

Pairing Phases of Polaritons

Jonathan Keeling



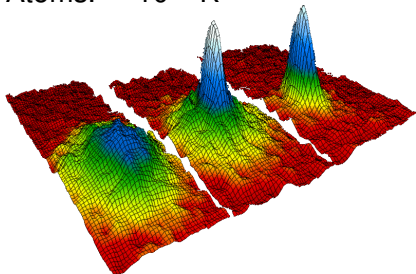
University of
St. Andrews

600
YEARS

Pisa, January 2014

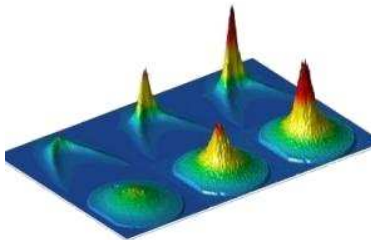
Bose-Einstein condensation: macroscopic occupation

Atoms. $\sim 10^{-7}\text{K}$



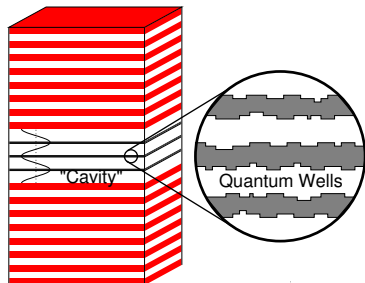
[Anderson *et al.* Science '95]

Polaritons. $\sim 20\text{K}$

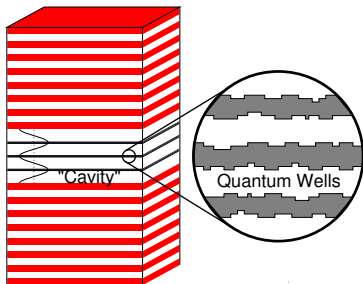


[Kasprzak *et al.* Nature, '06]

Microcavity polaritons

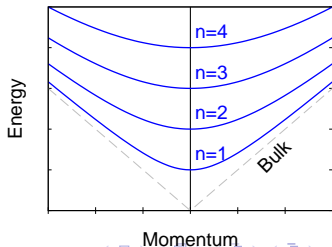


Microcavity polaritons

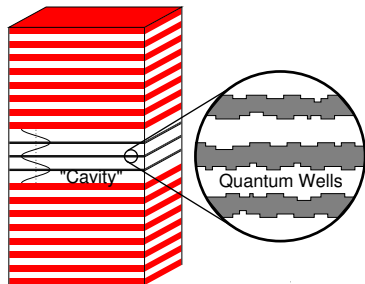


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2/2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$

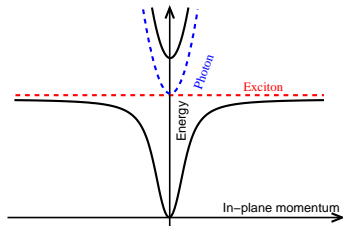


Microcavity polaritons

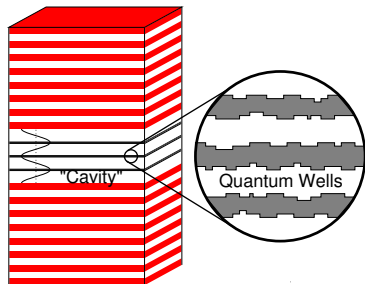


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2/2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$

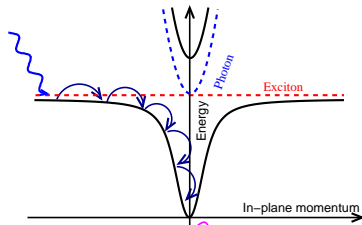


Microcavity polaritons

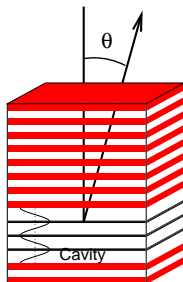
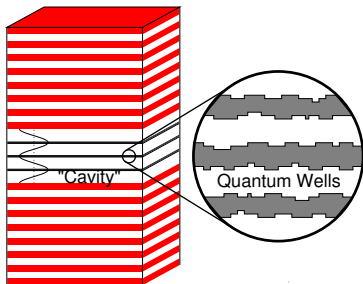


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + \frac{k^2}{2m^*} \\ m^* &\sim 10^{-4} m_e\end{aligned}$$

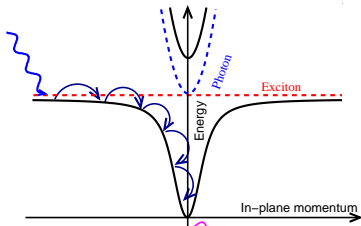


Microcavity polaritons

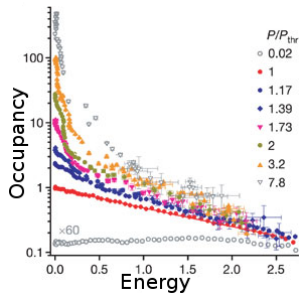
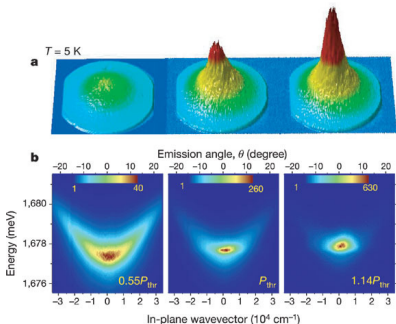


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + \frac{k^2}{2m^*} \\ m^* &\sim 10^{-4} m_e\end{aligned}$$

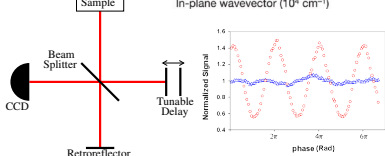
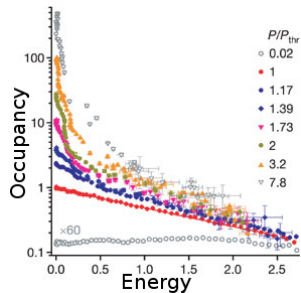
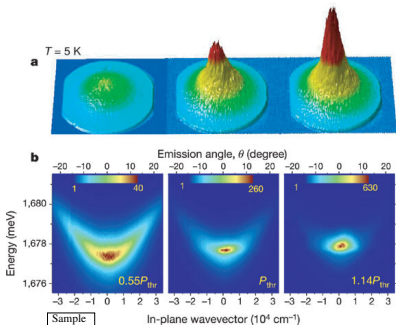


Polariton experiments: occupation and coherence

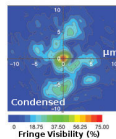
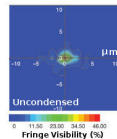
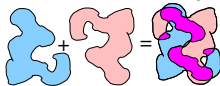


[Kasprzak, *et al.* Nature, '06]

Polariton experiments: occupation and coherence



Coherence map:

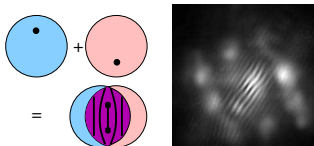


[Kasprzak, *et al.* Nature, '06]

(Some) other polariton condensation experiments

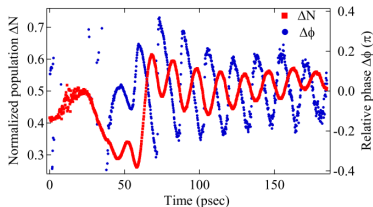
- Quantised vortices

[Lagoudakis *et al.* *Nat. Phys.* '08. *Science* '09, PRL '10; Sanvitto *et al.* *Nat. Phys.* '10; Roumpos *et al.* *Nat. Phys.* '10]



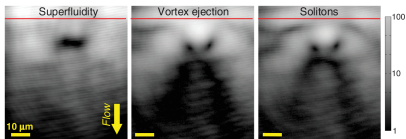
- Josephson oscillations

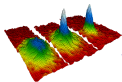
[Lagoudakis *et al.* PRL '10]



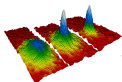
- Pattern formation/Hydrodynamics

[Amo *et al.* *Science* '11, *Nature* '09; Wertz *et al.* *Nat. Phys.* '10]



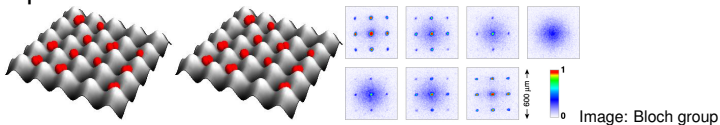


What can you do beyond BEC (atoms)

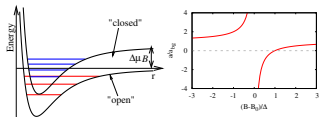


What can you do beyond BEC (atoms)

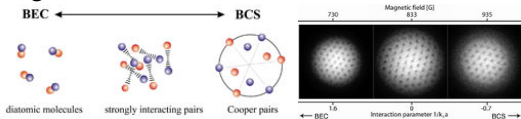
- Optical lattices



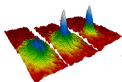
- Feshbach resonance — strong interactions



- Degenerate fermions — BEC BCS crossover

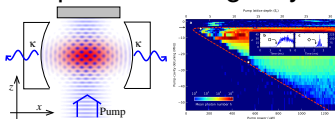


Images: Jin Group, Zweierlein group

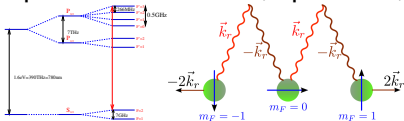


What can you do beyond BEC (atoms)

- Coupled matter-light systems



- Spinor condensates, spin-orbit, gauges



- Quenches, thermalisation, dynamics.

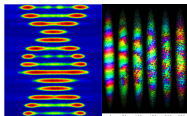
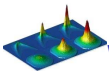


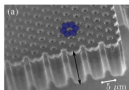
Image: Weiss group//Stamper Kurn group



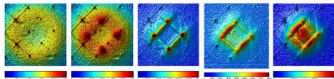
What of this can polaritons do?

- Engineered potentials; lattices

- ▶ Bloch — etched micropillars



- ▶ Yamamoto — metal lattices



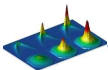
- ▶ Snoke — stress traps
- ▶ Krizhanovskii — Surface Acoustic Waves

- TE-TM interaction → Spin-Orbit

- ▶ Theory: Malpuech, Rubo

- “Feshbach” resonances

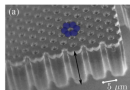
- ▶ Theory: Wouters, Carusotto
- ▶ Deveaud, Biexciton resonance



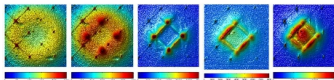
What of this can polaritons do?

- Engineered potentials; lattices

- ▶ Bloch — etched micropillars



- ▶ Yamamoto — metal lattices



- ▶ Snoke — stress traps
- ▶ Krizhanovskii — Surface Acoustic Waves

- TE-TM interaction → Spin-Orbit

- ▶ Theory: Malpuech, Rubo

- “Feshbach” resonances

- ▶ Theory: Wouters, Carusotto
- ▶ Deveaud, Biexciton resonance

This talk: *Spin structure, pairing phases of bosons*

Outline

- 1 Introduction
- 2 Pairing phases of atoms: review
- 3 Modelling polariton pairing
 - Exciton spin structure
 - Microscopic and effective Hamiltonian
- 4 Ground state phase diagram
 - Origin of multicritical behaviour
 - Signatures of phases
- 5 Finite temperature phase diagram
 - Variational MFT
 - Required temperatures, detuning
 - Candidate material systems

Acknowledgements

GROUP:



COLLABORATORS:



Francesca Marchetti, UAM

FUNDING:



EPSRC

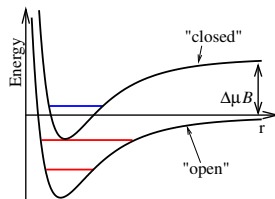
Engineering and Physical Sciences
Research Council

Outline

- 1 Introduction
- 2 Pairing phases of atoms: review
- 3 Modelling polariton pairing
 - Exciton spin structure
 - Microscopic and effective Hamiltonian
- 4 Ground state phase diagram
 - Origin of multicritical behaviour
 - Signatures of phases
- 5 Finite temperature phase diagram
 - Variational MFT
 - Required temperatures, detuning
 - Candidate material systems

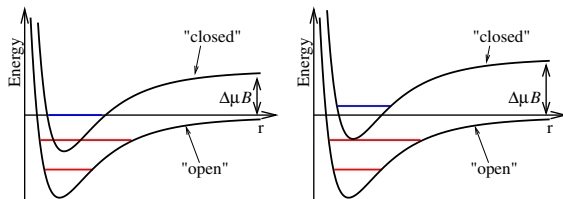
Tuning atomic interactions

- Open/closed channel — different F_z states ($F=J + I$)
- Magnetic field, $\Delta E(r \rightarrow \infty) = \text{const} + (\mu_a - \mu_b)B_z$
- Hybridisation by hyperfine coupling



Tuning atomic interactions

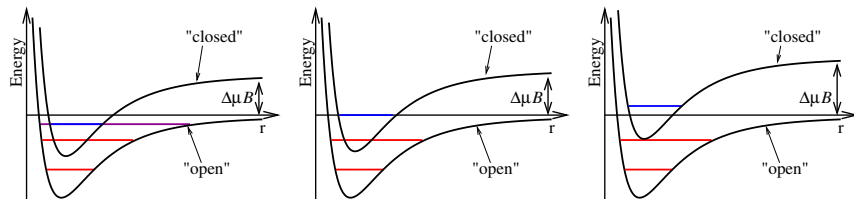
- Open/closed channel — different F_z states ($F=J+I$)
- Magnetic field, $\Delta E(r \rightarrow \infty) = \text{const} + (\mu_a - \mu_b)B_z$
- Hybridisation by hyperfine coupling



Tuning atomic interactions

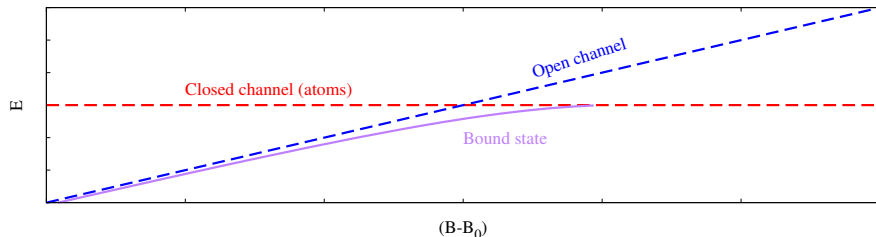
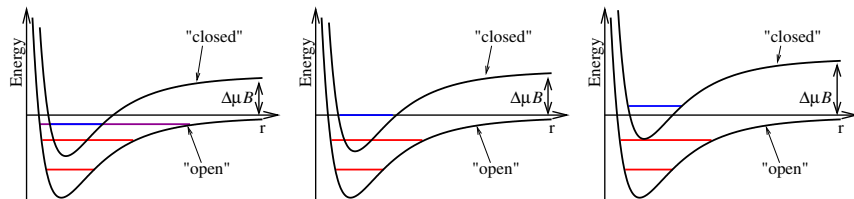
- Open/closed channel — different F_z states ($F=J+I$)
- Magnetic field, $\Delta E(r \rightarrow \infty) = \text{const} + (\mu_a - \mu_b)B_z$

• Hybridisation by hyperfine coupling

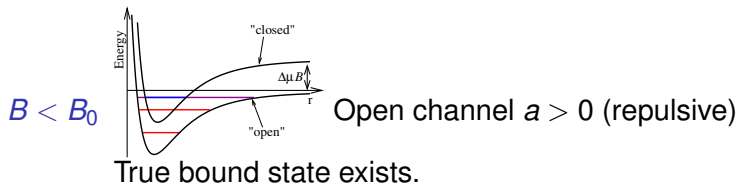


Tuning atomic interactions

- Open/closed channel — different F_z states ($F=J+I$)
- Magnetic field, $\Delta E(r \rightarrow \infty) = \text{const} + (\mu_a - \mu_b)B_z$
- Hybridisation by hyperfine coupling



Feshbach resonance

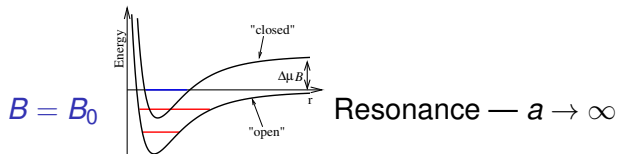
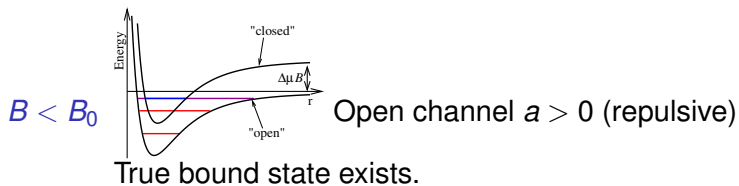


Resonance $\rightarrow a \rightarrow \infty$

Open channel $a < 0$ (attractive)

Hybridisation with continuum: Fano–Feshbach resonance

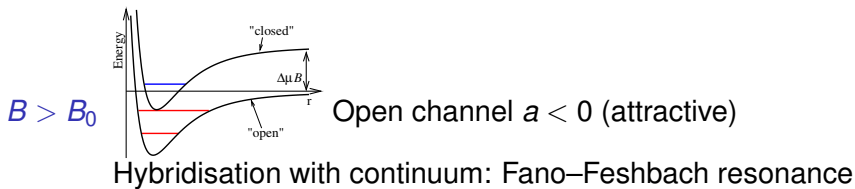
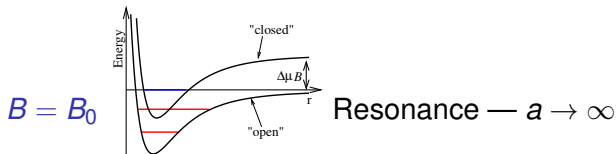
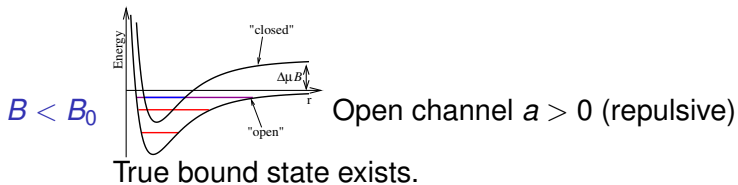
Feshbach resonance



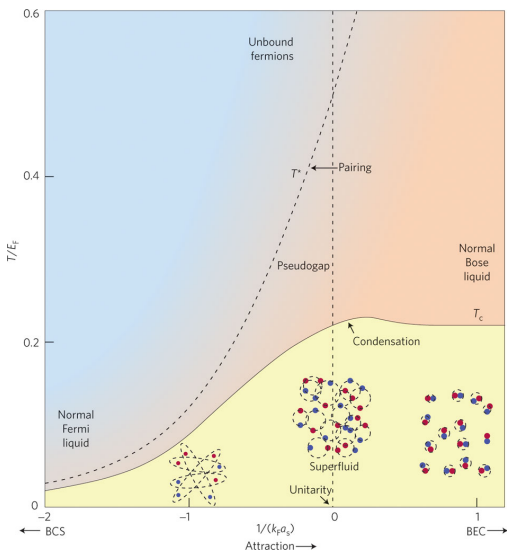
Open channel $a < 0$ (attractive)

Hybridisation with continuum: Fano–Feshbach resonance

Feshbach resonance



Pairing phases of Fermions

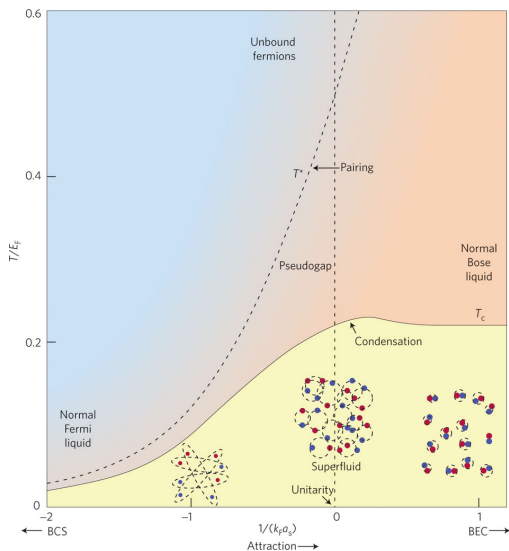


- Attractive interaction of open channel: BCS condensate.

- True bound state. Molecule BEC
- Always pair coherence: No phase transition (for s-wave).
- BEC-BCS crossover [Eagles, Leggett, Keldysh, Nozières, Randeria, ...]

From Randeria, Nat. Phys. News & Views '10

Pairing phases of Fermions

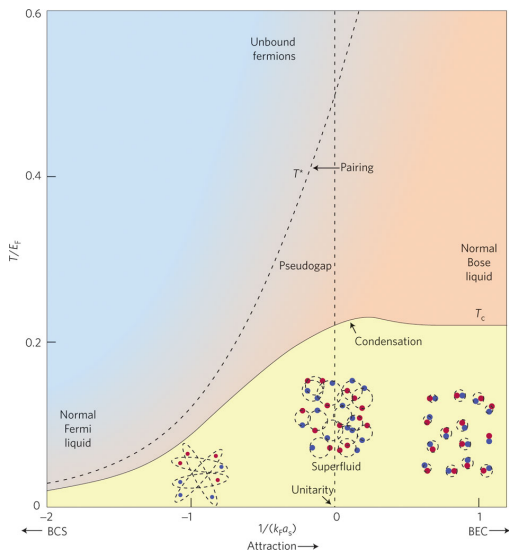


- Attractive interaction of open channel: BCS condensate.
- True bound state. Molecule BEC

- Always pair coherence: No phase transition (for s-wave).
- BEC-BCS crossover [Eagles, Leggett, Keldysh, Nozières, Randeria, ...]

From Randeria, Nat. Phys. News & Views '10

Pairing phases of Fermions

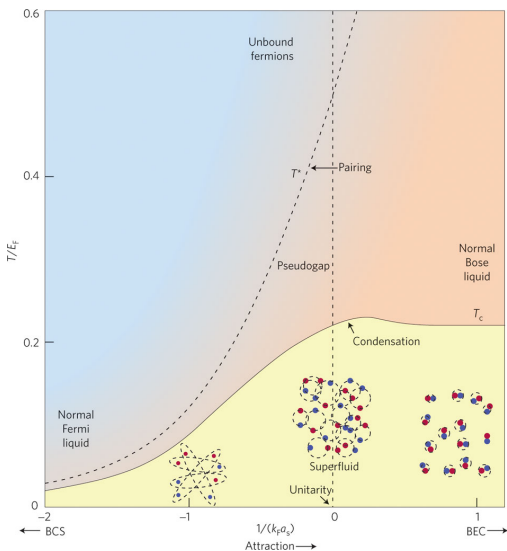


- Attractive interaction of open channel: BCS condensate.
- True bound state. Molecule BEC
- Always pair coherence: No phase transition (for s-wave).

• BEC-BCS crossover
[Eagles, Leggett, Keldysh, Nozières, Randeria, ...]

From Randeria, Nat. Phys. News & Views '10

Pairing phases of Fermions



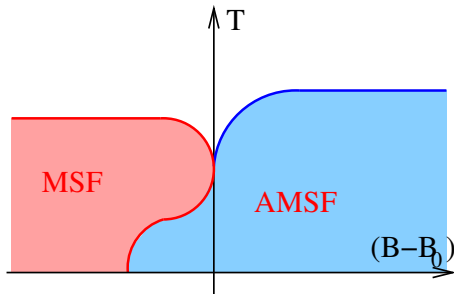
- Attractive interaction of open channel: BCS condensate.
- True bound state. Molecule BEC
- Always pair coherence: No phase transition (for s-wave).
- BEC-BCS crossover [Eagles, Leggett, Keldysh, Nozières, Randeria, ...]

From Randeria, Nat. Phys. News & Views '10

Pairing phases of Bosons

- Pair (molecule) coherence, MSF, vs atom coherence ASF, vs both AMSF.

● Homonuclear/heteronuclear cases distinct, symmetry $U(1) \times \mathbb{Z}_2$ vs $U(1) \times U(1)$.



● Heteronuclear: $\hat{H} = \dots + \hat{\psi}_m^\dagger \hat{\psi}_a \hat{\psi}_{a_2} + \text{h.c.}$

● If $\langle \hat{\psi}_m \rangle \neq 0$, free atomic phase, MSF

● If $\langle \hat{\psi}_a \rangle \neq 0, \langle \hat{\psi}_{a_2} \rangle \neq 0$ no free phase, AMSF

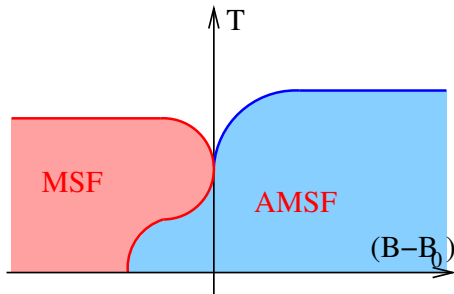
● Homonuclear: $\hat{H} = \dots + \hat{\psi}_m^\dagger \hat{\psi}_1 \hat{\psi}_2 + \text{h.c.}$

● If $\langle \hat{\psi}_m \rangle \neq 0$, free atomic sign $\rightarrow \mathbb{Z}_2$, MSF

[Radzihovsky *et al.* PRL '04, Romans *et al.* PRL '04, building on Nozières & St James, Timmermanns, Mueller, Thouless ...]

Pairing phases of Bosons

- Pair (molecule) coherence, MSF, vs atom coherence ASF, vs both AMSF.
- Homonuclear/heteronuclear cases distinct, symmetry $U(1) \times \mathbb{Z}_2$ vs $U(1) \times U(1)$.



• Heteronuclear: $\hat{H} = \dots + \hat{\psi}_m^\dagger \hat{\psi}_a \hat{\psi}_a^\dagger + \text{h.c.}$

• If $\langle \hat{\psi}_m \rangle \neq 0$, free atomic phase, MSF

• If $\langle \hat{\psi}_a \rangle \neq 0$, $\langle \hat{\psi}_m \rangle \neq 0$ no free phase, AMSF

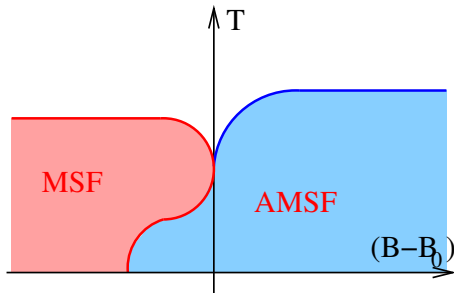
• Homonuclear: $\hat{H} = \dots + \hat{\psi}_m^\dagger \hat{\psi}_1 \hat{\psi}_2 + \text{h.c.}$

• If $\langle \hat{\psi}_m \rangle \neq 0$, free atomic sign $\rightarrow \mathbb{Z}_2$, MSF

[Radzihovsky *et al.* PRL '04, Romans *et al.* PRL '04, building on Nozières & St James, Timmermanns, Mueller, Thouless ...]

Pairing phases of Bosons

- Pair (molecule) coherence, MSF, vs atom coherence ASF, vs both AMSF.
- Homonuclear/heteronuclear cases distinct, symmetry $U(1) \times \mathbb{Z}_2$ vs $U(1) \times U(1)$.



- Heteronuclear: $\hat{H} = \dots + \hat{\psi}_m^\dagger \hat{\psi}_{a_1} \hat{\psi}_{a_2} + \text{h.c.}$
 - ▶ If $\langle \hat{\psi}_m \rangle \neq 0$, free atomic phase. MSF
 - ▶ If $\langle \hat{\psi}_{a_1} \rangle \neq 0, \langle \hat{\psi}_{a_2} \rangle \neq 0$ no free phase. AMSF

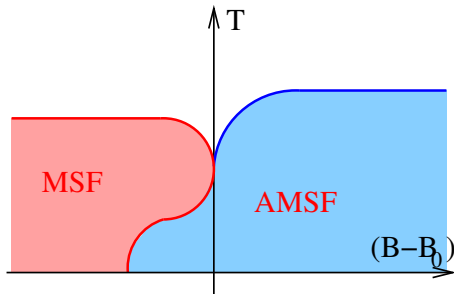
• Homonuclear: $\hat{H} = \dots + \hat{\psi}_m^\dagger \hat{\psi}_{a_1} \hat{\psi}_{a_2} + \text{h.c.}$

• If $\langle \hat{\psi}_m \rangle \neq 0$, free atomic sign $\rightarrow \mathbb{Z}_2$, MSF

[Radzihovsky *et al.* PRL '04, Romans *et al.* PRL '04, building on Nozières & St James, Timmermanns, Mueller, Thouless ...]

Pairing phases of Bosons

- Pair (molecule) coherence, MSF, vs atom coherence ASF, vs both AMSF.
- Homonuclear/heteronuclear cases distinct, symmetry $U(1) \times \mathbb{Z}_2$ vs $U(1) \times U(1)$.



- Heteronuclear: $\hat{H} = \dots + \hat{\psi}_m^\dagger \hat{\psi}_{a_1} \hat{\psi}_{a_2} + \text{h.c.}$
 - ▶ If $\langle \hat{\psi}_m \rangle \neq 0$, free atomic phase. MSF
 - ▶ If $\langle \hat{\psi}_{a_1} \rangle \neq 0, \langle \hat{\psi}_{a_2} \rangle \neq 0$ no free phase. AMSF
- Homonuclear: $\hat{H} = \dots + \hat{\psi}_m^\dagger \hat{\psi}_a \hat{\psi}_a + \text{h.c.}$
 - ▶ If $\langle \hat{\psi}_m \rangle \neq 0$, free atomic sign — \mathbb{Z}_2 . MSF

[Radzihovsky *et al.* PRL '04, Romans *et al.* PRL '04, building on Nozières & St James, Timmermanns, Mueller, Thouless ...]

Problems with cold atoms

- Signatures?

- ▶ Momentum distribution — indirect

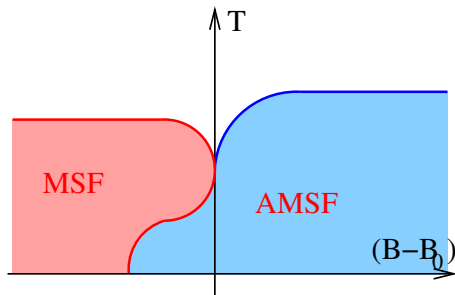
- ▶ Half vortices

- Cold atoms — (meta)stability issues

- ▶ Enhanced scattering at resonance

- ▶ First order transitions — high densities

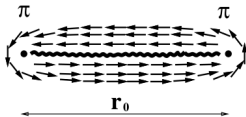
- ▶ Vibrationally hot molecules



Problems with cold atoms

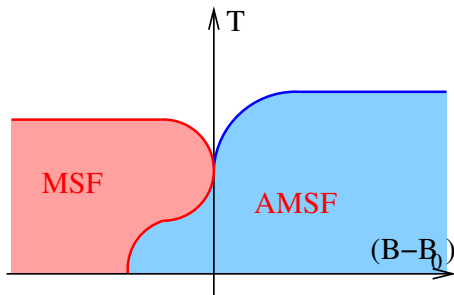
- Signatures?

- ▶ Momentum distribution — indirect
- ▶ Half vortices



- Cold atoms — (meta)stability issues

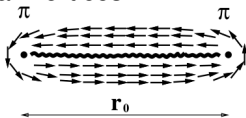
- ▶ Enhanced scattering at resonance
- ▶ First order transitions — high densities
- ▶ Vibrationally hot molecules



Problems with cold atoms

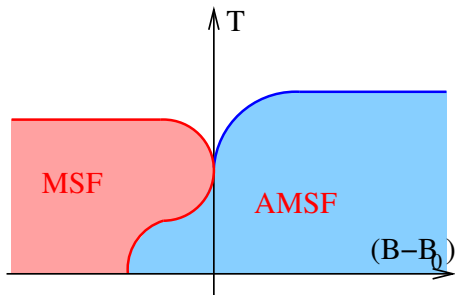
- Signatures?

- ▶ Momentum distribution — indirect
- ▶ Half vortices



- Cold atoms — (meta)stability issues

- ▶ Enhanced scattering at resonance
- ▶ First order transitions — high densities
- ▶ Vibrationally hot molecules



Outline

- 1 Introduction
- 2 Pairing phases of atoms: review
- 3 Modelling polariton pairing**
 - Exciton spin structure
 - Microscopic and effective Hamiltonian
- 4 Ground state phase diagram
 - Origin of multicritical behaviour
 - Signatures of phases
- 5 Finite temperature phase diagram
 - Variational MFT
 - Required temperatures, detuning
 - Candidate material systems

Exciton and polariton spin degrees of freedom

- Photon: two circular polarisation modes

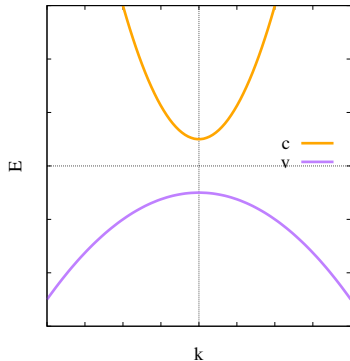
● Exciton: bound state of electron & hole

● Exciton spin states $J_z = +2, +1, -1, -2$

● Optically active states $J_z = \pm 1$

Exciton and polariton spin degrees of freedom

- Photon: two circular polarisation modes
- Exciton: bound state of electron & hole



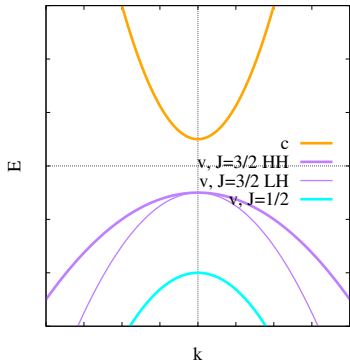
- ▶ $J = 1 \pm 1/2$ hole (p -orbital),
 $J = 1/2$ electron

- ▶ Spin orbit splits hole bands,
 4×2 states.
- ▶ Quantum well fixes k_z of hole
 2×2 states.

- Exciton spin states $J_z = +2, +1, -1, -2$
- Optically active states $J_z = \pm 1$

Exciton and polariton spin degrees of freedom

- Photon: two circular polarisation modes
- Exciton: bound state of electron & hole



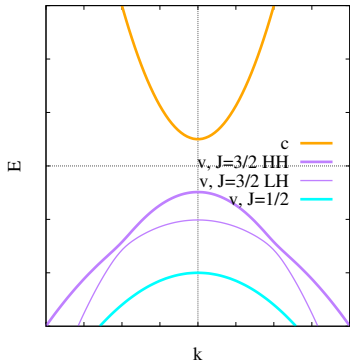
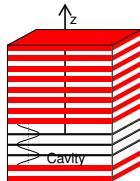
- ▶ $J = 1 \pm 1/2$ hole (p -orbital),
 $J = 1/2$ electron
- ▶ Spin orbit splits hole bands,
 4×2 states.

▶ Quantum well fixes k_y of hole
 2×2 states.

- Exciton spin states $J_z = +2, +1, -1, -2$
- Optically active states $J_z = \pm 1$

Exciton and polariton spin degrees of freedom

- Photon: two circular polarisation modes
- Exciton: bound state of electron & hole

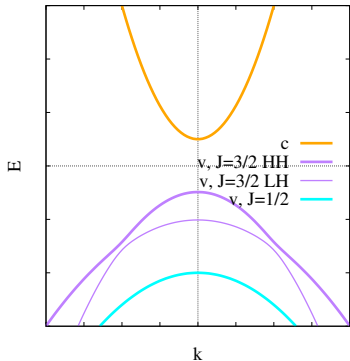
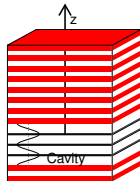


- ▶ $J = 1 \pm 1/2$ hole (p -orbital),
 $J = 1/2$ electron
- ▶ Spin orbit splits hole bands,
 4×2 states.
- ▶ Quantum well fixes k_z of hole
 2×2 states.

- Exciton spin states $J_x = +2, +1, -1, -2$
- Optically active states $J_x = \pm 1$

Exciton and polariton spin degrees of freedom

- Photon: two circular polarisation modes
- Exciton: bound state of electron & hole



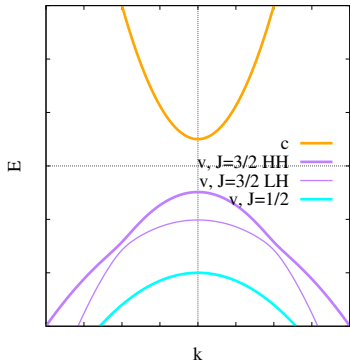
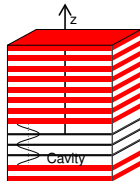
- ▶ $J = 1 \pm 1/2$ hole (p -orbital),
 $J = 1/2$ electron
- ▶ Spin orbit splits hole bands,
 4×2 states.
- ▶ Quantum well fixes k_z of hole
 2×2 states.

- Exciton spin states $J_z = +2, +1, -1, -2$

○ Optically active states $J_z = \pm 1$

Exciton and polariton spin degrees of freedom

- Photon: two circular polarisation modes
- Exciton: bound state of electron & hole



- ▶ $J = 1 \pm 1/2$ hole (p -orbital),
 $J = 1/2$ electron
- ▶ Spin orbit splits hole bands,
 4×2 states.
- ▶ Quantum well fixes k_z of hole
 2×2 states.

- Exciton spin states $J_z = +2, +1, -1, -2$
- Optically active states $J_z = \pm 1$

Exciton-photon model

- Microscopic model — coupled **exciton-photon** system

$$\begin{aligned}
 H = & \sum_k \left[\sum_{\sigma=\pm 2, \pm 1} \left(\frac{k^2}{2m_X} - \mu \right) \hat{X}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \sum_{\sigma=\pm 1} \left(\delta + \frac{k^2}{2m_C} - \mu \right) \hat{C}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right. \\
 & \left. + \sum_{\sigma=\pm 1} \Omega_R \left(\hat{C}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \hat{X}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right) \right] + \iint d^2 r d^2 R \sum_{\sigma, \sigma', \tau, \tau'=\pm 2, \pm 1} U_{\sigma'\tau'\tau\sigma}^{XX}(\mathbf{r}) \\
 & \times \hat{X}_{\sigma'}^\dagger \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau'}^\dagger \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau} \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\sigma} \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right)
 \end{aligned}$$

- Interaction U^{XX} has exchange structure
- For large Ω_R , neglect $\sigma = \pm 2$
- Interaction supports bound states in $U_{+1, -1, -1, +1}^{XX}$ channel — *bipolariton*
- NB, bipolariton, bound polaritons, *but* larger exciton fraction.

Exciton-photon model

- Microscopic model — coupled **exciton-photon** system

$$\begin{aligned}
 H = & \sum_k \left[\sum_{\sigma=\pm 2, \pm 1} \left(\frac{k^2}{2m_X} - \mu \right) \hat{X}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \sum_{\sigma=\pm 1} \left(\delta + \frac{k^2}{2m_C} - \mu \right) \hat{C}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right. \\
 & \left. + \sum_{\sigma=\pm 1} \Omega_R \left(\hat{C}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \hat{X}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right) \right] + \iint d^2 r d^2 R \sum_{\sigma, \sigma', \tau, \tau'=\pm 2, \pm 1} U_{\sigma'\tau'\tau\sigma}^{XX}(\mathbf{r}) \\
 & \times \hat{X}_{\sigma'}^\dagger \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau'}^\dagger \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau} \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\sigma} \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right)
 \end{aligned}$$

- Interaction U^{XX} has exchange structure

• For large Ω_R , neglect $\sigma = \pm 2$

• Interaction supports bound states in $U_{+1,-1,-1,+1}^{XX}$ channel — *bipolariton*

• NB, bipolariton, bound polaritons, but larger exciton fraction.

Exciton-photon model

- Microscopic model — coupled **exciton-photon** system

$$\begin{aligned}
 H = & \sum_k \left[\sum_{\sigma=\pm 2, \pm 1} \left(\frac{k^2}{2m_X} - \mu \right) \hat{X}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \sum_{\sigma=\pm 1} \left(\delta + \frac{k^2}{2m_C} - \mu \right) \hat{C}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right. \\
 & \left. + \sum_{\sigma=\pm 1} \Omega_R \left(\hat{C}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \hat{X}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right) \right] + \iint d^2 r d^2 R \sum_{\sigma, \sigma', \tau, \tau'=\pm 2, \pm 1} U_{\sigma' \tau' \tau \sigma}^{XX}(\mathbf{r}) \\
 & \times \hat{X}_{\sigma'}^\dagger \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau'}^\dagger \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau} \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\sigma} \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right)
 \end{aligned}$$

- Interaction U^{XX} has exchange structure
- For large Ω_R , neglect $\sigma = \pm 2$

• Interaction supports bound states in $U_{-1, -1, -1, +1}^{XX}$ channel — *bipolariton*

• NB, bipolariton, bound polaritons, but larger exciton fraction.

Exciton-photon model

- Microscopic model — coupled **exciton-photon** system

$$\begin{aligned}
 H = & \sum_k \left[\sum_{\sigma=\pm 2, \pm 1} \left(\frac{k^2}{2m_X} - \mu \right) \hat{X}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \sum_{\sigma=\pm 1} \left(\delta + \frac{k^2}{2m_C} - \mu \right) \hat{C}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right. \\
 & \left. + \sum_{\sigma=\pm 1} \Omega_R \left(\hat{C}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \hat{X}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right) \right] + \iint d^2 r d^2 R \sum_{\sigma, \sigma', \tau, \tau'=\pm 2, \pm 1} U_{\sigma'\tau'\tau\sigma}^{XX}(\mathbf{r}) \\
 & \times \hat{X}_{\sigma'}^\dagger \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau'}^\dagger \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau} \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\sigma} \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right)
 \end{aligned}$$

- Interaction U^{XX} has exchange structure
- For large Ω_R , neglect $\sigma = \pm 2$
- Interaction supports bound states in $U_{+1, -1, -1, +1}^{XX}$ channel — *bipolariton*

• NB, bipolariton, bound polaritons, but larger exciton fraction.

Exciton-photon model

- Microscopic model — coupled **exciton-photon** system

$$\begin{aligned}
 H = & \sum_k \left[\sum_{\sigma=\pm 2, \pm 1} \left(\frac{k^2}{2m_X} - \mu \right) \hat{X}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \sum_{\sigma=\pm 1} \left(\delta + \frac{k^2}{2m_C} - \mu \right) \hat{C}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right. \\
 & \left. + \sum_{\sigma=\pm 1} \Omega_R \left(\hat{C}_{k\sigma}^\dagger \hat{X}_{k\sigma} + \hat{X}_{k\sigma}^\dagger \hat{C}_{k\sigma} \right) \right] + \iint d^2 r d^2 R \sum_{\sigma, \sigma', \tau, \tau'=\pm 2, \pm 1} U_{\sigma'\tau'\tau\sigma}^{XX}(\mathbf{r}) \\
 & \times \hat{X}_{\sigma'}^\dagger \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau'}^\dagger \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\tau} \left(\mathbf{R} - \frac{\mathbf{r}}{2} \right) \hat{X}_{\sigma} \left(\mathbf{R} + \frac{\mathbf{r}}{2} \right)
 \end{aligned}$$

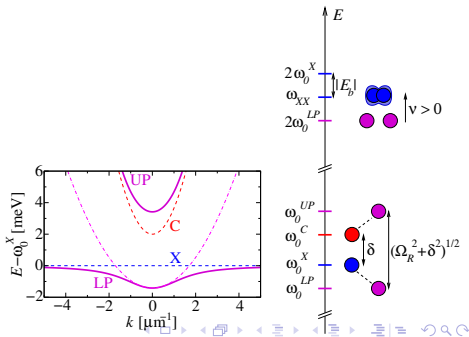
- Interaction U^{XX} has exchange structure
- For large Ω_R , neglect $\sigma = \pm 2$
- Interaction supports bound states in $U_{+1, -1, -1, +1}^{XX}$ channel — *bipolariton*
- NB, bipolariton, bound polaritons, *but* larger exciton fraction.

Polariton model

$$H = \sum_k \left[\sum_{\sigma=\uparrow,\downarrow} \left(\frac{k^2}{2m} - \mu \right) \hat{\psi}_{\sigma k}^\dagger \hat{\psi}_{\sigma k} + \left(\frac{k^2}{2m_m} + \nu - 2\mu \right) \hat{\psi}_{mk}^\dagger \hat{\psi}_{mk} \right]$$

$$+ \int d^2R \left[\sum_{\sigma=\uparrow,\downarrow,m} \frac{U_{\sigma\sigma}}{2} \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma \hat{\psi}_\sigma + U_{\uparrow\downarrow} \hat{\psi}_\uparrow^\dagger \hat{\psi}_\downarrow^\dagger \hat{\psi}_\uparrow \hat{\psi}_\downarrow + \frac{g}{2} \left(\hat{\psi}_\uparrow^\dagger \hat{\psi}_\downarrow^\dagger \hat{\psi}_m + \text{h.c.} \right) \right]$$

- Polariton dispersion m , detuning ν , interactions depend on δ
- Resonance width, dispersion derived from dressed exciton T matrix.



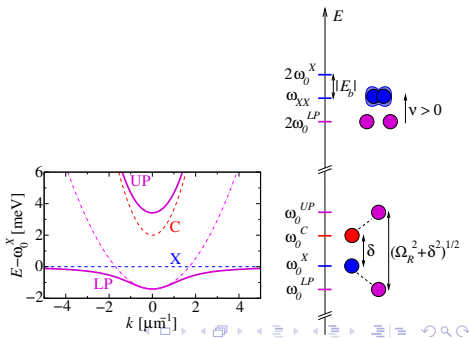
Polariton model

$$H = \sum_k \left[\sum_{\sigma=\uparrow,\downarrow} \left(\frac{k^2}{2m} - \mu \right) \hat{\psi}_{\sigma k}^\dagger \hat{\psi}_{\sigma k} + \left(\frac{k^2}{2m_m} + \nu - 2\mu \right) \hat{\psi}_{mk}^\dagger \hat{\psi}_{mk} \right]$$

$$+ \int d^2R \left[\sum_{\sigma=\uparrow,\downarrow,m} \frac{U_{\sigma\sigma}}{2} \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma \hat{\psi}_\sigma + U_{\uparrow\downarrow} \hat{\psi}_\downarrow^\dagger \hat{\psi}_\downarrow^\dagger \hat{\psi}_\uparrow \hat{\psi}_\downarrow + \frac{g}{2} \left(\hat{\psi}_\uparrow^\dagger \hat{\psi}_\downarrow^\dagger \hat{\psi}_m + \text{h.c.} \right) \right]$$

- Polariton dispersion m , detuning ν , interactions depend on δ

● Resonance width, dispersion derived from dressed exciton T matrix.



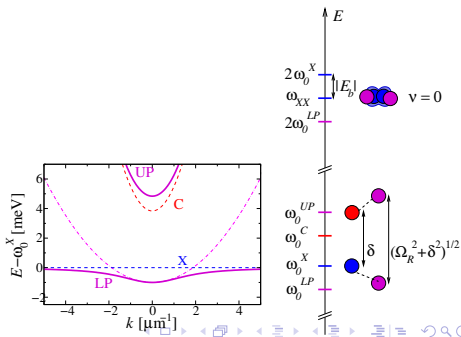
Polariton model

$$H = \sum_k \left[\sum_{\sigma=\uparrow,\downarrow} \left(\frac{k^2}{2m} - \mu \right) \hat{\psi}_{\sigma k}^\dagger \hat{\psi}_{\sigma k} + \left(\frac{k^2}{2m_m} + \nu - 2\mu \right) \hat{\psi}_{mk}^\dagger \hat{\psi}_{mk} \right]$$

$$+ \int d^2R \left[\sum_{\sigma=\uparrow,\downarrow,m} \frac{U_{\sigma\sigma}}{2} \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma \hat{\psi}_\sigma + U_{\uparrow\downarrow} \hat{\psi}_\downarrow^\dagger \hat{\psi}_\uparrow^\dagger \hat{\psi}_\uparrow \hat{\psi}_\downarrow + \frac{g}{2} \left(\hat{\psi}_\uparrow^\dagger \hat{\psi}_\downarrow^\dagger \hat{\psi}_m + \text{h.c.} \right) \right]$$

- Polariton dispersion m , detuning ν , interactions depend on δ

● Resonance width, dispersion derived from dressed exciton T matrix.

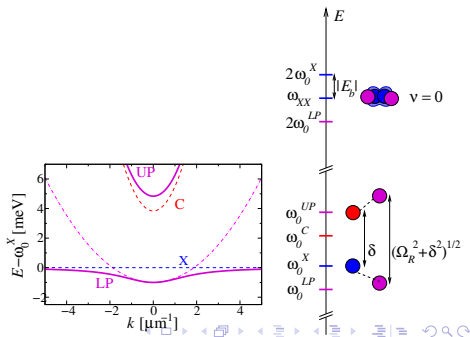


Polariton model

$$H = \sum_k \left[\sum_{\sigma=\uparrow,\downarrow} \left(\frac{k^2}{2m} - \mu \right) \hat{\psi}_{\sigma k}^\dagger \hat{\psi}_{\sigma k} + \left(\frac{k^2}{2m_m} + \nu - 2\mu \right) \hat{\psi}_{mk}^\dagger \hat{\psi}_{mk} \right]$$

$$+ \int d^2R \left[\sum_{\sigma=\uparrow,\downarrow,m} \frac{U_{\sigma\sigma}}{2} \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma \hat{\psi}_\sigma + U_{\uparrow\downarrow} \hat{\psi}_\downarrow^\dagger \hat{\psi}_\uparrow^\dagger \hat{\psi}_\uparrow \hat{\psi}_\downarrow + \frac{g}{2} \left(\hat{\psi}_\uparrow^\dagger \hat{\psi}_\downarrow^\dagger \hat{\psi}_m + \text{h.c.} \right) \right]$$

- Polariton dispersion m , detuning ν , interactions depend on δ
- Resonance width, dispersion derived from dressed exciton T matrix.

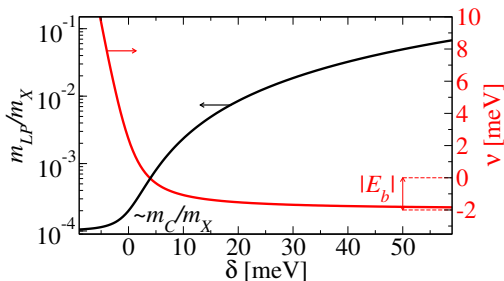


Polariton model

$$H = \sum_k \left[\sum_{\sigma=\uparrow,\downarrow} \left(\frac{k^2}{2m} - \mu \right) \hat{\psi}_{\sigma k}^\dagger \hat{\psi}_{\sigma k} + \left(\frac{k^2}{2m_m} + \nu - 2\mu \right) \hat{\psi}_{mk}^\dagger \hat{\psi}_{mk} \right]$$

$$+ \int d^2R \left[\sum_{\sigma=\uparrow,\downarrow,m} \frac{U_{\sigma\sigma}}{2} \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma^\dagger \hat{\psi}_\sigma \hat{\psi}_\sigma + U_{\uparrow\downarrow} \hat{\psi}_\downarrow^\dagger \hat{\psi}_\downarrow^\dagger \hat{\psi}_\uparrow \hat{\psi}_\downarrow + \frac{g}{2} \left(\hat{\psi}_\uparrow^\dagger \hat{\psi}_\downarrow^\dagger \hat{\psi}_m + \text{h.c.} \right) \right]$$

$\nu > 0$
Biexciton in
continuum

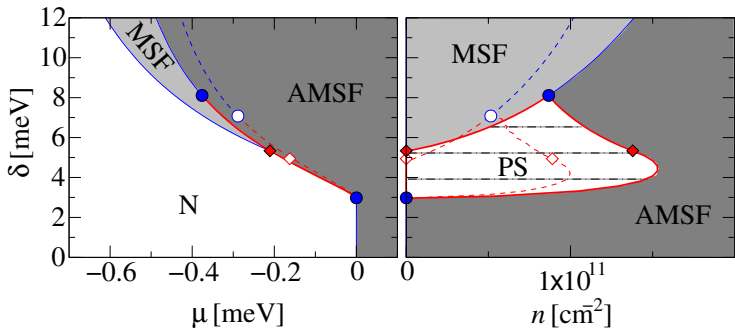


$\nu < 0$
Bound biexciton.
(Excitonic limit)

Outline

- 1 Introduction
- 2 Pairing phases of atoms: review
- 3 Modelling polariton pairing
 - Exciton spin structure
 - Microscopic and effective Hamiltonian
- 4 Ground state phase diagram**
 - Origin of multicritical behaviour
 - Signatures of phases
- 5 Finite temperature phase diagram
 - Variational MFT
 - Required temperatures, detuning
 - Candidate material systems

Phase diagram (ground state, $T = 0$)



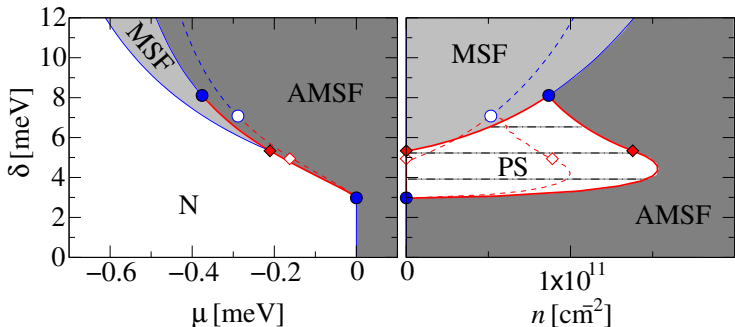
- Parameters for GaAs $\Omega_R = 4.4\text{meV}$, $E_b = 2\text{meV}$.
- Dashed: Mean-field. Solid: Next order fluctuations

• Two condensed phases:

MSF, $\langle \psi_f \rangle = \langle \psi_b \rangle = 0$, $\langle \psi_m \rangle \neq 0$. AMSF $\langle \psi_f \rangle \neq 0$.

• Small (or -ve) δ : $v \gg g$, no biexciton physics — “standard” WIDBG.

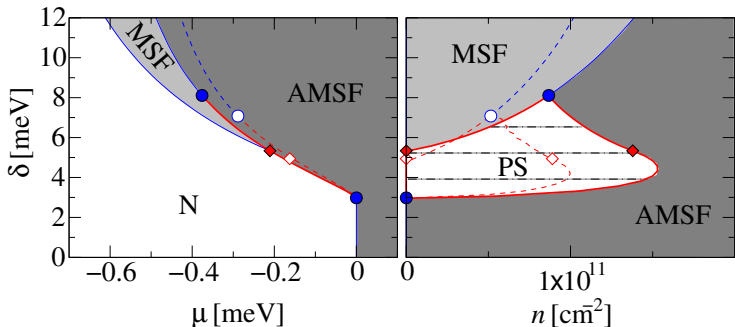
Phase diagram (ground state, $T = 0$)



- Parameters for GaAs $\Omega_R = 4.4\text{meV}$, $E_b = 2\text{meV}$.
- Dashed: Mean-field. Solid: Next order fluctuations
- Two condensed phases:
 $\text{MSF}, \langle \hat{\psi}_\uparrow \rangle = \langle \hat{\psi}_\downarrow \rangle = 0, \langle \hat{\psi}_m \rangle \neq 0.$ $\text{AMSF} \forall_\sigma : \langle \hat{\psi}_\sigma \rangle \neq 0.$

• Small (or -ve) δ & $\nu \gg g$, no bion physics — “standard” WIDBG.

Phase diagram (ground state, $T = 0$)



- Parameters for GaAs $\Omega_R = 4.4\text{meV}$, $E_b = 2\text{meV}$.
- Dashed: Mean-field. Solid: Next order fluctuations
- Two condensed phases:
MSF, $\langle \hat{\psi}_\uparrow \rangle = \langle \hat{\psi}_\downarrow \rangle = 0$, $\langle \hat{\psi}_m \rangle \neq 0$. AMSF $\forall_\sigma : \langle \hat{\psi}_\sigma \rangle \neq 0$.
- Small (or -ve) δ : $\nu \gg g$, no biexciton physics — “standard” WIDBG.

Mean field ansatz

- Hamiltonian: Weakly Interacting Bose Gas
- First ansatz: mean-field theory [Zhou *et al.* PRA '08]

$$|\Psi\rangle \propto \exp\left(-\sum_{\sigma=\uparrow,\downarrow,m} \psi_{\sigma} \hat{\psi}_{k=0,\sigma}\right).$$

- No magnetic field, take $\psi_{\uparrow} = \psi_{\downarrow} = \psi_0$ (real):

$$E = \langle \Psi | H | \Psi \rangle = -2\mu\psi_0^2 + \left(\frac{U_{\uparrow\uparrow} + U_{\downarrow\downarrow} + U_{\uparrow\downarrow}}{2}\right)\psi_0^4 + (\nu - 2\mu)\psi_m^2 + \frac{U_{mm}}{2}\psi_m^4 + g\nu_0^2\psi_m$$

- Minimise $E(\psi_0, \psi_m)$: first order transitions exist.

Mean field ansatz

- Hamiltonian: Weakly Interacting Bose Gas
- First ansatz: mean-field theory [Zhou *et al.* PRA '08]

$$|\Psi\rangle \propto \exp\left(-\sum_{\sigma=\uparrow,\downarrow,m} \psi_{\sigma} \hat{\psi}_{k=0,\sigma}\right).$$

- No magnetic field, take $\psi_{\uparrow} = \psi_{\downarrow} = \psi_0$ (real):

$$E = \langle \Psi | H | \Psi \rangle = -2\mu\psi_0^2 + \left(\frac{U_{\uparrow\uparrow} + U_{\downarrow\downarrow}}{2} + U_{\uparrow\downarrow}\right)\psi_0^4 \\ + (\nu - 2\mu)\psi_m^2 + \frac{U_{mm}}{2}\psi_m^4 + g\psi_0^2\psi_m$$

• Minimise $E(\psi_0, \psi_m)$: first order transitions exist.

Mean field ansatz

- Hamiltonian: Weakly Interacting Bose Gas
- First ansatz: mean-field theory [Zhou *et al.* PRA '08]

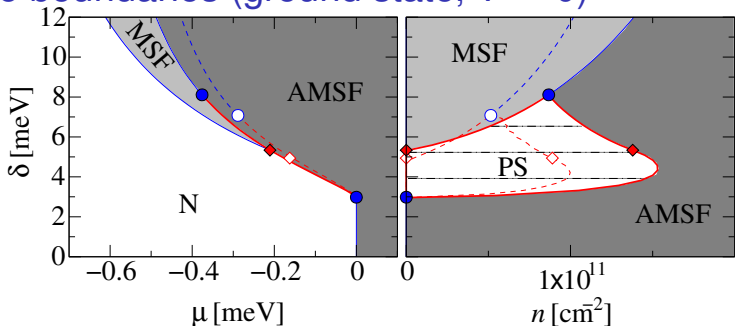
$$|\Psi\rangle \propto \exp\left(-\sum_{\sigma=\uparrow,\downarrow,m} \psi_{\sigma} \hat{\psi}_{k=0,\sigma}\right).$$

- No magnetic field, take $\psi_{\uparrow} = \psi_{\downarrow} = \psi_0$ (real):

$$E = \langle \Psi | H | \Psi \rangle = -2\mu\psi_0^2 + \left(\frac{U_{\uparrow\uparrow} + U_{\downarrow\downarrow}}{2} + U_{\uparrow\downarrow}\right)\psi_0^4 \\ + (\nu - 2\mu)\psi_m^2 + \frac{U_{mm}}{2}\psi_m^4 + g\psi_0^2\psi_m$$

- Minimise $E(\psi_0, \psi_m)$: first order transitions exist.

Phase boundaries (ground state, $T = 0$)



- Small δ : Fluctuations drive N–AMSF 1st order
 - ▶ (Larkin-Pikin mechanism, coupling $g\psi_0^2\psi_m$)

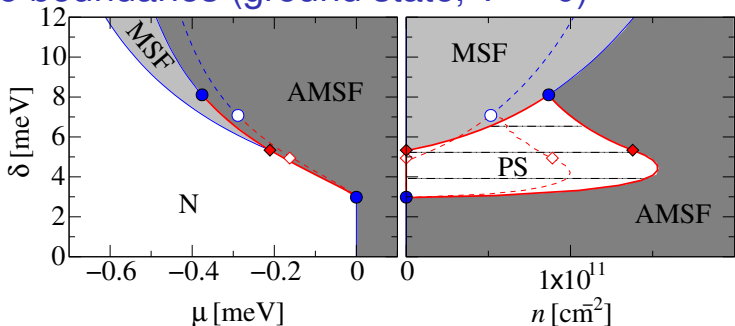
● Large phase-separation region — high densities.

● First-order BEC transition!

- Naive resonance $\delta_r = 3.84$, tricritical point at $(\delta, \mu) = (\delta_r, 0)$ replaced by critical end-point at $\delta > \delta_r$.

- Large δ : MSF-AMSF transition always occurs as U_{mm} renormalises energy.

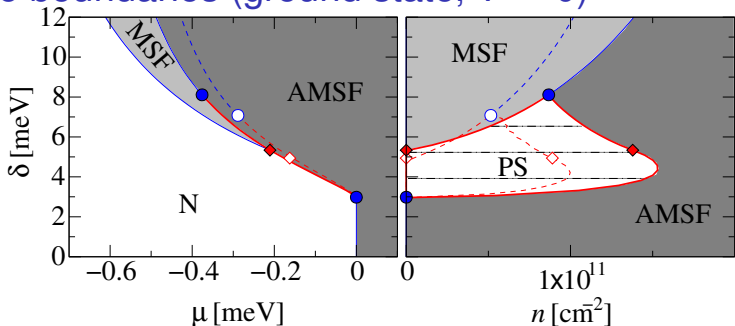
Phase boundaries (ground state, $T = 0$)



- Small δ : Fluctuations drive N–AMSF 1st order
 - ▶ (Larkin-Pikin mechanism, coupling $g\psi_0^2\psi_m$)
 - ▶ Large phase-separation region — high densities. First-order BEC transition!

- Naive resonance $\delta_r = 3.84$, tricritical point at $(\delta, \mu) = (\delta_r, 0)$ replaced by critical end-point at $\delta > \delta_r$
- Large δ : MSF-AMSF transition always occurs as U_{mm} renormalises energy

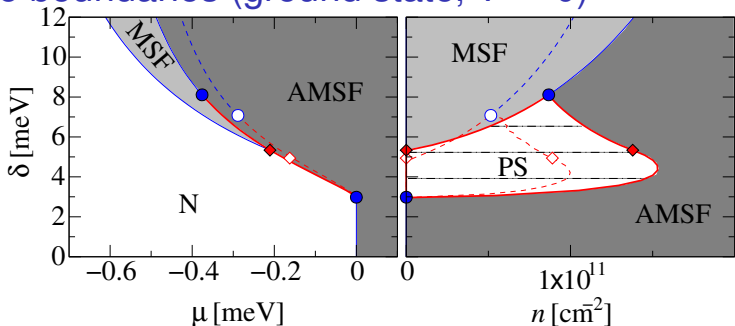
Phase boundaries (ground state, $T = 0$)



- Small δ : Fluctuations drive N–AMSF 1st order
 - ▶ (Larkin-Pikin mechanism, coupling $g\psi_0^2\psi_m$)
 - ▶ Large phase-separation region — high densities.
First-order BEC transition!
- Naive resonance $\delta_r = 3.84$, tricritical point at $(\delta, \mu) = (\delta_r, 0)$
replaced by critical end-point at $\delta > \delta_r$

● Large δ : MSF-AMSF transition always occurs as U_{res}
renormalises energy

Phase boundaries (ground state, $T = 0$)



- Small δ : Fluctuations drive N–AMSF 1st order
 - ▶ (Larkin-Pikin mechanism, coupling $g\psi_0^2\psi_m$)
 - ▶ Large phase-separation region — high densities.
First-order BEC transition!
- Naive resonance $\delta_r = 3.84$, tricritical point at $(\delta, \mu) = (\delta_r, 0)$ replaced by critical end-point at $\delta > \delta_r$
- Large δ : MSF-AMSF transition always occurs as U_{mm} renormalises energy.

Beyond mean-field

- Fluctuation effects?
 - ▶ Polariton fluctuations irrelevant: $mU \sim 10^{-4}$.
 - ▶ Exciton fluctuations important: $m_m U \sim 1$.
- Next order theory: [Nozières & St James, J. Phys '82]

$$|\Psi\rangle \propto \exp \left(- \sum_{\sigma=\uparrow,\downarrow,m} \psi_\sigma \hat{\psi}_{k=0,\sigma} + \sum_{k,\gamma=a,b,m} \tanh(\theta_{k\gamma}) \hat{b}_{k\gamma}^\dagger \hat{b}_{-k\gamma}^\dagger \right).$$

where $\hat{b}_{km}^\dagger = \hat{\psi}_{km}^\dagger$ and $\begin{pmatrix} \hat{\psi}_{\mathbf{k}\uparrow}^\dagger \\ \hat{\psi}_{-\mathbf{k}\downarrow}^\dagger \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}a}^\dagger \\ \hat{b}_{-\mathbf{k}b}^\dagger \end{pmatrix}$,

• Variational functional $E[\psi_0, \psi_m, \theta_{k\gamma}]$

• Can show minimum $\theta_{k\gamma}$ has form $\tanh(2\theta_{k\gamma}) = \frac{\alpha_\gamma}{\beta_\gamma + k^2/2m_\gamma}$

• Finite only if $|\alpha_\gamma| < \beta_\gamma$

Beyond mean-field

- Fluctuation effects?
 - ▶ Polariton fluctuations irrelevant: $mU \sim 10^{-4}$.
 - ▶ Exciton fluctuations important: $m_m U \sim 1$.
- Next order theory: [Nozières & St James, J. Phys '82]

$$|\Psi\rangle \propto \exp \left(- \sum_{\sigma=\uparrow,\downarrow,m} \psi_\sigma \hat{\psi}_{k=0,\sigma} + \sum_{k,\gamma=a,b,m} \tanh(\theta_{k\gamma}) \hat{b}_{k\gamma}^\dagger \hat{b}_{-k\gamma}^\dagger \right).$$

where $\hat{b}_{km}^\dagger = \hat{\psi}_{km}^\dagger$ and $\begin{pmatrix} \hat{\psi}_{\mathbf{k}\uparrow}^\dagger \\ \hat{\psi}_{-\mathbf{k}\downarrow}^\dagger \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}a}^\dagger \\ \hat{b}_{-\mathbf{k}b}^\dagger \end{pmatrix}$,

- Variational functional $E[\psi_0, \psi_m, \theta_{k\gamma}]$

- Can show minimum $\theta_{k\gamma}$ has form $\tanh(2\theta_{k\gamma}) = \frac{\alpha_\gamma}{\beta_\gamma + k^2/2m_\gamma}$
 - ★ Finite only if $|\alpha_\gamma| < \beta_\gamma$

Beyond mean-field

- Fluctuation effects?
 - ▶ Polariton fluctuations irrelevant: $mU \sim 10^{-4}$.
 - ▶ Exciton fluctuations important: $m_m U \sim 1$.
- Next order theory: [Nozières & St James, J. Phys '82]

$$|\Psi\rangle \propto \exp \left(- \sum_{\sigma=\uparrow,\downarrow,m} \psi_\sigma \hat{\psi}_{\mathbf{k}=0,\sigma} + \sum_{k,\gamma=a,b,m} \tanh(\theta_{k\gamma}) \hat{b}_{k\gamma}^\dagger \hat{b}_{-k\gamma}^\dagger \right).$$

where $\hat{b}_{km}^\dagger = \hat{\psi}_{km}^\dagger$ and $\begin{pmatrix} \hat{\psi}_{\mathbf{k}\uparrow}^\dagger \\ \hat{\psi}_{-\mathbf{k}\downarrow}^\dagger \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}a}^\dagger \\ \hat{b}_{-\mathbf{k}b}^\dagger \end{pmatrix}$,

- Variational functional $E[\psi_0, \psi_m, \theta_{k\gamma}]$
- Can show minimum $\theta_{k\gamma}$ has form $\tanh(2\theta_{k\gamma}) = \frac{\alpha_\gamma}{\beta_\gamma + k^2/2m_\gamma}$
 - ▶ Finite only if $|\alpha_\gamma| < \beta_\gamma$

Variational function $E(\psi_0, \psi_m, \theta_{k\gamma}, \beta_\gamma)$

Beyond mean-field

- Fluctuation effects?
 - ▶ Polariton fluctuations irrelevant: $mU \sim 10^{-4}$.
 - ▶ Exciton fluctuations important: $m_m U \sim 1$.
- Next order theory: [Nozières & St James, J. Phys '82]

$$|\Psi\rangle \propto \exp \left(- \sum_{\sigma=\uparrow,\downarrow,m} \psi_\sigma \hat{\psi}_{\mathbf{k}=0,\sigma} + \sum_{k,\gamma=a,b,m} \tanh(\theta_{k\gamma}) \hat{b}_{k\gamma}^\dagger \hat{b}_{-k\gamma}^\dagger \right).$$

where $\hat{b}_{km}^\dagger = \hat{\psi}_{km}^\dagger$ and $\begin{pmatrix} \hat{\psi}_{\mathbf{k}\uparrow}^\dagger \\ \hat{\psi}_{-\mathbf{k}\downarrow}^\dagger \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}a}^\dagger \\ \hat{b}_{-\mathbf{k}b}^\dagger \end{pmatrix}$,

- Variational functional $E[\psi_0, \psi_m, \theta_{k\gamma}]$
- Can show minimum $\theta_{k\gamma}$ has form $\tanh(2\theta_{k\gamma}) = \frac{\alpha_\gamma}{\beta_\gamma + k^2/2m_\gamma}$
 - ▶ Finite only if $|\alpha_\gamma| < \beta_\gamma$
 - ▶ Variational function $E(\psi_0, \psi_m, \alpha_\gamma, \beta_\gamma)$

Outline

- 1 Introduction
- 2 Pairing phases of atoms: review
- 3 Modelling polariton pairing
 - Exciton spin structure
 - Microscopic and effective Hamiltonian
- 4 Ground state phase diagram**
 - Origin of multicritical behaviour
 - Signatures of phases**
- 5 Finite temperature phase diagram
 - Variational MFT
 - Required temperatures, detuning
 - Candidate material systems

Detecting MSF vs AMSF

- Phase separation

- Direct access to polariton phase coherence

- Vortex structure — half vortices, $\psi_{\uparrow} = e^{im_{\uparrow}\theta}$, $\psi_{\downarrow} = e^{im_{\downarrow}\theta}$,
MSF has $(m_{\uparrow}, m_{\downarrow}) = (1/2, 1/2)$

Polariton "half-vortex" $(m_{\uparrow}, m_{\downarrow}) = (1, 0)$ [Lagoudakis *et al.* Science '09]

Detecting MSF vs AMSF

- Phase separation
- Direct access to polariton phase coherence

→ AMSF: standard $g_o^{(1)} = \langle \psi_o^\dagger(r, t) \psi_o(0, 0) \rangle$. (2D, QLRO)

- Vortex structure → half vortices, $\psi_\uparrow = e^{im_\uparrow\theta}$, $\psi_\downarrow = e^{im_\downarrow\theta}$,
MSF has $(m_\uparrow, m_\downarrow) = (1/2, 1/2)$

Polariton "half-vortex" $(m_\uparrow, m_\downarrow) = (1, 0)$ [Lagoudakis *et al.* Science '09]

Detecting MSF vs AMSF

- Phase separation
- Direct access to polariton phase coherence
 - ▶ AMSF: standard $g_{\sigma}^{(1)} = \langle \psi_{\sigma}^{\dagger}(r, t) \psi_{\sigma}(0, 0) \rangle$. (2D, QLRO)

● Vortex structure \rightarrow half vortices, $\psi_{\uparrow} = e^{im_{\uparrow}\theta}$, $\psi_{\downarrow} = e^{im_{\downarrow}\theta}$,
MSF has $(m_{\uparrow}, m_{\downarrow}) = (1/2, 1/2)$

Polariton "half-vortex" $(m_{\uparrow}, m_{\downarrow}) = (1, 0)$ [Lagoudakis *et al.* Science '09]

Detecting MSF vs AMSF

- Phase separation
- Direct access to polariton phase coherence
 - ▶ AMSF: standard $g_{\sigma}^{(1)} = \langle \psi_{\sigma}^{\dagger}(r, t) \psi_{\sigma}(0, 0) \rangle$. (2D, QLRO)

● Vortex structure \rightarrow half vortices, $\psi_{\uparrow} = e^{im_{\uparrow}\theta}$, $\psi_{\downarrow} = e^{im_{\downarrow}\theta}$,
MSF has $(m_{\uparrow}, m_{\downarrow}) = (1/2, 1/2)$

Polariton "half-vortex" $(m_{\uparrow}, m_{\downarrow}) = (1, 0)$ [Lagoudakis *et al.* Science '09]

Detecting MSF vs AMSF

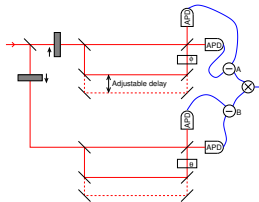
- Phase separation
- Direct access to polariton phase coherence

▶ AMSF: standard $g_{\sigma}^{(1)} = \langle \psi_{\sigma}^{\dagger}(\mathbf{r}, t) \psi_{\sigma}(\mathbf{0}, 0) \rangle$. (2D, QLRO)

▶ MSF: no “atomic” coherence, but:

$$g_m^{(1)} = \langle \psi_{\uparrow}^{\dagger}(\mathbf{r}, t) \psi_{\downarrow}^{\dagger}(\mathbf{r}, t) \psi_{\downarrow}(\mathbf{0}, 0) \psi_{\uparrow}(\mathbf{0}, 0) \rangle$$

not g_2 – see time/space labels



• Vortex structure \rightarrow half vortices, $\psi_{\uparrow} = e^{im_{\uparrow}\theta}$, $\psi_{\downarrow} = e^{im_{\downarrow}\theta}$

MSF has $(m_{\uparrow}, m_{\downarrow}) = (1/2, 1/2)$

Polariton “half-vortex” $(m_{\uparrow}, m_{\downarrow}) = (1, 0)$ [Lagoudakis *et al.* Science '09]

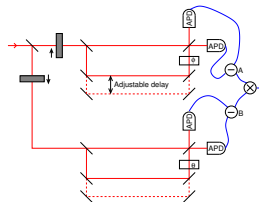
Detecting MSF vs AMSF

- Phase separation
- Direct access to polariton phase coherence
 - ▶ AMSF: standard $g_{\sigma}^{(1)} = \langle \psi_{\sigma}^{\dagger}(\mathbf{r}, t) \psi_{\sigma}(0, 0) \rangle$. (2D, QLRO)

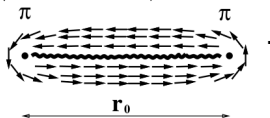
- ▶ MSF: no “atomic” coherence, but:

$$g_m^{(1)} = \langle \psi_{\uparrow}^{\dagger}(\mathbf{r}, t) \psi_{\downarrow}^{\dagger}(\mathbf{r}, t) \psi_{\downarrow}(0, 0) \psi_{\uparrow}(0, 0) \rangle$$

not g_2 – see time/space labels



- Vortex structure — half vortices, $\psi_{\uparrow} = e^{im_{\uparrow}\theta}$, $\psi_{\downarrow} = e^{im_{\downarrow}\theta}$,
MSF has $(m_{\uparrow}, m_{\downarrow}) = (1/2, 1/2)$



Polariton “half-vortex” $(m_{\uparrow}, m_{\downarrow}) = (1, 0)$ [Lagoudakis *et al.* Science '09]

Outline

- 1 Introduction
- 2 Pairing phases of atoms: review
- 3 Modelling polariton pairing
 - Exciton spin structure
 - Microscopic and effective Hamiltonian
- 4 Ground state phase diagram
 - Origin of multicritical behaviour
 - Signatures of phases
- 5 **Finite temperature phase diagram**
 - **Variational MFT**
 - **Required temperatures, detuning**
 - **Candidate material systems**

Finite T calculation

- Finite T — minimize free energy

- Use Feynman-Jensen inequality:

$$F = -k_B T \ln \left[\text{Tr} e^{-H/k_B T} \right] \leq F_{\text{MF}} + \langle H - \hat{H}_{\text{MF}} \rangle_{\text{MF}}$$

Where $\langle \dots \rangle_{\text{MF}}$ calculated using $\rho = e^{(F_{\text{MF}} - \hat{H}_{\text{MF}})/k_B T}$

- Ansatz $\hat{H}_{\text{MF}} \rightarrow$ Variational $F(\psi_0, \psi_m, \alpha_\gamma, \beta_\gamma)$.

$$H_{\text{MF}} = \sum_{\gamma} \left\{ -\sqrt{\lambda} \psi_{\gamma} (\alpha_{\gamma} + \beta_{\gamma}) (\hat{b}_{0\gamma}^{\dagger} + \hat{b}_{0\gamma}) + \frac{1}{2} \sum_{\mathbf{k}} (\hat{b}_{\mathbf{k}\gamma}^{\dagger} - \hat{b}_{-\mathbf{k}\gamma}) \begin{pmatrix} \alpha_{\mathbf{k}\gamma} + \beta_{\mathbf{k}\gamma} & \alpha_{\mathbf{k}\gamma} \\ \alpha_{\mathbf{k}\gamma} & \alpha_{\mathbf{k}\gamma} + \beta_{\mathbf{k}\gamma} \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}\gamma} \\ \hat{b}_{-\mathbf{k}\gamma}^{\dagger} \end{pmatrix} \right\}.$$

Finite T calculation

- Finite T — minimize free energy
- Use Feynman-Jensen inequality:

$$F = -k_B T \ln \left[\text{Tr} e^{-\hat{H}/k_B T} \right] \leq F_{\text{MF}} + \langle \hat{H} - \hat{H}_{\text{MF}} \rangle_{\text{MF}}$$

Where $\langle \dots \rangle_{\text{MF}}$ calculated using $\rho = e^{(F_{\text{MF}} - \hat{H}_{\text{MF}})/k_B T}$

• Ansatz $\hat{H}_{\text{MF}} \rightarrow$ Variational $F(\psi_0, \psi_m, \alpha_\gamma, \beta_\gamma)$.

$$F_{\text{MF}} = \sum_\gamma \left\{ -\sqrt{\lambda} \psi_\gamma (\alpha_\gamma + \beta_\gamma) (b_{0\gamma}^\dagger + b_{0\gamma}) + \frac{1}{2} \sum_{\mathbf{k}} (b_{\mathbf{k}\gamma}^\dagger \ b_{-\mathbf{k}\gamma}) \begin{pmatrix} \alpha_\gamma + \beta_\gamma & \alpha_\gamma \\ \alpha_\gamma & \alpha_\gamma + \beta_\gamma \end{pmatrix} \begin{pmatrix} b_{\mathbf{k}\gamma} \\ b_{-\mathbf{k}\gamma}^\dagger \end{pmatrix} \right\}.$$

Finite T calculation

- Finite T — minimize free energy
- Use Feynman-Jensen inequality:

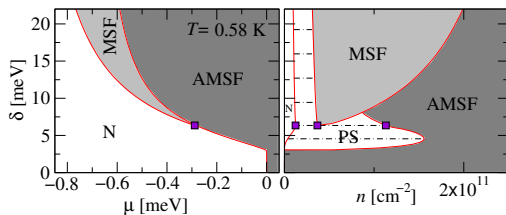
$$F = -k_B T \ln \left[\text{Tr} e^{-\hat{H}/k_B T} \right] \leq F_{\text{MF}} + \langle \hat{H} - \hat{H}_{\text{MF}} \rangle_{\text{MF}}$$

Where $\langle \dots \rangle_{\text{MF}}$ calculated using $\rho = e^{(F_{\text{MF}} - \hat{H}_{\text{MF}})/k_B T}$

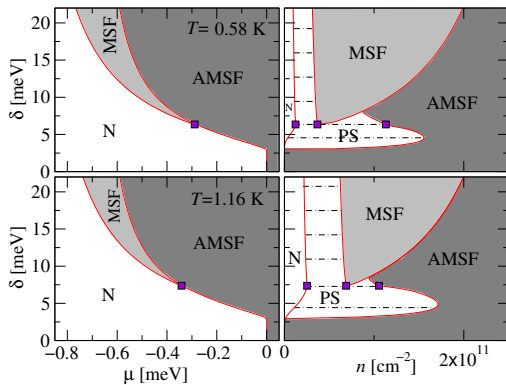
- Ansatz $\hat{H}_{\text{MF}} \rightarrow$ Variational $F(\psi_0, \psi_m, \alpha_\gamma, \beta_\gamma)$.

$$\hat{H}_{\text{MF}} = \sum_{\gamma} \left\{ -\sqrt{\mathcal{A}} \psi_{\gamma} (\alpha_{\gamma} + \beta_{\gamma}) (\hat{b}_{0\gamma}^{\dagger} + \hat{b}_{0\gamma}) \right. \\ \left. + \frac{1}{2} \sum_{\mathbf{k}} \begin{pmatrix} \hat{b}_{\mathbf{k}\gamma}^{\dagger} & \hat{b}_{-\mathbf{k}\gamma} \end{pmatrix} \begin{pmatrix} \epsilon_{\mathbf{k}\gamma} + \beta_{\gamma} & \alpha_{\gamma} \\ \alpha_{\gamma} & \epsilon_{\mathbf{k}\gamma} + \beta_{\gamma} \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}\gamma} \\ \hat{b}_{-\mathbf{k}\gamma}^{\dagger} \end{pmatrix} \right\}.$$

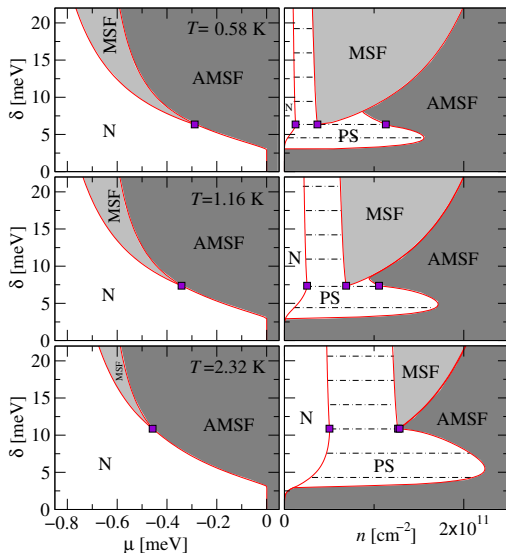
Phase diagram, finite temperature



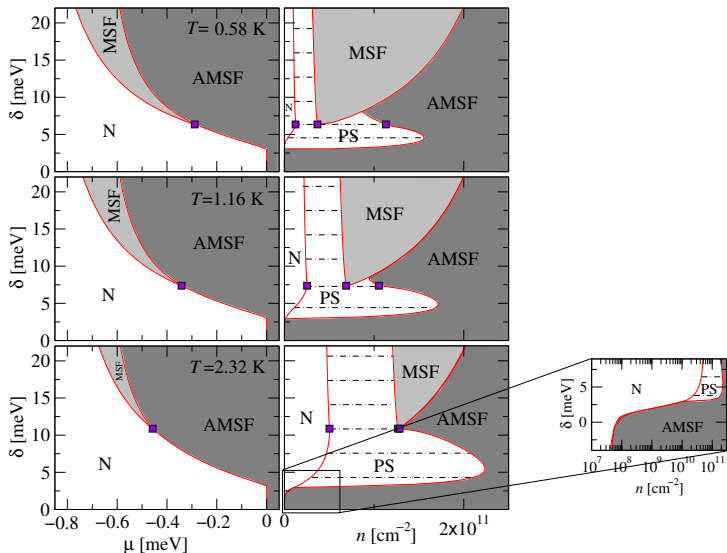
Phase diagram, finite temperature



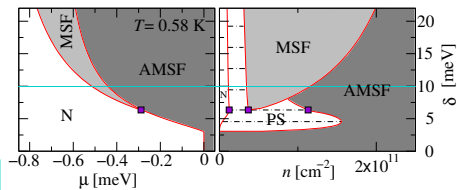
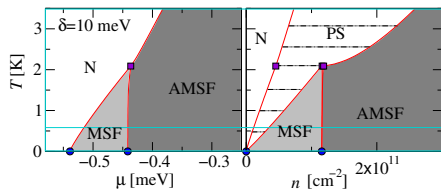
Phase diagram, finite temperature



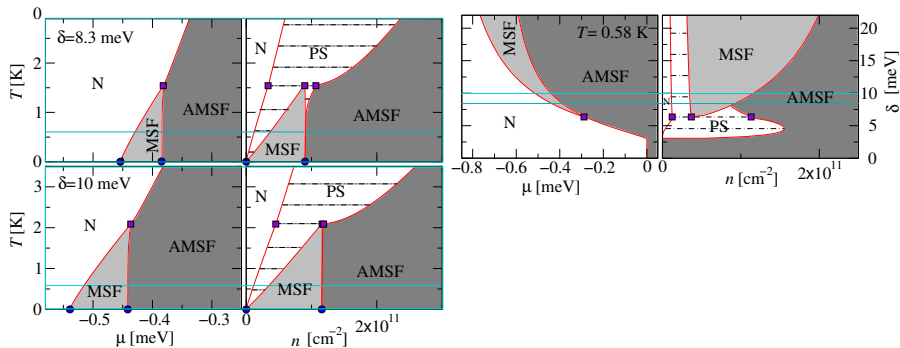
Phase diagram, finite temperature



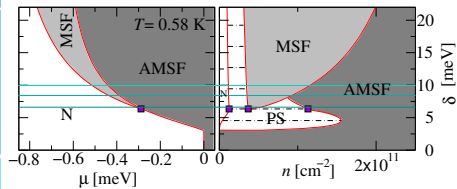
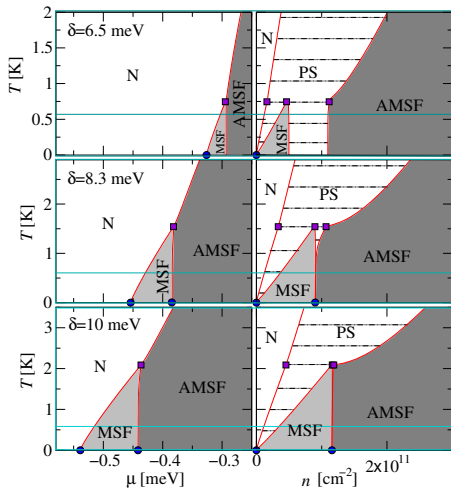
Phase diagram, vs temperature



Phase diagram, vs temperature



Phase diagram, vs temperature



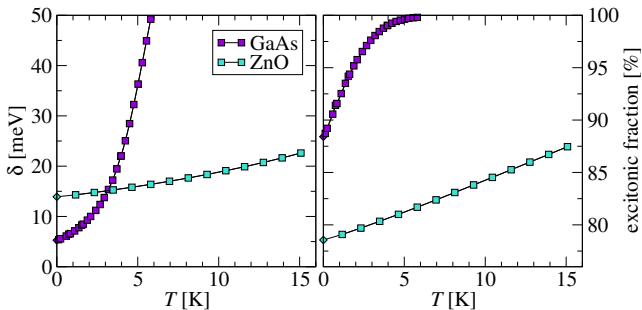
Evolution of triple point

- Required T, δ for MSF: Triple point
- Excitonic fraction $c_0^2 = \frac{1}{2} \left[1 + \frac{\delta}{\sqrt{\delta^2 + \Omega^2}} \right]$. $\delta \gg \Omega$: Pure exciton

- GaAs, need low T /high exciton fraction
- ZnO, easy to attain MSF

Evolution of triple point

- Required T, δ for MSF: Triple point
- Excitonic fraction $c_0^2 = \frac{1}{2} \left[1 + \frac{\delta}{\sqrt{\delta^2 + \Omega^2}} \right]$. $\delta \gg \Omega$: Pure exciton

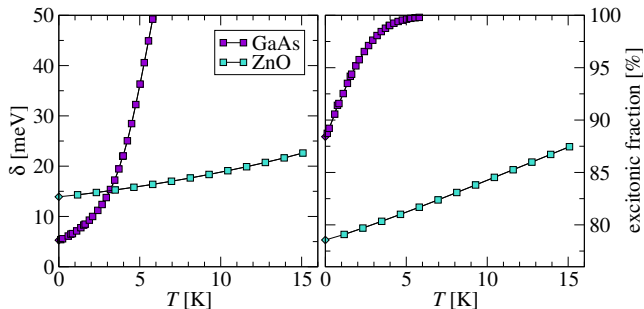


• GaAs, need low T /high exciton fraction

• ZnO, easy to attain MSF

Evolution of triple point

- Required T, δ for MSF: Triple point
- Excitonic fraction $c_0^2 = \frac{1}{2} \left[1 + \frac{\delta}{\sqrt{\delta^2 + \Omega^2}} \right]$. $\delta \gg \Omega$: Pure exciton

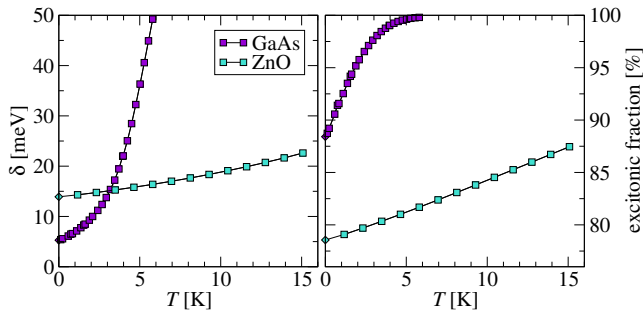


- GaAs, need low T /high exciton fraction

• ZnO, easy to attain MSF

Evolution of triple point

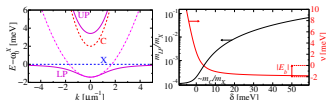
- Required T, δ for MSF: Triple point
- Excitonic fraction $c_0^2 = \frac{1}{2} \left[1 + \frac{\delta}{\sqrt{\delta^2 + \Omega^2}} \right]$. $\delta \gg \Omega$: Pure exciton



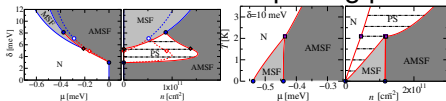
- GaAs, need low T /high exciton fraction
- ZnO, easy to attain MSF

Summary

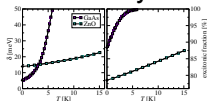
- Polariton — biexciton detuning \rightarrow Feshbach resonance



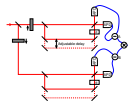
- Polaritons can show pairing phase



- MSF easily attainable for ZnO; "possible" for GaA

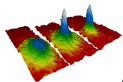


- Possible signatures in coherence and vortices



[Marchetti and Keeling, arXiv:1308.1032]

Differences & new opportunities



Atoms

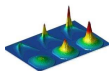
Metastable, isolated

Fixed mass $\sim 10^3 m_0$

Couples to light

Can be 3D

Spin: $S_z = -S, -S + 1, \dots, S$



Polaritons

Finite lifetime, in semiconductor

Tunable mass $\sim 10^{-4} m_0$

Is part light

Always 2D or less.

Light polarisations $S_z = \pm 1$

Variational MFT for WIDBG

- Test validity. WIDBG $\hat{H} = \sum_{\mathbf{k}} \frac{k^2}{2m} \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \frac{U}{2} \int d^2r \psi^\dagger \psi^\dagger \psi \psi$
- VMFT for WIDBG:

$$\hat{H}_{\text{MF}} = -\sqrt{\mathcal{A}}\psi(\alpha + \beta)(\hat{b}_0^\dagger + \hat{b}_0) + \frac{1}{2} \sum_{\mathbf{k}} \begin{pmatrix} \hat{b}_{\mathbf{k}}^\dagger & \hat{b}_{-\mathbf{k}} \end{pmatrix} \begin{pmatrix} \epsilon_{\mathbf{k}} + \beta & \alpha \\ \alpha & \epsilon_{\mathbf{k}} + \beta \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}} \\ \hat{b}_{-\mathbf{k}}^\dagger \end{pmatrix}.$$

- Compare to 2D EOS, $\rho(\mu) = T(\mu/T)$
- CF Hartree-Fock-Popov-Bogoliubov method, include $U\rho$ in Σ

Variational MFT for WIDBG

- Test validity. WIDBG $\hat{H} = \sum_{\mathbf{k}} \frac{k^2}{2m} \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \frac{U}{2} \int d^2r \psi^\dagger \psi^\dagger \psi \psi$
- VMFT for WIDBG:

$$\hat{H}_{\text{MF}} = -\sqrt{\mathcal{A}}\psi(\alpha + \beta)(\hat{b}_0^\dagger + \hat{b}_0) + \frac{1}{2} \sum_{\mathbf{k}} \begin{pmatrix} \hat{b}_{\mathbf{k}}^\dagger & \hat{b}_{-\mathbf{k}} \end{pmatrix} \begin{pmatrix} \epsilon_{\mathbf{k}} + \beta & \alpha \\ \alpha & \epsilon_{\mathbf{k}} + \beta \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}} \\ \hat{b}_{-\mathbf{k}}^\dagger \end{pmatrix}.$$

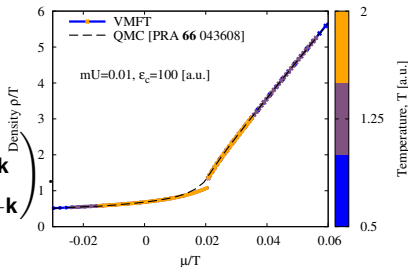
- Compare to 2D EOS, $\rho(\mu) = T(\mu/T)$
- CF Hartree-Fock-Popov-Bogoliubov method, include $U\rho$ in Σ

Variational MFT for WIDBG

- Test validity. WIDBG $\hat{H} = \sum_{\mathbf{k}} \frac{k^2}{2m} \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \frac{U}{2} \int d^2r \psi^\dagger \psi^\dagger \psi \psi$

- VMFT for WIDBG:

$$\hat{H}_{\text{MF}} = -\sqrt{\mathcal{A}} \psi(\alpha + \beta)(\hat{b}_0^\dagger + \hat{b}_0) + \frac{1}{2} \sum_{\mathbf{k}} \begin{pmatrix} \hat{b}_{\mathbf{k}}^\dagger & \hat{b}_{-\mathbf{k}} \end{pmatrix} \begin{pmatrix} \epsilon_{\mathbf{k}} + \beta & \alpha \\ \alpha & \epsilon_{\mathbf{k}} + \beta \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}} \\ \hat{b}_{-\mathbf{k}}^\dagger \end{pmatrix}$$



- Compare to 2D EOS, $\rho(\mu) = Tf(\mu/T)$

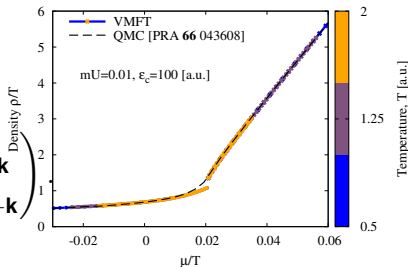
• CF Hartree-Fock-Popov-Bogoluibov method, include U_p in Σ

Variational MFT for WIDBG

- Test validity. WIDBG $\hat{H} = \sum_k \frac{k^2}{2m} \psi_k^\dagger \psi_k + \frac{U}{2} \int d^2r \psi^\dagger \psi^\dagger \psi \psi$

- VMFT for WIDBG:

$$\hat{H}_{MF} = -\sqrt{\mathcal{A}} \psi(\alpha + \beta)(\hat{b}_0^\dagger + \hat{b}_0) + \frac{1}{2} \sum_{\mathbf{k}} \begin{pmatrix} \hat{b}_{\mathbf{k}}^\dagger & \hat{b}_{-\mathbf{k}} \end{pmatrix} \begin{pmatrix} \epsilon_{\mathbf{k}} + \beta & \alpha \\ \alpha & \epsilon_{\mathbf{k}} + \beta \end{pmatrix} \begin{pmatrix} \hat{b}_{\mathbf{k}} \\ \hat{b}_{-\mathbf{k}}^\dagger \end{pmatrix}$$



- Compare to 2D EOS, $\rho(\mu) = Tf(\mu/T)$
- CF Hartree-Fock-Popov-Bogoluibov method, include $U\rho$ in Σ