

Superradiance and self organisation with fermionic atoms

Jonathan Keeling



University of
St Andrews

600
YEARS



QSOE, August 2013

Dicke effect: Superradiance

PHYSICAL REVIEW

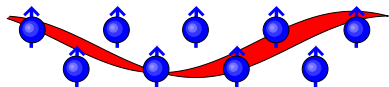
VOLUME 93, NUMBER 1

JANUARY 1, 1954

Coherence in Spontaneous Radiation Processes

R. H. DICKE

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey



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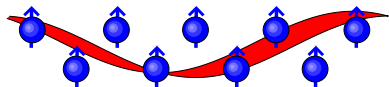
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Collective decay:

$$\frac{d\rho}{dt} = -\frac{\Gamma}{2} [S^+ S^- \rho - S^- \rho S^+ + \rho S^+ S^-]$$

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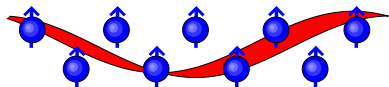
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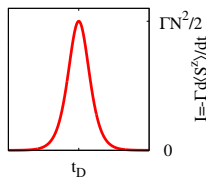
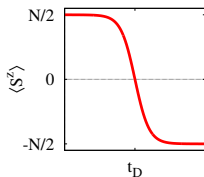
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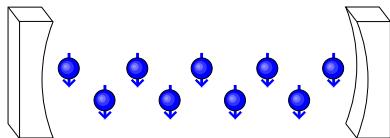
$$\frac{d\rho}{dt} = -\frac{\Gamma}{2} [S^+ S^- \rho - S^- \rho S^+ + \rho S^+ S^-]$$

If $S^z = |S| = N/2$ initially:

$$I \propto -\Gamma \frac{d\langle S^z \rangle}{dt} = \frac{\Gamma N^2}{4} \text{sech}^2 \left[\frac{\Gamma N}{2} t \right]$$



Dicke model and Dicke-Hepp-Lieb transition



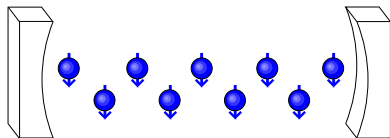
$$H = \omega \psi^\dagger \psi + \sum_i \omega_0 \mathbf{S}_i^z + g(\psi + \psi^\dagger)(\mathbf{S}_i^+ + \mathbf{S}_i^-)$$

Rotating wave approx \rightarrow Tavis-Cummings model

- Coherent state: $|\psi\rangle \rightarrow e^{\lambda \psi^\dagger + \gamma \mathbf{S}^+} |\Omega\rangle$
- Small g , min at $\lambda, \gamma = 0$

[Hepp, Lieb, Ann. Phys. '73]

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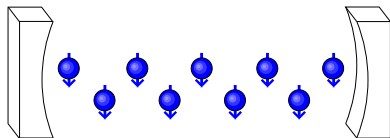
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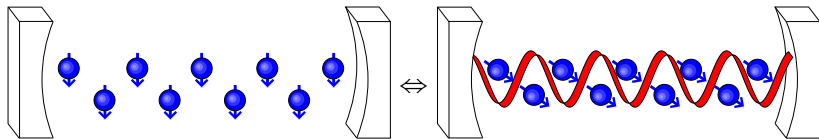
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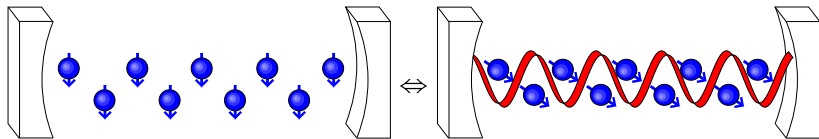
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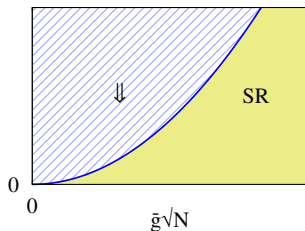


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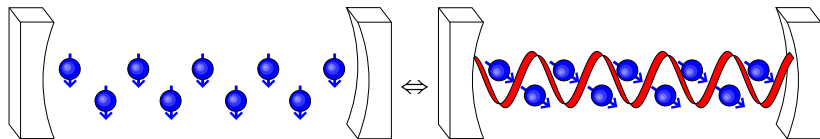
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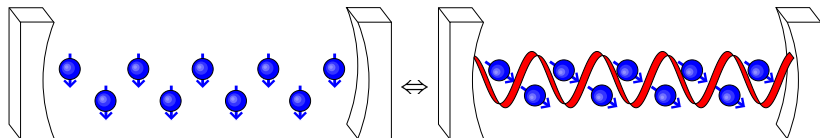
No go theorem and transition



Spontaneous polarisation if: $4Ng^2 > \omega\omega_0$

[Rzazewski *et al* PRL '75]

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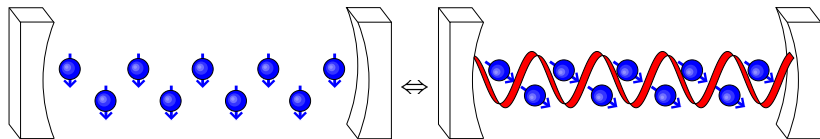
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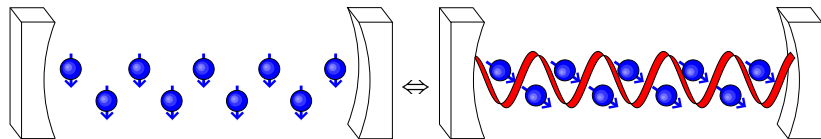
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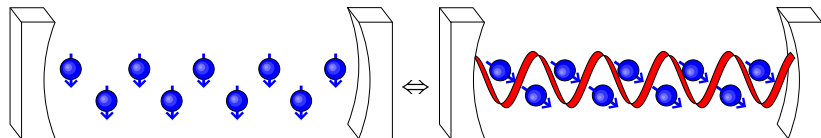
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But f -sum rule states: $g^2/\omega_0 < \zeta$. **No transition**

[Rzazewski *et al* PRL '75]

Dicke phase transition: ways out

Problem: $g^2/\omega_0 < \zeta$ for intrinsic parameters. **Solutions:**

- ④ Gauge/interpretation of "photon"
 - Ferroelectric transition in $\mathbf{D} \cdot \mathbf{r}$ gauge.
[JK JPCM '07, Vukics & Domokos PRA 2012]
 - Circuit QED [Nataf and Cluti, Nat. Comm. '10; Viehmann et al. PRL '11]
- ④ Grand canonical ensemble:
 - If $H \rightarrow H - \mu(S^z + \psi^\dagger \psi)$, need only:
 $g^2 N > (\omega - \mu)(\omega_0 - \mu)$
 - Incoherent pumping — polariton condensation.
- ④ Dissociate g, ω_0 ,
 - e.g. Raman scheme: $\omega_0 \ll \omega$.
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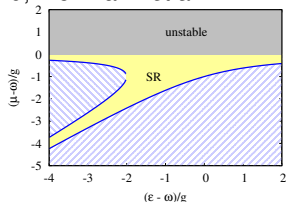
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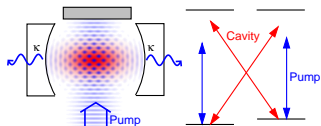
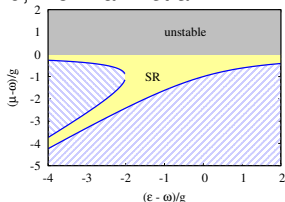
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1 Dicke model, superradiance and no-go theorem

2 Superradiance and self-organisation

- Raman scheme
- Rayleigh scheme and hierarchies of H_{eff}
- Generalized Dicke equilibrium theory

3 Fermionic self organisation

- Equilibrium phase diagrams
- Landau theory and microscopics
- Evolution with filling

4 Open system dynamics

- Linear stability with losses
- Attractors of the Dicke model phases

Acknowledgements

PEOPLE:



FUNDING:



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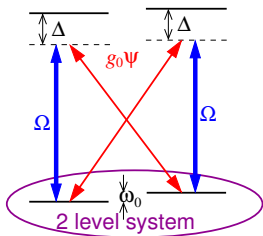
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Raman scheme, decoupling g, ω_0



$$H = \omega_0 S^Z + g(\psi + \psi^\dagger)(S^- + S^+) + \omega \psi^\dagger \psi$$

- 2 Level system, $|\downarrow\rangle, |\uparrow\rangle$

- Coupling $g = \frac{g_0 \Omega}{2\Delta}$

- Rotating frame of pump, $\omega = \omega_{\text{cavity}} - \omega_{\text{pump}}$

- Imbalanced case (internal states):

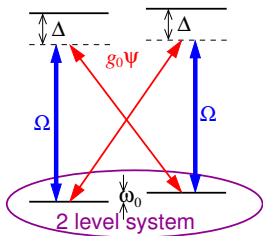
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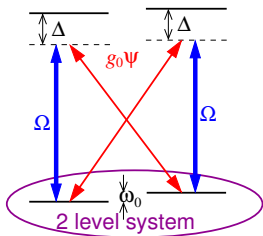
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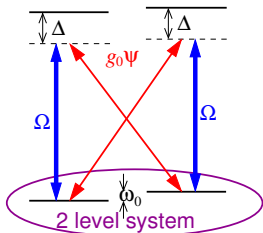
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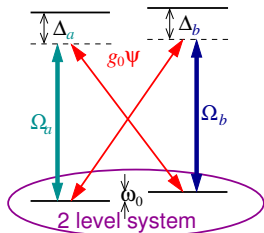
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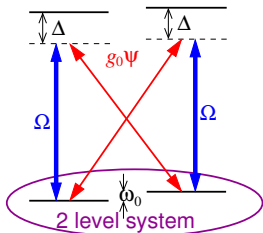
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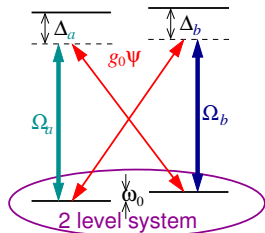
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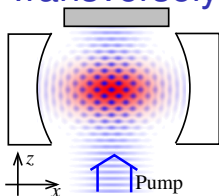
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Transversely pumped cavity

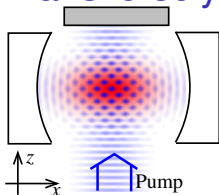


Internal state \rightarrow momentum states

- 1 Full description

$$H_0 = \omega_{\text{cavity}} \psi^\dagger \psi + \int d^2r \left[\sum_{\alpha=e,g} c_\alpha^\dagger \left(\frac{-\nabla^2}{2m} \right) c_\alpha \right. \\ \left. + \omega_{\text{atom}} c_e^\dagger c_e + E(x, z, t) (c_e^\dagger c_g + c_g^\dagger c_e) \right]$$

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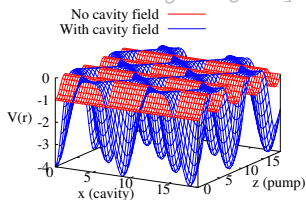
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2 Eliminate e state

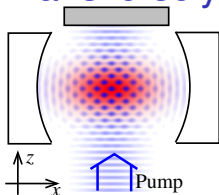
Rotating frame $\omega = \omega_{\text{cavity}} - \omega_{\text{pump}} - N\delta$

$$H = \omega \psi^\dagger \psi + \int d^2r c^\dagger(\mathbf{r}) \left(-\frac{\nabla^2}{2m} - V(\mathbf{r}) \right) c(\mathbf{r})$$

$$V(\mathbf{r}) = \frac{g^2}{2\Delta} \psi^\dagger \psi \cos(2qx) + \frac{g\Omega}{\Delta} (\psi + \psi^\dagger) \cos(qx) \cos(qz) + \frac{\Omega^2}{2\Delta} \cos(2qz)$$



Transversely pumped cavity



Internal state \rightarrow momentum states

1 Full description

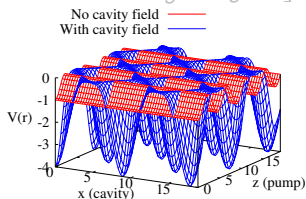
$$H_0 = \omega_{\text{cavity}} \psi^\dagger \psi + \int d^2r \left[\sum_{\alpha=e,g} c_\alpha^\dagger \left(\frac{-\nabla^2}{2m} \right) c_\alpha + \omega_{\text{atom}} c_e^\dagger c_e + E(x, z, t) (c_e^\dagger c_g + c_g^\dagger c_e) \right]$$

2 Eliminate e state

Rotating frame $\omega = \omega_{\text{cavity}} - \omega_{\text{pump}} - N\delta$

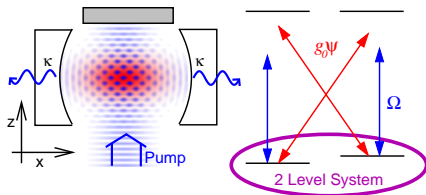
$$H = \omega \psi^\dagger \psi + \int d^2r c^\dagger(\mathbf{r}) \left(-\frac{\nabla^2}{2m} - V(\mathbf{r}) \right) c(\mathbf{r})$$

$$V(\mathbf{r}) = \frac{g^2}{2\Delta} \psi^\dagger \psi \cos(2qx) + \frac{g\Omega}{\Delta} (\psi + \psi^\dagger) \cos(qx) \cos(qz) + \frac{\Omega^2}{2\Delta} \cos(2qz)$$



3 Dicke: project to atomic states $\phi(x, z) \propto \begin{cases} 1 \\ \cos(qz) \cos(qz) \end{cases}$

Extended Dicke model



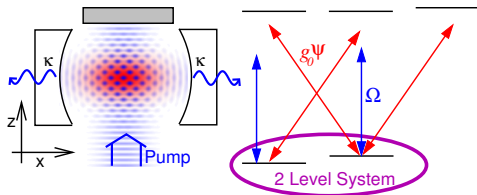
Reduced basis:

$$\phi(x, z) \propto \begin{cases} 1 & \Downarrow \\ \cos(qz) \cos(qz) & \Uparrow \end{cases}$$

$$H = \omega \psi^\dagger \psi + \omega_0 S^z + g(\psi + \psi^\dagger)(S^- + S^+)$$

[Baumann *et al* Nature '10]

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“Feedback” due to extra states $U = -\frac{g_0^2}{4\Delta}$

[Baumann *et al* Nature '10]

Extended Dicke phase diagram

Ground state energy, $\lambda = \langle \psi \rangle / \sqrt{N}$:

$$\frac{E}{N} = \omega \lambda^2 - \frac{N}{2} \sqrt{(\omega_0 + UN\lambda^2)^2 + 16g^2N\lambda^2}$$

• Superradiant transition:

$$4g^2N > \left(\omega - \frac{UN}{2}\right) \omega_0$$

• Stability, $\lambda \rightarrow \infty$

$$E \sim \left(\omega - \frac{UN}{2}\right) \lambda^2 + \dots$$

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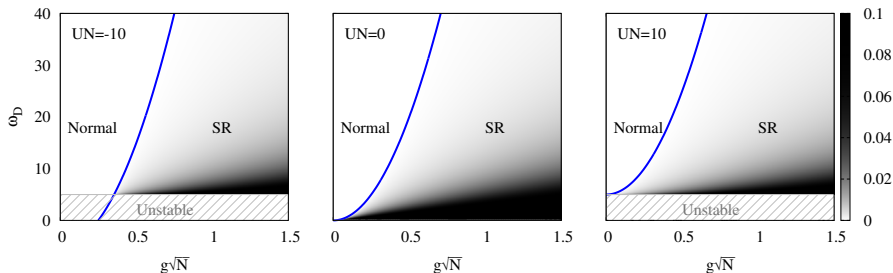
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Outline

1 Dicke model, superradiance and no-go theorem

2 Superradiance and self-organisation

- Raman scheme
- Rayleigh scheme and hierarchies of H_{eff}
- Generalized Dicke equilibrium theory

3 Fermionic self organisation

- Equilibrium phase diagrams
- Landau theory and microscopics
- Evolution with filling

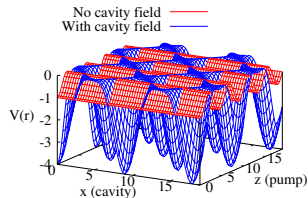
4 Open system dynamics

- Linear stability with losses
- Attractors of the Dicke model phases

Fermions in optical cavities

$$H = \omega \psi^\dagger \psi + \int d^2 \mathbf{r} c^\dagger(\mathbf{r}) \left(-\frac{\nabla^2}{2m} - V(\mathbf{r}) \right) c(\mathbf{r})$$

$$V(\mathbf{r}) = \frac{g^2}{2\Delta} \psi^\dagger \psi \cos(2qx) + \frac{g\Omega}{\Delta} (\psi + \psi^\dagger) \cos(qx) \cos(qz) + \frac{\Omega^2}{2\Delta} \cos(2qz)$$



[Domokos & Ritsch, PRL '03; Black *et al.* PRL '03; Piazza, Strack, Zweger
arXiv:1305.2928]

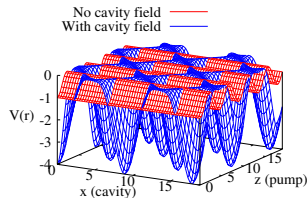
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- Commensurability effects
- Rescale with $\sqrt{2q}$, $\omega_r = \hbar^2 q^2 / 2m$
 - $n_r = N/N_s$, $\omega \rightarrow \tilde{\omega}$
 - $\Omega \rightarrow P$, $\psi \rightarrow \lambda$

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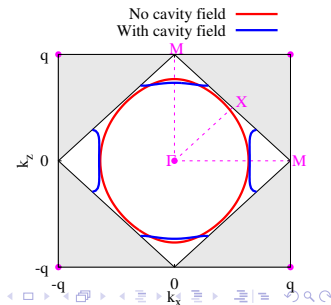
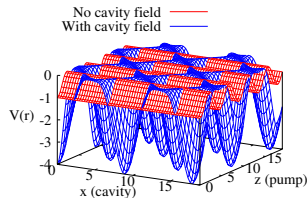
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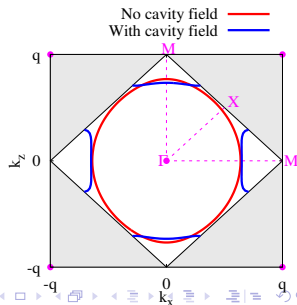
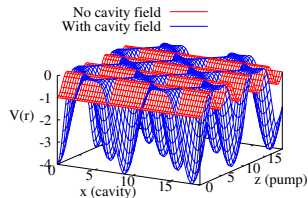
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Free energy to phase diagram

- Free energy $f = F/N_L\omega_r$

$$f(\tilde{\omega}, P, n_F \rightarrow \mu; \lambda) = \tilde{\omega}\lambda^2 + \mu n_F - \frac{1}{\beta} \int_{BZ} d^2k \sum_n \ln \left[1 + e^{-\beta(\epsilon_{\mathbf{k},n} - \mu)} \right]$$

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• Instability, $\lambda \rightarrow \infty$,

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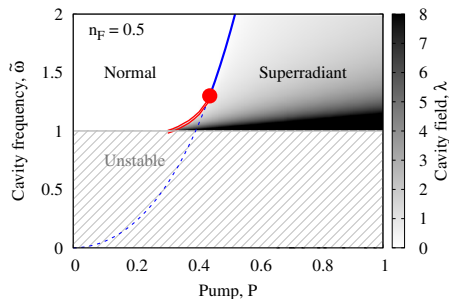
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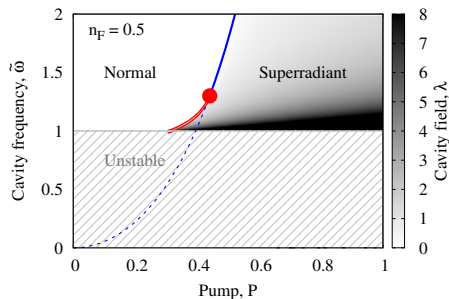
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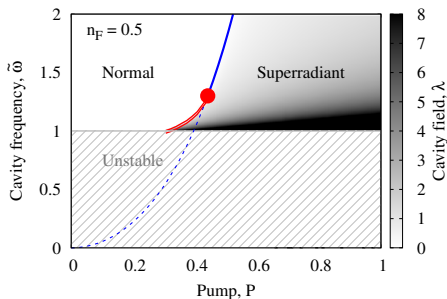
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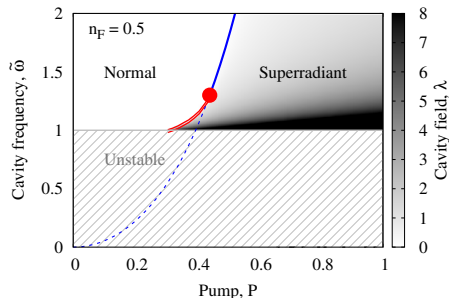
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$$f = a\lambda^2 + b\lambda^4 + c\lambda^6$$

$$b < 0 \text{ at small } P.$$

Origin of first order transition



- $\epsilon_{\mathbf{k},n}$ from $\hat{h} = k^2 \delta_{\mathbf{k},\mathbf{k}'} - V_{\mathbf{k},\mathbf{k}'}$

$$V_{\mathbf{k},\mathbf{k}'} = \lambda^2 \sum_{\mathbf{s}} \delta_{\mathbf{k},\mathbf{k}'+\mathbf{s}\sqrt{2}\hat{\mathbf{x}}} + P\lambda \sum_{\mathbf{s},\mathbf{s}'} \delta_{\mathbf{k},\mathbf{k}'+\frac{\mathbf{s}}{\sqrt{2}}\hat{\mathbf{x}}+\frac{\mathbf{s}'}{\sqrt{2}}\hat{\mathbf{z}}} + P^2 \sum_{\mathbf{s}} \delta_{\mathbf{k},\mathbf{k}'+\mathbf{s}\sqrt{2}\hat{\mathbf{z}}}$$

Landau expansion: $f = a(\tilde{\omega}, P, n_F)\lambda^2 + b(P, n_F)\lambda^4 + c(P, n_F)\lambda^6$

• Second order perturbation theory,

$$= \lambda^2 m_{\mathbf{k},\mathbf{k}'}^2 (\epsilon_{\mathbf{k}} - \epsilon_{\mathbf{k}'})$$

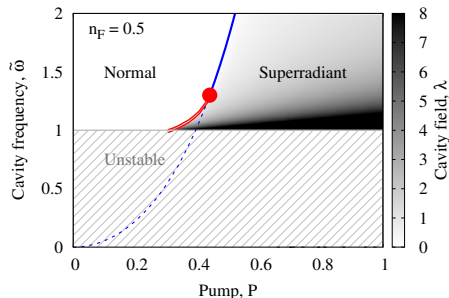
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• But needs state $\phi(x, z) = \cos(\sqrt{2}x)$

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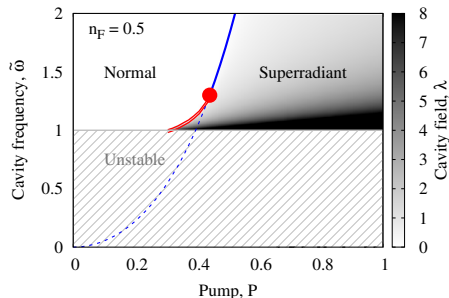
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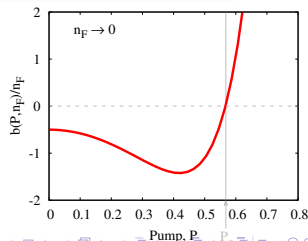
$$+ P\lambda \sum_{s,s'} \delta_{\mathbf{k},\mathbf{k}'+\frac{s}{\sqrt{2}}\hat{x}+\frac{s'}{\sqrt{2}}\hat{z}}$$

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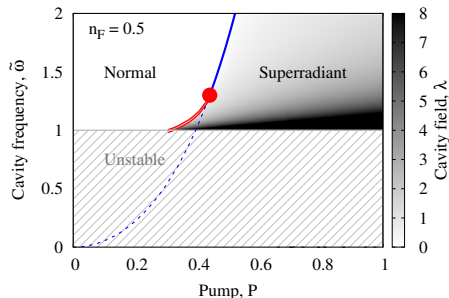
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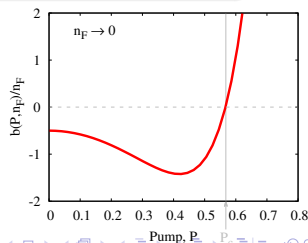
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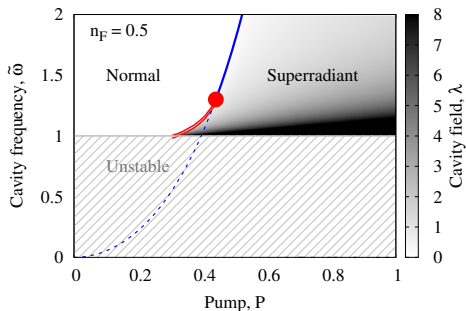
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Higher fillings

$$f = a\lambda^2 + b\lambda^4 + c\lambda^6$$

- Phase diagram unchanged for $n_F < 1$
- 2nd order line $a = 0$
- Tricritical **•** at $a = b = 0$

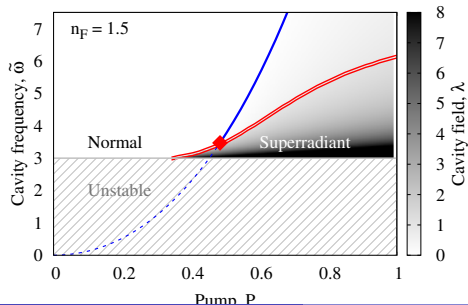
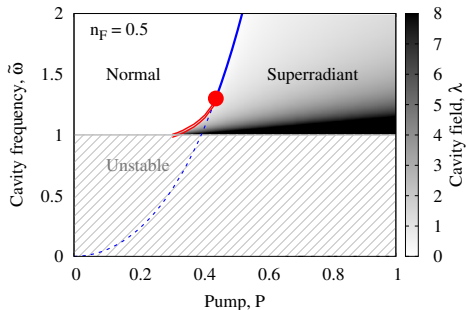


- 2nd band, new structure.
- Critical end-point **•**
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- SR-SR phase boundary
- No symmetry breaking
- Liquid-gas type (metamagnetic)

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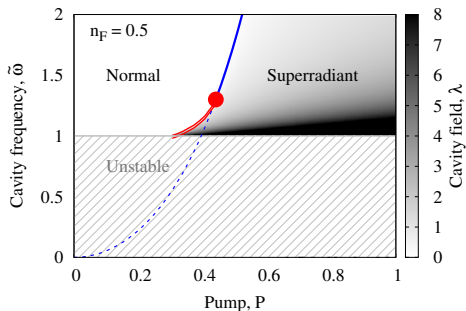
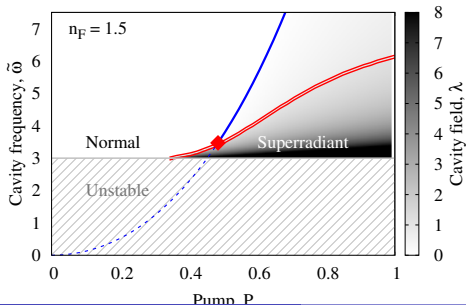
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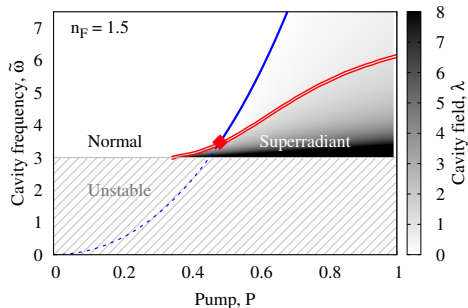
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Why liquid gas?



Liquid-gas transition

- $f(\lambda) \rightarrow$ multiple minima
- Higher orders in λ

• Plot bands $\text{Im}[\epsilon_k(\omega_k, P)]$

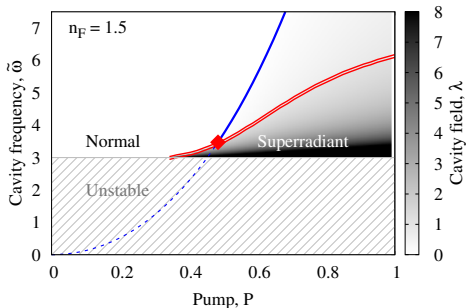
• Contribution of 2nd band

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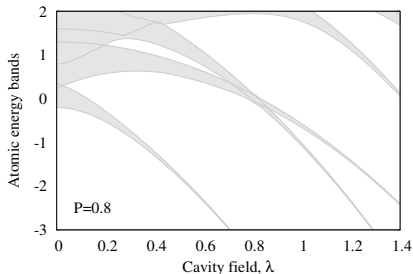
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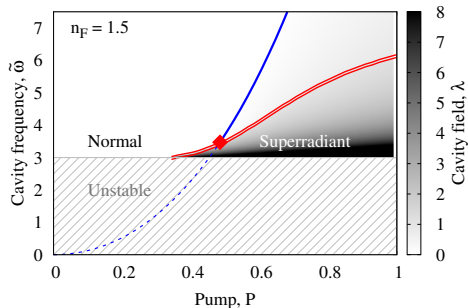
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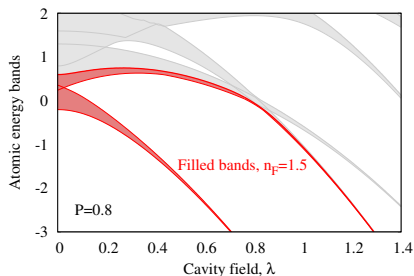
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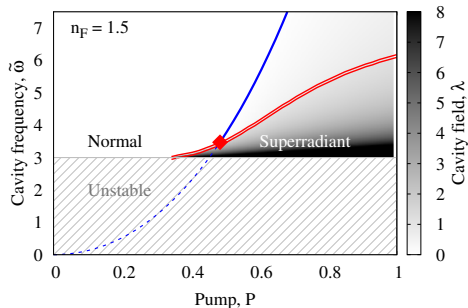
- p_1, p_2 orbitals cross at $P = \lambda$
- $n > 1$ bands initially go up

Liquid-gas transition

- $f(\lambda) \rightarrow$ multiple minima
- Higher orders in λ



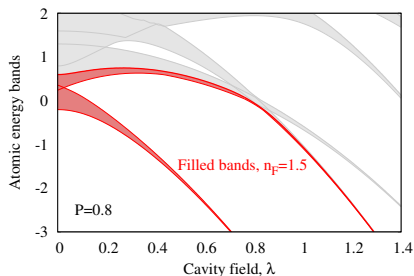
Why liquid gas?



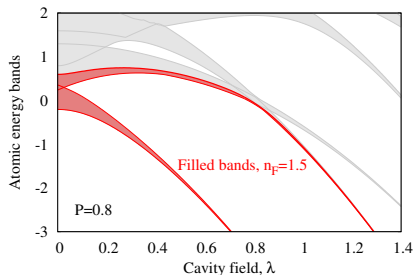
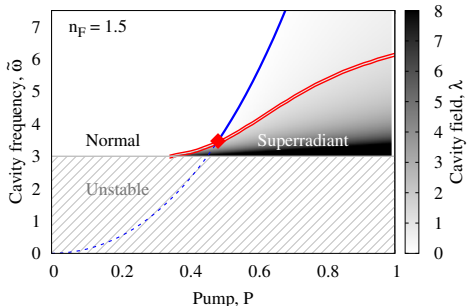
- Plot bands $\inf_k [\epsilon_{\mathbf{k}, n}]$
- Contribution of 2nd band
- Non-trivial form:
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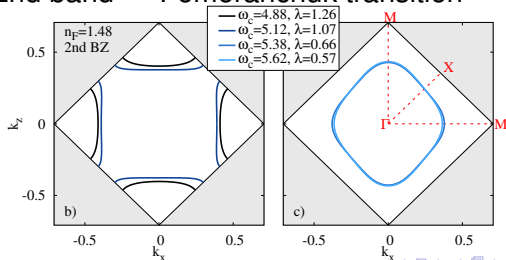
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Signatures: Changing FS topology

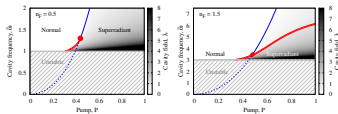


Near half-filled 2nd band — Pomeranchuk transition



Evolution with density

- Phase topology change:

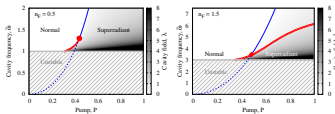


- Fix P , plot vs n_F
- SR–SR alter critical point \odot
- Peak in 2nd order line $0 = a(\tilde{\omega}, n_F, P) = \tilde{\omega} + \xi(P, n_F)$
- Susceptibility ξ asymptote $P \rightarrow \infty$

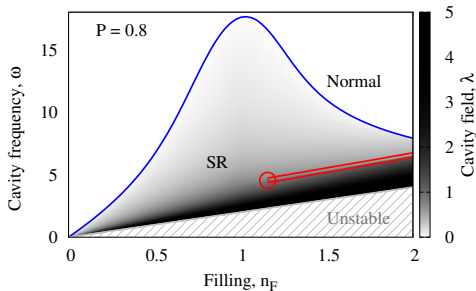
$$\xi \simeq 16P^2 \ln \left| \frac{1 - n_F}{1 + n_F} \right|$$

Evolution with density

- Phase topology change:



- Fix P , plot vs n_F



- SR-SR after critical point

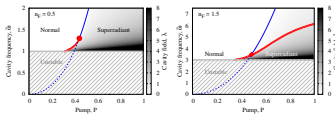
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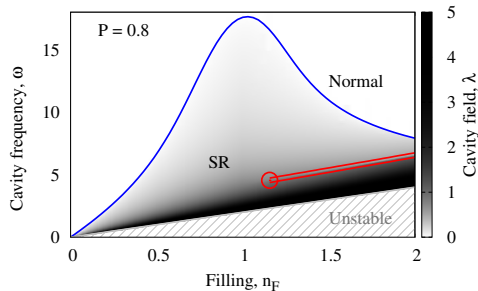
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Evolution with density

- Phase topology change:



- Fix P , plot vs n_F
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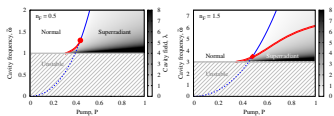


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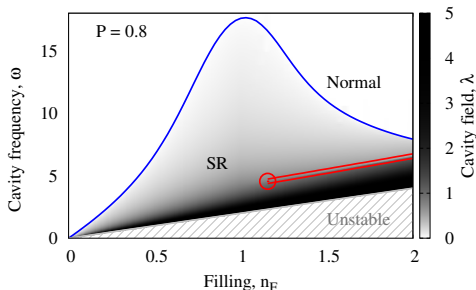
- Phase topology change:



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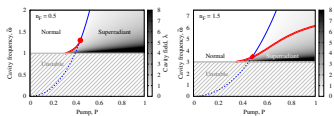
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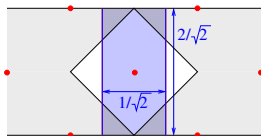
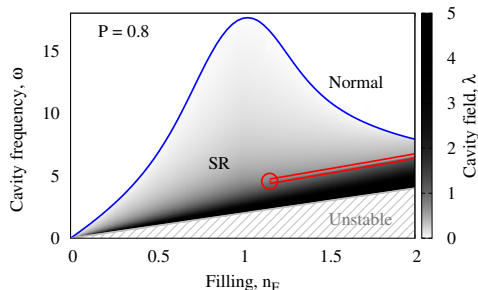


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- Nesting of $\cos\left(\frac{x}{\sqrt{2}}\right) \cos\left(\frac{z}{\sqrt{2}}\right)$ at $n_F = 1$



Outline

1 Dicke model, superradiance and no-go theorem

2 Superradiance and self-organisation

- Raman scheme
- Rayleigh scheme and hierarchies of H_{eff}
- Generalized Dicke equilibrium theory

3 Fermionic self organisation

- Equilibrium phase diagrams
- Landau theory and microscopics
- Evolution with filling

4 Open system dynamics

- Linear stability with losses
- Attractors of the Dicke model phases

Open system: Stable attractors

- Open system, $\dot{\rho} = -i[H, \rho] - \kappa\mathcal{L}[\psi]$
- Instead of $\min(F) \rightarrow$, stable attractors

- What survives

- Normal-SR boundary: Consider $\psi = ue^{-i\nu t} + v^* e^{i\nu^* t}$

- What must change

- Unstable region \rightarrow new attractors

- Known unknowns

- First order transitions/multistability?
 - Spinodal lines?
 - Limit cycles?

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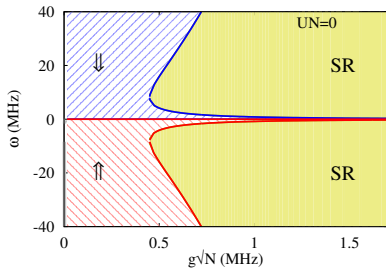
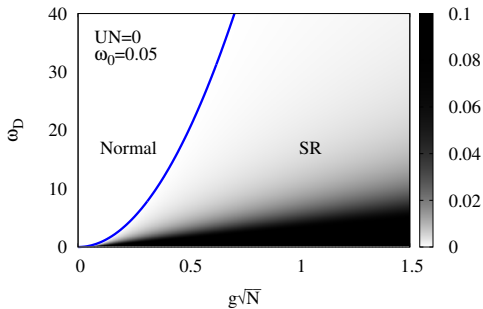
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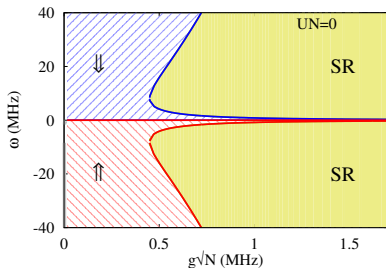
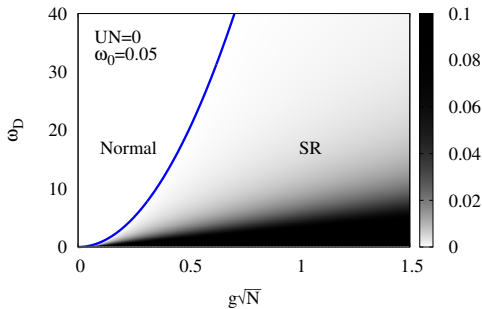
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Equilibrium Dicke vs open phase diagram, $UN = 0$



- Shift boundary $(\kappa^2 + \omega^2)/\omega = -\zeta(\omega_0)$
- Allow negative $\omega \rightarrow$ inverted

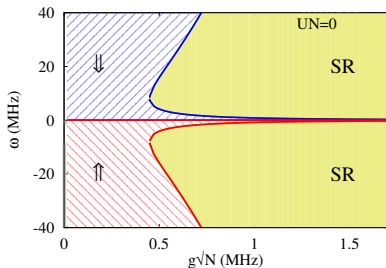
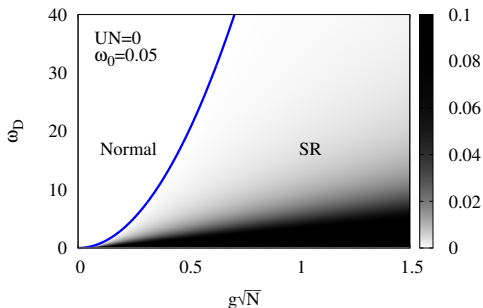
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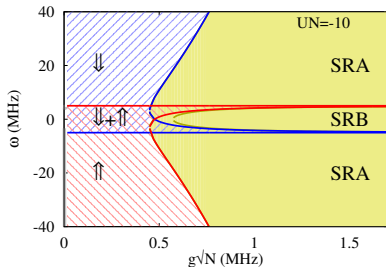
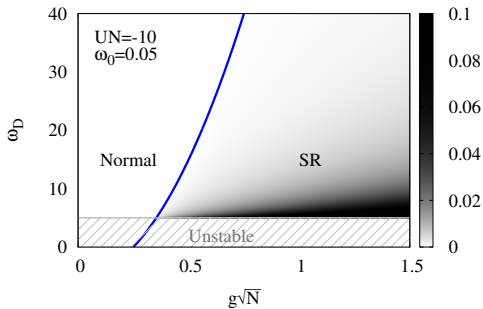
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Equilibrium Dicke vs open phase diagram, $UN = 0$



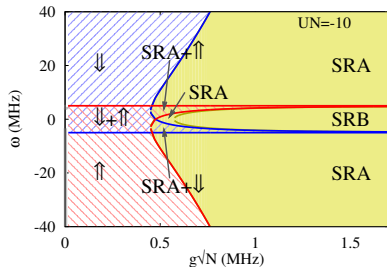
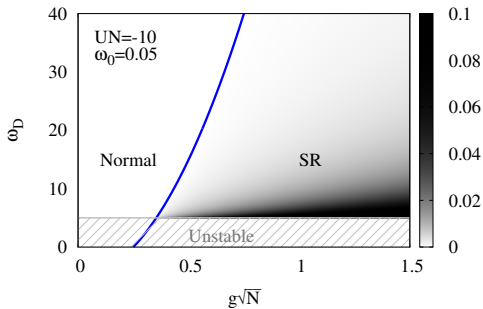
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... Dicke ... $UN = -10\text{MHz}$



- Coexistence regions
- Unstable \rightarrow SRB

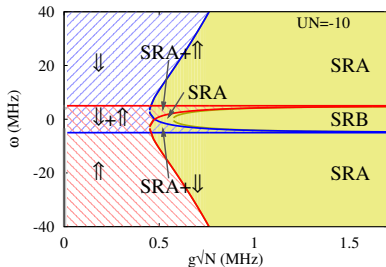
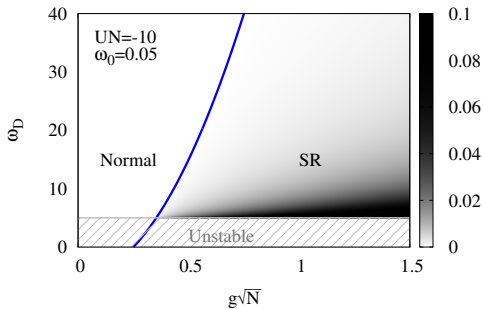
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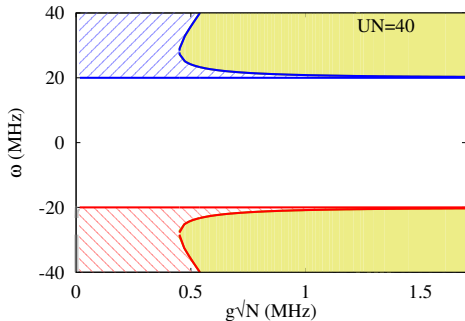
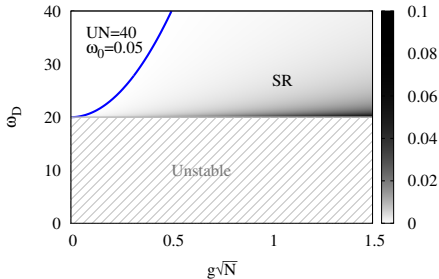
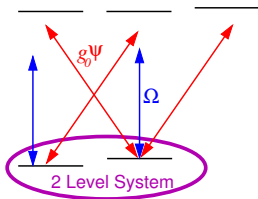
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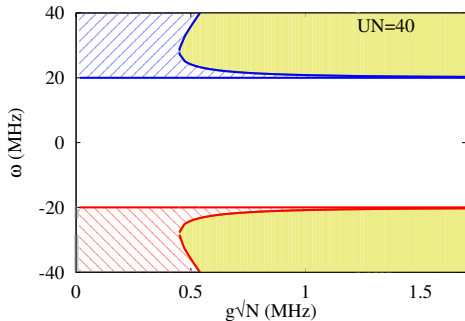
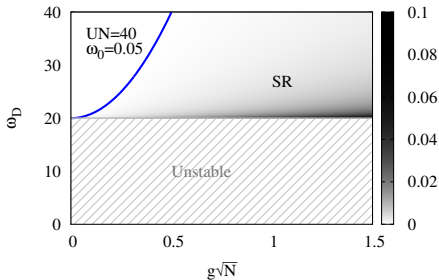
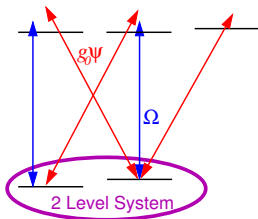
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Changing U :



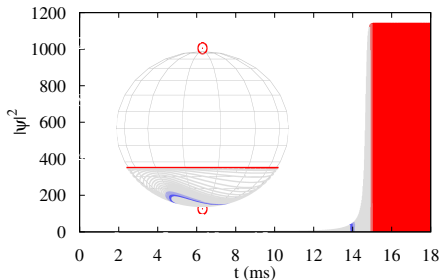
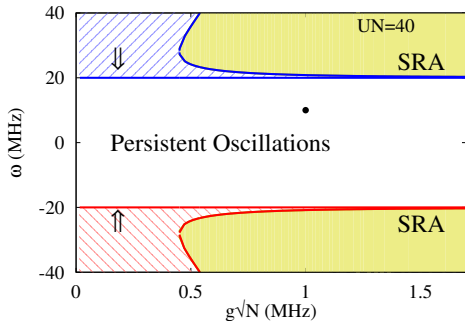
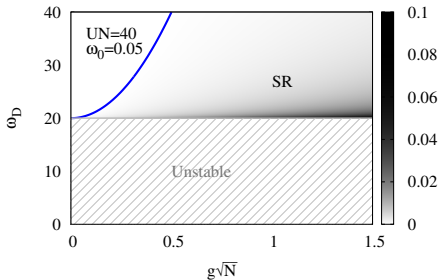
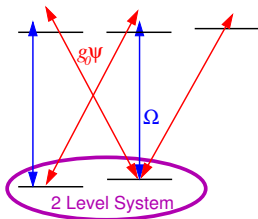
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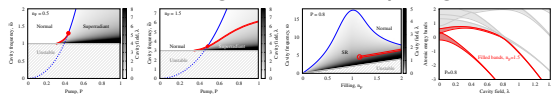
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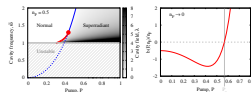


Summary

- Fermions self organisation, liquid gas, and multicritical points

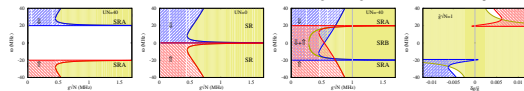


- First order transitions for bosons, outside Dicke model



JK, Bhassen, Simons *et al.* arXiv:1308.????

- Dicke model shows many dynamical phases



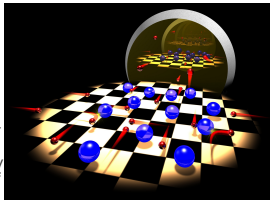
JK *et al.* PRL '10, Bhaseen *et al.* PRA '12

Many body quantum optics and correlated states of light

9:00 am on Monday 28 October 2013 – 5:00 pm on Tuesday 29 October 2013

at: **The Royal Society at Chicheley Hall, home of the Kavli Royal Society International Centre, Buckinghamshire**

Theo Murphy international scientific meeting organised by Dr Jonathan Keeling, Professor Steven Girvin, Dr Michael Hartmann and Professor Peter Littlewood FRS.



List of speakers and chairs

Professor Iacopo Carusotto, Professor Andrew Cleland, Professor Hui Deng, Professor Tilman Esslinger, Professor Rosario Fazio, Professor Ed Hinds, Professor Andrew Houck, Professor Ataç İmamoğlu, Professor Jens Koch, Professor Misha Lukin, Professor Martin Plenio, Professor Arno Rauschenbeutel, Professor Timothy Spiller, Professor Jacob Taylor, Professor Hakan Türeci, Professor Andreas Wallraff

Attending this event

This is a residential conference which allows for increased discussion and networking. It is free to attend, however participants need to cover their accommodation and catering costs if required. Places are limited and therefore pre-registration is essential.

5 Classical dynamics

6 Ferroelectric transition

7 Grand canonical

Classical dynamics of the extended Dicke model

Open dynamical system:

$$H = \omega\psi^\dagger\psi + \omega_0\mathbf{S}^z + g(\psi + \psi^\dagger)(\mathbf{S}^- + \mathbf{S}^+) + U\mathbf{S}_z\psi^\dagger\psi.$$
$$\partial_t\rho = -i[H, \rho] - \kappa(\psi^\dagger\psi\rho - 2\psi\rho\psi^\dagger + \rho\psi^\dagger\psi)$$

- Neglects quantum fluctuations
- Linearisation about fixed point \rightarrow stability, spectrum

[JK *et al.* PRL '10, Bhaseen *et al.* PRA '12]

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Classical EOM
($|\mathbf{S}| = N/2 \gg 1$)

$$\dot{\mathbf{S}}^- = -i(\omega_0 + U|\psi|^2)\mathbf{S}^- + 2ig(\psi + \psi^*)\mathbf{S}^z$$

$$\dot{\mathbf{S}}^z = ig(\psi + \psi^*)(\mathbf{S}^- - \mathbf{S}^+)$$

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[JK *et al.* PRL '10, Bhaseen *et al.* PRA '12]

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[JK *et al.* PRL '10, Bhaseen *et al.* PRA '12]

Fixed points (steady states)

$$0 = i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z$$

$$0 = ig(\psi + \psi^*)(S^- - S^+)$$

$$0 = -[\kappa + i(\omega + US^z)]\psi - ig(S^- + S^+)$$

- $\psi = 0, S = (0, 0, \pm N/2)$ always a solution.
- If $g > g_c, \psi \neq 0$ too
 - $S^z = -g[S^-] = 0$
 - $\psi' = \Re[\psi] = 0$

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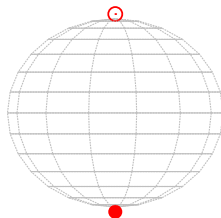
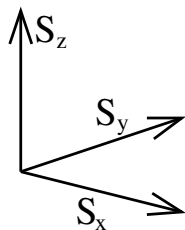
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$$\Rightarrow S^z = -g[S^-] = 0$$

$$\Rightarrow \psi = \Re[\psi] = 0$$



Small g : \uparrow, \downarrow only.
 $(\omega = 30\text{MHz}, UN = -40\text{MHz})$

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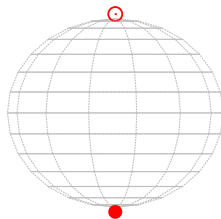
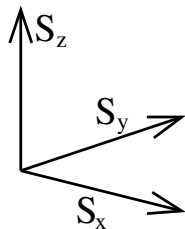
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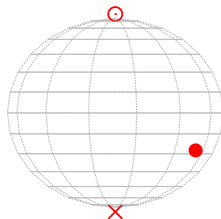
A $S^y = -\Im[S^-] = 0$

B $\psi' = \Re[\psi] = 0$



Small g : \uparrow, \downarrow only.

($\omega = 30\text{MHz}, UN = -40\text{MHz}$)



Larger g : SR too.

Ferroelectric transition

Atoms in Coulomb gauge

$$H = \sum \omega_k a_k^\dagger a_k + \sum_i [p_i - eA(r_i)]^2 + V_{\text{coul}}$$

Ferroelectric transition

Atoms in **Coulomb gauge**

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Two-level systems — dipole-dipole coupling

$$H = \omega_0 S^z + \omega \psi^\dagger \psi + g(S^+ + S^-)(\psi + \psi^\dagger) + N\zeta(\psi + \psi^\dagger)^2 - \eta(S^+ - S^-)^2$$

(nb $g^2, \zeta, \eta \propto 1/V$).

Ferroelectric transition

Atoms in **Coulomb gauge**

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Gauge transform to dipole gauge $\mathbf{D} \cdot \mathbf{r}$

$$H = \omega_0 S^z + \omega \psi^\dagger \psi + \bar{g}(S^+ - S^-)(\psi - \psi^\dagger)$$

“Dicke” transition at $\omega_0 < N\bar{g}^2/\omega \equiv 2\eta N$

But, ψ describes **electric displacement**

Grand canonical ensemble

Grand canonical ensemble:

- If $H \rightarrow H - \mu(S^z + \psi^\dagger\psi)$, need only: $g^2 N > (\omega - \mu)|\omega_0 - \mu|$

- Fix density / fix $\mu > 0$ — pumping

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 - $g^2 N > (\omega - \mu)(\omega_0 - \mu)$
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[Eastham and Littlewood, PRB '01]

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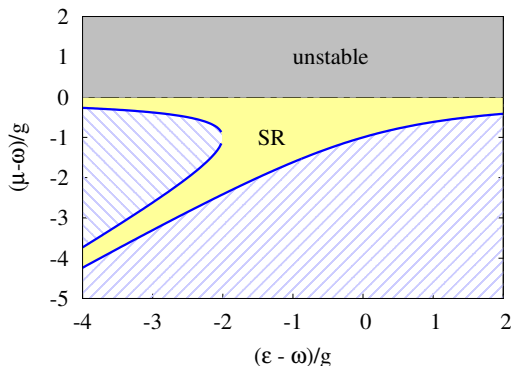
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