

Polariton and photon condensates in organic materials

Jonathan Keeling



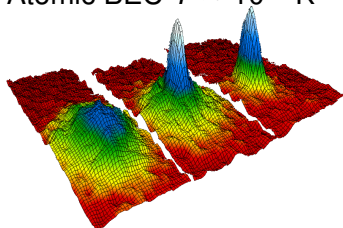
University of
St Andrews

1413-2013

Sheffield, October 2014

Coherent states of matter and light

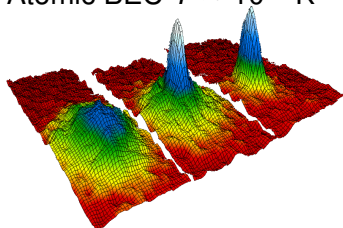
Atomic BEC $T \sim 10^{-7}\text{K}$



[Anderson *et al.* Science '95]

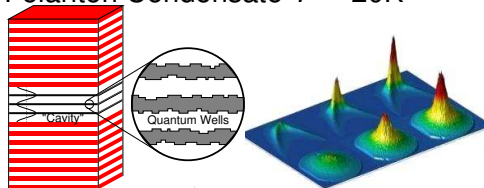
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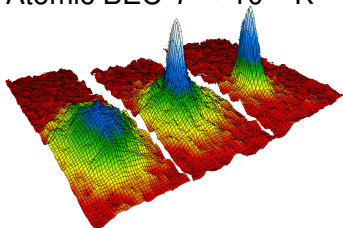
Polariton Condensate $T \sim 20\text{K}$



[Kasprzak *et al.* Nature, '06]

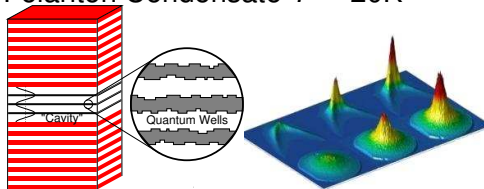
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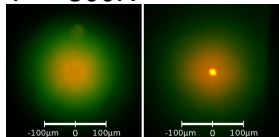
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Photon Condensate

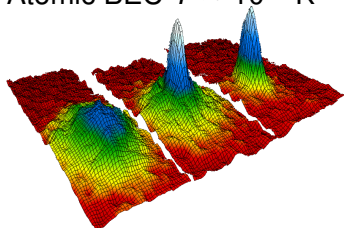
$T \sim 300\text{K}$



[Klaers *et al.* Nature, '10]

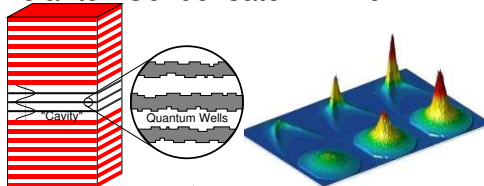
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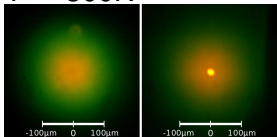
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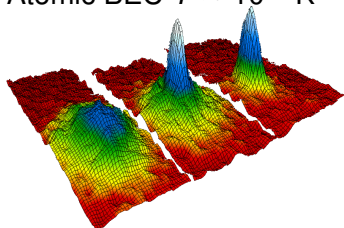
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Laser
 $T \sim ?, < 0, \infty$



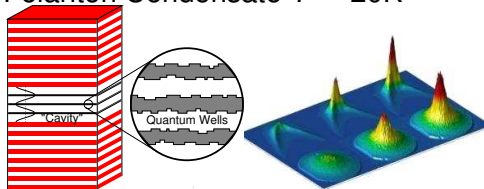
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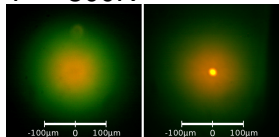
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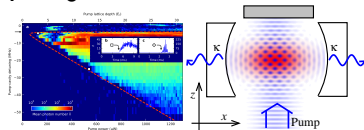


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Laser
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Superradiance transition
 $T \sim 0$



[Baumann *et al.* Nature, '10]

Outline

- 1 Condensation, superradiance, lasing
 - Polariton condensation and Dicke model
 - Condensation vs superradiance transition
 - Non-equilibrium condensation vs lasing
- 2 Room temperature condensates: Organic polaritons
 - Dicke phase diagram with phonons
 - Condensation of phonon replicas?
- 3 Room temperature condensates: Photons
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 - Time evolution
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Acknowledgements

GROUP:



COLLABORATORS: Szymanska (UCL), Reja (MPI-PKS), Littlewood (ANL), De Liberato (Soton)

FUNDING:

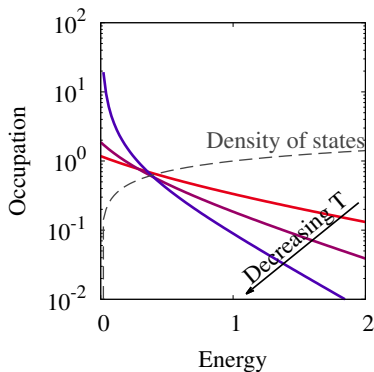
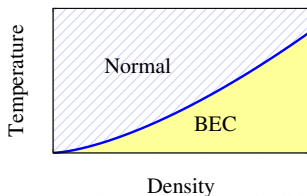


“Textbook” BEC

• Non-interacting viewpoint

▶ BE distribution: $\mu < \omega_0$

▶ $T_c = \frac{2\pi\hbar^2}{m} \left(\frac{n}{\xi_d} \right)^{2/d}$



• Interacting approach (WIDBG)

$$H = \sum_k \omega_k \psi_k^\dagger \psi_k + \frac{g}{2V} \sum_{k,k',q} \psi_{k+q}^\dagger \psi_{k-q}^\dagger \psi_k \psi_{k'}$$

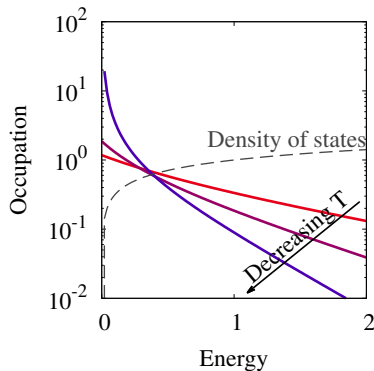
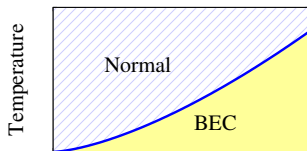
• Mean field: $|\psi_0|^2 = (\mu - \omega_0)/V$

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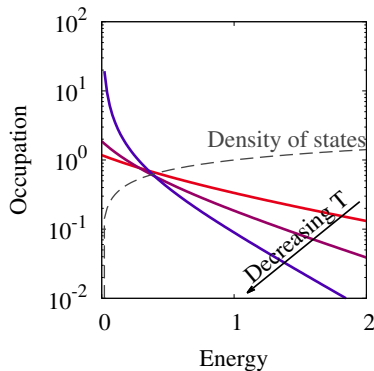
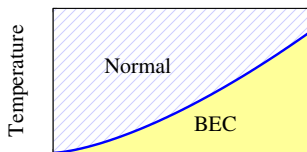
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“Textbook” Laser: Maxwell Bloch equations

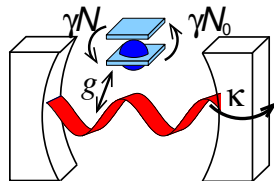
$$H = \omega_0 \psi^\dagger \psi + \sum_{\alpha} \epsilon_{\alpha} S_{\alpha}^Z + g_{\alpha, \mathbf{k}} \left(\psi S_{\alpha}^{+} + \psi^{\dagger} S_{\alpha}^{-} \right)$$

$$\text{Maxwell-Bloch eqns: } P = -i \langle S^{-} \rangle, N = 2 \langle S^Z \rangle$$

$$\partial_t \psi = -i \omega_0 \psi - \kappa \psi + \sum_{\alpha} g_{\alpha} P_{\alpha}$$

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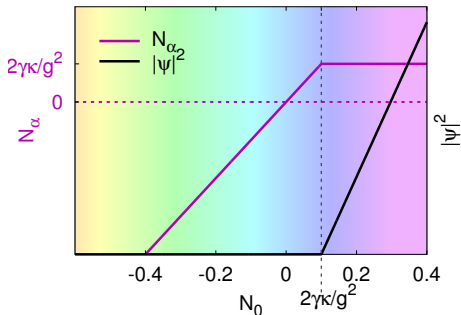
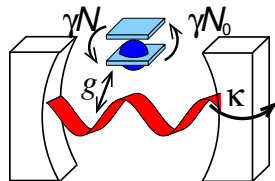
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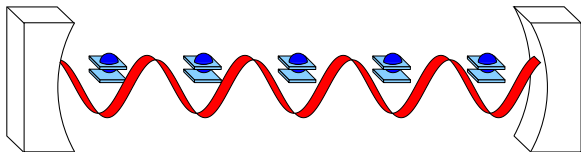
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$$|\psi|^2 > 0 \text{ if } N_0 g^2 > 2\gamma\kappa$$

“Textbook” Dicke-Hepp-Lieb superradiance

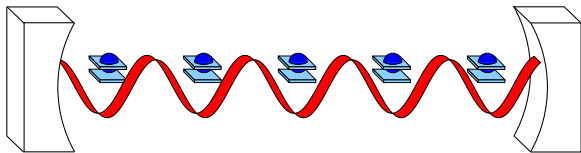


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- Coherent state: $|\Psi\rangle \rightarrow e^{\lambda \psi^\dagger + \eta S^+} |\Omega\rangle$
- Small g , min at $\lambda, \eta = 0$

[Hepp, Lieb, Ann. Phys. '73]

“Textbook” Dicke-Hepp-Lieb superradiance



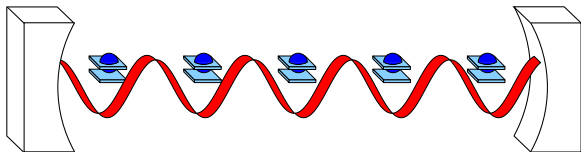
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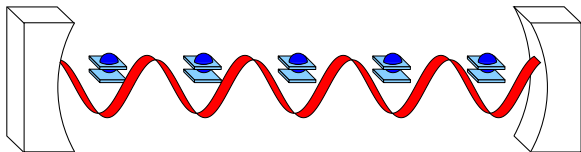
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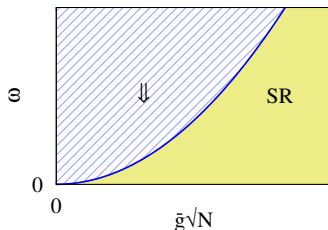
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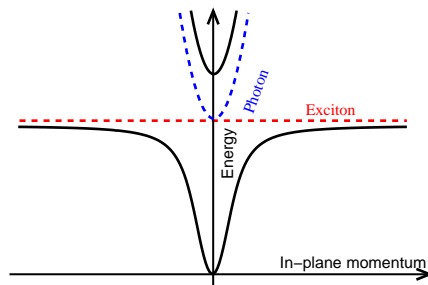
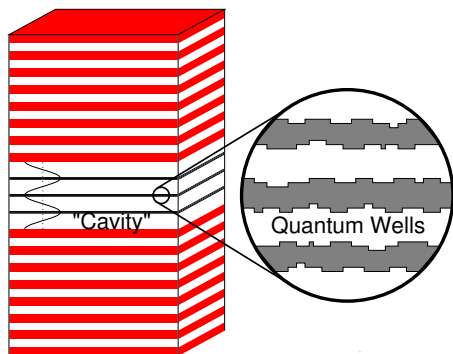
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- Dicke phase diagram with phonons
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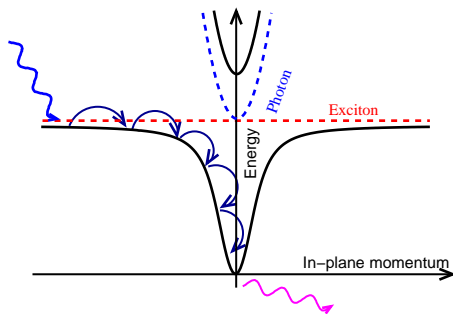
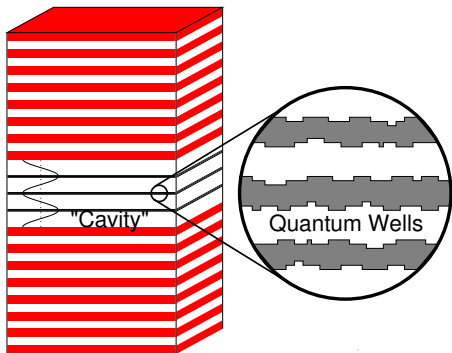
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Microcavity polaritons

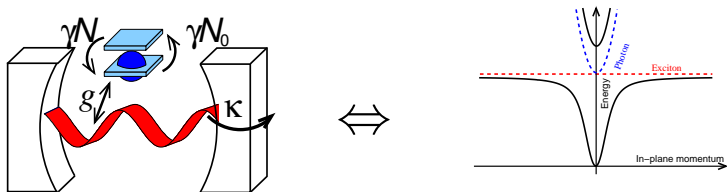


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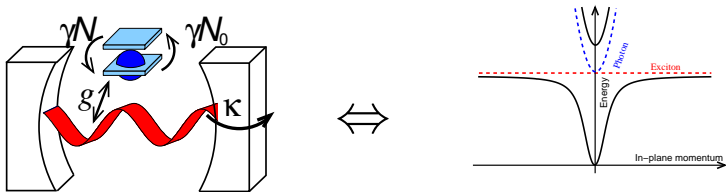
Lasing-condensation crossover model

- Use model that can show lasing and condensation:



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Dicke model:

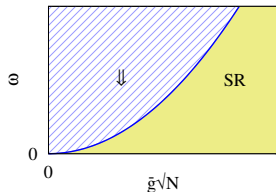
$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} [\epsilon S_{\alpha}^Z + g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} S_{\alpha}^{+} + \text{H.c.}]$$

Dicke-Hepp-Lieb superradiance and modes

$$H = \omega \psi^\dagger \psi + \epsilon S^z + g \left(\psi^\dagger S^- + \psi S^+ \right)$$

Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

- Normal state, $S^z = -N/2 + B^\dagger B$
 $H = \omega \psi^\dagger \psi + \epsilon B^\dagger B + g\sqrt{N} \left(\psi^\dagger B + \psi B^\dagger \right)$
- Excitation cost E :
 $(E - \omega)(E - \epsilon) = g^2 N$
- Transition when $E = 0$



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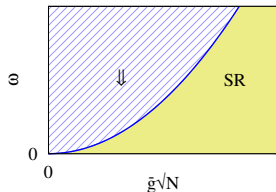
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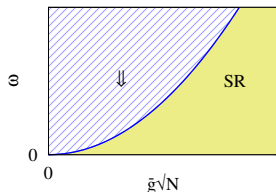
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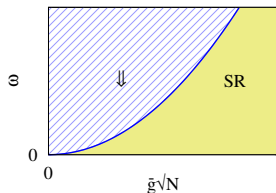
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Grand canonical ensemble

Grand canonical ensemble:

- If $H \rightarrow H - \mu(\mathbf{S}^z + \psi^\dagger\psi)$, need only: $g^2 N > (\omega - \mu)|\epsilon - \mu|$

- Fix density / fix $\mu > 0$ — pumping

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[Eastham and Littlewood, PRB '01]

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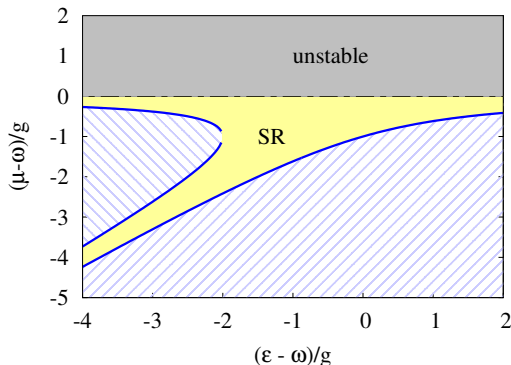
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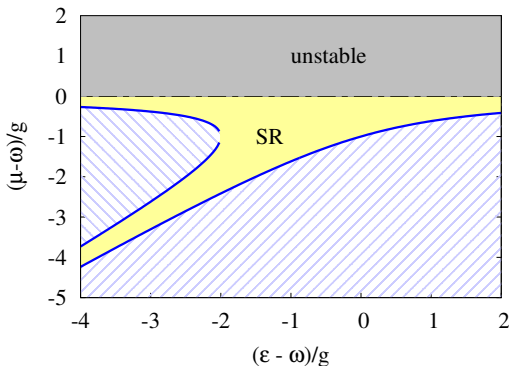
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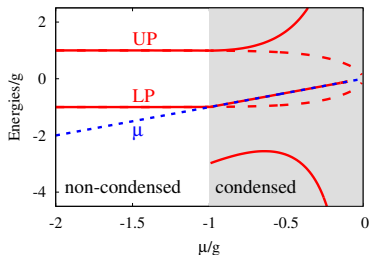
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Plot for $\omega = \epsilon = 0$

[Eastham and Littlewood, PRB '01]

Polariton model and equilibrium results

- Localised excitons, propagating photons

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- Self-consistent polarisation and field

$$(\omega - \mu) \psi = \sum_{\alpha} \frac{g_{\alpha}^2 \psi}{E_{\alpha}} \tanh(\beta E_{\alpha}/2), \quad E_{\alpha}^2 = (\epsilon_{\alpha} - \mu)^2 + 4g_{\alpha}^2 |\psi|^2$$

→ Phase diagram:

Polariton model and equilibrium results

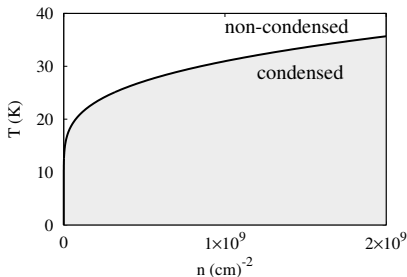
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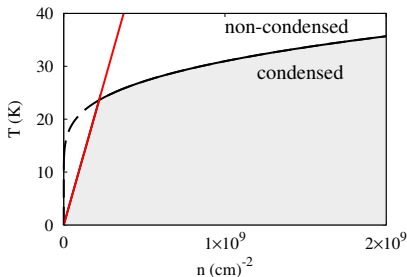
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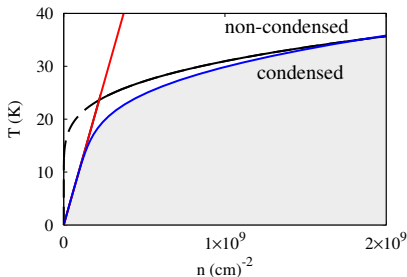
- Localised excitons, propagating photons

$$H - \mu N = \sum_{\mathbf{k}} (\omega_{\mathbf{k}} - \mu) \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} (\epsilon_{\alpha} - \mu) S_{\alpha}^Z + g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} S_{\alpha}^+ + \text{H.c.}$$

- Self-consistent polarisation and field

$$(\omega - \mu) \psi = \sum_{\alpha} \frac{g_{\alpha}^2 \psi}{E_{\alpha}} \tanh(\beta E_{\alpha} / 2), \quad E_{\alpha}^2 = (\epsilon_{\alpha} - \mu)^2 + 4g_{\alpha}^2 |\psi|^2$$

- Phase diagram:



Non-equilibrium condensation vs lasing

1 Condensation, superradiance, lasing

- Polariton condensation and Dicke model
- Condensation vs superradiance transition
- **Non-equilibrium condensation vs lasing**

2 Room temperature condensates: Organic polaritons

- Dicke phase diagram with phonons
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- Phase diagram
- Time evolution
- Linewidth

Simple Laser: Maxwell Bloch equations

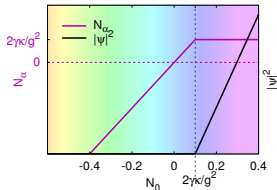
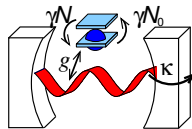
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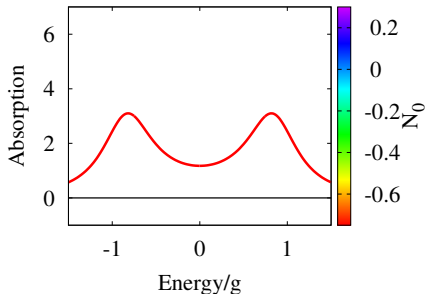
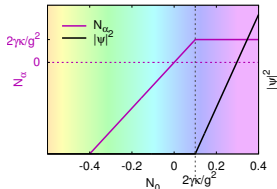
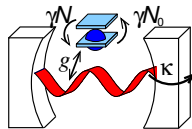
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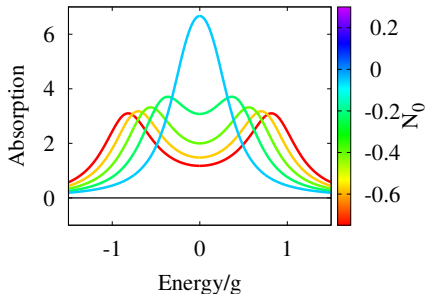
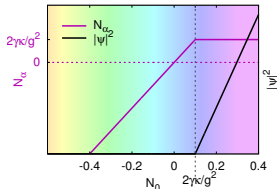
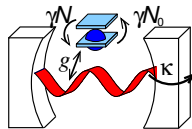
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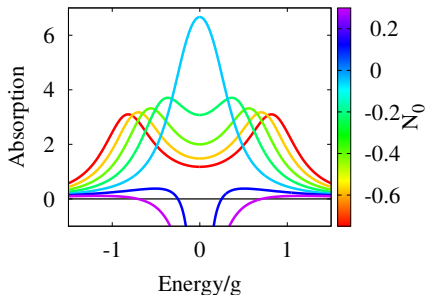
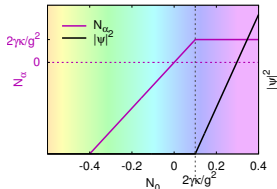
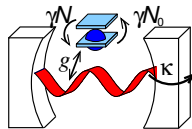
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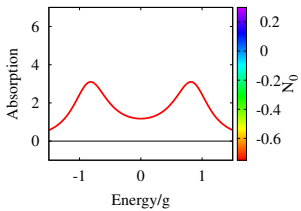
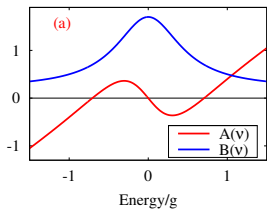
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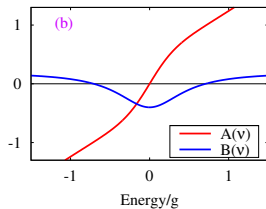
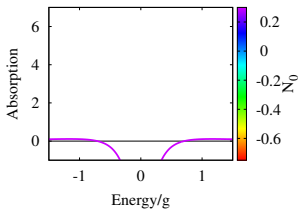
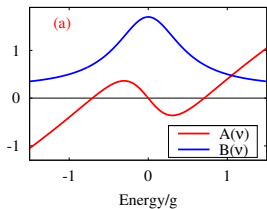
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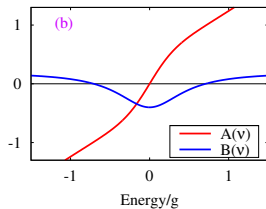
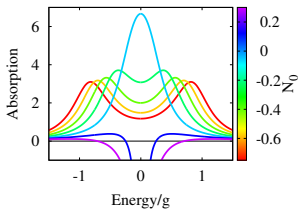
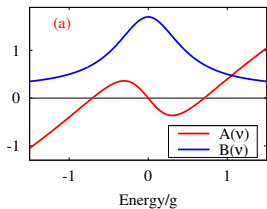
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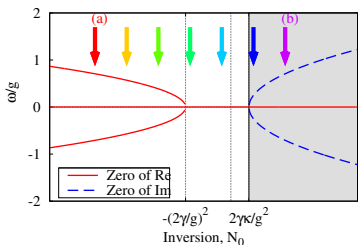


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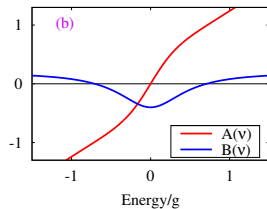
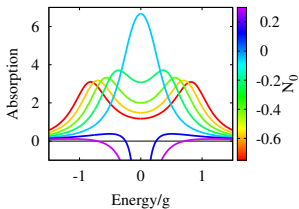
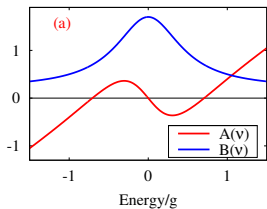


Laser:

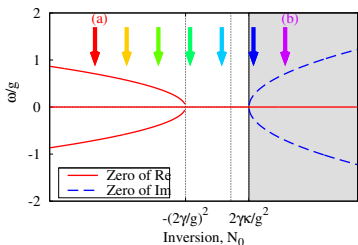


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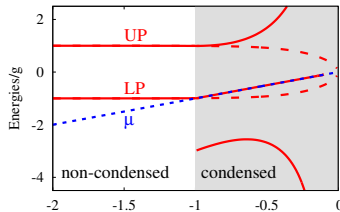
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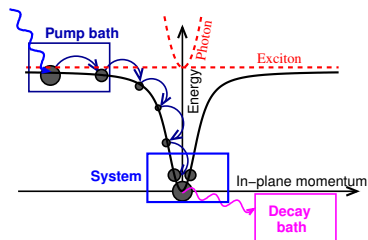
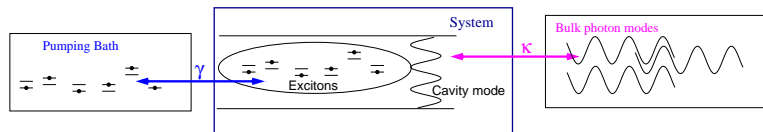
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Equilibrium:



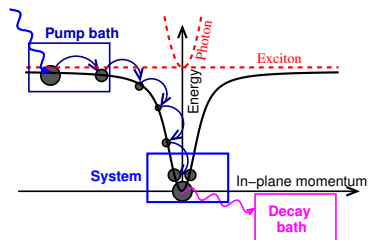
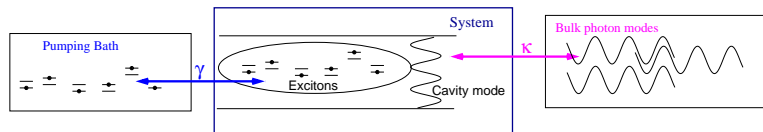
Non-equilibrium description: baths



$$H = H_{\text{sys}} + H_{\text{sys,bath}} + H_{\text{bath}}$$

- Decay bath: Empty ($\mu \rightarrow -\infty$)
- Pump bath: Thermal μ_B, T_B

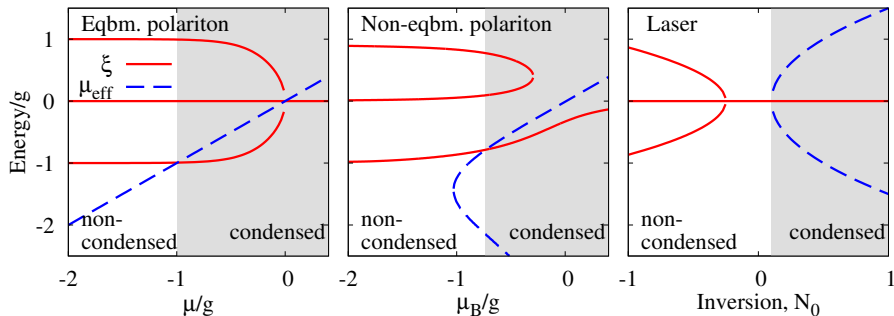
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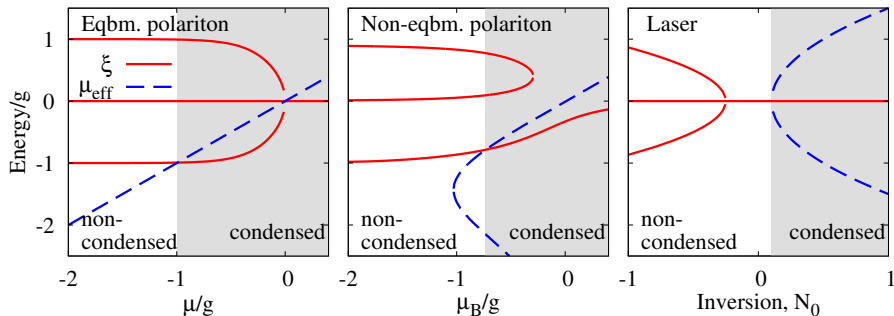
Strong coupling and lasing — low temperature phenomenon



- inversionless
- allows strong coupling
- requires low $T \leftrightarrow$ condensation
- Related weak-coupling inversionless lasing

[Szymanska *et al.* PRL '06; Keeling *et al.* book chapter 1010.3338]

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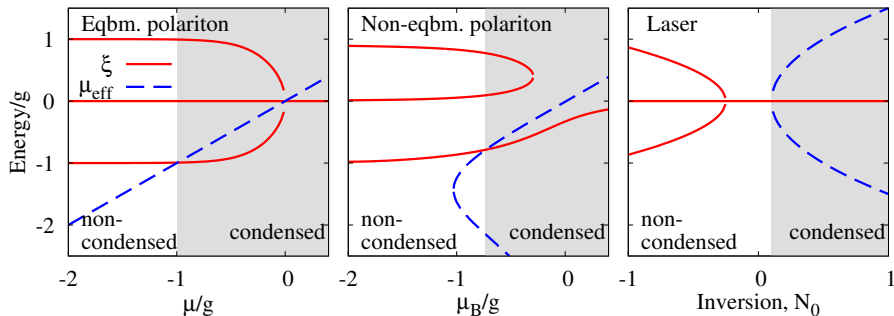


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Room temperature condensates: Organic polaritons

1 Condensation, superradiance, lasing

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2 Room temperature condensates: Organic polaritons

- Dicke phase diagram with phonons
- Condensation of phonon replicas?

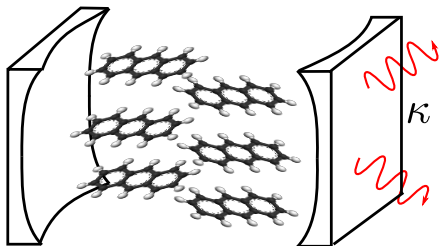
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Organic materials in microcavities

- What?

- Why?

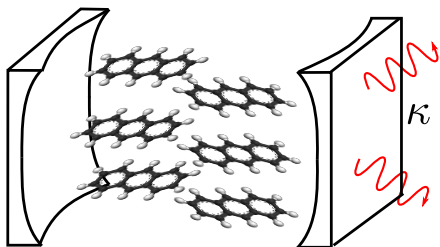


- Lasing threshold at room T

[Lidzey *et al.*, Nat. '98, Kena Cohen and Forrest, Nat. Photon '10; Plumhoff *et al.* Nat. Materials '14, Daskalakis *et al.* *ibid* '14]

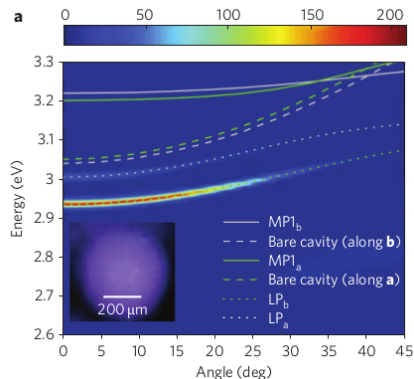
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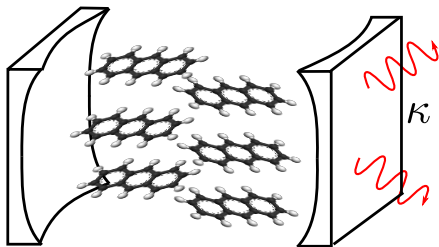


Polariton splitting: 0.1 eV \leftrightarrow 1000K.

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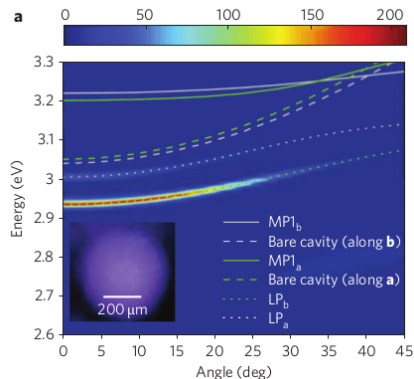
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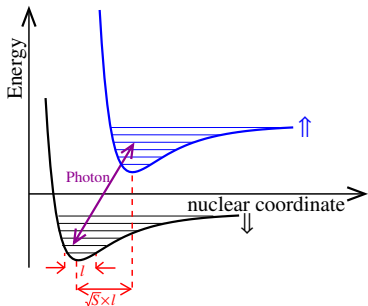
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Dicke Holstein Model

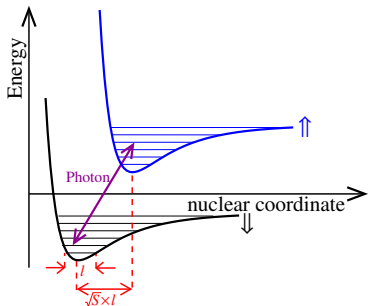


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- Phonon frequency Ω
- Huang-Rhys parameter S — phonon coupling

- Phase diagram with $S \neq 0$
 - 2LS energy $\epsilon - n\Omega$
- Polariton spectrum, phonon replicas
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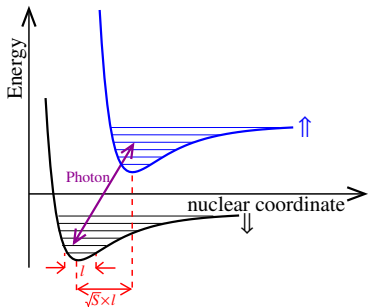


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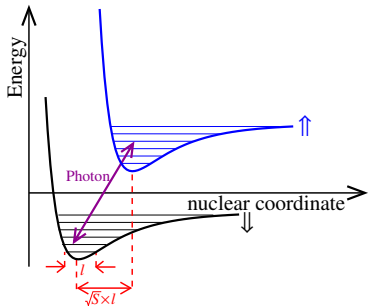
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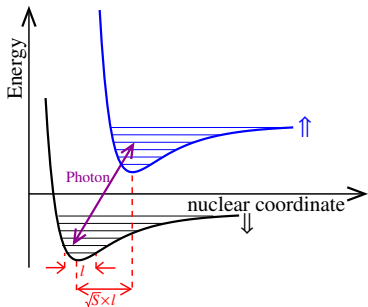
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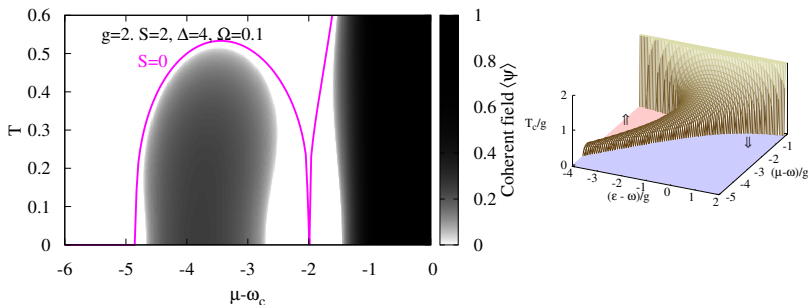
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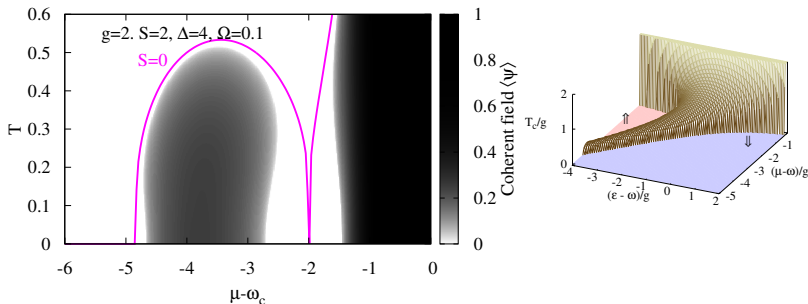


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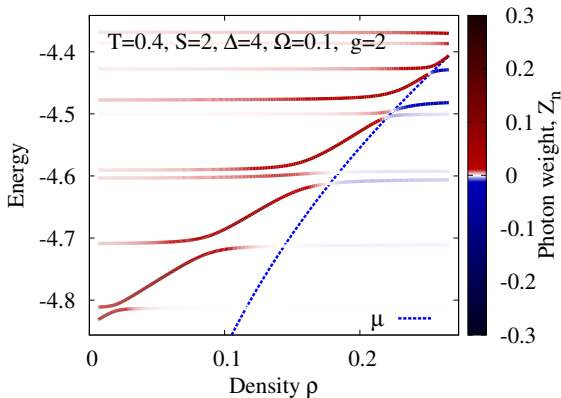
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Polariton spectrum: photon weight



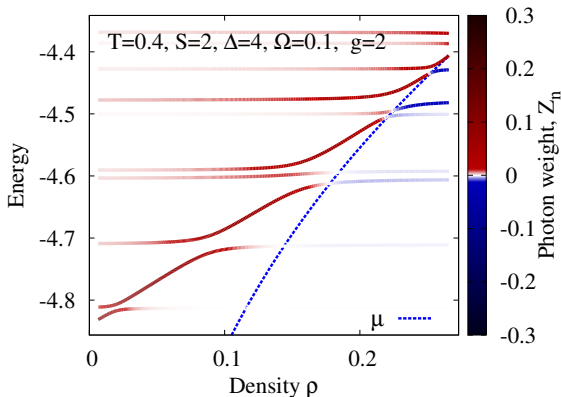
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• What is nature of polariton mode?

• $D(t) = -\langle \psi^\dagger(t)\psi(0) \rangle$, $D(\omega) = \sum_n \frac{Z_n}{\omega - \omega_n}$

[Cwik *et al.* EPL '14]

Polariton spectrum: photon weight

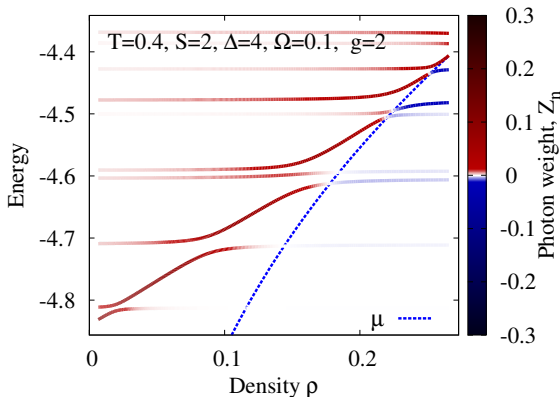


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- What is nature of polariton mode?

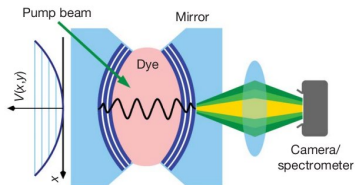
- $\mathcal{D}(t) = -i\langle \psi^\dagger(t)\psi(0) \rangle, \quad \mathcal{D}(\omega) = \sum_n \frac{Z_n}{\omega - \omega_n}$

[Cwik *et al.* EPL '14]

Room temperature condensates: Photons

- 1 Condensation, superradiance, lasing
 - Polariton condensation and Dicke model
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 - Time evolution
 - Linewidth

Photon BEC experiments

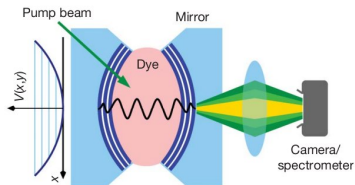


- Dye filled microcavity

➤ No strong coupling

[Klaers et al, Nature, 2010]

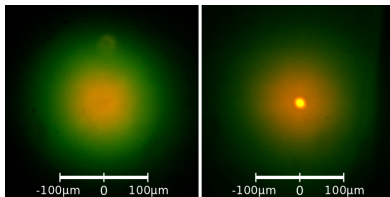
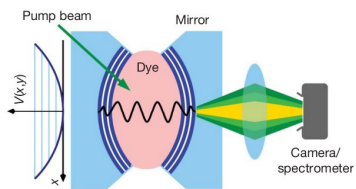
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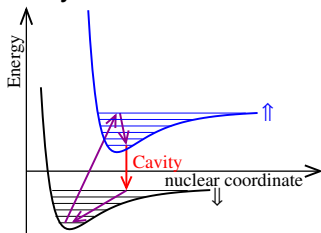
Relation to dye laser

- No electronic inversion
- No strong coupling
- No single cavity mode
 - ▶ Condensate mode is not maximum gain
 - ▶ Gain/Absorption in balance
- Thermalised many-mode system

Relation to dye laser

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4 Level Dye Laser

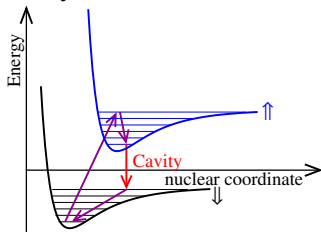


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But:

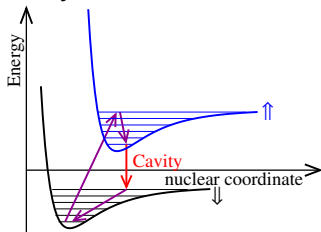
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Modelling

$$H_{\text{sys}} = \sum_m \omega_m \psi_m^\dagger \psi_m + \sum_\alpha \left[\frac{\epsilon}{2} \sigma_\alpha^z + g (\psi_m \sigma_\alpha^+ + \text{H.c.}) \right]$$

- 2D harmonic cavity

$$\omega_m = \omega_{\text{cutoff}} + m\omega_{H.O.}$$

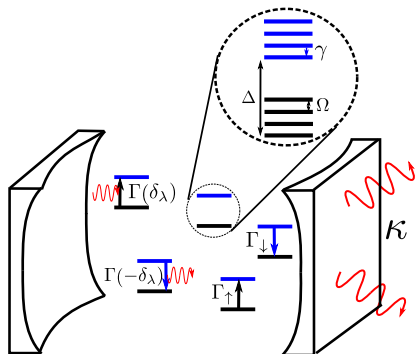
$$\text{Degeneracies } g_m = m + 1$$

- Local vibrational mode

 - Phonon frequency Ω

 - Huang-Rhys parameter S

 - phonon coupling



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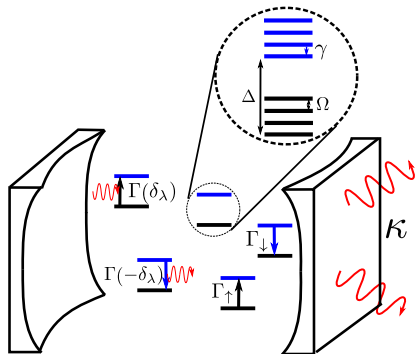
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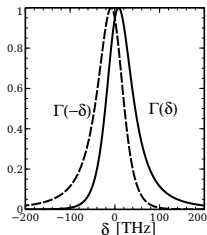
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Modelling

Master equation

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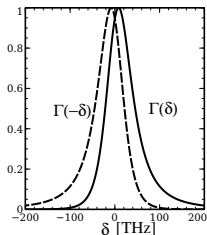
[Marthaler et al PRL '11, Kirton & JK PRL '13]

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 $\Gamma(+\delta) \simeq \Gamma(-\delta) e^{\beta\delta}$
- Expt: $\omega_0 < \epsilon$
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Distribution $g_m n_m$

- Master equation \rightarrow Rate equation

$$\partial_t n_m = -\kappa n_m + N [\Gamma(-\delta_m)(n_m + 1)\langle\sigma^{ee}\rangle - \Gamma(\delta_m)n_m\langle\sigma^{gg}\rangle]$$

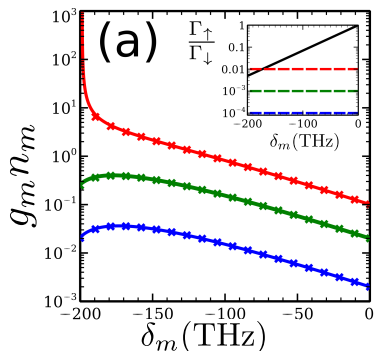
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Low loss: Thermal

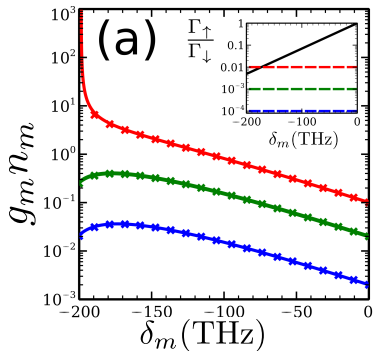
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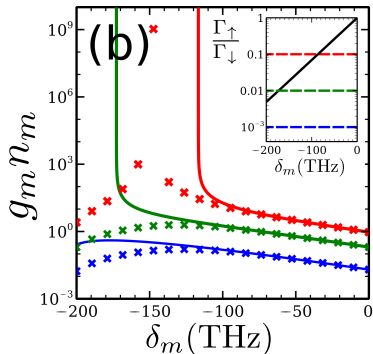
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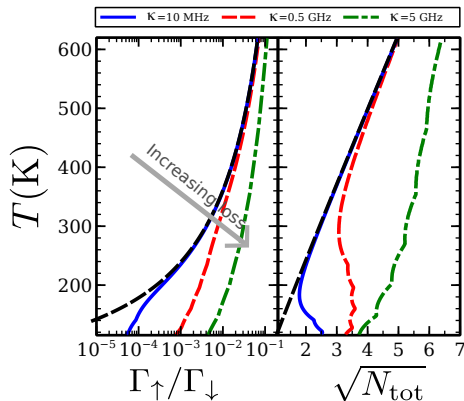
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[Kirton & JK PRL '13]



High loss \rightarrow Laser

Threshold condition



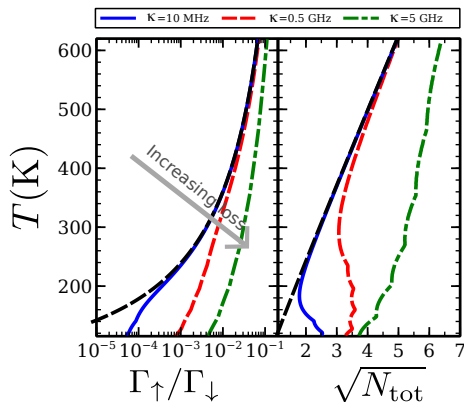
Compare threshold:

- Pump rate (Laser)
- Critical density (condensate)

- Thermal at low/high temperature
- High loss, κ competes with $\Gamma(\pm\delta_0)$
- Low temperature, $\Gamma(\pm\delta_0)$ shrinks

[Kirton & JK PRL '13]

Threshold condition



Compare threshold:

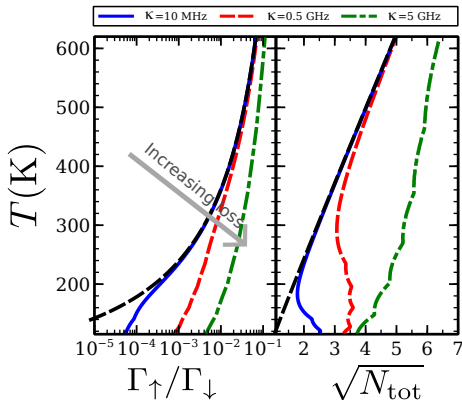
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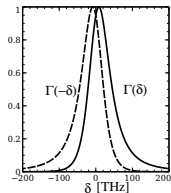
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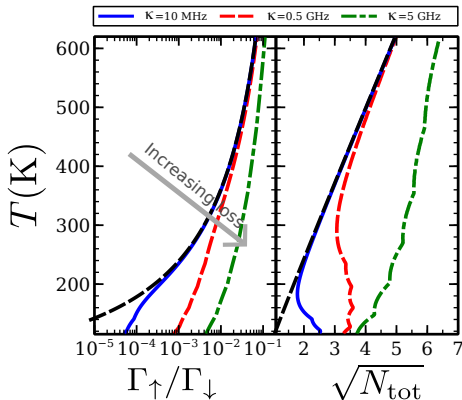
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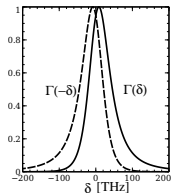


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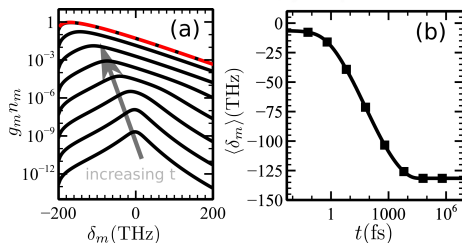
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Time evolution

- Initial state: excited molecules
- Initial emission, follows gain peak
- Thermalisation by repeated absorption
- Above threshold, jump to condensate



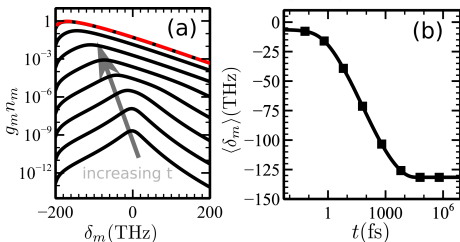
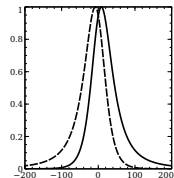
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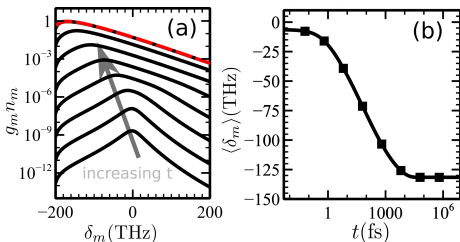
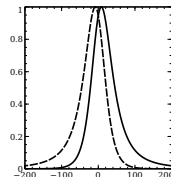


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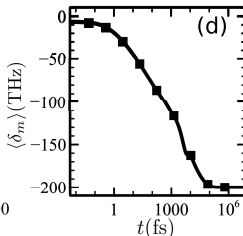
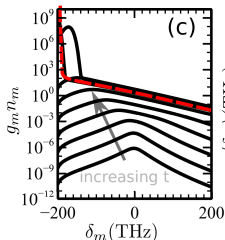
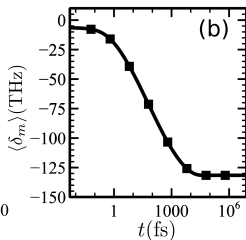
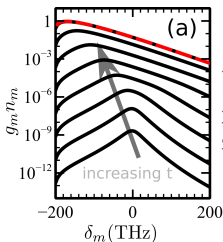
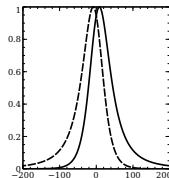
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Linewidth

- 1 Condensation, superradiance, lasing
 - Polariton condensation and Dicke model
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 - Dicke phase diagram with phonons
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Quantum model, linewidth

Full Master equation:

$$\dot{\rho} = -i[H_0, \rho] - \frac{\kappa}{2} \mathcal{L}[\psi] - \sum_{\alpha} \left[\frac{\Gamma_{\uparrow}}{2} \mathcal{L}[\sigma_{\alpha}^{+}] + \frac{\Gamma_{\downarrow}}{2} \mathcal{L}[\sigma_{\alpha}^{-}] \right] \\ - \sum_{\alpha} \left[\frac{\Gamma(\delta = \omega - \epsilon)}{2} \mathcal{L}[\sigma_{\alpha}^{+} \psi] + \frac{\Gamma(-\delta = \epsilon - \omega)}{2} \mathcal{L}[\sigma_{\alpha}^{-} \psi^{\dagger}] \right]$$

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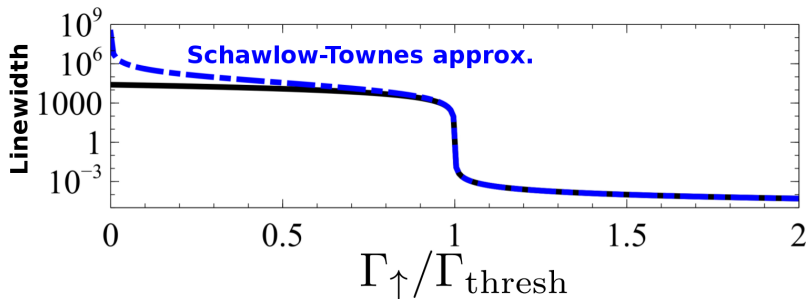
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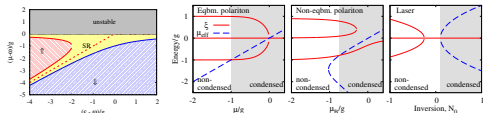
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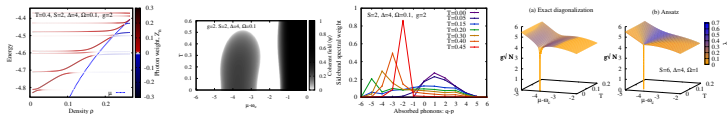


Summary

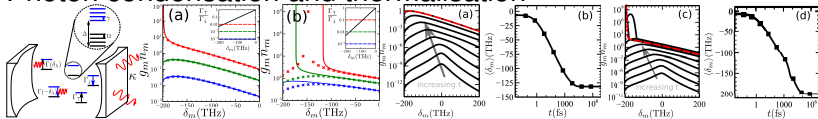
● Polariton condensation vs lasing; superradiance



● Reentrance, phonon assisted transition, 1st order at $S \gg 1$



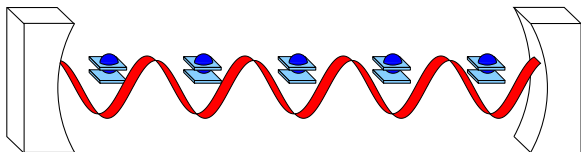
● Photon condensation and thermalisation



Extra slides

- 4 No go theorem
- 5 Dicke finite T
- 6 Retarded Green's function for laser
- 7 Organic properties
- 8 Ultra-strong phonon coupling?
- 9 Anticrossing vs ρ

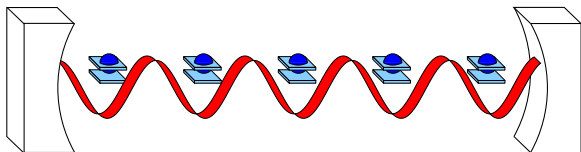
No go theorem and transition



Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

[Rzazewski *et al* PRL '75]

No go theorem and transition



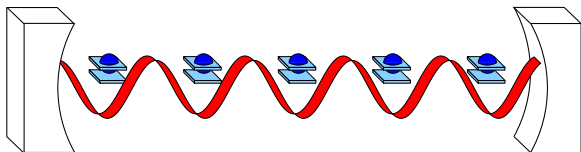
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No go theorem: Minimal coupling $(p - eA)^2/2m$

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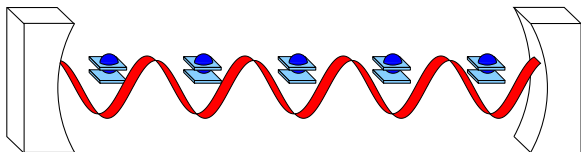
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For large N , $\omega \rightarrow \omega + 2N\zeta$. (RWA)

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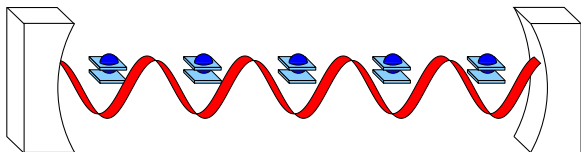
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But Thomas-Reiche-Kuhn sum rule states: $g^2/\epsilon < 2\zeta$. **No transition**

[Rzazewski *et al* PRL '75]

Dicke phase transition: ways out

Problem: $g^2/\omega_0 < 2\zeta$ for intrinsic parameters. **Solutions:**

- Interpretation
 - Ferroelectric transition in $\mathbf{D} \cdot \mathbf{r}$ gauge.
[JK JPCM '07, Vukics & Domokos PRA 2012]
 - Circuit QED [Nataf and Cluzet, Nat. Comm. '10; Viehmann *et al.* PRL '11]
- Grand canonical ensemble:
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e.g. Raman scheme: $\omega_0 \ll \omega$.
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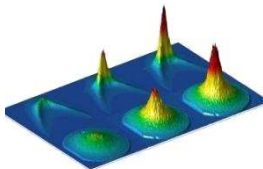
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- Circuit QED [Nataf and Ciuti, Nat. Comm. '10; Viehmann *et al.* PRL '11]

- Grand canonical ensemble:

- ▶ If $H \rightarrow H - \mu(\mathbf{S}^z + \psi^\dagger\psi)$, need only:
 $g^2 N > (\omega - \mu)(\omega_0 - \mu)$
- ▶ Incoherent pumping — polariton condensation.



- Dissociate g, ω_0 ,

e.g. Raman scheme: $\omega_0 \ll \omega$.

[Dimer *et al.* PRA '07; Baumann *et al.* Nature

'10. Also, Black *et al.* PRL '03]

Dicke phase transition: ways out

Problem: $g^2/\omega_0 < 2\zeta$ for intrinsic parameters. **Solutions:**

- **Interpretation**

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[JK JPCM '07, Vukics & Domokos PRA 2012]

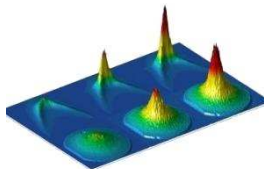
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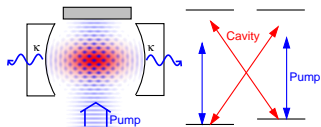
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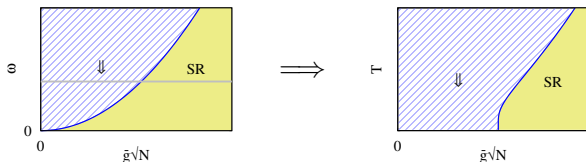
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Grand canonical Dicke, finite temperature

- Finite temperature:

$$Ng^2 \tanh(\beta\epsilon/2) > \omega\epsilon$$



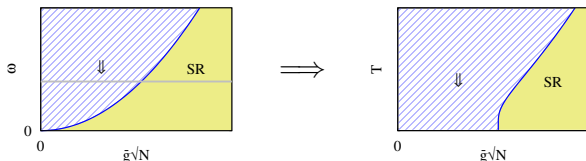
[Hepp, Lieb, Ann. Phys. '73]

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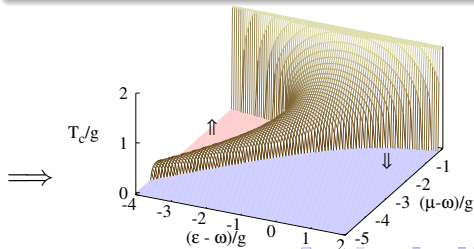
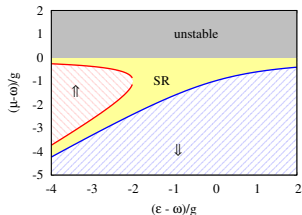
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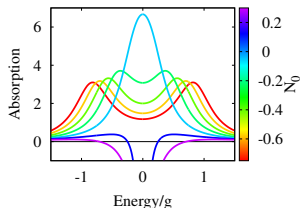
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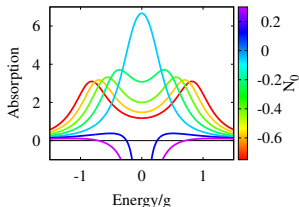


Maxwell-Bloch Equations: Retarded Green's function



- Introduce $D^R(\omega)$:
Response to perturbation
- Absorption = $-2\Im[D^R(\omega)]$

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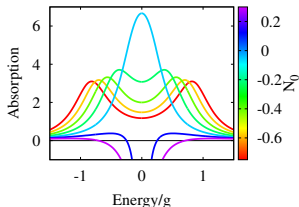
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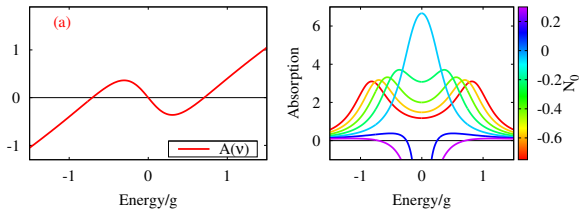
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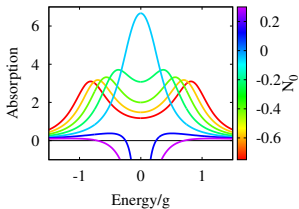
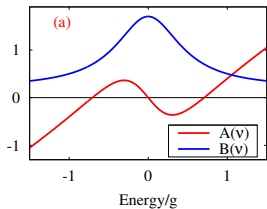
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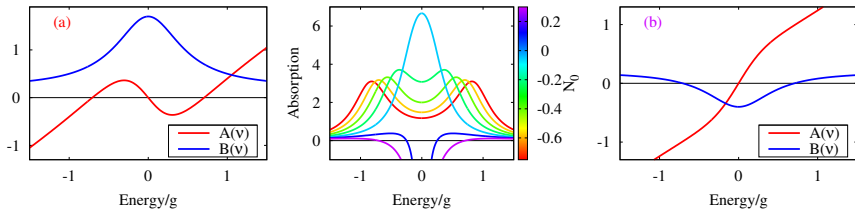
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Organic materials in microcavities

- State of art:
 - ▶ Strong coupling:
 - ★ J aggregates [Bulovic *et al.*]
 - ★ Crystalline anthracene [Forrest *et al.*]

▶ Threshold: Anthracene

[Kena Cohen and Forrest, Nat. Photon 2010]

● Differences

▶ Stronger coupling

Organic materials in microcavities

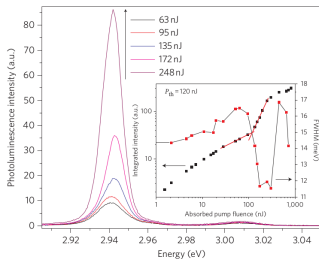
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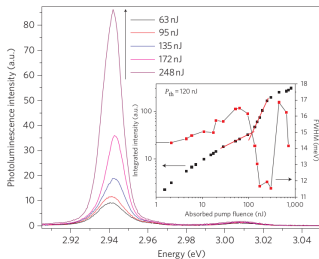
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- ▶ Singlet-Triplet conversion — dark states

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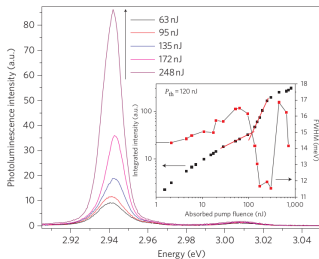
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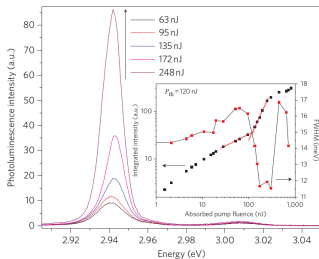
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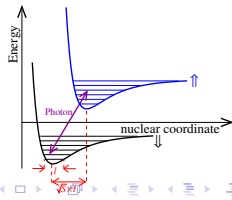
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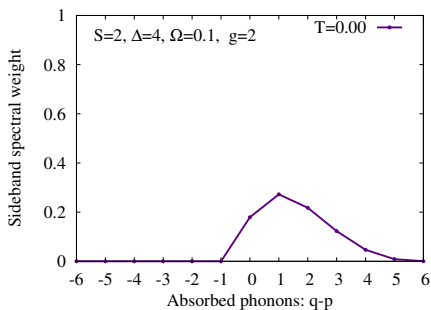
Polariton spectrum: what condensed

- Repeat weight for n -phonon channel
- Eigenvector that is macroscopically occupied
- Optimal $T \sim 2\Omega$

[Cwik *et al.* EPL '14]

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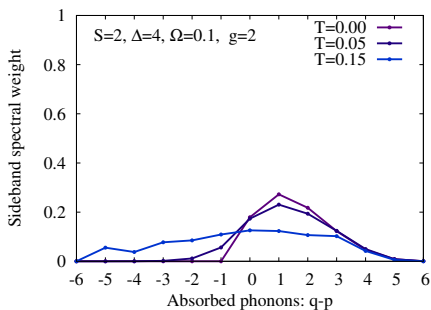


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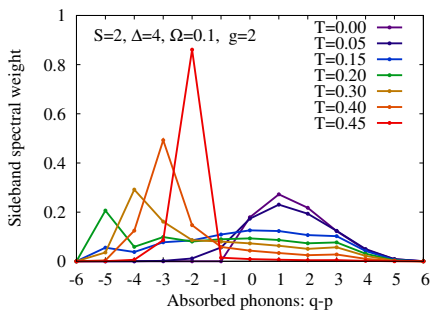


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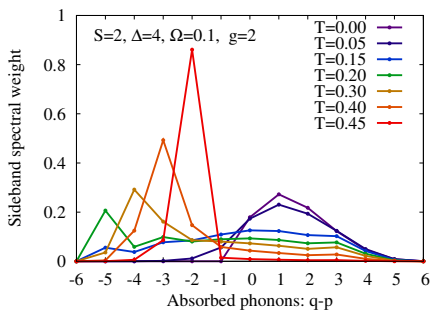


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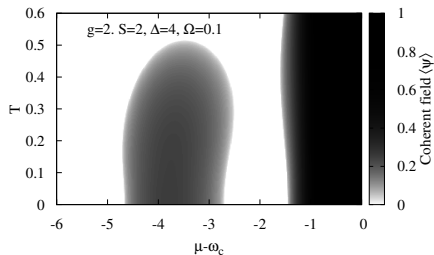
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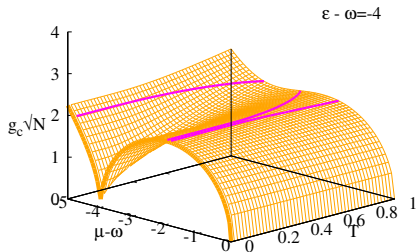
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Organic polaritons

- 4 No go theorem
- 5 Dicke finite T
- 6 Retarded Green's function for laser
- 7 Organic properties
- 8 Ultra-strong phonon coupling?
- 9 Anticrossing vs ρ**

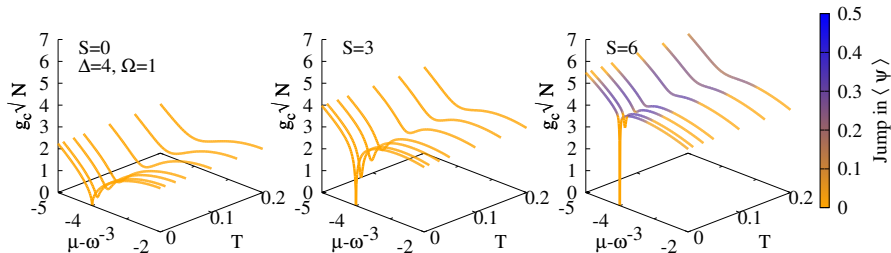
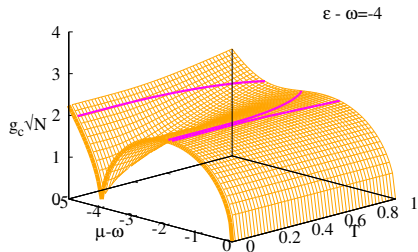
Critical coupling with increasing S

- Re-orient phase diagram
- g vs μ, T
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Explanation: Polaron formation

- Unitary transform

$$H_\alpha \rightarrow \tilde{H}_\alpha = e^{K_\alpha} H_\alpha e^{-K_\alpha} \quad K = \sqrt{S} S_\alpha^z (b_\alpha^\dagger - b_\alpha)$$

- Coupling moves to S^\pm

$$\tilde{H}_\alpha = \text{const.} + \epsilon S_\alpha^z + \Omega b_\alpha^\dagger b_\alpha + g \left[\psi S_\alpha^+ e^{\sqrt{S}(b_\alpha^\dagger - b_\alpha)} + \text{H.c.} \right]$$

- Optimal phonon displacements, $\sim \sqrt{S}$

- Reduced $g_{\text{eff}} \sim g \times \exp(-S/2)$

- For $\psi \neq 0$, competition

$$\text{Variational MFT } |\psi\rangle_\alpha \sim \exp(-\eta K_\alpha - \zeta (b_\alpha^\dagger)) |0, S\rangle_\alpha$$

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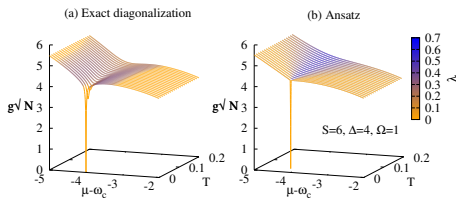
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Collective polaron formation

- Compares well at $S \gg 1$
- Coherent bosonic state



- Feedback: Large/small $g_{\text{eff}} \leftrightarrow \lambda = \langle \psi \rangle$
- Variational free energy

$$F = (\omega_c - \mu)\lambda^2 + N \left\{ \Omega \left[\zeta^2 - S \frac{\eta(2 - \eta)}{4} \right] - T \ln \left[2 \cosh \left(\frac{\xi}{T} \right) \right] \right\}$$

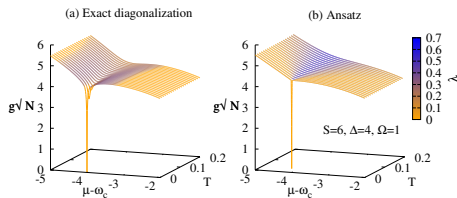
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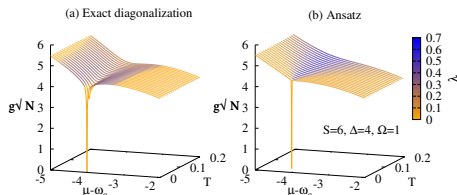
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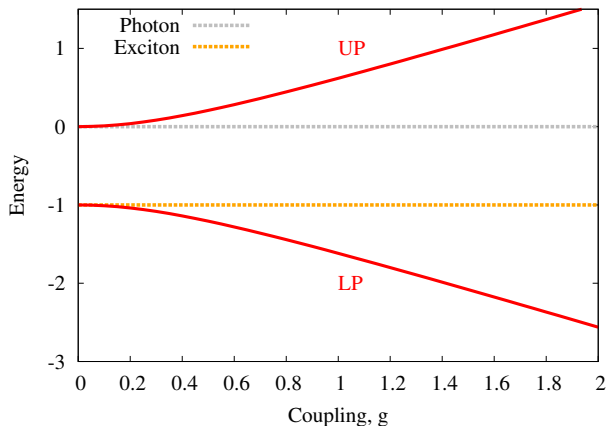
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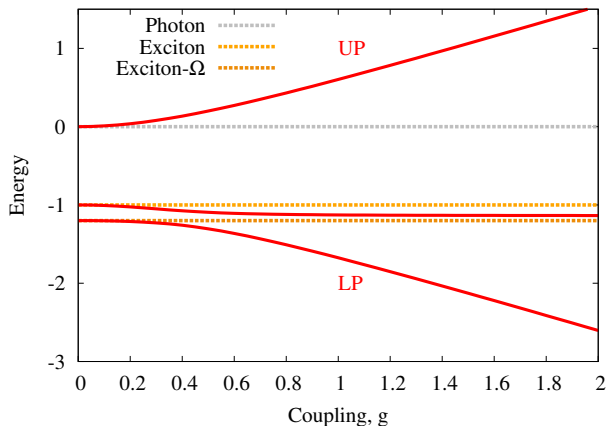
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Polariton spectrum — coupled oscillators

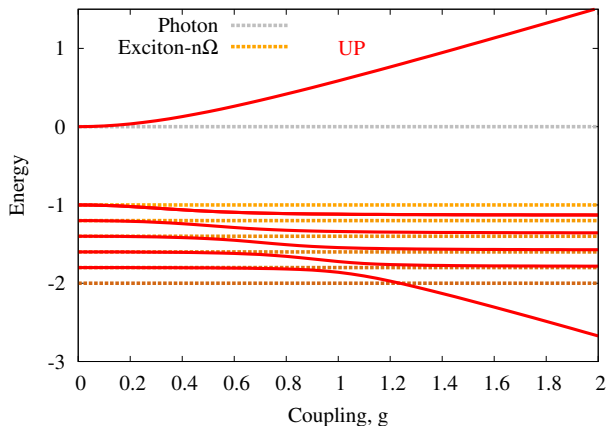
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