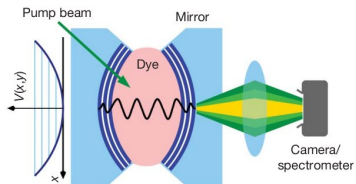
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Photon BEC experiments

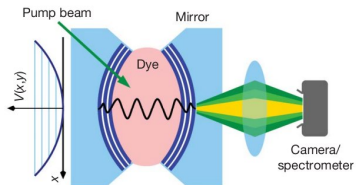


- Dye filled microcavity

➤ No strong coupling

[Klaers et al, Nature, 2010]

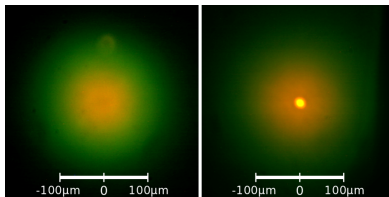
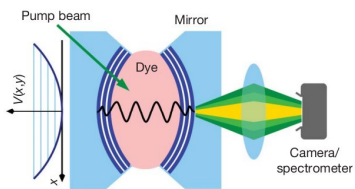
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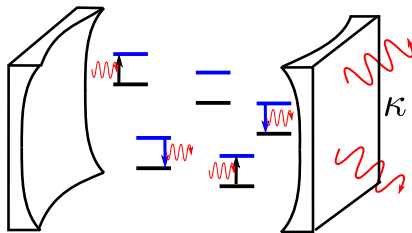
Dicke Holstein Model

- Dicke model: 2LS \leftrightarrow photons

- Molecular vibrational mode

- Phonon frequency Ω

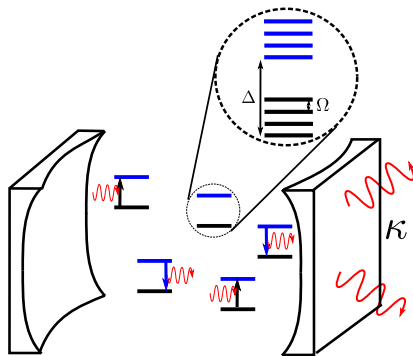
- Huang-Rhys parameter S —
coupling strength



$$H_{\text{sys}} = \omega \psi^\dagger \psi + \sum_{\alpha} \left[\frac{\epsilon}{2} \sigma_{\alpha}^z + g (\psi + \psi^\dagger) (\sigma_{\alpha}^+ + \sigma_{\alpha}^-) \right]$$

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$$H_{\text{sys}} = \omega \psi^\dagger \psi + \sum_{\alpha} \left[\frac{\epsilon}{2} \sigma_{\alpha}^z + g (\psi + \psi^\dagger) (\sigma_{\alpha}^+ + \sigma_{\alpha}^-) \right] + \sum_{\alpha} \Omega \left\{ b_{\alpha}^\dagger b_{\alpha} + \sqrt{S} \sigma_{\alpha}^z (b_{\alpha}^\dagger + b_{\alpha}) \right\}$$

Overview

1 Introduction: organic molecules

2 Modelling photon BEC

- Threshold behaviour
- Time evolution
- Pump-spot size dependence

3 Strong coupling: polaritons

- Polariton spectrum nature

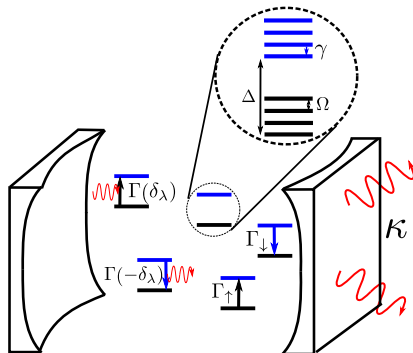
Modelling

$$H_{\text{sys}} = \sum_m \omega_m \psi_m^\dagger \psi_m + \sum_\alpha \left[\frac{\epsilon}{2} \sigma_\alpha^z + g (\psi_m \sigma_\alpha^+ + \text{H.c.}) \right] \\ + \sum_\alpha \Omega \left\{ b_\alpha^\dagger b_\alpha + \sqrt{S} \sigma_\alpha^z (b_\alpha^\dagger + b_\alpha) \right\}$$

- **2D** harmonic oscillator

$$\omega_m = \omega_{\text{cutoff}} + m\omega_{H.O.}$$

- Incoherent processes: excitation, decay, loss, vibrational thermalisation.
- Weak coupling, perturbative in g

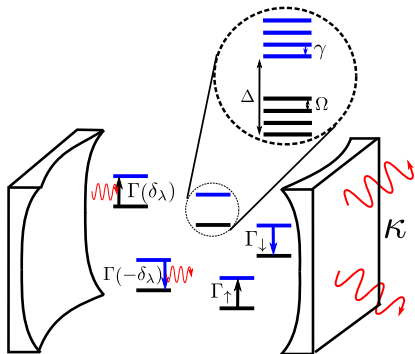


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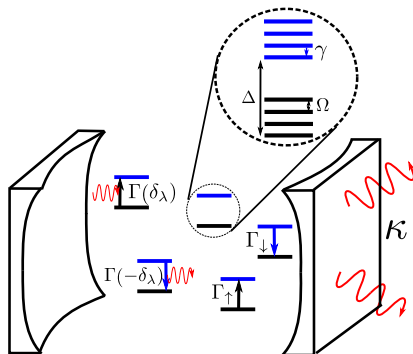
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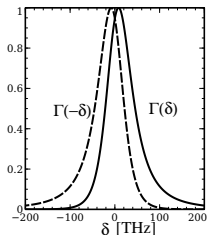
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Modelling

Master equation

$$\dot{\rho} = -i[H_0, \rho] - \sum_m \frac{\kappa}{2} \mathcal{L}[\psi_m] - \sum_\alpha \left[\frac{\Gamma_\uparrow}{2} \mathcal{L}[\sigma_\alpha^+] + \frac{\Gamma_\downarrow}{2} \mathcal{L}[\sigma_\alpha^-] \right] \\ - \sum_{m,\alpha} \left[\frac{\Gamma(\delta_m = \omega_m - \epsilon)}{2} \mathcal{L}[\sigma_\alpha^+ \psi_m] + \frac{\Gamma(-\delta_m = \epsilon - \omega_m)}{2} \mathcal{L}[\sigma_\alpha^- \psi_m^\dagger] \right]$$



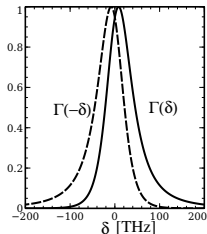
- Kennard-Stepanov
 $\Gamma(+\delta) \simeq \Gamma(-\delta) e^{\beta\delta}$
- Expt: $\omega_0 < \epsilon$
- $\Gamma \rightarrow 0$ at large δ

[Marthaler et al PRL '11, Kirton & JK PRL '13]

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Microscopic model – calculating $\Gamma(\delta)$

How to calculate $\Gamma(\delta)$

- Polaron transform (exact)

$$h_\alpha = \frac{\epsilon}{2} \sigma_\alpha^z + g (\psi_m \sigma_\alpha^+ D_\alpha + \text{H.c.}) + \Omega b_\alpha^\dagger b_\alpha,$$

$$D_\alpha = \exp \left[2\sqrt{S}(b_\alpha^\dagger - b_\alpha) \right]$$

- Correlation function:

$$\Gamma(\delta) = 2g^2 \Re \int dt \langle D_\alpha^\dagger(t) D_\alpha(0) \rangle \exp \left[-(\Gamma_\uparrow + \Gamma_\downarrow) \frac{t}{2} \right] e^{-i\delta t}$$

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Microscopic model – requirements for Kennard-Stepanov

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$$\langle D_\alpha^\dagger(t) D_\alpha(0) \rangle = \langle D_\alpha^\dagger(-t - i\beta) D_\alpha(0) \rangle$$

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Distribution $g_m n_m$

- Master equation \rightarrow Rate equation

$$\partial_t n_m = -\kappa n_m + \Gamma(-\delta_m)(n_m + 1)N_{\uparrow} - \Gamma(\delta_m)n_m N_{\downarrow}$$

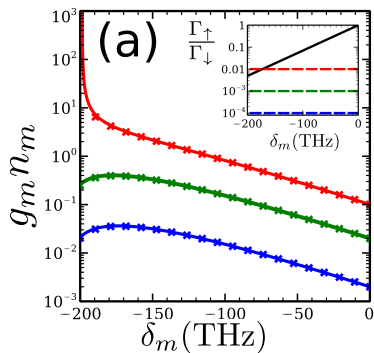
- Bose-Einstein distribution without losses

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Low loss: Thermal

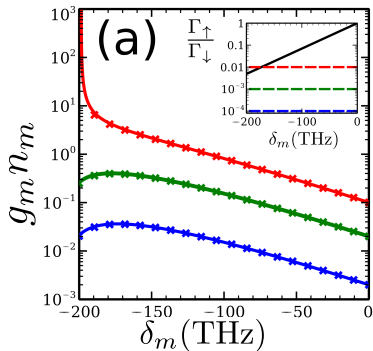
[Kirton & JK PRL '13]

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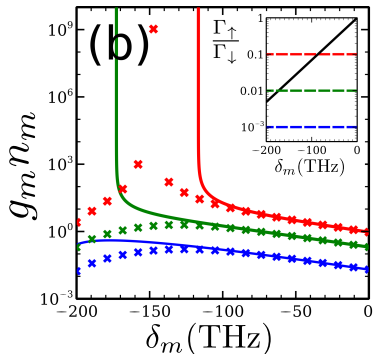
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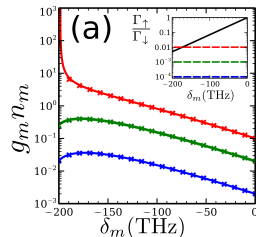


High loss \rightarrow Laser

Chemical potential?

- Steady state distribution:

$$\frac{n_m}{n_m + 1} = \frac{\Gamma(-\delta_m) N_{\uparrow}}{\kappa + \Gamma(\delta_m) N_{\downarrow}}$$



- $\kappa \ll N\Gamma(\delta)$, Kennard-Stepanov

$$\frac{n_m}{n_m + 1} = e^{-\beta \epsilon_m + \beta \mu}, \quad e^{\beta \mu} = \frac{N_{\uparrow}}{N_{\downarrow}} = \frac{\Gamma_{\uparrow} + \sum_m \Gamma(\delta_m) n_m}{\Gamma_{\downarrow} + \sum_m \Gamma(-\delta_m) (n_m + 1)}$$

- Below threshold, $\mu = k_B T \ln[\Gamma_{\uparrow}/\Gamma_{\downarrow}]$
- Above threshold, $\mu \rightarrow \delta_0$

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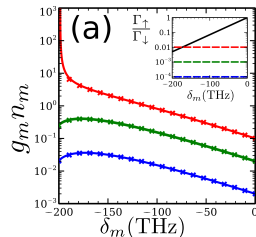
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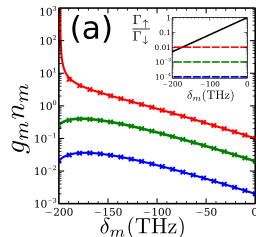
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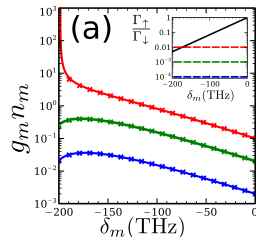
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Threshold condition

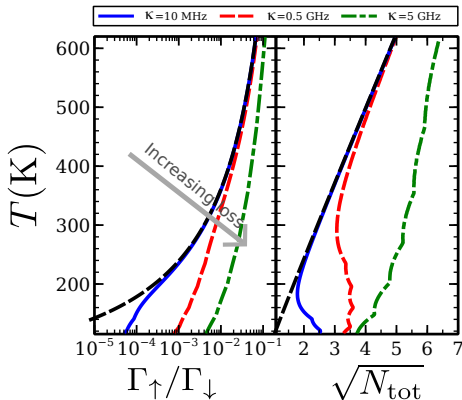
Use: $\max[n_m] = 1/(\beta\epsilon) \rightarrow k_B T_c = \sqrt{6/\pi^2\epsilon}\sqrt{N}$.

- Pump rate (Laser)
- Critical density (condensate)

- Thermal at low κ /high temperature
- High loss, κ competes with $\Gamma(\pm\delta_0)$
- Low temperature, $\Gamma(\pm\delta_0)$ shrinks

Threshold condition

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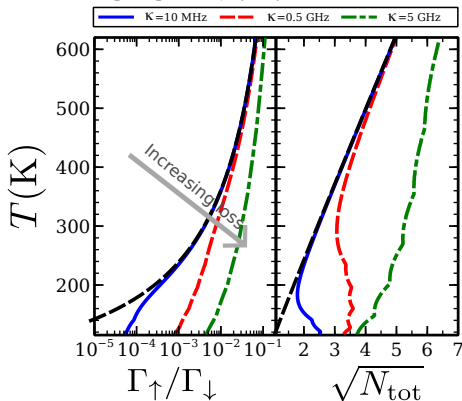
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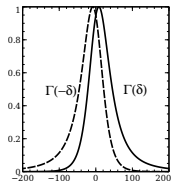


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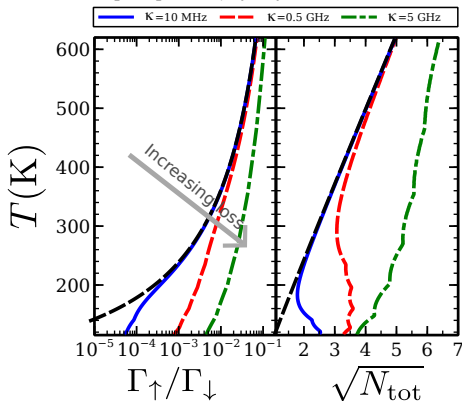
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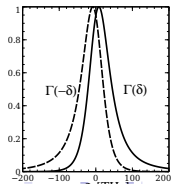
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Time evolution

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2 Modelling photon BEC

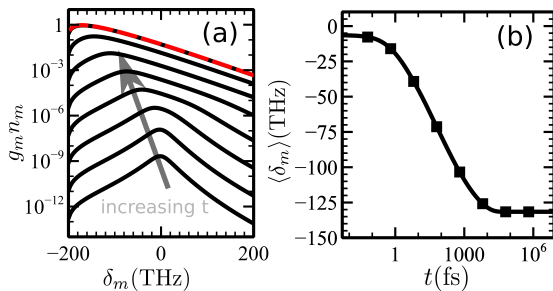
- Threshold behaviour
- **Time evolution**
- Pump-spot size dependence

3 Strong coupling: polaritons

- Polariton spectrum nature

Time evolution

- Initial state: excited molecules
 - Initial emission, follows gain peak
 - Thermalisation by repeated absorption

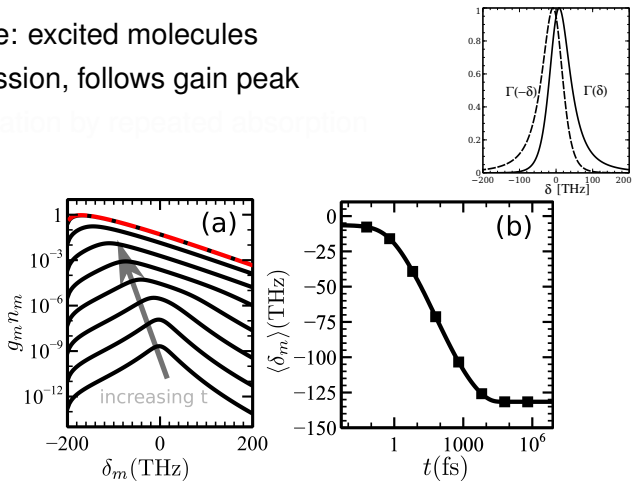


[Kirton & JK arXiv:1410.6632]

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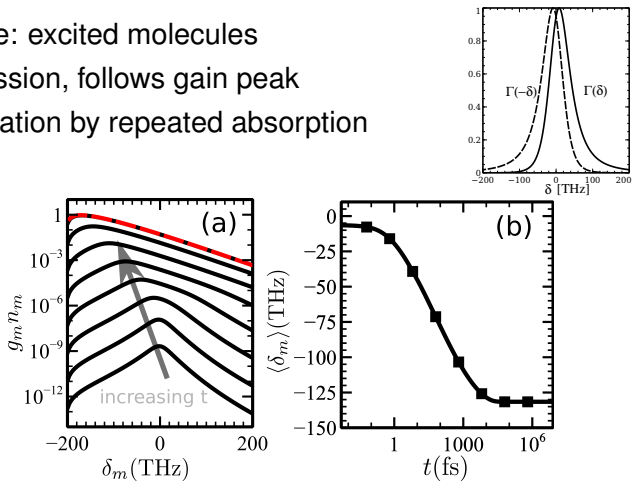
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Pump spot size

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Spatially varying pump intensity

- Consider effects of pump profile, $\Gamma_{\uparrow}(\mathbf{r}) = \Gamma_{\uparrow} \exp(-r^2/2r_{\text{spot}}^2)$

$$\partial_t \rho_{\uparrow}(\mathbf{r}) = -\tilde{\Gamma}_{\downarrow}(\mathbf{r})\rho_{\uparrow}(\mathbf{r}) + \tilde{\Gamma}_{\uparrow}(\mathbf{r})\rho_{\downarrow}(\mathbf{r})$$

- Varying excited density – differential coupling to modes

$$\partial_t n_m = \Gamma(-\delta_m) Q_m (n_m + 1) - [\kappa + \Gamma(\delta_m)(\rho_m - Q_m)] n_m$$

$$Q_m = \int dr \rho_{\uparrow}(r) |\psi_m(r)|^2, \quad \rho_{\uparrow} + \rho_{\downarrow} = \rho_m$$

- Experiments: [Marell & Nyman, arXiv:1410.6822]

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NB $\Gamma(\delta)$ differs by area factor

Spatially varying pump intensity

- Consider effects of pump profile, $\Gamma_{\uparrow}(\mathbf{r}) = \Gamma_{\uparrow} \exp(-r^2/2r_{\text{spot}}^2)$

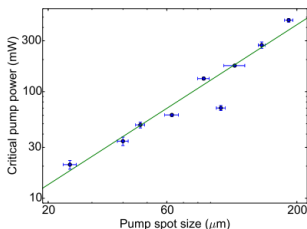
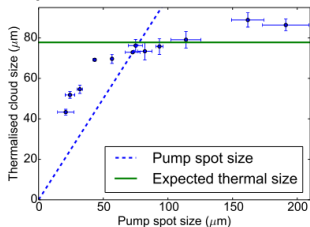
$$\partial_t \rho_{\uparrow}(\mathbf{r}) = -\tilde{\Gamma}_{\downarrow}(\mathbf{r})\rho_{\uparrow}(\mathbf{r}) + \tilde{\Gamma}_{\uparrow}(\mathbf{r})\rho_{\downarrow}(\mathbf{r})$$

- Varying excited density – differential coupling to modes

$$\partial_t n_m = \Gamma(-\delta_m) O_m (n_m + 1) - [\kappa + \Gamma(\delta_m)(\rho_m - O_m)] n_m$$

$$O_m = \int d\mathbf{r} \rho_{\uparrow}(\mathbf{r}) |\psi_m(\mathbf{r})|^2, \quad \rho_{\uparrow} + \rho_{\downarrow} = \rho_m$$

- Experiments: [Marelic & Nyman, arXiv:1410.6822]



NB $\Gamma(\delta)$ differs by area factor

Spatially varying pump: below threshold

- Far below threshold:

- ▶ Excitation: $\rho_{\uparrow}(\mathbf{r}) \simeq \rho_m \frac{\Gamma_{\uparrow}(\mathbf{r})}{\Gamma_{\downarrow}} \ll \rho_m$

- ▶ If $\kappa \ll \rho_m \Gamma(\delta_m)$, $\frac{n_m}{n_m + 1} \simeq e^{-\beta \delta_m} \times \int d\mathbf{r} \frac{\Gamma_{\uparrow}(\mathbf{r})}{\Gamma_{\downarrow}} |\psi_m(\mathbf{r})|^2$

• Resulting profile, $I(\mathbf{r}) = \sum_m \rho_m |\psi_m(\mathbf{r})|^2$

Spatially varying pump: below threshold

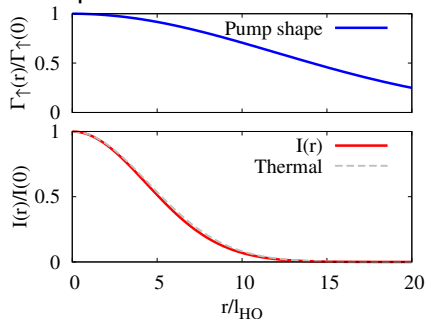
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- Resulting profile, $I(\mathbf{r}) = \sum_m n_m |\psi_m(\mathbf{r})|^2$

Cloud profile



▶ If $\kappa_{\text{trap}} \rightarrow 0$, $\Omega_m \propto \delta_{m,0}$ so $I(r) \propto |\psi_0(r)|^2$.

Spatially varying pump: below threshold

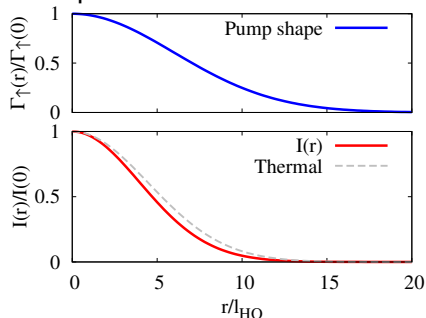
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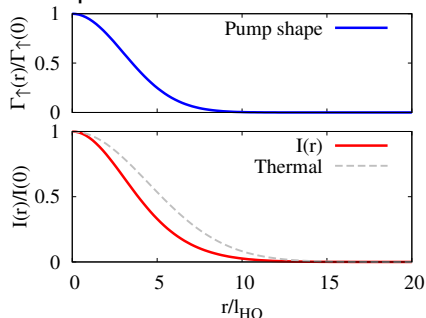
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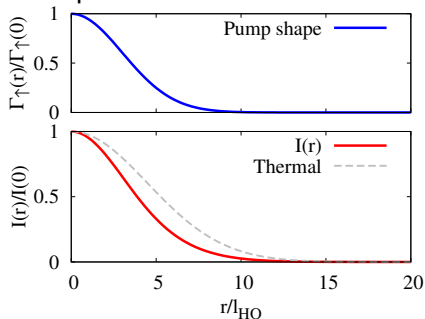
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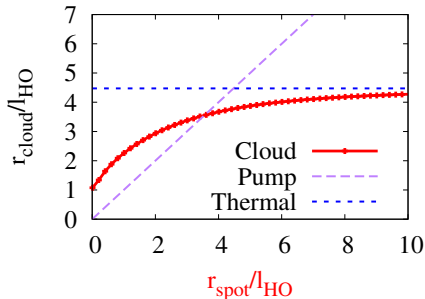
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Cloud profile



Cloud size:

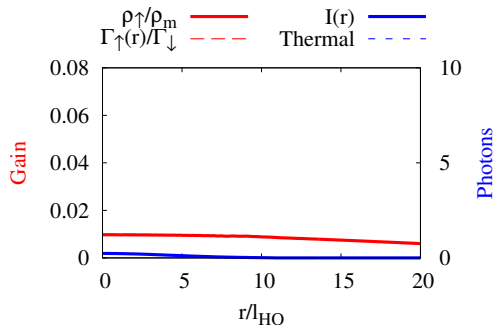


- If $r_{\text{spot}} \rightarrow 0$, $O_m \propto \delta_{m,0}$ so $I(r) \propto |\psi_0(r)|^2$.

Near threshold behaviour

- Large spot, $r_{\text{spot}}/l_{\text{HO}} = 20$

- “Gain saturation” at centre
- Non Boltzmann peak — BEC

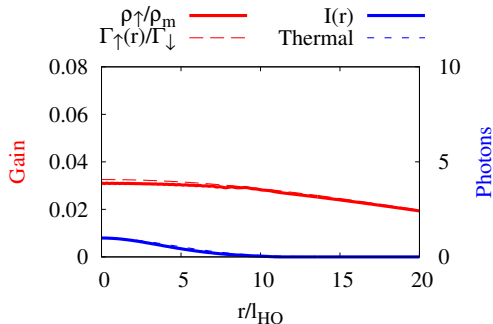


- If thermal: $(n_m + 1)\Gamma(-\delta_m) = n_m\Gamma(\delta_m)\zeta^{-1}$
- Saturation of $\rho_{\uparrow} = \rho_m/(1 + \zeta^{-1})$ — spatial equilibration

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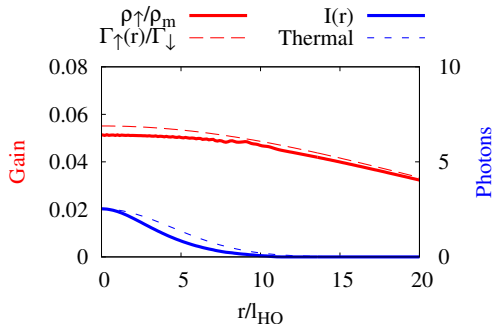


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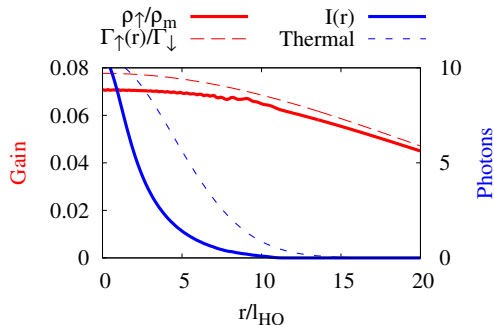
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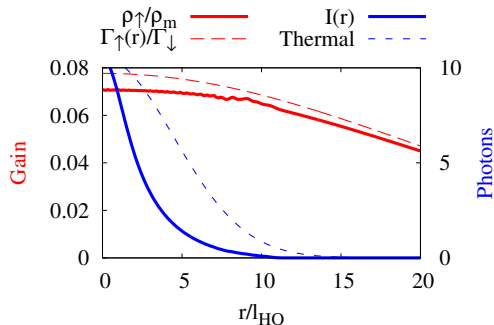
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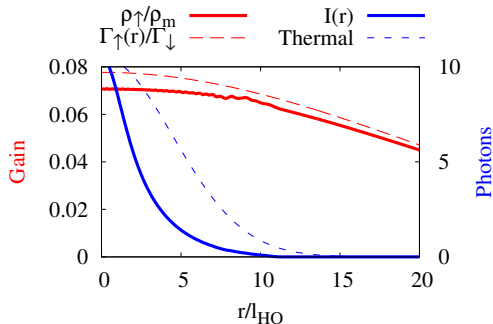


$$\frac{\rho_{\uparrow}(r)}{\rho_m} = \frac{\tilde{\Gamma}_{\uparrow}(r)}{\tilde{\Gamma}_{\uparrow}(r) + \tilde{\Gamma}_{\downarrow}(r)}$$

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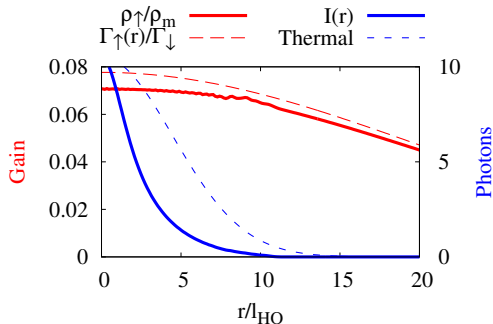


$$\frac{\rho_{\uparrow}(r)}{\rho_m} = \frac{\Gamma_{\uparrow}(r) + \sum_m n_m \Gamma(\delta_m) |\psi_m(r)|^2}{\Gamma_{\uparrow}(r) + \Gamma_{\downarrow} + \sum_m [n_m \Gamma(\delta_m) + (n_m + 1) \Gamma(-\delta_m)] |\psi_m(r)|^2}$$

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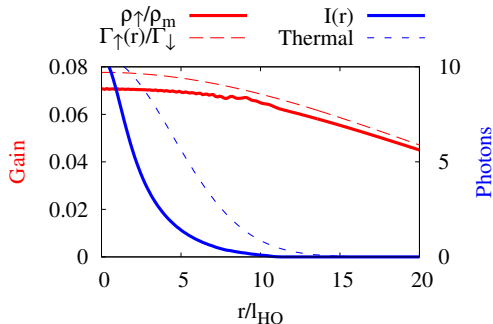
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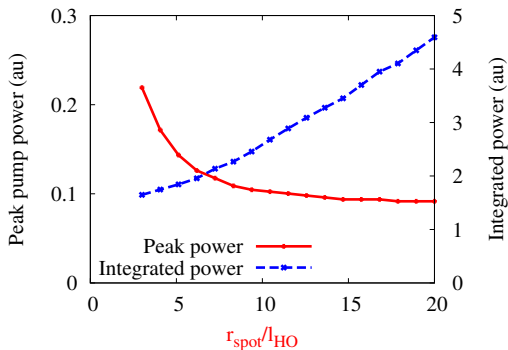


$$\frac{\rho_{\uparrow}(r)}{\rho_m} = \frac{\Gamma_{\uparrow}(r) + \Gamma_{\text{photon}}(r)}{\Gamma_{\uparrow}(r) + \Gamma_{\downarrow} + (1 + \zeta^{-1})\Gamma_{\text{photon}}(r)}$$

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Effect of spot size on threshold

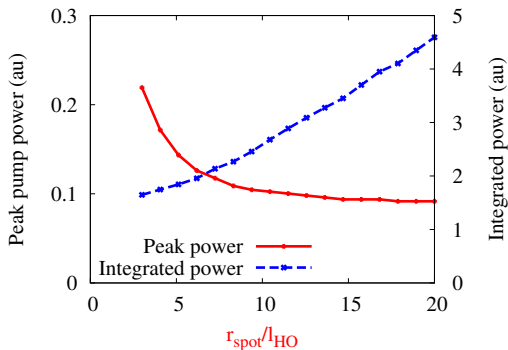
Threshold power:



- Small spot, integrated power saturates
- Large spot, peak power saturates

Effect of spot size on threshold

Threshold power:

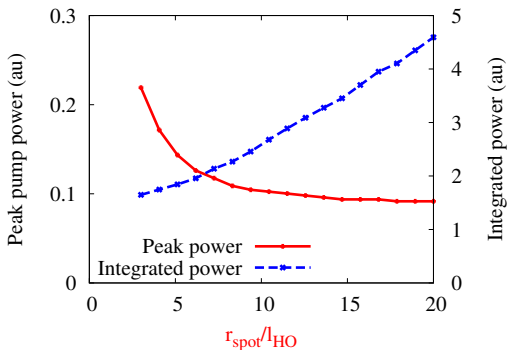


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Strong coupling: polaritons

1 Introduction: organic molecules

2 Modelling photon BEC

- Threshold behaviour
- Time evolution
- Pump-spot size dependence

3 Strong coupling: polaritons

- Polariton spectrum nature

Strong coupling phase diagram — mean field

- Mean field — single photon mode

$$H = \omega \psi^\dagger \psi + \sum_{\alpha} \left[\epsilon \mathbf{S}_{\alpha}^z + g \left(\psi \mathbf{S}_{\alpha}^+ + \psi^\dagger \mathbf{S}_{\alpha}^- \right) + \Omega \left\{ b_{\alpha}^\dagger b_{\alpha} + \sqrt{S} \left(b_{\alpha}^\dagger + b_{\alpha} \right) \mathbf{S}_{\alpha}^z \right\} \right]$$

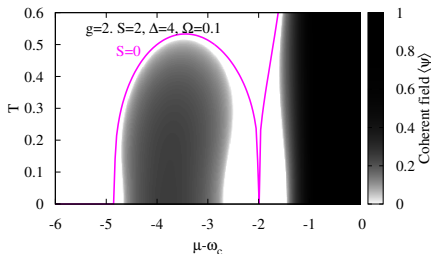
- $\epsilon = \omega - \Delta$
- Mot lobes if $\epsilon < \omega - 2g$
- S reduces g_{eff}

- Reentrant behaviour — Min μ at $k_B T \sim 0.1 \Omega$

Strong coupling phase diagram — mean field

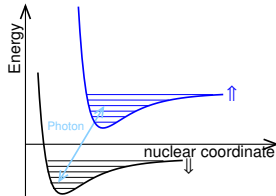
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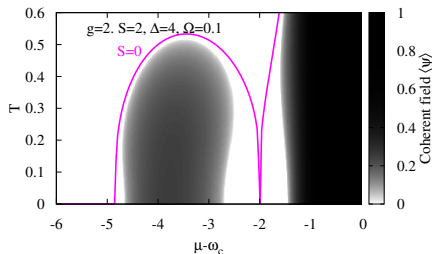


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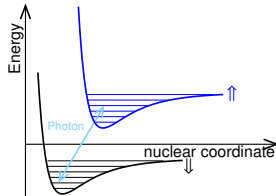
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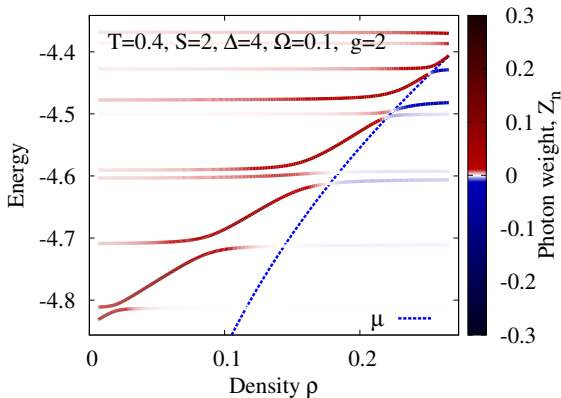
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Polariton spectrum: photon weight



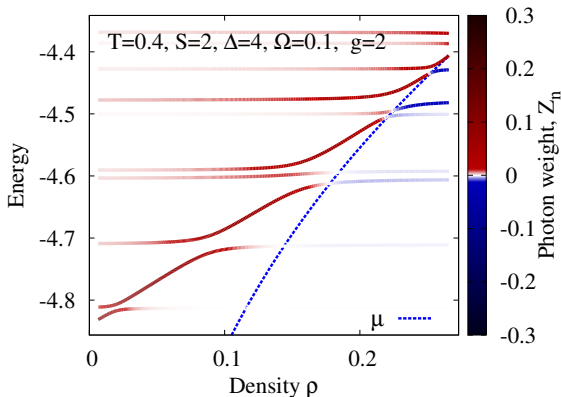
- Saturating 2LS: $g_{\text{eff}}^2 \sim g^2(1 - 2\rho)$

• What is nature of polariton mode?

• $G^{\mu}(\tau) = -i \langle \psi^{\dagger}(\tau) \psi(0) \rangle$, $G^{\mu}(\omega) = \sum_n \frac{Z_n}{\omega - \omega_n}$

[Cwik *et al.* EPL '14]

Polariton spectrum: photon weight

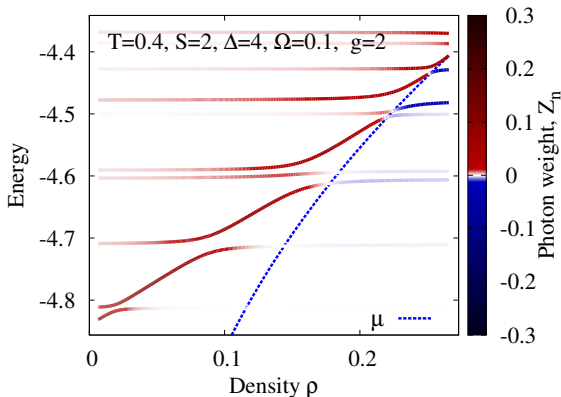


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$$g^2(\rho) = \langle \psi(\rho) | \psi(0) \rangle, \quad g^2(\rho) = \sum_n \frac{Z_n}{\omega_n - \omega_0}$$

[Cwik *et al.* EPL '14]

Polariton spectrum: photon weight



- Saturating 2LS: $g_{\text{eff}}^2 \sim g^2(1 - 2\rho)$
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- $G^R(t) = -i\langle \psi^\dagger(t)\psi(0) \rangle, \quad G^R(\nu) = \sum_n \frac{Z_n}{\nu - \omega_n}$

[Cwik *et al.* EPL '14]

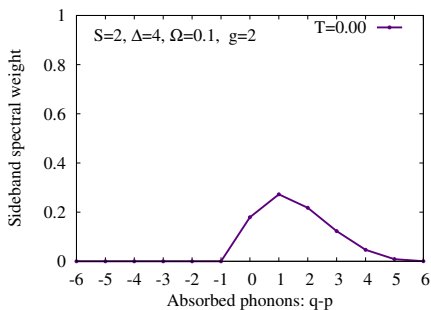
Polariton spectrum: what condensed

- Repeat weight for n -phonon channel
- Eigenvector that is macroscopically occupied
- Optimal $T \sim 2\Omega$

[Cwik *et al.* EPL '14]

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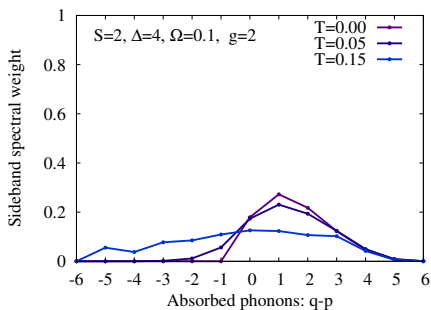


[Cwik *et al.* EPL '14]

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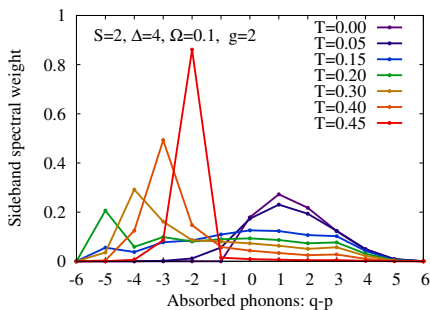


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[Cwik *et al.* EPL '14]

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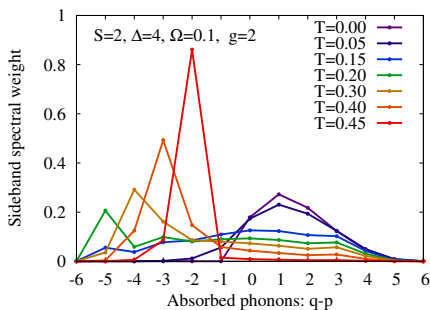


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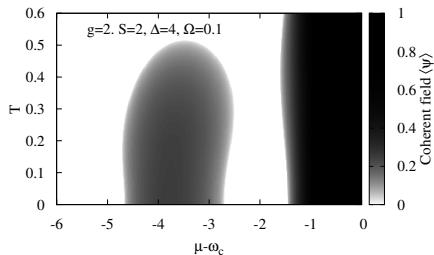
[Cwik *et al.* EPL '14]

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[Cwik *et al.* EPL '14]

Acknowledgements

GROUP:



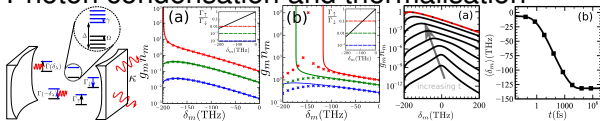
COLLABORATORS: Reja (MPI-PKS), Littlewood (ANL & Chicago)

FUNDING:



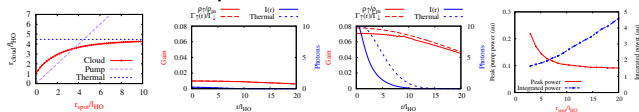
Summary

● Photon condensation and thermalisation

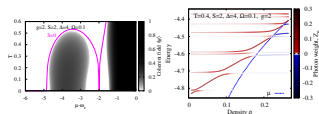


[Kirton & Keeling, PRL '13, arXiv:1410.6632]

● Effects of finite spot size



● Reentrance, phonon assisted transition, 1st order at $S \gg 1$



[Cwik *et al.* EPL '14]

Extra Slides

4 Ultra-strong phonon coupling?

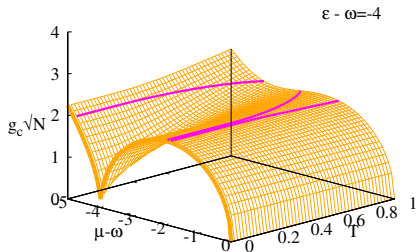
5 Anticrossing vs ρ

6 Vibrational reconfiguration

Critical coupling with increasing S

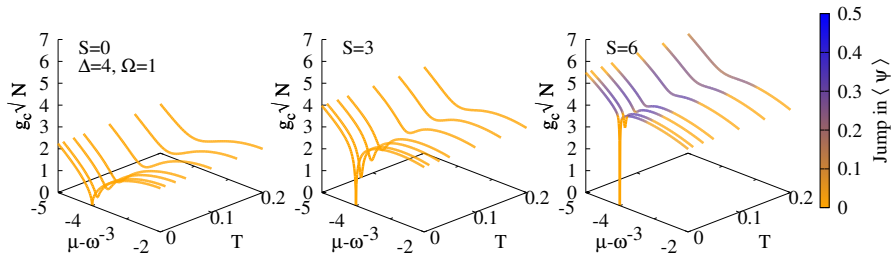
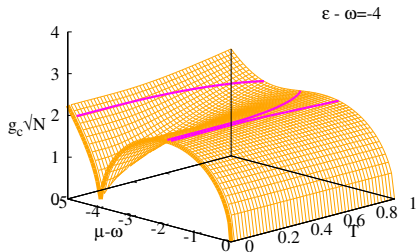
- Re-orient phase diagram
- g vs μ, T

• Colors \rightarrow Jump of $\langle \psi \rangle$



Critical coupling with increasing S

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- g vs μ, T
- Colors \rightarrow Jump of $\langle \psi \rangle$



Explanation: Polaron formation

- Unitary transform

$$H_\alpha \rightarrow \tilde{H}_\alpha = e^{K_\alpha} H_\alpha e^{-K_\alpha} \quad K = \sqrt{S} S_\alpha^z (b_\alpha^\dagger - b_\alpha)$$

- Coupling moves to S^\pm

$$\tilde{H}_\alpha = \text{const.} + \epsilon S_\alpha^z + \Omega b_\alpha^\dagger b_\alpha + g \left[\psi S_\alpha^+ e^{\sqrt{S}(b_\alpha^\dagger - b_\alpha)} + \text{H.c.} \right]$$

- Optimal phonon displacements, $\sim \sqrt{S}$

- Reduced $g_{\text{eff}} \sim g \times \exp(-S/2)$

- For $\psi \neq 0$, competition

$$\text{Variational MFT } |\psi\rangle_\alpha \sim \exp(-\eta K_\alpha - \zeta (b_\alpha^\dagger)) |0, S\rangle_\alpha$$

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Explanation: Polaron formation

- Unitary transform

$$H_\alpha \rightarrow \tilde{H}_\alpha = e^{K_\alpha} H_\alpha e^{-K_\alpha} \quad K = \sqrt{S} S_\alpha^z (b_\alpha^\dagger - b_\alpha)$$

- Coupling moves to S^\pm

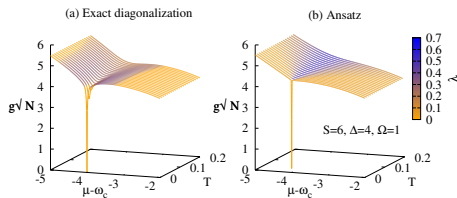
$$\tilde{H}_\alpha = \text{const.} + \epsilon S_\alpha^z + \Omega b_\alpha^\dagger b_\alpha + g \left[\psi S_\alpha^+ e^{\sqrt{S}(b_\alpha^\dagger - b_\alpha)} + \text{H.c.} \right]$$

- Optimal phonon displacements, $\sim \sqrt{S}$
- Reduced $g_{\text{eff}} \sim g \times \exp(-S/2)$
- For $\psi \neq 0$, competition

Variational MFT $|\psi\rangle_\alpha \sim \exp(-\eta K_\alpha - \zeta b_\alpha^\dagger) |0, \mathbf{S}\rangle_\alpha$

Collective polaron formation

- Compares well at $S \gg 1$
- Coherent bosonic state



- Feedback: Large/small $g_{\text{eff}} \leftrightarrow \lambda = \langle \psi \rangle$
- Variational free energy

$$F = (\omega_c - \mu)\lambda^2 + N \left\{ \Omega \left[\zeta^2 - S \frac{\eta(2 - \eta)}{4} \right] - T \ln \left[2 \cosh \left(\frac{\xi}{T} \right) \right] \right\}$$

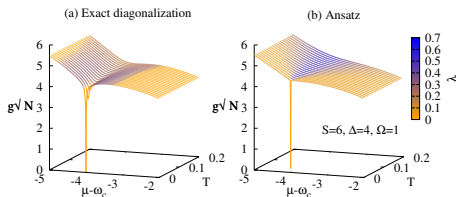
Effective 2LS energy in field:

$$\zeta^2 = \left(\frac{\epsilon - \mu}{2} + \Omega \sqrt{S(1 - \eta)} \zeta \right)^2 + g^2 \lambda^2 e^{-2\eta/T}$$

[Cwik *et al.* EPL '14]

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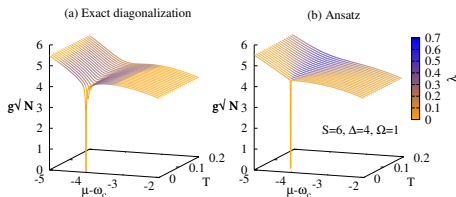
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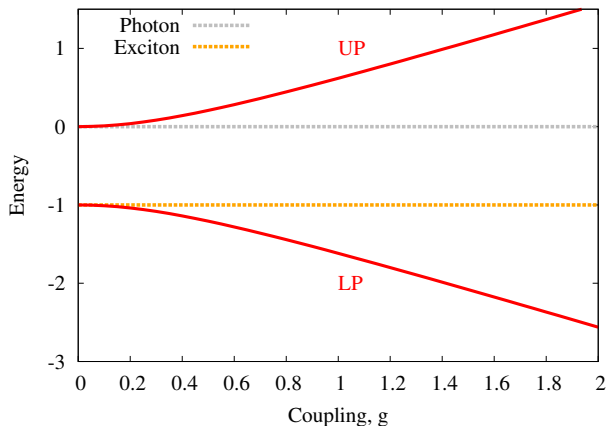
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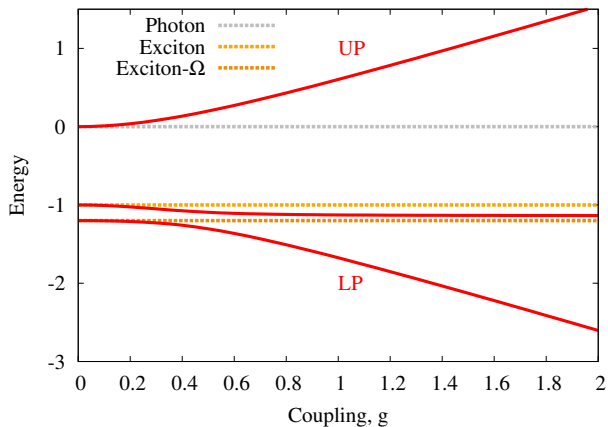
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Polariton spectrum — coupled oscillators

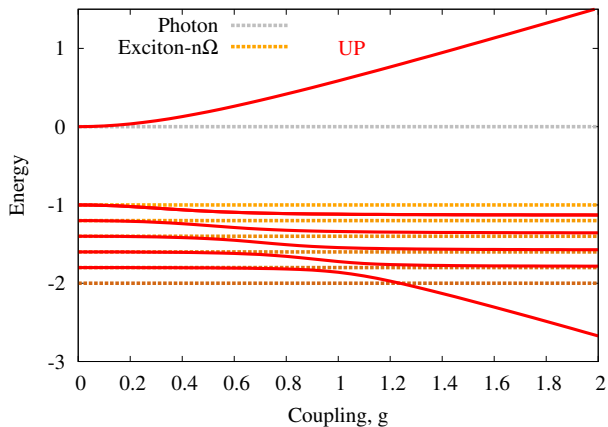
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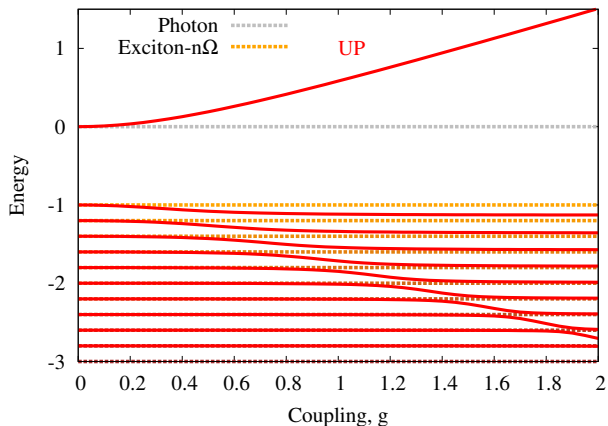
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Polariton spectrum — coupled oscillators



Vibrational reconfiguration

- $H = H_0 + H_1, H_1 = \sum_{n,k} g_{n,k} (\psi_k^\dagger \sigma_n^+ + \text{H.c.})$
- Schrieffer-Wolff: admixture of excited state

$$H_{\text{eff,vacuum}} = H_0 - \frac{g^2 N}{2(\epsilon + \omega)} \left\{ 1 - \frac{\Omega \sqrt{S}(b + b^\dagger)}{\epsilon + \omega} + \mathcal{O} \left[\left(\frac{\Omega}{\epsilon} \right)^2, \frac{g\sqrt{N}}{\epsilon} \right] \right\}$$

Reduced vibrational offset

$$\sqrt{S} \rightarrow \sqrt{S} \left(1 - \frac{g^2 N}{\epsilon + \omega} \right)$$

Increased effective coupling:

$$g_{\text{eff}}^2 = g^2 \exp(-S)$$

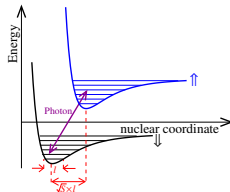
Numerically tiny effect, $\Omega \ll \epsilon$

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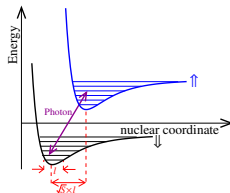
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