

Polariton and photon condensates with organic molecules

Jonathan Keeling



University of
St Andrews

600
YEARS



Telluride, July 2013

Polariton and photon condensates with organic molecules: The “Dicke-Holstein” model

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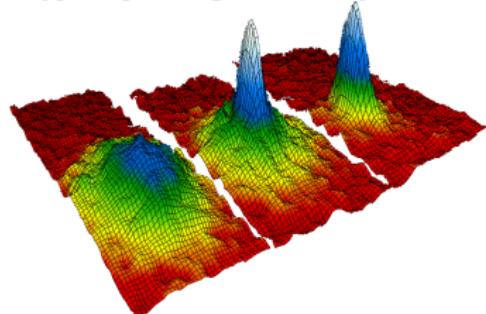
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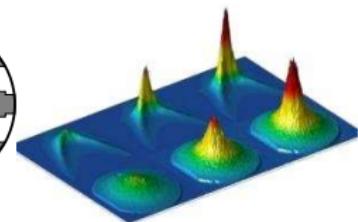
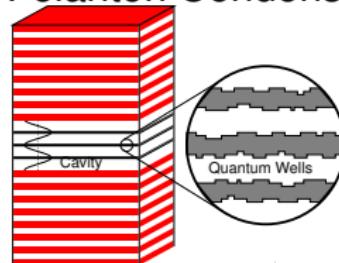
Coherent states of matter and light

Atomic BEC $T \sim 10^{-7}$ K



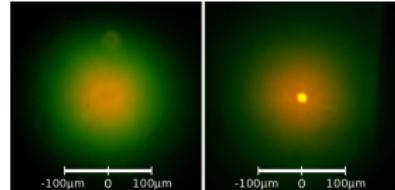
[Anderson *et al.* Science '95]

Polariton Condensate $T \sim 20$ K



[Kasprzak *et al.* Nature, '06]

Photon Condensate
 $T \sim 300$ K

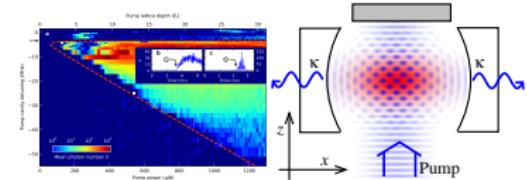


[Klaers *et al.* Nature, '10]

Laser
 $T \sim ?, < 0, \infty$



Superradiance transition
 $T \sim 0$



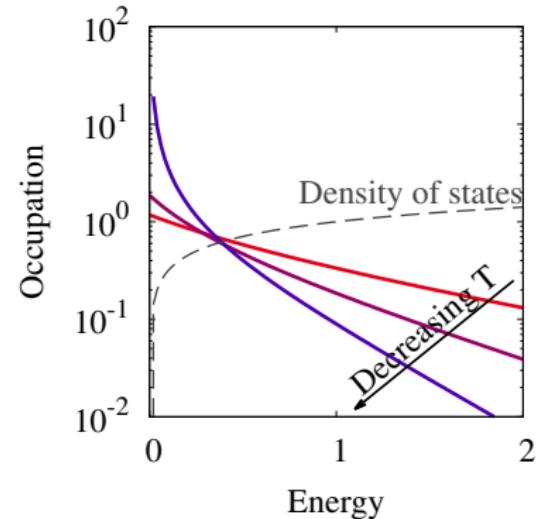
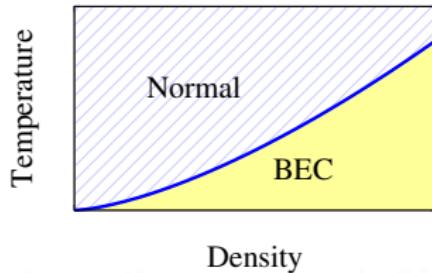
[Baumann *et al.* Nature, '10]

“Textbook” BEC

- **Non-interacting** viewpoint

- ▶ BE distribution: $\mu < \omega_0$

- ▶ $T_c = \frac{2\pi\hbar^2}{m} \left(\frac{n}{\xi_d}\right)^{2/d}$



- Interacting approach (MBG)

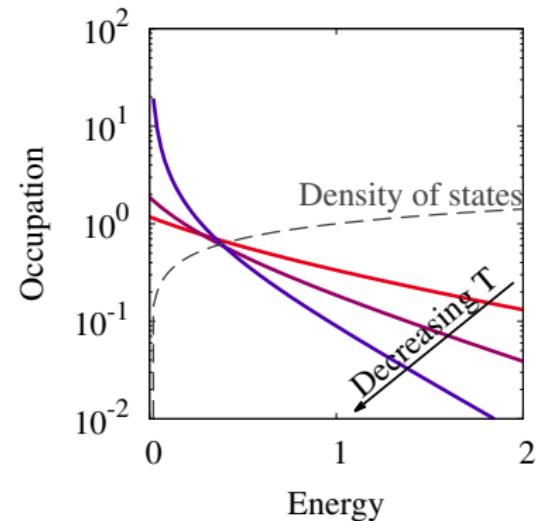
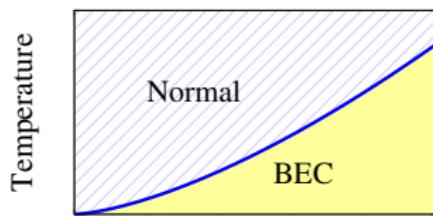
- Mean field approach (MF)

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- **Interacting** approach (WIDBG)

$$H = \sum_k \omega_k \psi_k^\dagger \psi_k + \frac{g}{2V} \sum_{k,k',q} \psi_{k+q}^\dagger \psi_{k'-q}^\dagger \psi_{k+q} \psi_k$$

- ▶ Mean field: $|\psi|^2 = (\mu - \omega_0)/V$

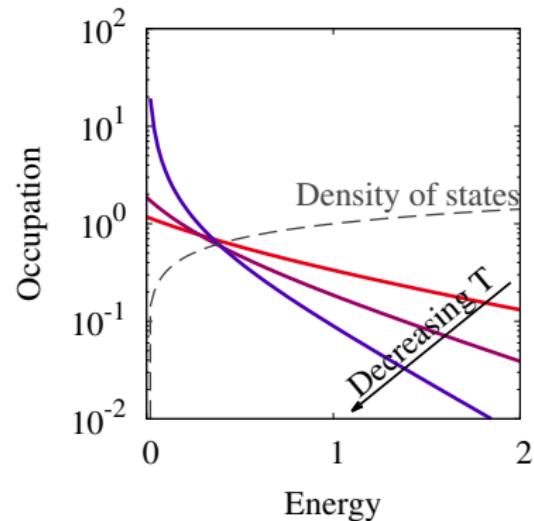
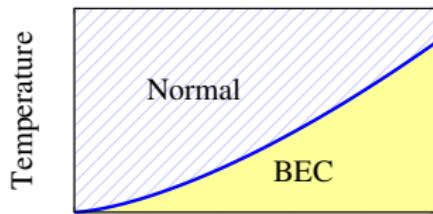
For $\mu < \omega_0$, the mean field energy vanishes at $T=0$.

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- ▶ Mean field: $|\psi|^2 = (\mu - \omega_0)/V$
- ▶ Fluctuations deplete condensate, vanishes at $T > T_c$

“Textbook” Laser: Maxwell Bloch equations

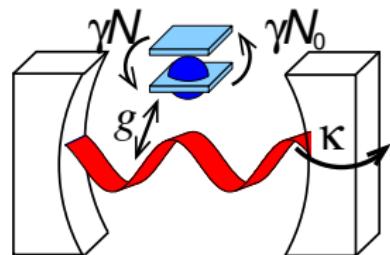
$$H = \omega_0 \psi^\dagger \psi + \sum_{\alpha} \epsilon_{\alpha} S_{\alpha}^z + g_{\alpha, \mathbf{k}} (\psi S_{\alpha}^+ + \psi^\dagger S_{\alpha}^-)$$

Maxwell-Bloch eqns: $P = -i\langle S^- \rangle$, $N = 2\langle S^z \rangle$

$$\partial_t \psi = -i\omega_0 \psi - \kappa \psi + \sum_{\alpha} g_{\alpha} P_{\alpha}$$

$$\partial_t P_{\alpha} = -2i\epsilon_{\alpha} P_{\alpha} - 2\gamma P + g_{\alpha} \psi N_{\alpha}$$

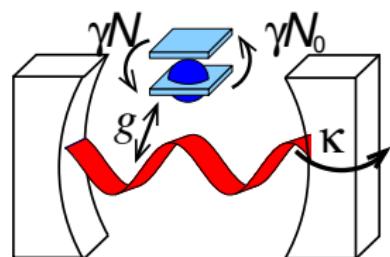
$$\partial_t N_{\alpha} = 2\gamma(N_0 - N_{\alpha}) - 2g_{\alpha}(\psi^* P_{\alpha} + P_{\alpha}^* \psi)$$



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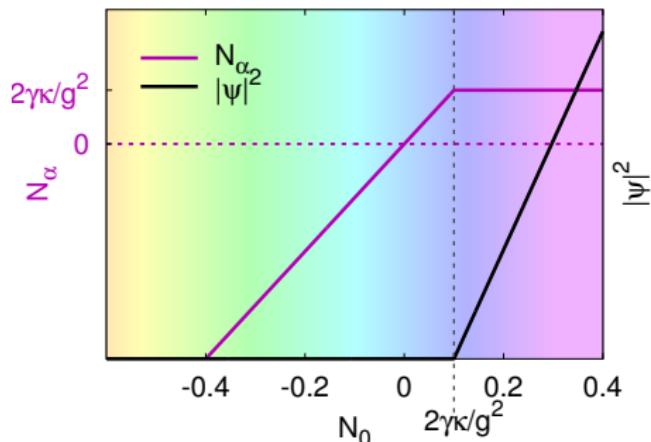
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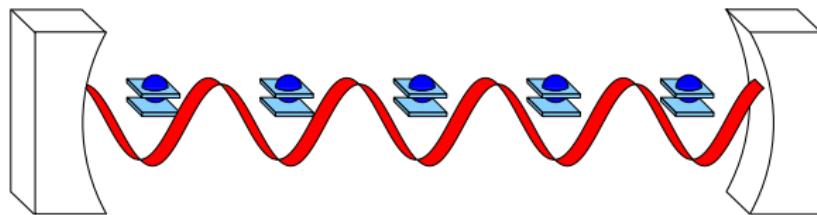
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$|\psi|^2 > 0$ if $N_0 g^2 > 2\gamma\kappa$

“Textbook” Dicke-Hepp-Lieb superradiance

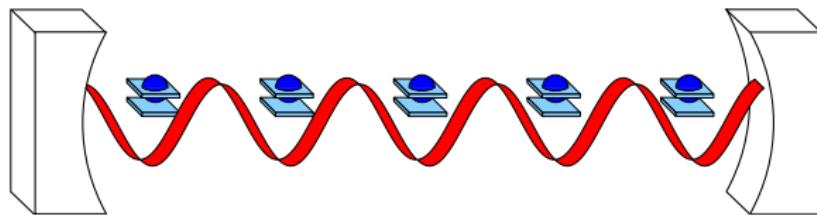


$$H = \omega \psi^\dagger \psi + \sum_{\alpha} \epsilon S_{\alpha}^z + g (\psi^\dagger S_{\alpha}^- + \psi S_{\alpha}^+)$$

- Coherent state: $|\Psi\rangle \rightarrow e^{i\phi_1^{\dagger} + i\phi_2^{\dagger}} |\Omega\rangle$
- Small g , min at $\lambda, \eta = 0$

[Hepp, Lieb, Ann. Phys. '73]

“Textbook” Dicke-Hepp-Lieb superradiance

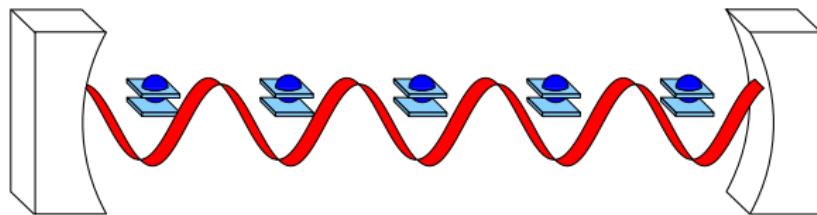


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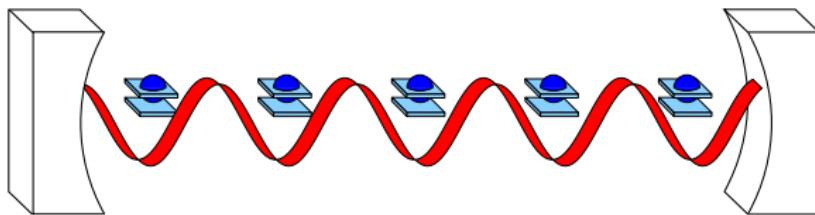
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[Hepp, Lieb, Ann. Phys. '73]

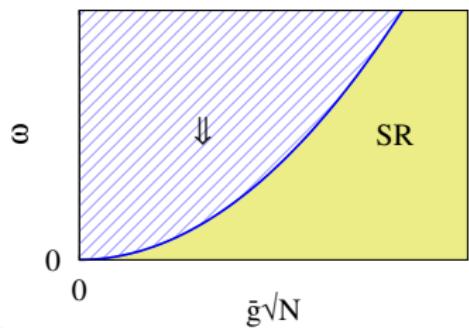
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[Hepp, Lieb, Ann. Phys. '73]

Outline

1 Introduction

- Polariton condensation
- Condensation, superradiance, lasing
 - Condensation vs superradiance transition
 - Non-equilibrium condensation vs lasing

2 Room temperature condensates: Organic polaritons

- Dicke phase diagram with phonons
- Condensation of phonon replicas?
- Ultra-strong phonon coupling?

3 Room temperature condensates: Photons

- Lasing model and thermalisation
- Critical properties

Acknowledgements

GROUP:

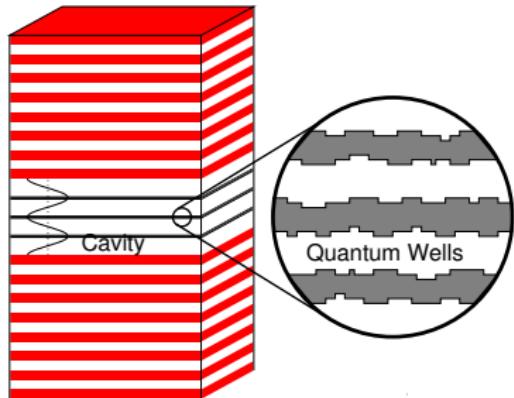


COLLABORATORS: Szymanska (Warwick), Reja (Cam.), Littlewood (ANL)

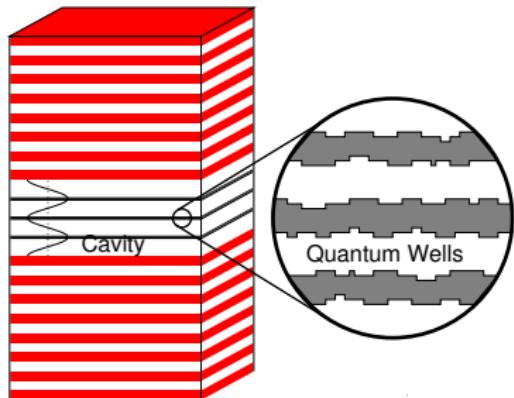
FUNDING:



Microcavity polaritons

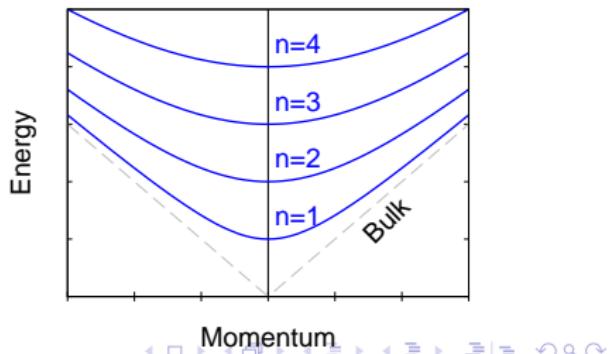


Microcavity polaritons

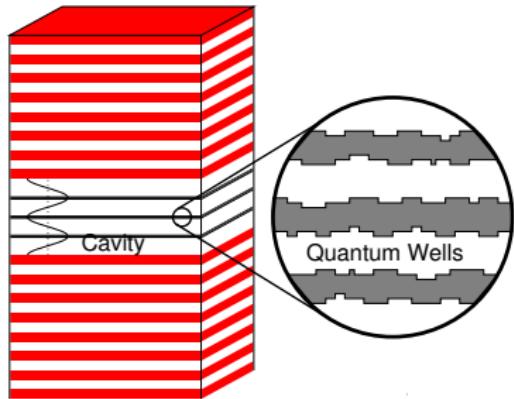


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2 / 2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$

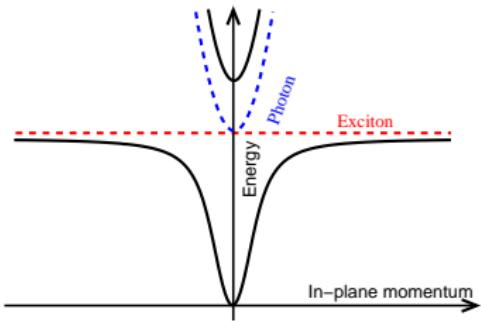


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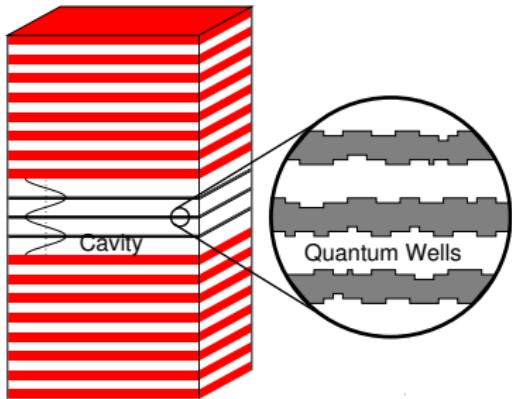


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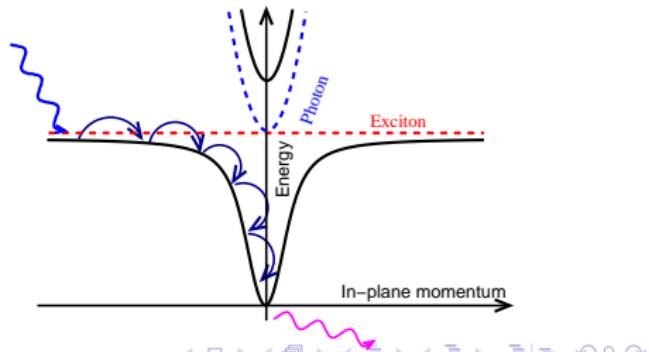


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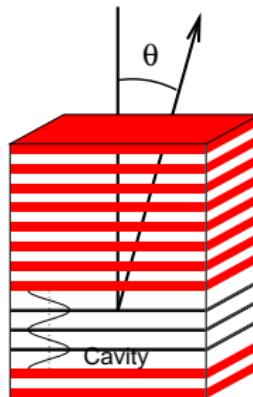
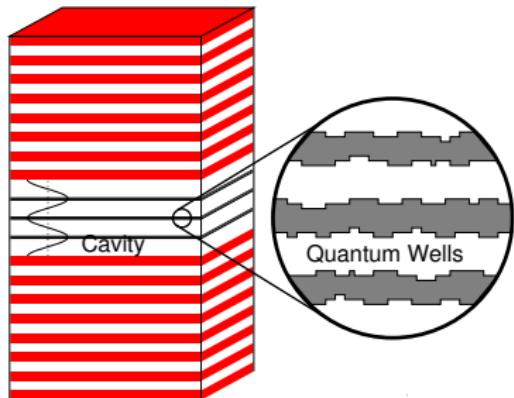
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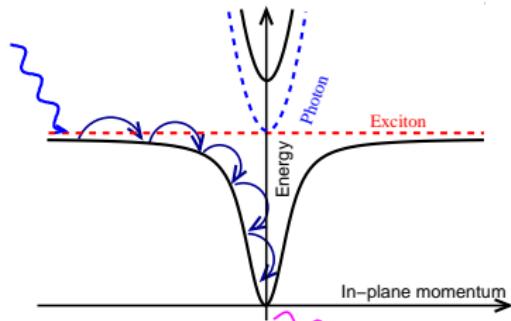


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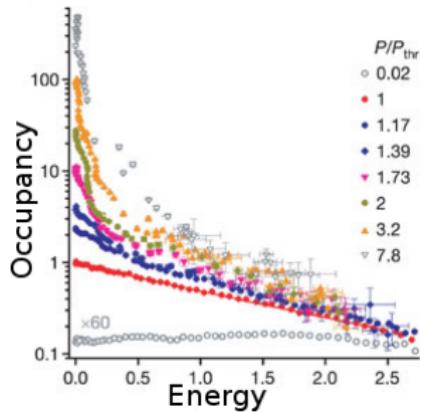
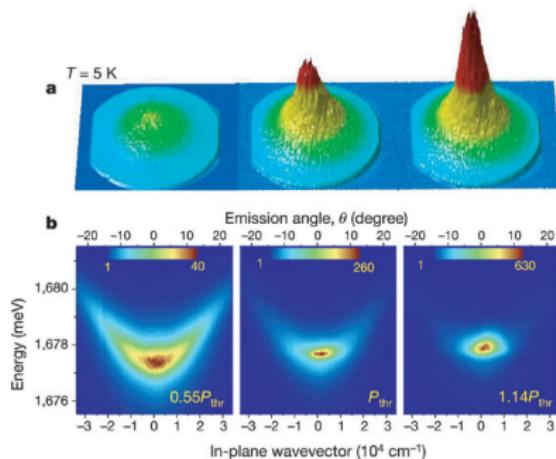
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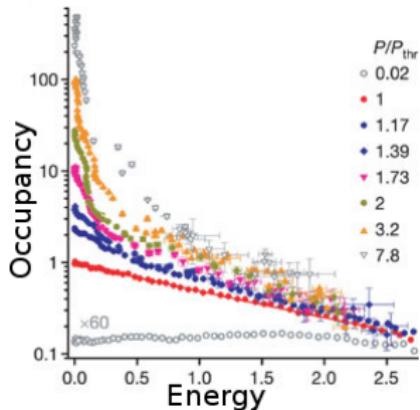
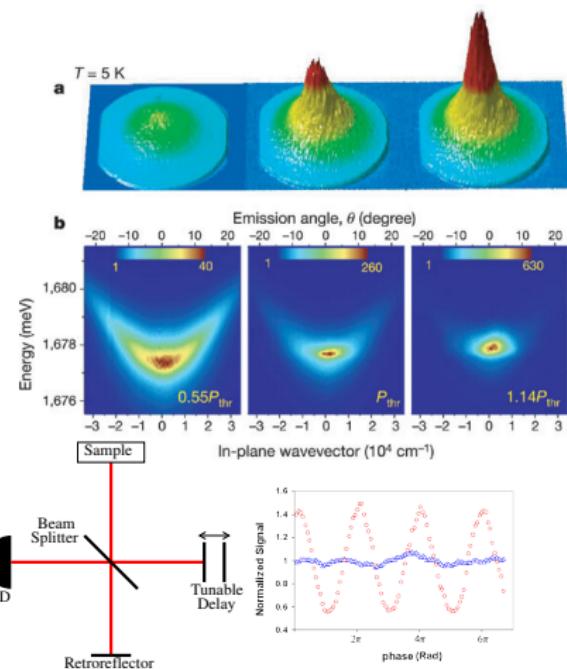


Polariton experiments: occupation and coherence

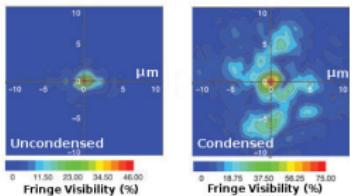


[Kasprzak, *et al.* Nature, '06]

Polariton experiments: occupation and coherence



Coherence map:

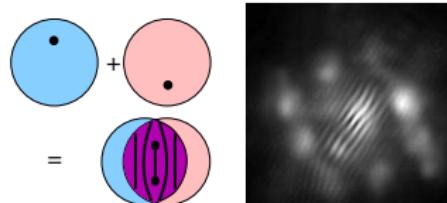


[Kasprzak, et al. Nature, '06]

(Some) other polariton condensation experiments

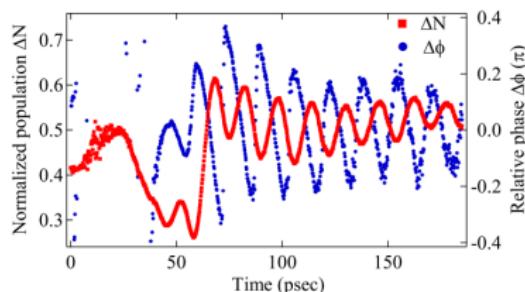
- Quantised vortices

[Lagoudakis *et al.* Nat. Phys. '08; Science '09, PRL '10; Sanvitto *et al.* Nat. Phys. '10; Roumpos *et al.* Nat. Phys. '10]



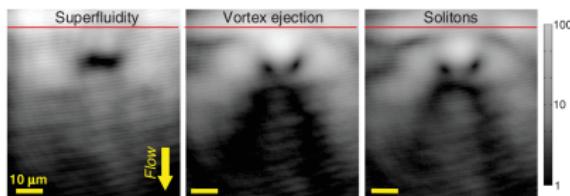
- Josephson oscillations

[Lagoudakis *et al.* PRL '10]



- Pattern formation/Hydrodynamics

[Amo *et al.* Science '11, Nature '09; Wertz *et al.* Nat. Phys '10]



Non-equilibrium condensation vs lasing

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2 Room temperature condensates: Organic polaritons

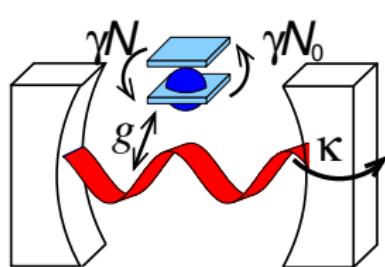
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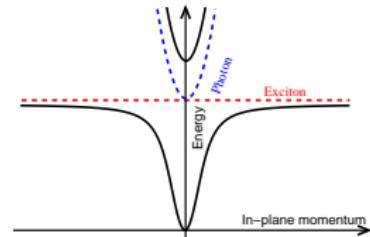
- Lasing model and thermalisation
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Lasing-condensation crossover model

- Use model that can show lasing and condensation:



\Leftrightarrow



Lasing-condensation crossover model

- Use model that can show lasing and condensation:



Dicke model:

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \sum_{\alpha} [\epsilon S_\alpha^z + g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} S_\alpha^+ + \text{H.c.}]$$

Dicke-Hepp-Lieb superradiance and modes

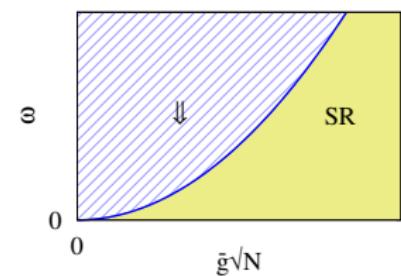
$$H = \omega\psi^\dagger\psi + \epsilon S^z + g(\psi^\dagger S^- + \psi S^+)$$

Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

Normal state, $S^z = -N/2 + \bar{B}B$

$$H = \omega\psi^\dagger\psi + \epsilon B^\dagger B + g\sqrt{N}(\phi^\dagger B + \phi B^\dagger)$$

Excitation cost E



[Hepp, Lieb, Ann. Phys. '73]

$$(E-\omega)(E-\epsilon) = g^2 N$$

Transition when $E = 0$

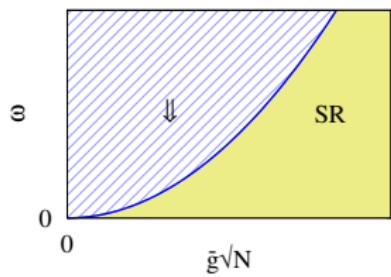
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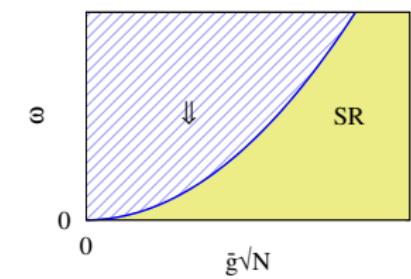
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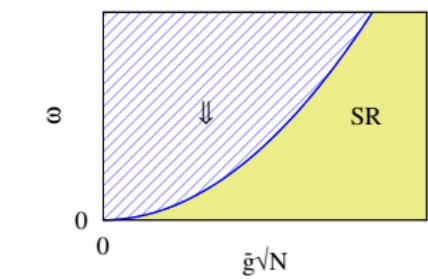
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- Transition when $E = 0$

Grand canonical ensemble

Grand canonical ensemble:

- If $H \rightarrow H - \mu(S^z + \psi^\dagger\psi)$, need only: $g^2N > (\omega - \mu)|\epsilon - \mu|$

Fix density, then add damping

→ Transition at:
 $g^2N > (\omega - \mu)(\epsilon - \mu)$
 γ hits lowest mode

[Eastham and Littlewood, PRB '01]

Grand canonical ensemble

Grand canonical ensemble:

- If $H \rightarrow H - \mu(S^z + \psi^\dagger\psi)$, need only: $g^2N > (\omega - \mu)|\epsilon - \mu|$
- Fix density / fix $\mu > 0$ — pumping

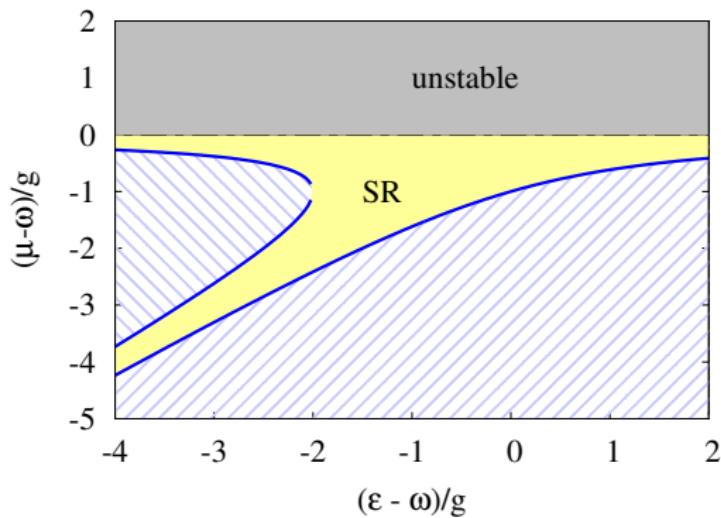
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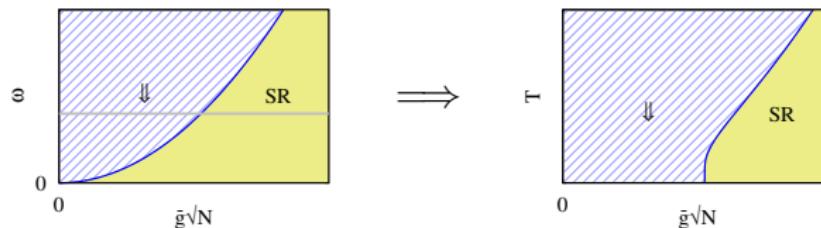
[Eastham and Littlewood, PRB '01]

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Grand canonical Dicke, finite temperature

- Finite temperature:

$$Ng^2 \tanh(\beta\epsilon) > \omega\epsilon$$



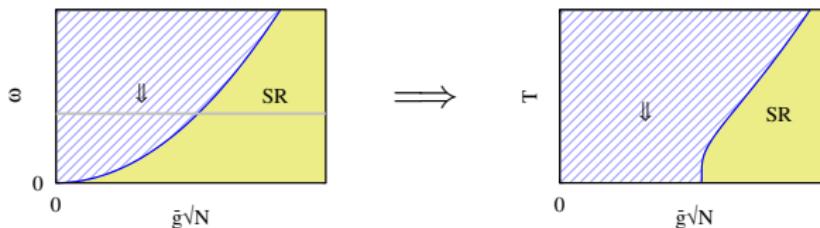
[Hepp, Lieb, Ann. Phys. '73]

With chemical potential $Ng^2 \tanh(\beta(\epsilon - \mu)) > (\omega - \mu)(\epsilon - \mu)$

Grand canonical Dicke, finite temperature

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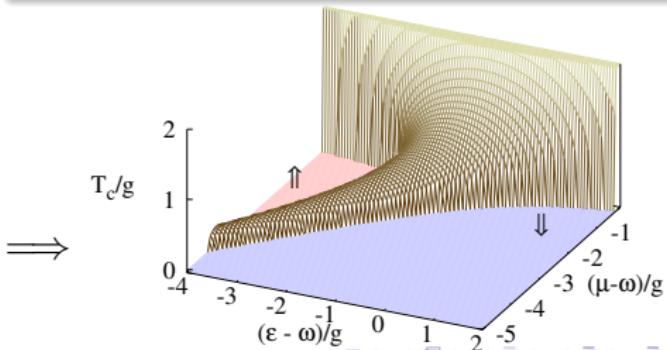
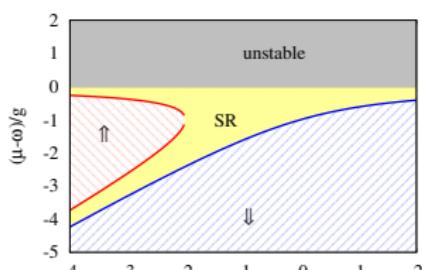
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[Hepp, Lieb, Ann. Phys. '73]

- With chemical potential

$$Ng^2 \tanh(\beta(\epsilon - \mu)) > (\omega - \mu)(\epsilon - \mu)$$



Polariton model and equilibrium results

Localised excitons, propagating photons

$$H - \mu N = \sum_{\mathbf{k}} (\omega_{\mathbf{k}} - \mu) \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \sum_{\alpha} (\epsilon_{\alpha} - \mu) S_{\alpha}^z + g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} S_{\alpha}^+ + \text{H.c.}$$

Self-consistent polarisation and field

$$(\omega - \mu) \psi = \sum_{\alpha} \frac{g_{\alpha}^2 \psi}{2 E_{\alpha}} \tanh(\beta E_{\alpha}), \quad E_{\alpha}^2 = \left(\frac{\epsilon_{\alpha} - \mu}{2} \right)^2 + g_{\alpha}^2 |\psi|^2$$

Polariton model and equilibrium results

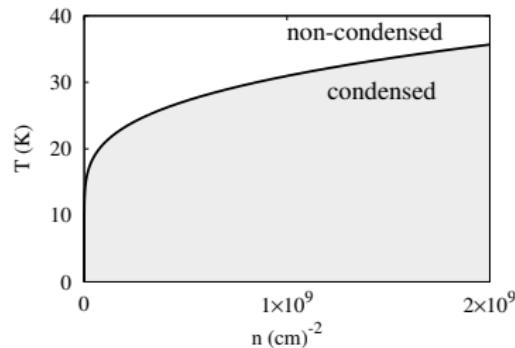
Localised excitons, propagating photons

$$H - \mu N = \sum_{\mathbf{k}} (\omega_{\mathbf{k}} - \mu) \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \sum_{\alpha} (\epsilon_{\alpha} - \mu) S_{\alpha}^z + g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} S_{\alpha}^+ + \text{H.c.}$$

Self-consistent polarisation and field

$$(\omega - \mu) \psi = \sum_{\alpha} \frac{g_{\alpha}^2 \psi}{2 E_{\alpha}} \tanh(\beta E_{\alpha}), \quad E_{\alpha}^2 = \left(\frac{\epsilon_{\alpha} - \mu}{2} \right)^2 + g_{\alpha}^2 |\psi|^2$$

Phase diagram:



Polariton model and equilibrium results

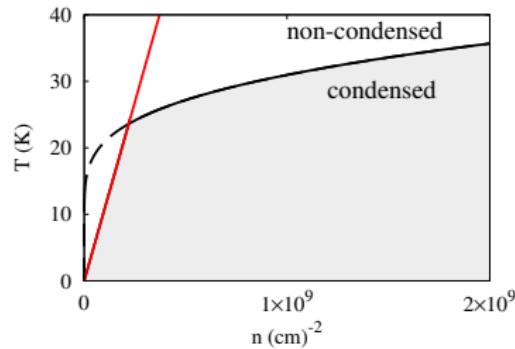
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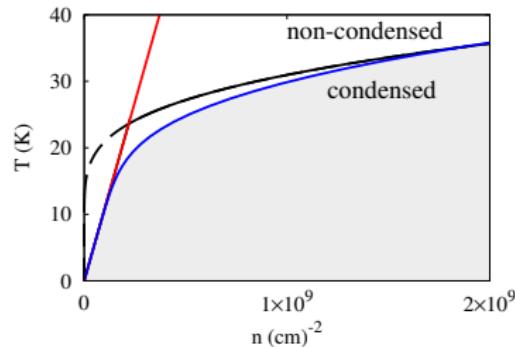
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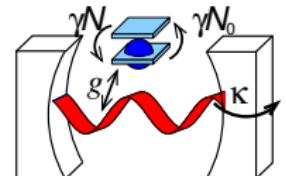
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Simple Laser: Maxwell Bloch equations

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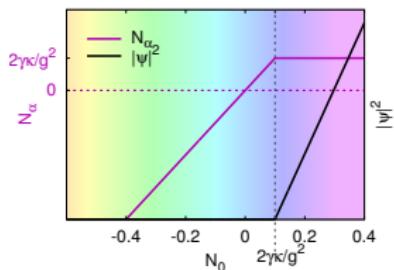


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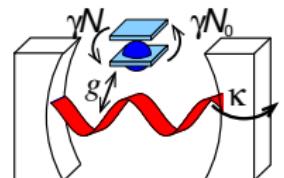
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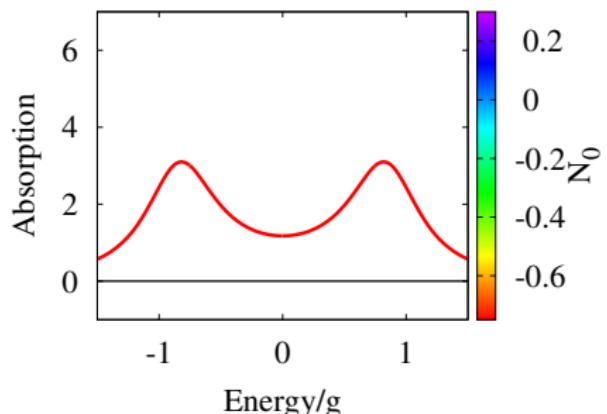
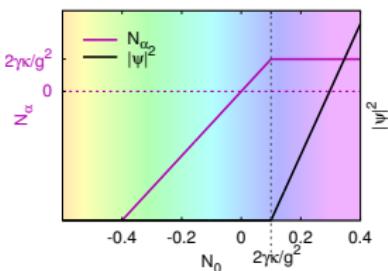


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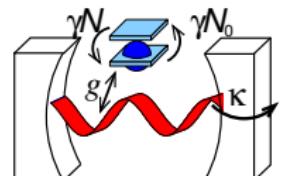


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Inversion causes collapse before lasing @ $g^2 N_0 = 2\gamma\kappa$

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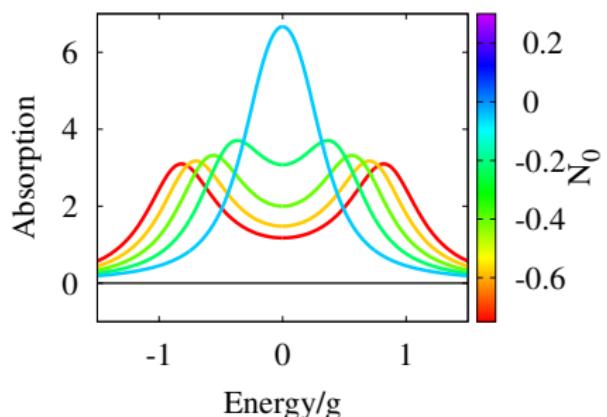
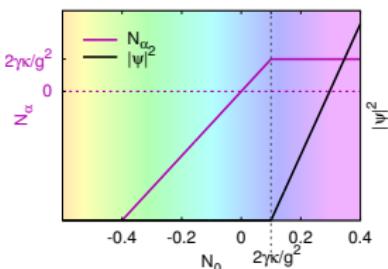


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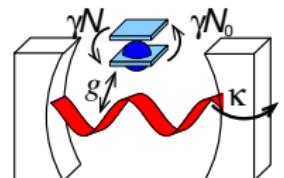
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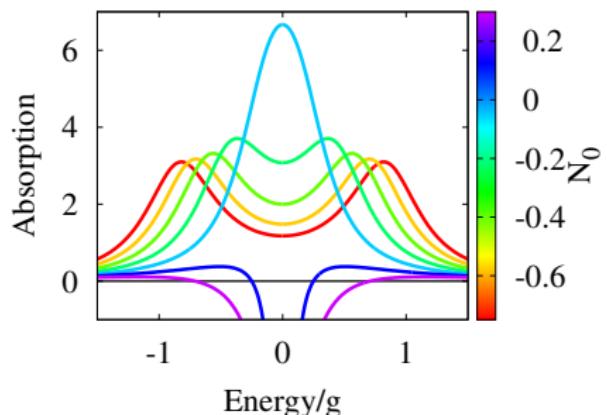
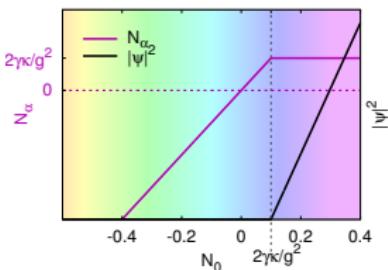


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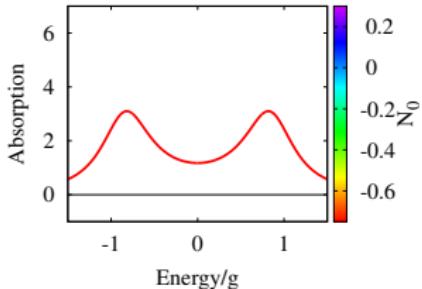
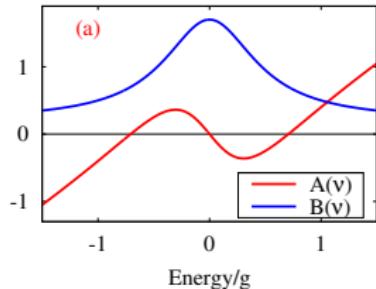
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Poles of Retarded Green's function and gain

$$\left[D^R(\nu) \right]^{-1} = \nu - \omega_k + i\kappa + \frac{g^2 N_0}{\nu - 2\epsilon + i2\gamma}$$

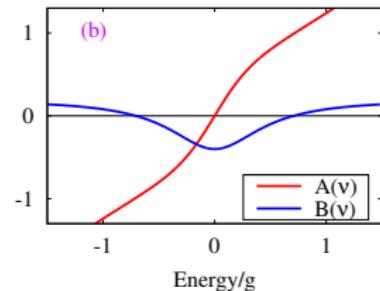
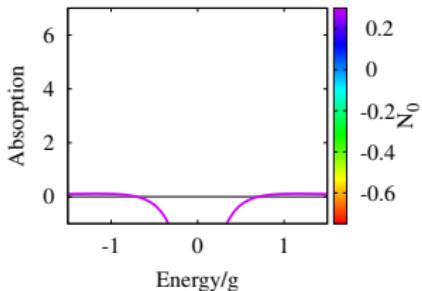
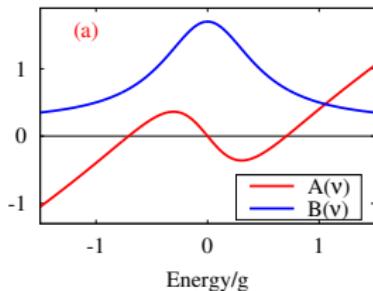
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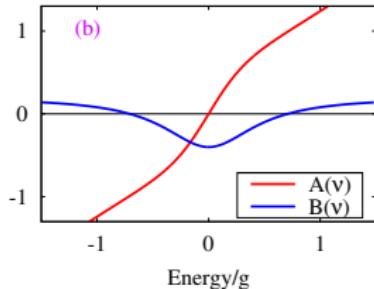
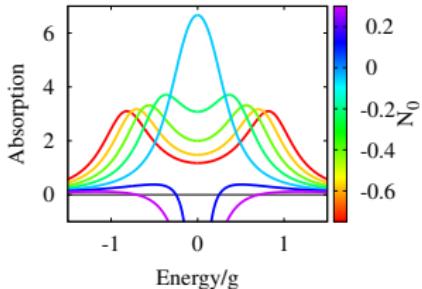
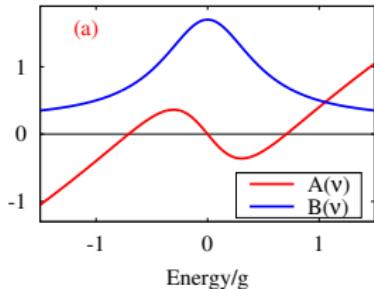
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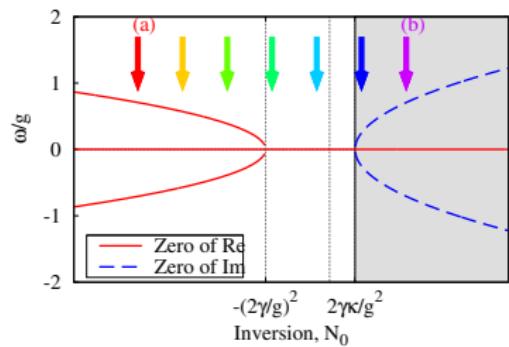


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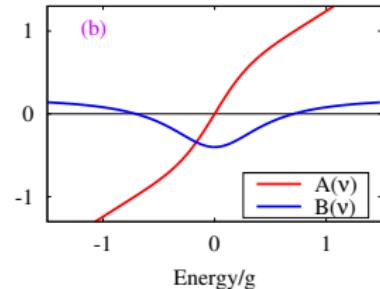
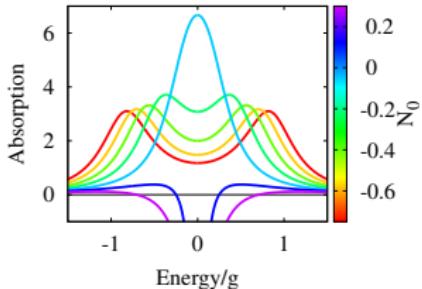
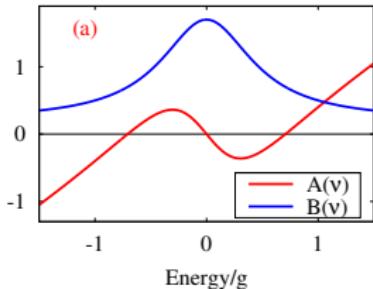


Laser:

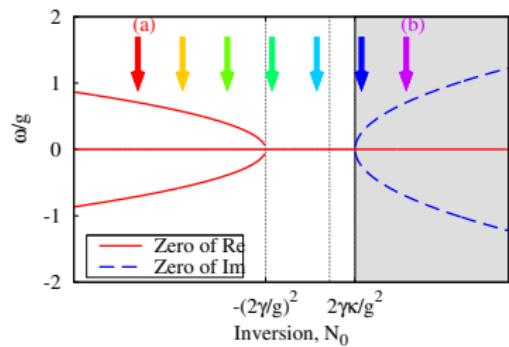


Poles of Retarded Green's function and gain

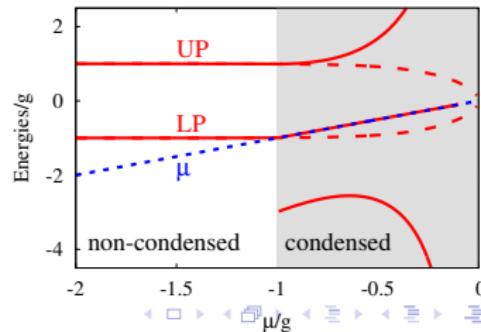
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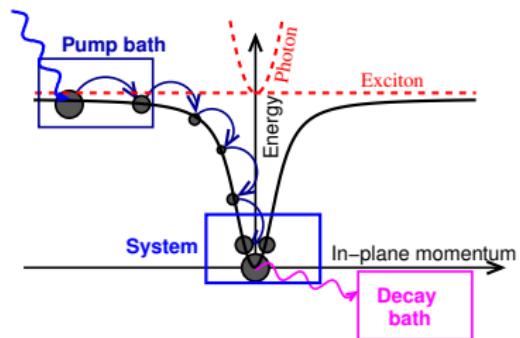
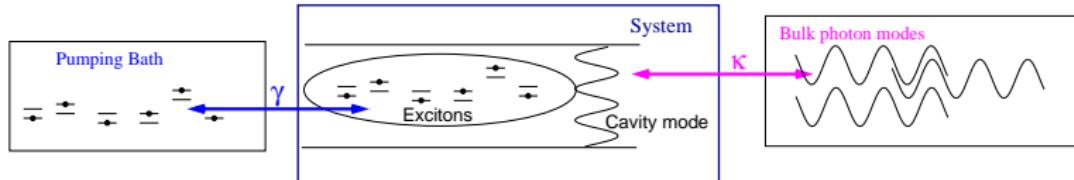
Laser:



Equilibrium:



Non-equilibrium description: baths

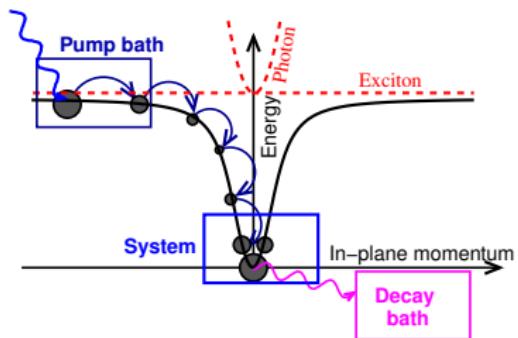
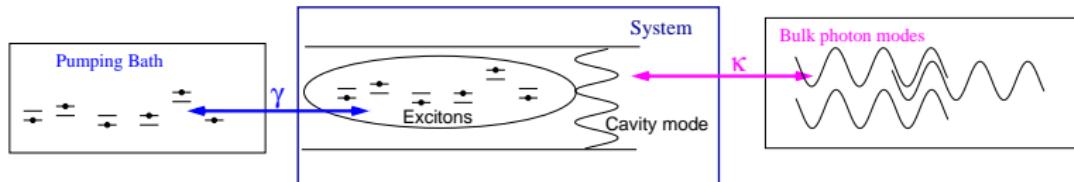


$$H = H_{\text{sys}} + H_{\text{sys,bath}} + H_{\text{bath}}$$

→ Decay bath: Empty ($\mu \rightarrow -\infty$)

→ Pump bath: Thermal μ_p, T_p

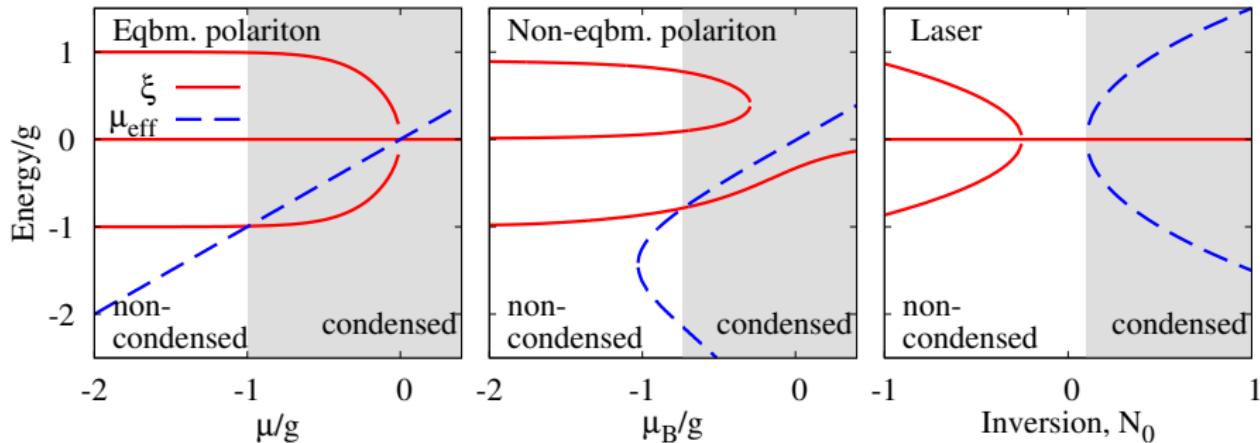
Non-equilibrium description: baths



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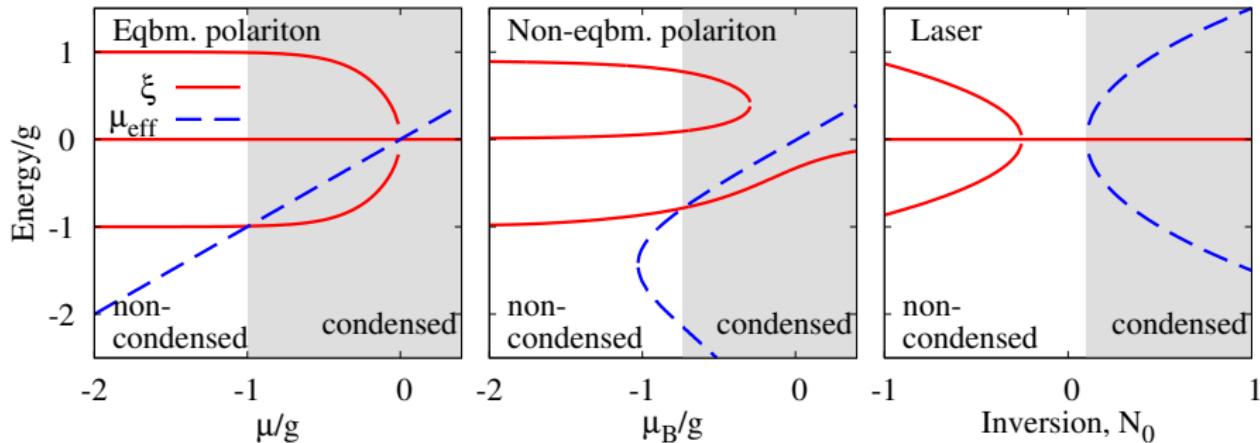
Strong coupling and lasing — low temperature phenomenon



- inversionless
- allows strong coupling
- requires low $T \rightarrow$ condensation
- Related weak-coupling inversionless lasing

[Szymanska *et al.* PRL '06; Keeling *et al.* book chapter 1010.3338]

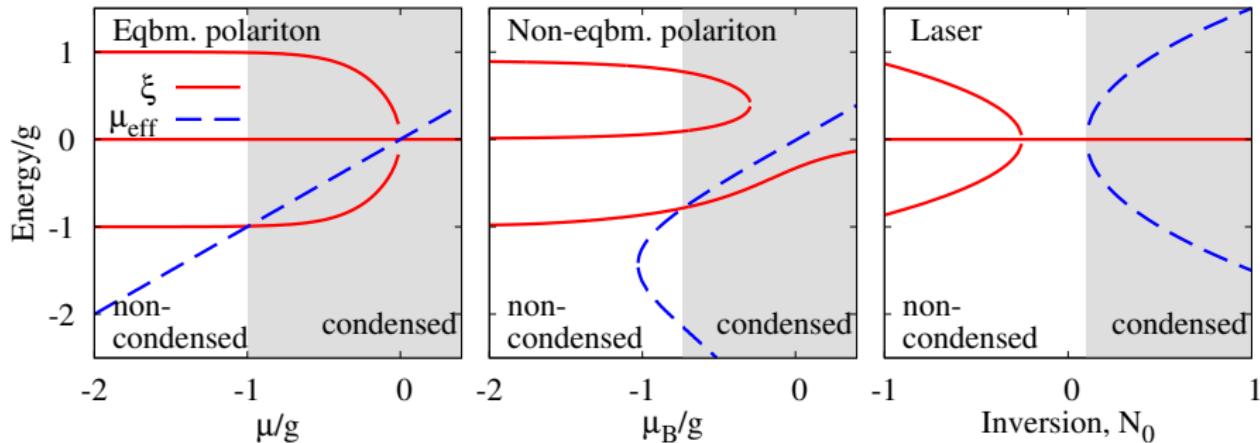
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Organic polaritons: photon-exciton-phonon coupling

1 Introduction

- Polariton condensation
- Condensation, superradiance, lasing
 - Condensation vs superradiance transition
 - Non-equilibrium condensation vs lasing

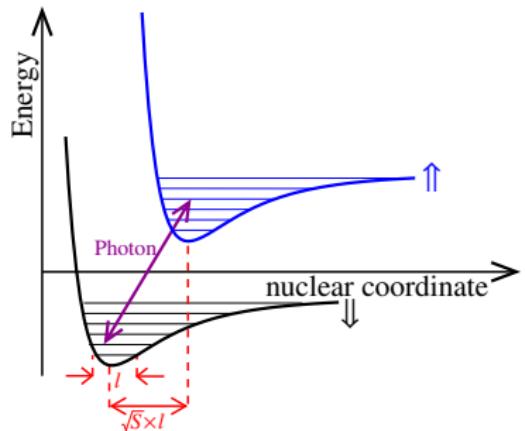
2 Room temperature condensates: Organic polaritons

- Dicke phase diagram with phonons
- Condensation of phonon replicas?
- Ultra-strong phonon coupling?

3 Room temperature condensates: Photons

- Lasing model and thermalisation
- Critical properties

Dicke Holstein Model

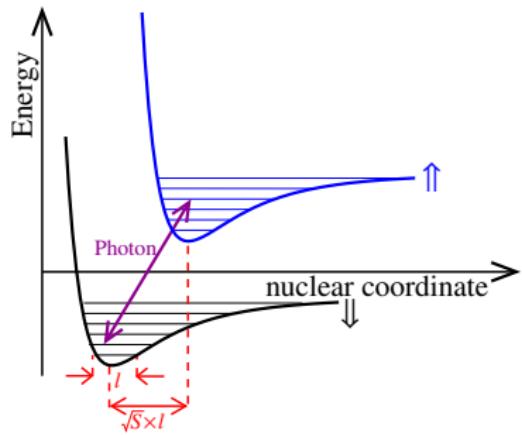


$$H = \omega \psi^\dagger \psi + \sum_{\alpha} \left[\epsilon S_{\alpha}^z + g \left(\psi S_{\alpha}^+ + \psi^\dagger S_{\alpha}^- \right) \right. \\ \left. + \Omega \left\{ b_{\alpha}^\dagger b_{\alpha} + \sqrt{S} \left(b_{\alpha}^\dagger + b_{\alpha} \right) S_{\alpha}^z \right\} \right]$$

- Photon frequency Ω
- Huang-Rhys parameter S — phonon coupling

- Phase diagram with $S \neq 0$
 - 2LS energy $\epsilon - n\Omega$
- Polariton spectrum, phonon replicas
- Ultra-strong phonon coupling?

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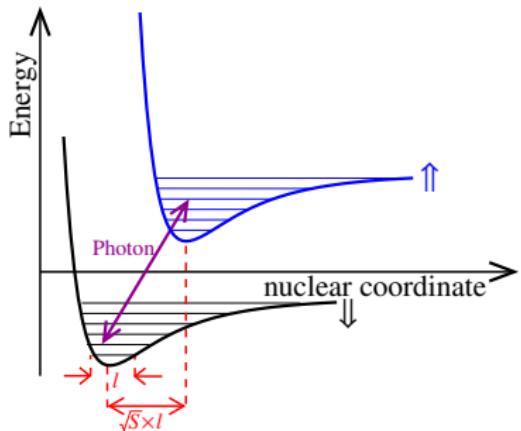


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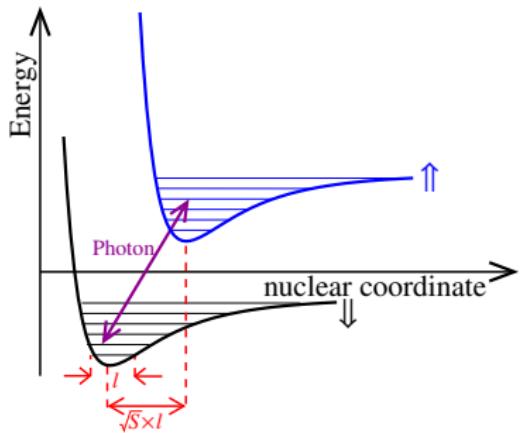
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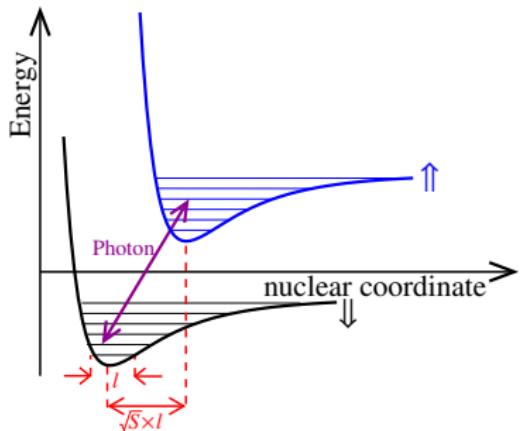
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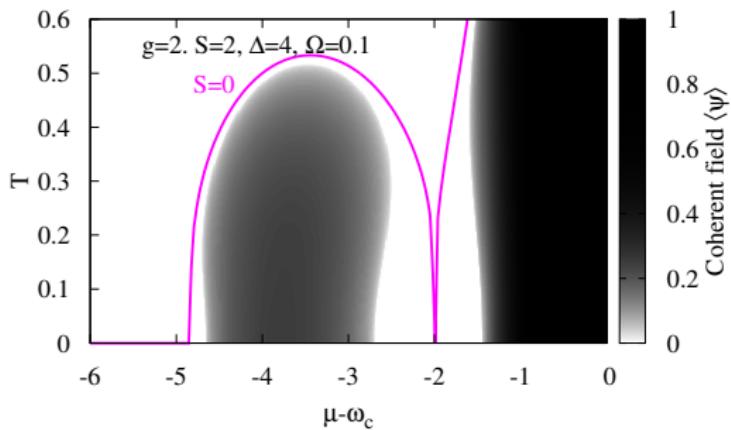
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- Superradiance condensation - reduces overlap
- Reentrant behaviour - min p at $T \sim 0.2$

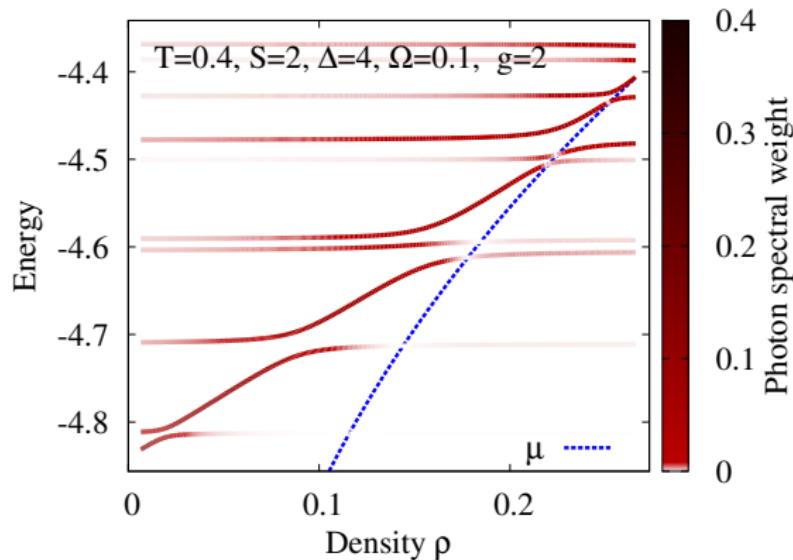
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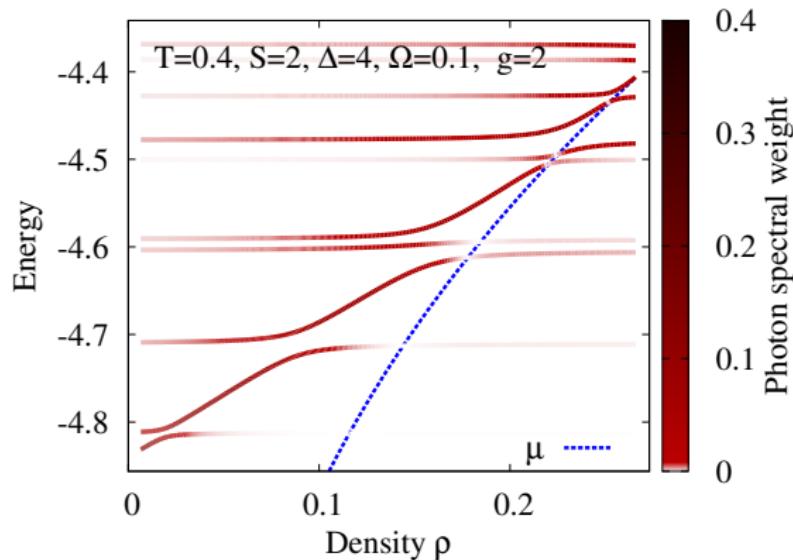
- S suppresses condensation — reduces overlap
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Polariton spectrum: photon weight



- Saturating 2LS: $g_{\text{eff}}^2 \sim g^2(1 - 2\rho)$

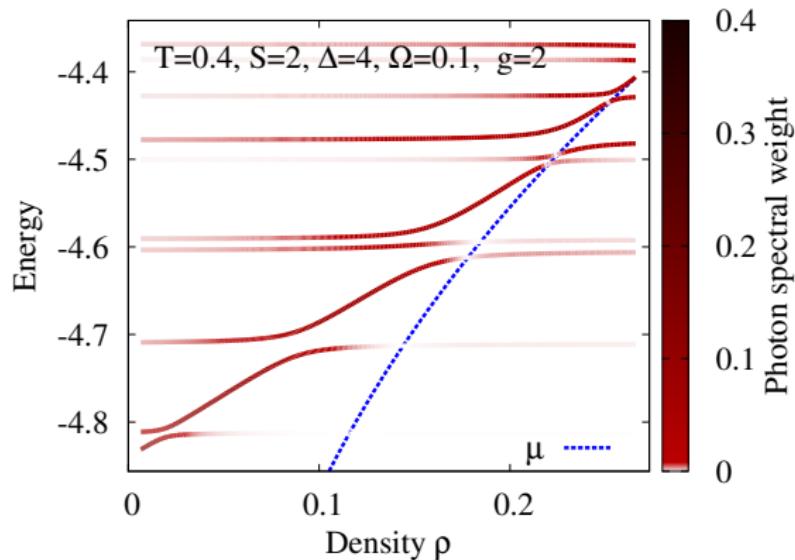
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- What is nature of polariton mode?

[Cwik *et al.* arXiv:1303.3702]

Polariton spectrum: photon weight



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- What is nature of polariton mode?
- $\mathcal{D}(t) = -i\langle \psi^\dagger(t)\psi(0) \rangle$, $\mathcal{D}(\omega) = \sum_n \frac{Z_n}{\omega - \omega_n}$

[Cwik *et al.* arXiv:1303.3702]

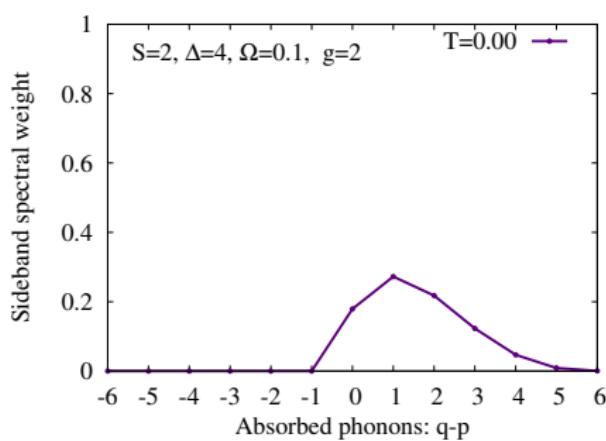
Polariton spectrum: what condensed

- Repeat weight for n -phonon channel
 - Eigenvector that is macroscopically occupied
 - Optimal $T \sim 20$

[Cwik *et al.* arXiv:1303.3702]

Polariton spectrum: what condensed

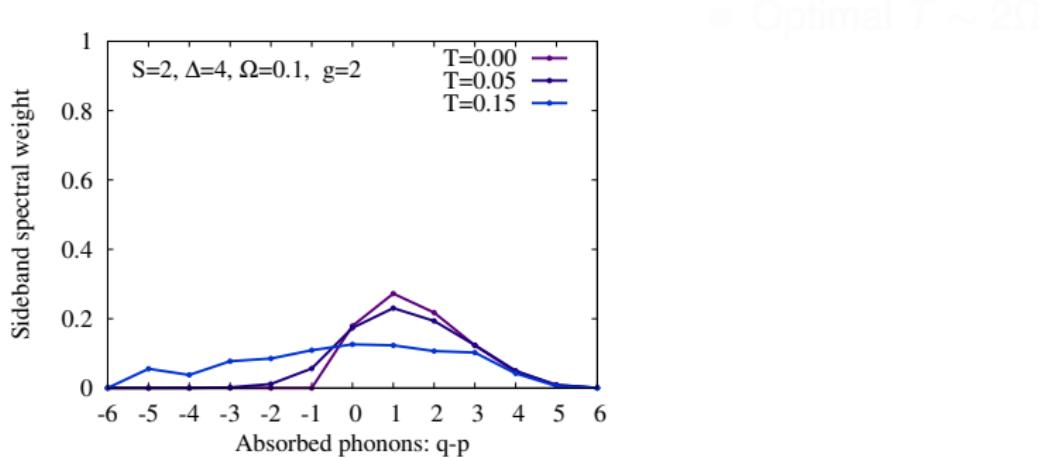
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[Cwik *et al.* arXiv:1303.3702]

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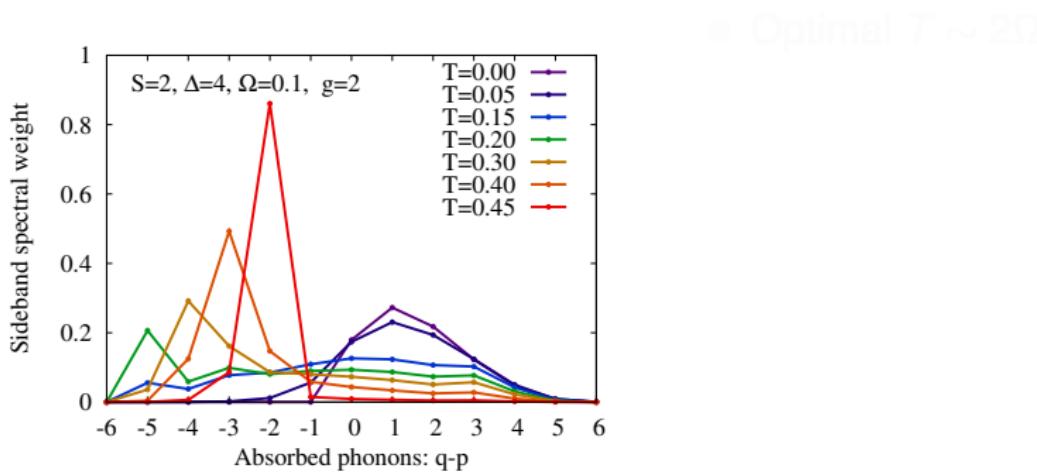
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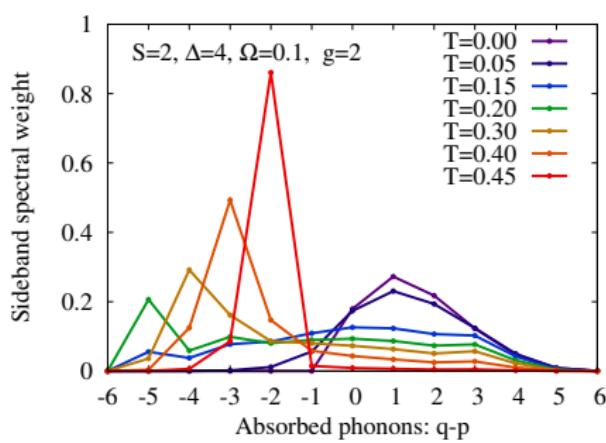
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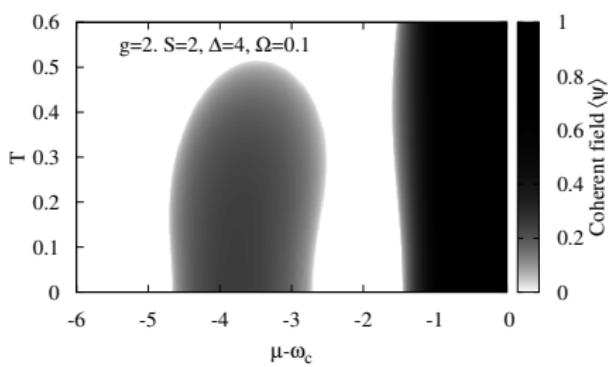
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- Repeat weight for n -phonon channel
- Eigenvector that is macroscopically occupied



- Optimal $T \sim 2\Omega$



[Cwik *et al.* arXiv:1303.3702]

Organic polaritons

1 Introduction

- Polariton condensation
- Condensation, superradiance, lasing
 - Condensation vs superradiance transition
 - Non-equilibrium condensation vs lasing

2 Room temperature condensates: Organic polaritons

- Dicke phase diagram with phonons
- Condensation of phonon replicas?
- Ultra-strong phonon coupling?

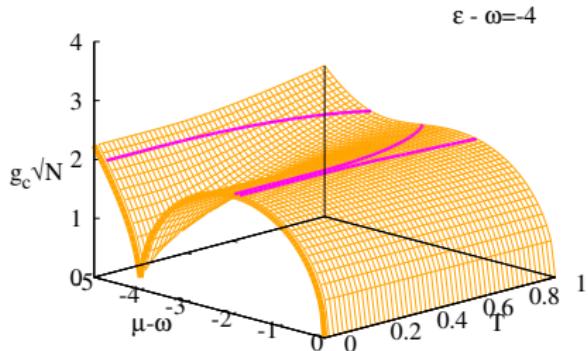
3 Room temperature condensates: Photons

- Lasing model and thermalisation
- Critical properties

Critical coupling with increasing S

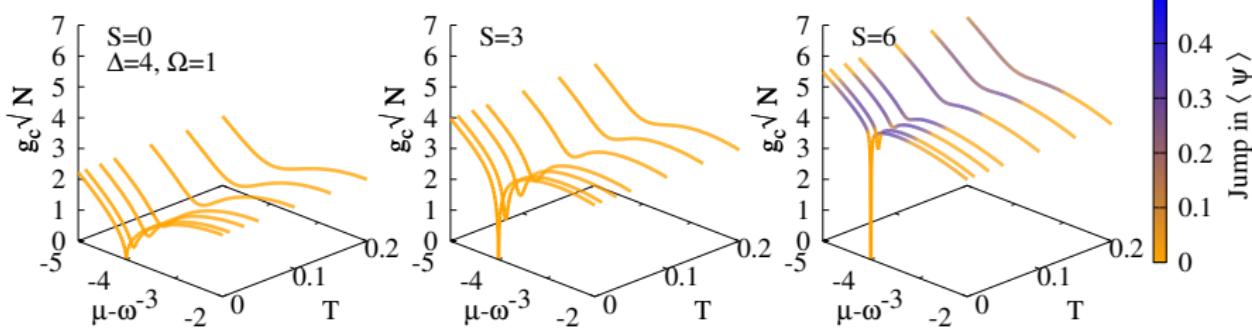
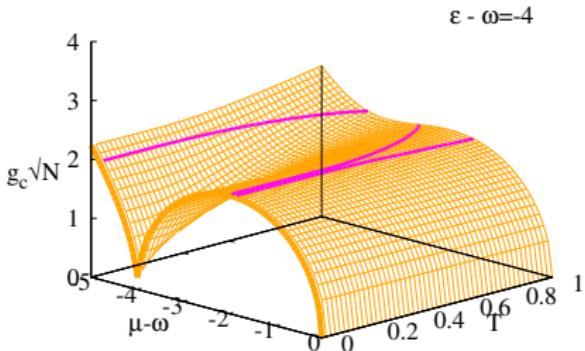
- Re-orient phase diagram
- g vs μ, T

\rightarrow $\text{collapse} \rightarrow \text{Jump of } \langle \hat{n} \rangle$



Critical coupling with increasing S

- Re-orient phase diagram
- g vs μ, T
- Colors \rightarrow Jump of $\langle \psi \rangle$



Explanation: Polaron formation

- Unitary transform

$$H_\alpha \rightarrow \tilde{H}_\alpha = e^{K_\alpha} H_\alpha e^{-K_\alpha} \quad K = \sqrt{S} S_\alpha^z (b_\alpha^\dagger - b_\alpha)$$

- Coupling moves to S^z

$$\tilde{H}_\alpha = \text{const.} + \epsilon S_\alpha^x + g b_\alpha^\dagger b_\alpha + g [g s t_z e^{iS(b_\alpha^\dagger - b_\alpha)} + \text{H.c.}]$$

- Optimal phonon displacements, $\sim \sqrt{S}$

- Reduced $g_{eff} \sim g \times \exp(-S/2)$

- For $\eta \neq 0$, competition

$$\text{Variational MFT } |\phi\rangle_\alpha \sim \exp(-\eta K_\alpha - \langle b_\alpha^\dagger \rangle) |0, S\rangle_\alpha$$

Explanation: Polaron formation

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- Optimal phonon displacements, $\sim \sqrt{S}$

- Reduced $g_{\text{eff}} \sim g \times \exp(-S/2)$

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Variational MFT $|\phi\rangle_v \sim \exp(-\eta K_\alpha - \langle b_\alpha^\dagger \rangle) |0, S\rangle$

Explanation: Polaron formation

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- Optimal phonon displacements, $\sim \sqrt{S}$
- Reduced $g_{\text{eff}} \sim g \times \exp(-S/2)$

QUESTION

Variational MFT $|\phi\rangle_c \sim \exp(-\eta K_c - \langle b_\alpha^\dagger \rangle) |0, S\rangle_c$

Explanation: Polaron formation

- Unitary transform

$$H_\alpha \rightarrow \tilde{H}_\alpha = e^{K_\alpha} H_\alpha e^{-K_\alpha} \quad K = \sqrt{S} S_\alpha^z (b_\alpha^\dagger - b_\alpha)$$

- Coupling moves to S^\pm

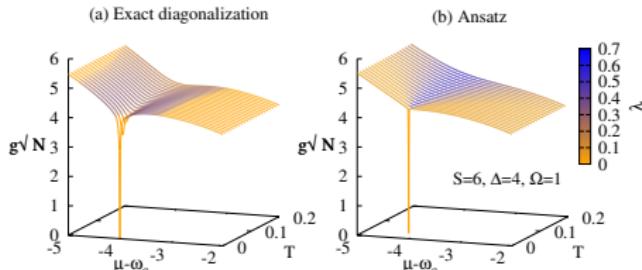
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- For $\psi \neq 0$, competition

Variational MFT $|\psi\rangle_\alpha \sim \exp(-\eta K_\alpha - \zeta b_\alpha^\dagger) |0, \mathbf{S}\rangle_\alpha$

Collective polaron formation

- Compares well at $S \gg 1$
- Coherent bosonic state



- Feedback: Large/small $\beta g\omega \rightarrow \lambda = (\beta)$
- Variational free energy

$$F = (\omega_c - \mu)\lambda^2 + N \left\{ \Omega \left[\xi^2 - S^2 \frac{(2-\eta)}{\eta} \right] - T \ln \left[2 \cosh \left(\frac{\xi}{T} \right) \right] \right\}$$

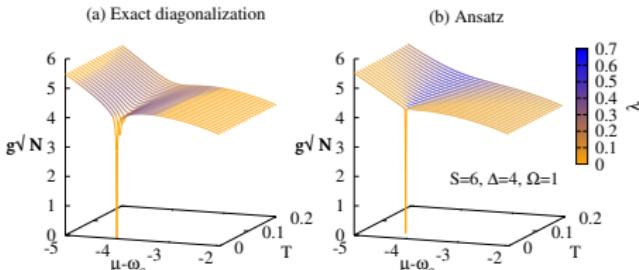
Effective 2LS energy in field:

$$\xi^2 = \left(\frac{\xi - \mu}{2} + \alpha \sqrt{S} (1 - \eta) \right)^2 + g^2 \lambda^2 e^{-S\beta}$$

[Cwik *et al.* arXiv:1303.3702]

Collective polaron formation

- Compares well at $S \gg 1$
- Coherent bosonic state
- Feedback: Large/small g_{eff} $\leftrightarrow \lambda = \langle \psi \rangle$



Effective 2LS energy:

$$E = (\omega_c - \mu)^2 + N \left[\Omega \left[\xi^2 - \frac{\Omega^2(2-\eta)}{2} \right] - T \ln \left[2 \cosh \left(\frac{\Omega}{T} \right) \right] \right]$$

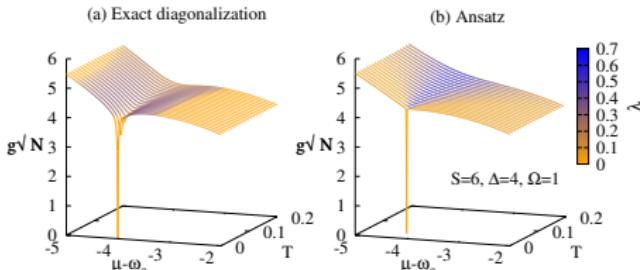
Effective 2LS energy in field:

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[Cwik *et al.* arXiv:1303.3702]

Collective polaron formation

- Compares well at $S \gg 1$
- Coherent bosonic state



- Feedback: Large/small g_{eff} $\leftrightarrow \lambda = \langle \psi \rangle$
- Variational free energy

$$F = (\omega_c - \mu)\lambda^2 + N \left\{ \Omega \left[\zeta^2 - S \frac{\eta(2-\eta)}{4} \right] - T \ln \left[2 \cosh \left(\frac{\xi}{T} \right) \right] \right\}$$

Effective 2LS energy in field:

$$\xi^2 = \left(\frac{\epsilon - \mu}{2} + \Omega \sqrt{S} (1 - \eta) \zeta \right)^2 + g^2 \lambda^2 e^{-S\eta^2}$$

[Cwik *et al.* arXiv:1303.3702]

Polariton and photon Condensation

1 Introduction

- Polariton condensation
- Condensation, superradiance, lasing
 - Condensation vs superradiance transition
 - Non-equilibrium condensation vs lasing

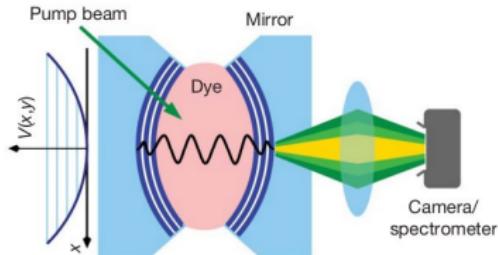
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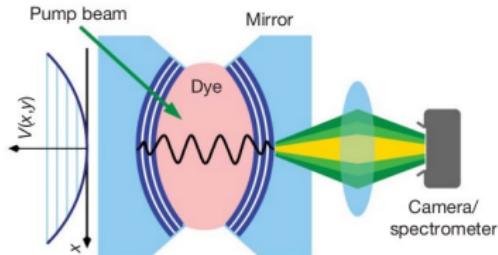
Photon BEC experiments



- Dye filled microcavity

[Klaers et al, Nature, 2010]

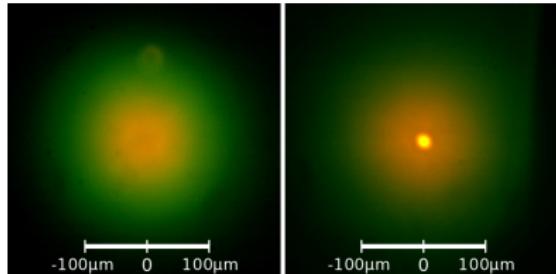
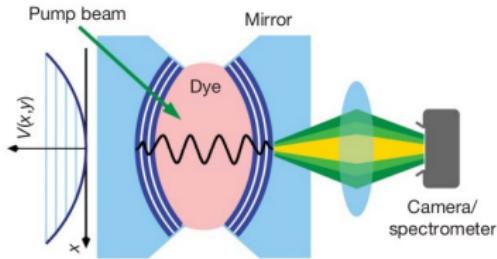
Photon BEC experiments



- Dye filled microcavity
- No strong coupling

[Klaers et al, Nature, 2010]

Photon BEC experiments



- Dye filled microcavity
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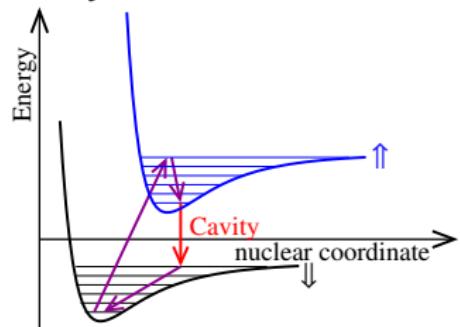
Relation to dye laser

- No electronic inversion
- No strong coupling
 - No single cavity mode
 - Condensate mode is not maximum gain
 - Gain/Absorption in balance
 - Thermalised many-mode system

Relation to dye laser

- No electronic inversion
- No strong coupling

4 Level Dye Laser

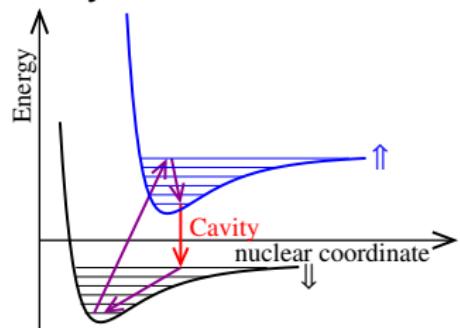


- No single cavity mode
- Condensate mode is not maximum gain
- Gain/Absorption in balance
- Thermalised many-mode system

Relation to dye laser

- No electronic inversion
- No strong coupling

4 Level Dye Laser



But:

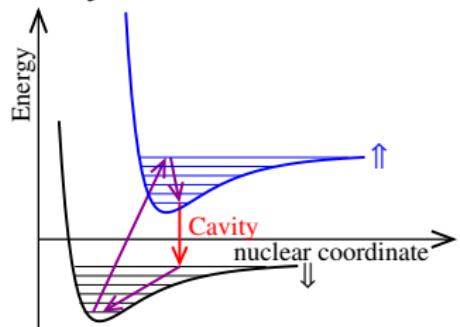
- No single cavity mode
 - ▶ Condensate mode is not maximum gain
 - ▶ Gain/Absorption in balance

• Thermalised many-mode system

Relation to dye laser

- No electronic inversion
- No strong coupling

4 Level Dye Laser



But:

- No single cavity mode
 - ▶ Condensate mode is not maximum gain
 - ▶ Gain/Absorption in balance
- Thermalised many-mode system

Modelling

$$H_{\text{sys}} = \sum_m \omega_m \psi_m^\dagger \psi_m + \sum_\alpha [\epsilon S_\alpha^z + g (\psi_m S_\alpha^+ + \text{H.c.})]$$

]

- 2D harmonic cavity

$$\omega_m = \omega_{\text{cutoff}} + m\omega_{H.O.}$$

$$\text{Degeneracies } g_m = m + 1$$

↳ [http://arxiv.org/abs/1303.0003](#)

Modelling

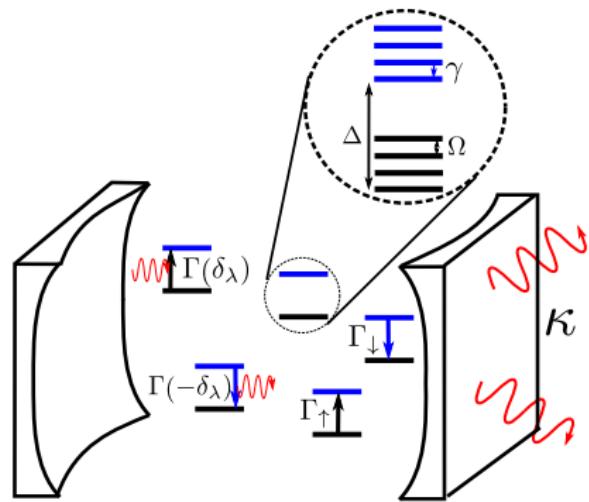
$$H_{\text{sys}} = \sum_m \omega_m \psi_m^\dagger \psi_m + \sum_\alpha [\epsilon S_\alpha^z + g (\psi_m S_\alpha^+ + \text{H.c.}) + \Omega \{ b_\alpha^\dagger b_\alpha + 2\sqrt{\epsilon} S_\alpha^z (b_\alpha^\dagger + b_\alpha) \}]$$

- 2D harmonic cavity

$$\omega_m = \omega_{\text{cutoff}} + m\omega_{\text{H.O.}}$$

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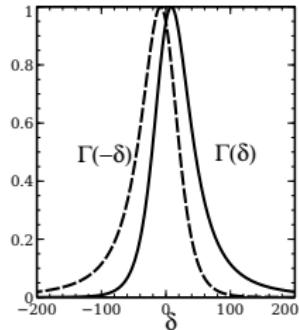
- Local vibrational mode



Modelling

Rate equation

$$\dot{\rho} = -i[H_0, \rho] - \sum_m \frac{\kappa}{2} \mathcal{L}[\psi_m] - \sum_{\alpha} \left[\frac{\Gamma_{\uparrow}}{2} \mathcal{L}[S_{\alpha}^{+}] + \frac{\Gamma_{\downarrow}}{2} \mathcal{L}[S_{\alpha}^{-}] \right] \\ - \sum_{m,\alpha} \left[\frac{\Gamma(\delta_m = \omega_m - \epsilon)}{2} \mathcal{L}[S_{\alpha}^{+} \psi_m] + \frac{\Gamma(-\delta_m = \epsilon - \omega_m)}{2} \mathcal{L}[S_{\alpha}^{-} \psi_m^{\dagger}] \right]$$



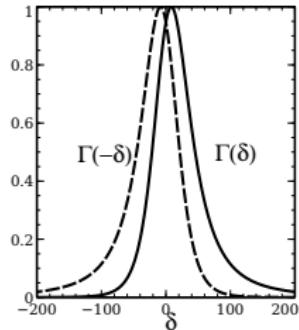
$$\Rightarrow \Gamma(-\delta) \approx \Gamma(-\delta) e^{i\delta}$$
$$\Rightarrow \Gamma \rightarrow 0 \text{ at large } \delta$$

[Marthaler et al PRL '11, Kirton & JK arXiv:1303.3459]

Modelling

Rate equation

$$\dot{\rho} = -i[H_0, \rho] - \sum_m \frac{\kappa}{2} \mathcal{L}[\psi_m] - \sum_{\alpha} \left[\frac{\Gamma_{\uparrow}}{2} \mathcal{L}[S_{\alpha}^{+}] + \frac{\Gamma_{\downarrow}}{2} \mathcal{L}[S_{\alpha}^{-}] \right] \\ - \sum_{m,\alpha} \left[\frac{\Gamma(\delta_m = \omega_m - \epsilon)}{2} \mathcal{L}[S_{\alpha}^{+} \psi_m] + \frac{\Gamma(-\delta_m = \epsilon - \omega_m)}{2} \mathcal{L}[S_{\alpha}^{-} \psi_m^{\dagger}] \right]$$



- $\Gamma(+\delta) \simeq \Gamma(-\delta) e^{\beta\delta}$
- $\Gamma \rightarrow 0$ at large δ

[Marthaler et al PRL '11, Kirton & JK arXiv:1303.3459]

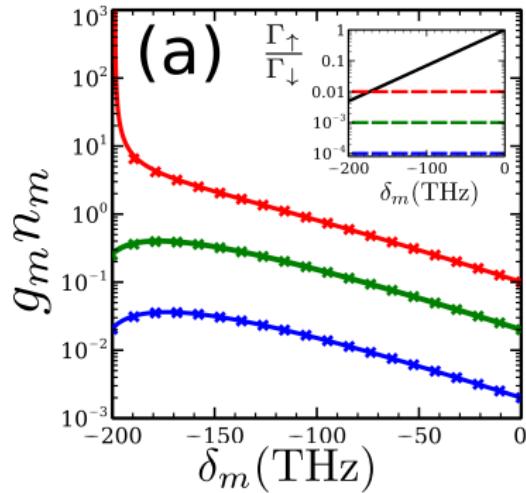
Distribution $g_m n_m$

- Rate equation — include spontaneous emission
- Bose-Einstein distribution without losses

[Kirton & JK arXiv:1303.3459]

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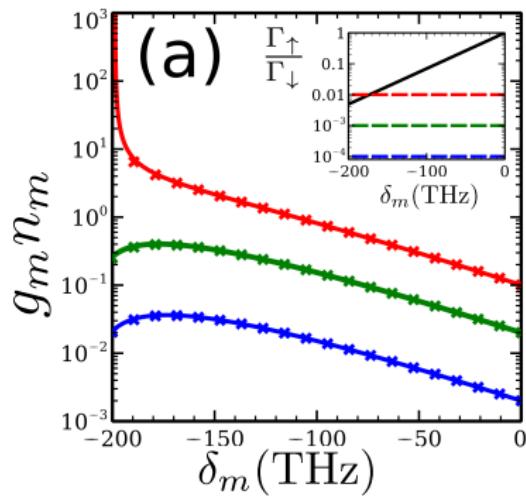


Low loss: Thermal

[Kirton & JK arXiv:1303.3459]

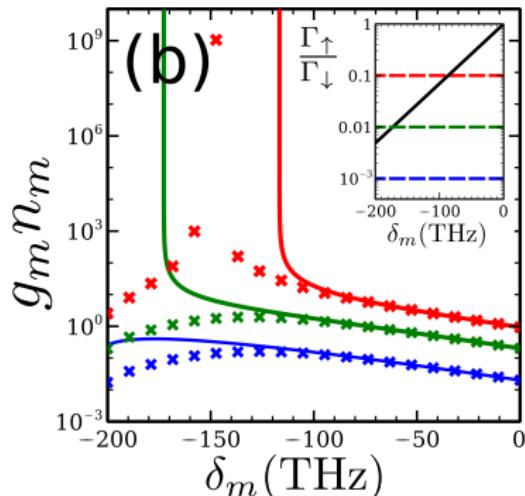
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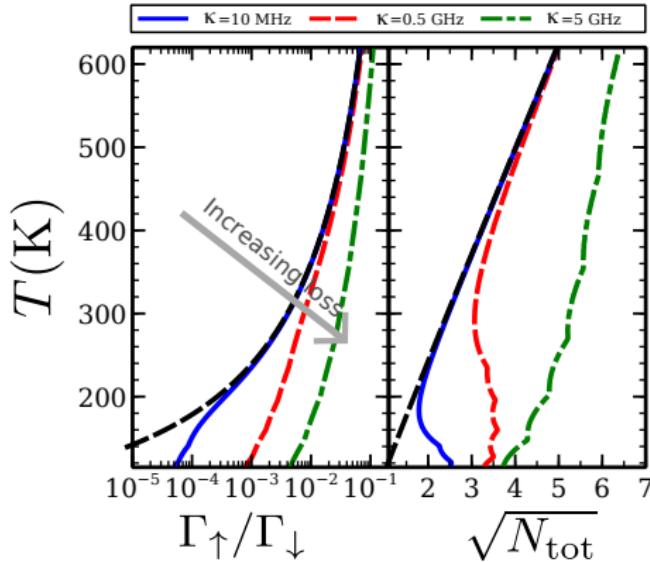
Low loss: Thermal

[Kirton & JK arXiv:1303.3459]



High loss \rightarrow Laser

Threshold condition



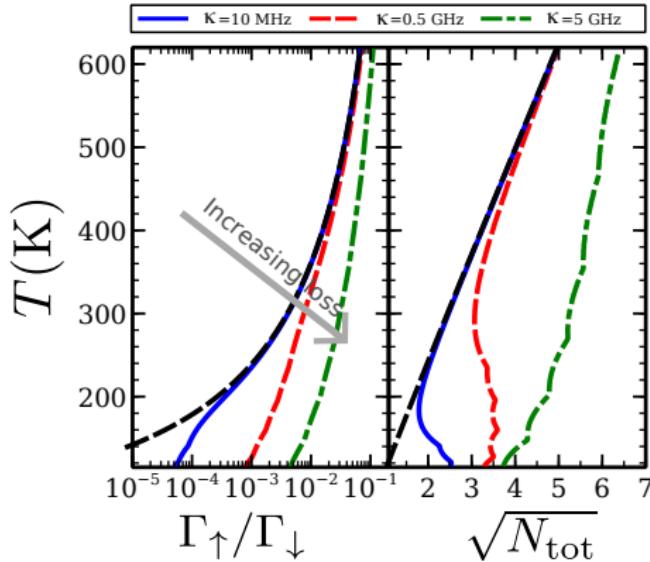
Compare threshold:

- Pump rate (Laser)
- Critical density (condensate)

- Thermal at low γ /high temperature
- High loss, γ competes with $\Gamma(\pm\delta_0)$
- Low temperature, $\Gamma(\pm\delta_0)$ shrinks

[Kirton & JK arXiv:1303.3459]

Threshold condition



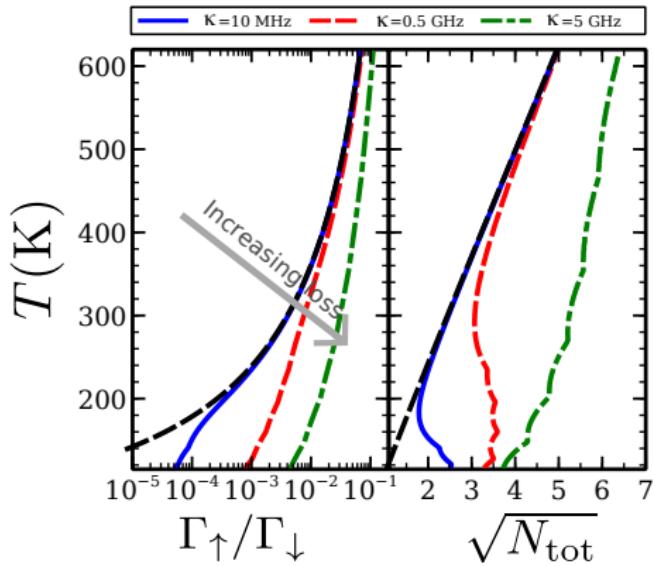
Compare threshold:

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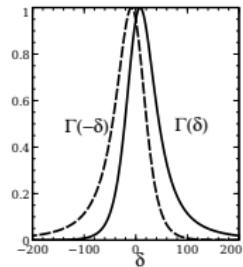
Threshold condition



Compare threshold:

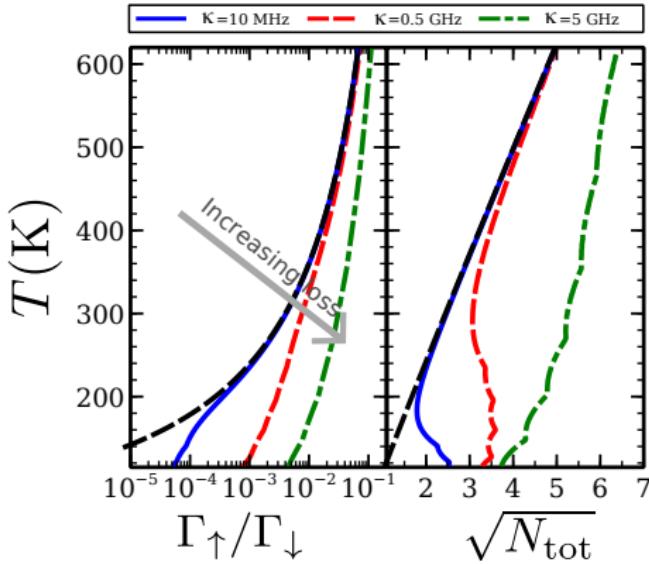
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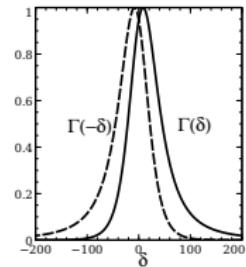


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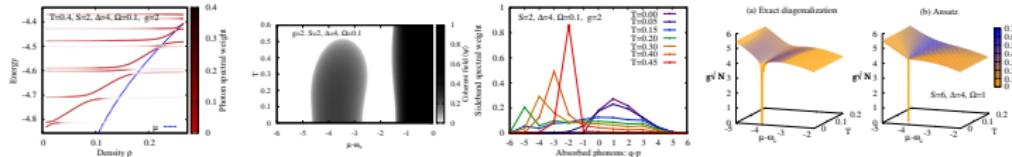
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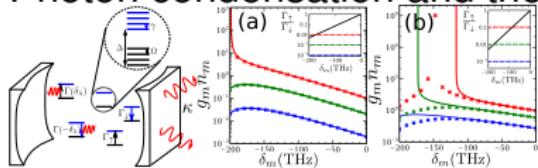


Summary

- Reentrance, phonon assisted transition, 1st order at $S \gg 1$



- Photon condensation and thermalisation



Many body quantum optics and correlated states of light

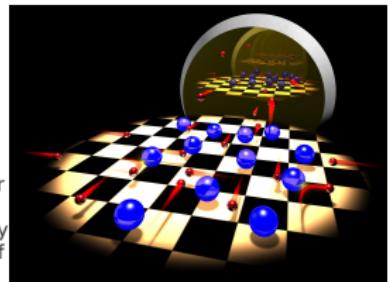
9:00 am on Monday 28 October 2013 – 5:00 pm on Tuesday 29 October 2013

at:[The Royal Society at Chicheley Hall, home of the Kavli Royal Society International Centre, Buckinghamshire](#)

Theo Murphy international scientific meeting organised by Dr Jonathan Keeling, Professor Steven Girvin, Dr Michael Hartmann and Professor Peter Littlewood FRS.

List of speakers and chairs

Professor Iacopo Carusotto, Professor Andrew Cleland, Professor Hui Deng, Professor Tilman Esslinger, Professor Rosario Fazio, Professor Ed Hinds, Professor Andrew Houck, Professor Ataç İmamoğlu, Professor Jens Koch, Professor Misha Lukin, Professor Martin Plenio, Professor Arno Rauschenbeutel, Professor Timothy Spiller, Professor Jacob Taylor, Professor Hakan Tureci, Professor Andreas Wallraff



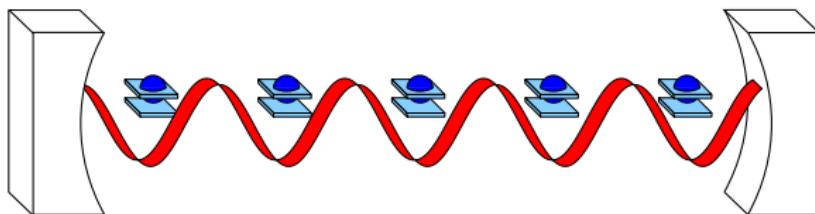
Attending this event

This is a residential conference which allows for increased discussion and networking. It is free to attend, however participants need to cover their accommodation and catering costs if required. Places are limited and therefore pre-registration is essential.

Extra slides

- 4 No go theorem
- 5 Retarded Green's function for laser
- 6 Organic properties
- 7 Anticrossing vs ρ

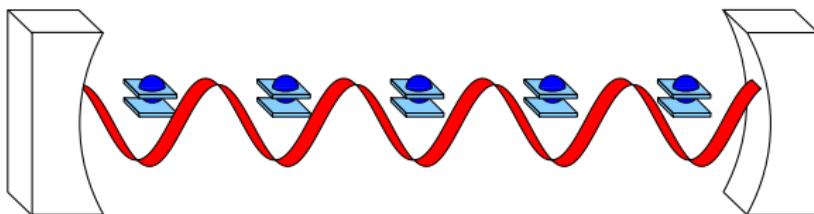
No go theorem and transition



Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

[Rzazewski *et al* PRL '75]

No go theorem and transition



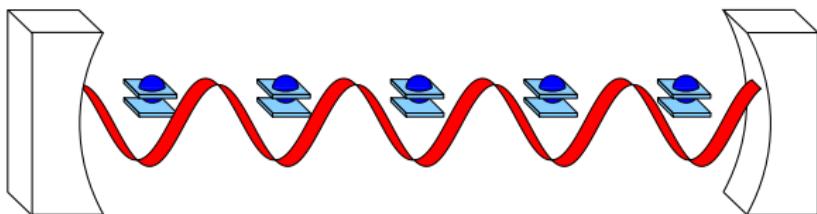
Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

No go theorem: Minimal coupling $(p - eA)^2/2m$

$$-\sum_i \frac{e}{m} A \cdot p_i \Leftrightarrow g(\psi^\dagger S^- + \psi S^+), \quad \sum_i \frac{A^2}{2m} \Leftrightarrow N\zeta(\psi + \psi^\dagger)^2$$

[Rzazewski *et al* PRL '75]

No go theorem and transition



Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

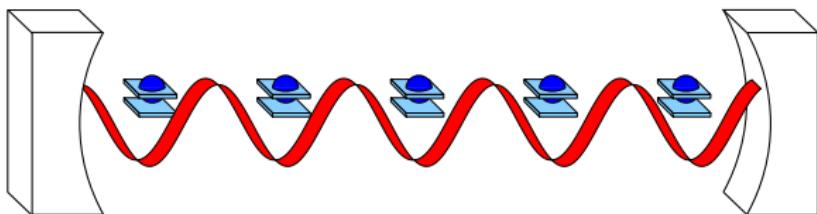
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For large N , $\omega \rightarrow \omega + 2N\zeta$. (RWA)

[Rzazewski *et al* PRL '75]

No go theorem and transition



Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

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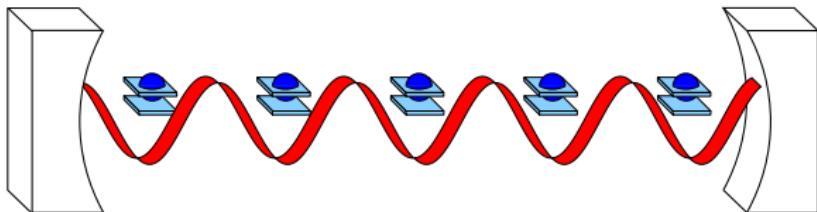
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For large N , $\omega \rightarrow \omega + 2N\zeta$. (RWA)

Need $Ng^2 > \epsilon(\omega + 2N\zeta)$.

[Rzazewski *et al* PRL '75]

No go theorem and transition



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No go theorem: Minimal coupling $(p - eA)^2/2m$

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For large N , $\omega \rightarrow \omega + 2N\zeta$. (RWA)

Need $Ng^2 > \epsilon(\omega + 2N\zeta)$.

But Thomas-Reiche-Kuhn sum rule states: $g^2/\epsilon < 2\zeta$. **No transition**
[Rzazewski *et al* PRL '75]

Dicke phase transition: ways out

Problem: $g^2/\omega_0 < 2\zeta$ for intrinsic parameters. **Solutions:**

- Interpretation:
Pseudoelectric transition in D-r gauge.
[Knap et al., Vukics & Demokos PRA 2012]
- Circuit QED [Nataf and Ciuti, Nat. Comm. '10; Viehmann et al. PRL '11]
- Grand canonical ensemble:
 - If $H \rightarrow H - \mu(57 + 97\beta)$, need only:
 $g^2 N > (\omega - \mu)(\omega_0 - \mu)$
 - Incoherent pumping — polariton condensation.
- Dissociate g, ω_0 ,
e.g. Raman scheme: $\omega_R \ll \omega$.
(Dimer et al. PRA '07; Baumann et al. Nature '10; Also, Black et al. PRL '03)

Dicke phase transition: ways out

Problem: $g^2/\omega_0 < 2\zeta$ for intrinsic parameters. **Solutions:**

- **Interpretation**

Ferroelectric transition in $\mathbf{D} \cdot \mathbf{r}$ gauge.

[JK JPCM '07, Vukics & Domokos PRA 2012]

→ Dicke phase transition and condensate. [H. Heimann et al. PRL 111]

• Grand canonical ensemble:

→ If $H \rightarrow H - \mu(57 + 97\beta)$, need only:

$$g^2 N > (\omega - \mu)(\omega_0 - \mu)$$

Incoherent pumping — polariton condensation.

• Dissociate g, ω_0 ,

e.g. Raman scheme: $\omega_R \ll \omega$.

(Dimer et al. PRA 107; Baumann et al. Nature

10. Also, Black et al. PRL 103).

Dicke phase transition: ways out

Problem: $g^2/\omega_0 < 2\zeta$ for intrinsic parameters. **Solutions:**

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Ferroelectric transition in $\mathbf{D} \cdot \mathbf{r}$ gauge.

[JK JPCM '07, Vukics & Domokos PRA 2012]

- Circuit QED [Nataf and Ciuti, Nat. Comm. '10; Viehmann *et al.* PRL '11]

- Grand canonical ensemble:

- If $H \rightarrow H - \mu(57 + 67\beta)$, need only:

$$g^2 N > (\omega - \mu)(\omega_0 - \mu)$$

- Incoherent pumping — polariton condensation.

- Dissociate g, ω_g ,

- e.g. Raman scheme: $\omega_g \ll \omega$.

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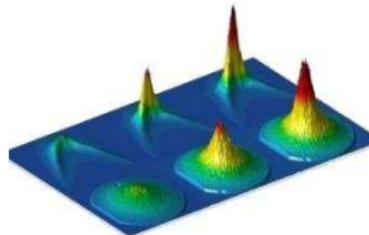
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• Dicke phase transition

$\omega_0/\hbar\omega_{\text{cav}} \ll \omega_0$

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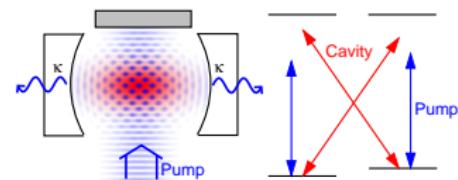
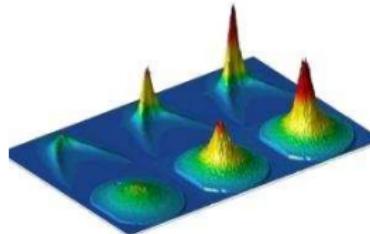
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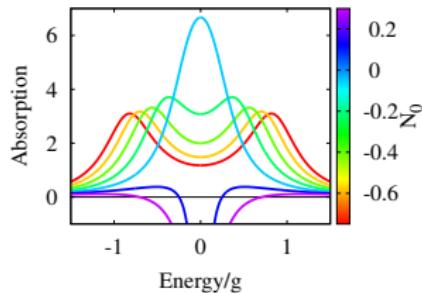
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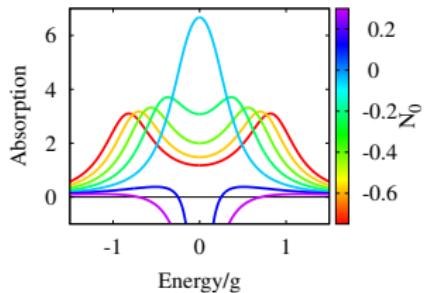


Maxwell-Bloch Equations: Retarded Green's function



- Introduce $D^R(\omega)$:
Response to perturbation
- Absorption = $-2\Im[D^R(\omega)]$

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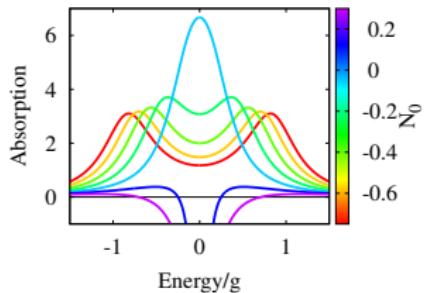
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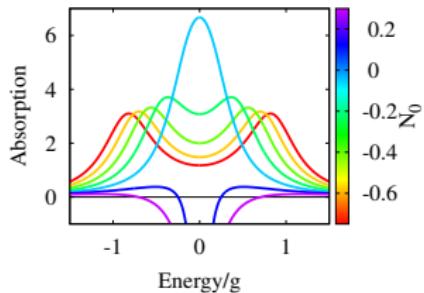
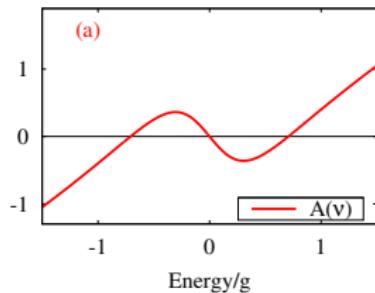
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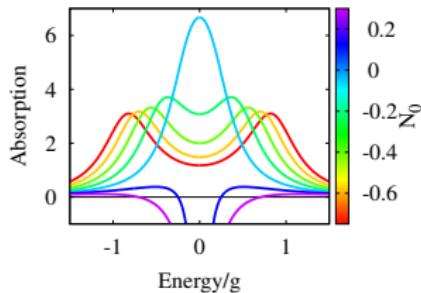
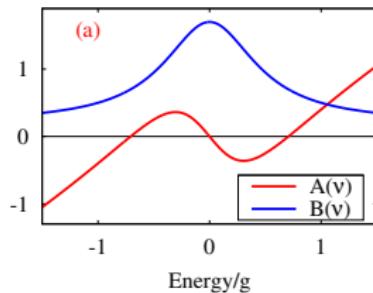
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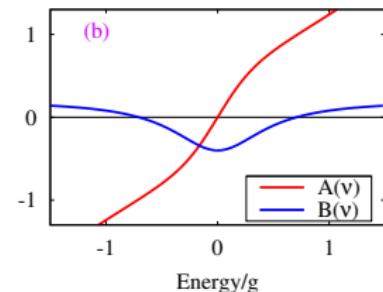
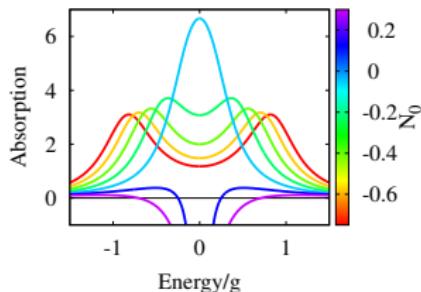
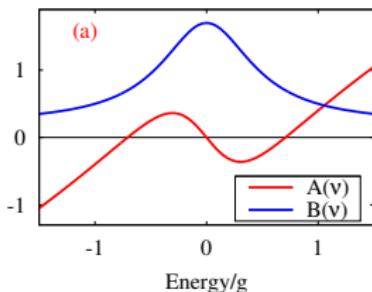
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Organic materials in microcavities

- State of art:

- ▶ Strong coupling:
 - ★ J aggregates [Bulovic *et al.*]
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- ▶ Threshold: Anthracene

[Kena Cohen and Forrest, Nat. Photon 2010]

- Differences

- ▶ Stronger coupling

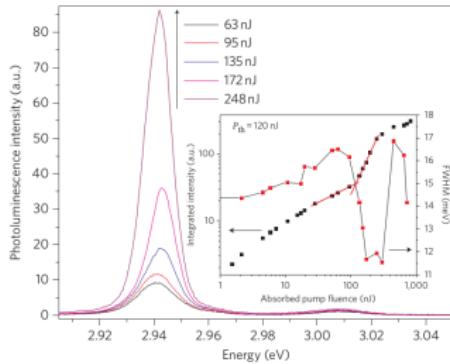
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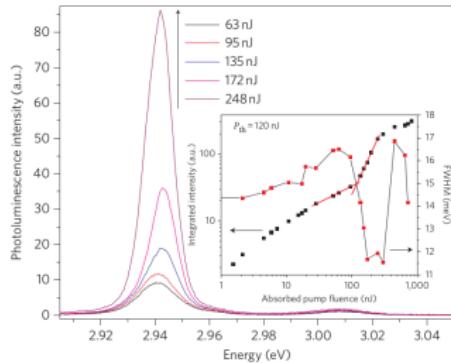
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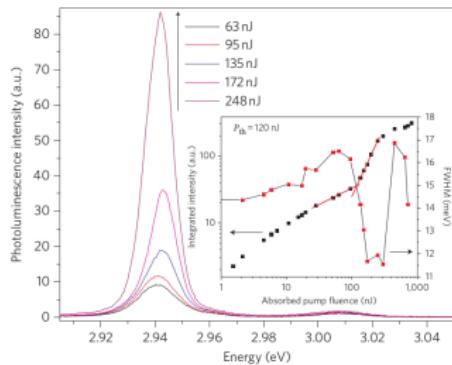
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→ Singlet-Triplet conversion — dark states

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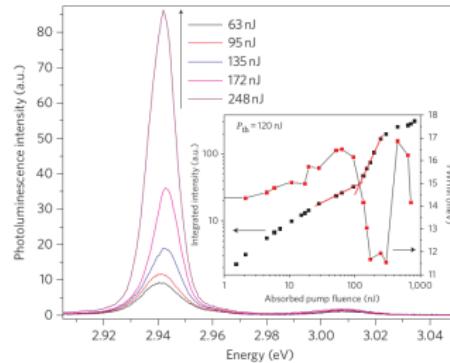
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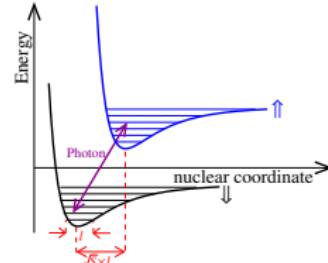
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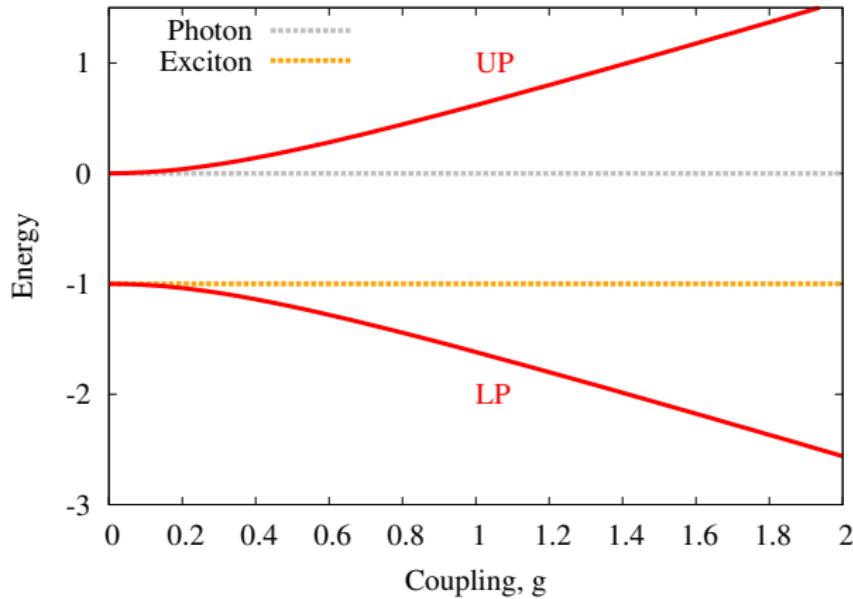
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 - ▶ Vibrational sidebands

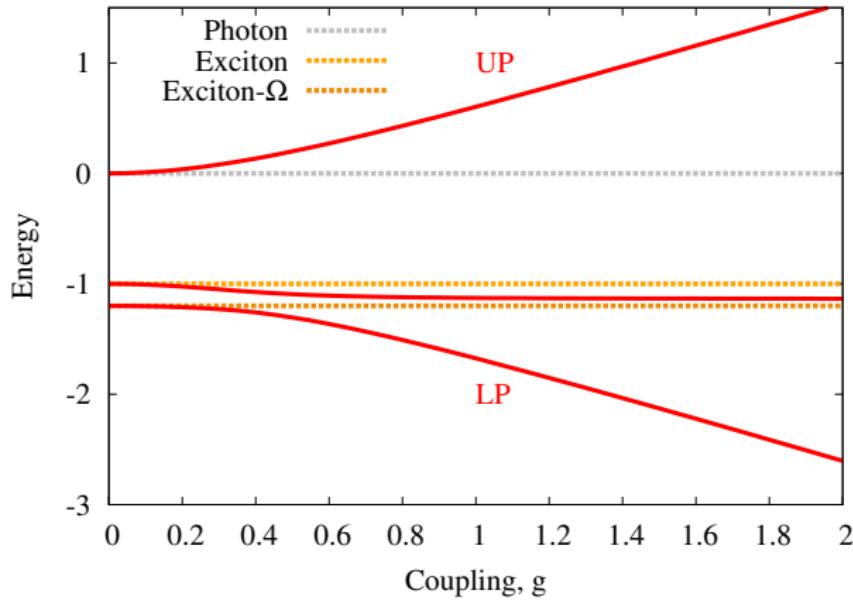


Polariton spectrum — coupled oscillators

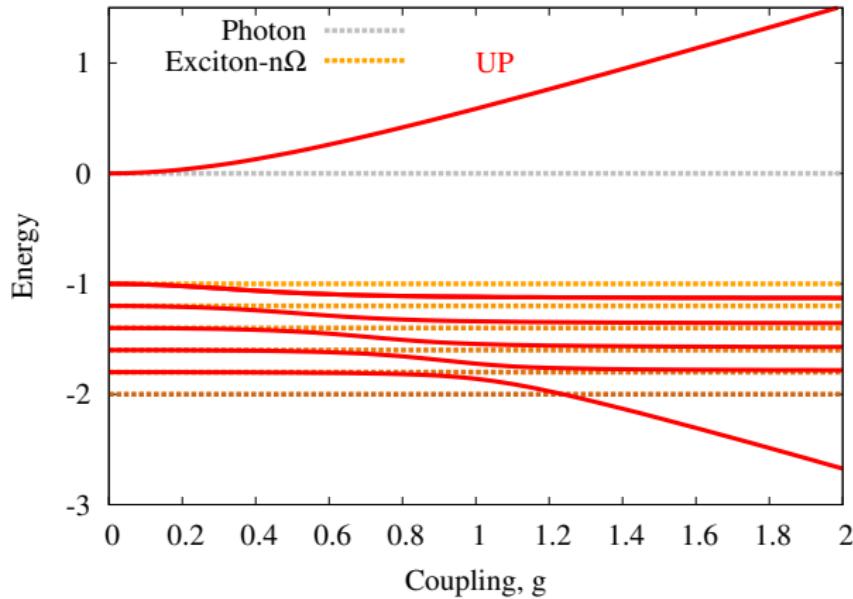
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