

Polariton and photon condensates in organic materials

Jonathan Keeling



University of
St Andrews

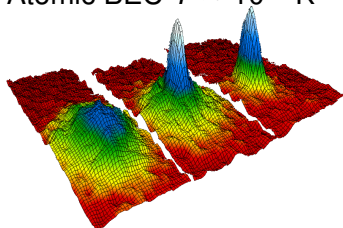
600
YEARS



Oxford, June 2013

Coherent states of matter and light

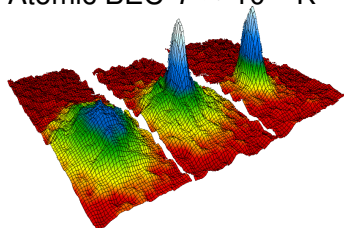
Atomic BEC $T \sim 10^{-7}\text{K}$



[Anderson *et al.* Science '95]

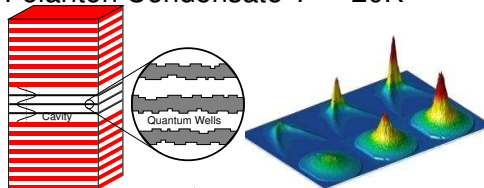
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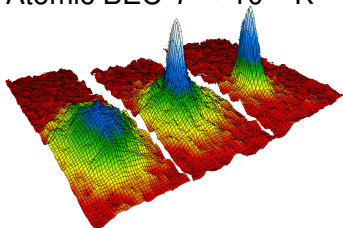
Polariton Condensate $T \sim 20\text{K}$



[Kasprzak *et al.* Nature, '06]

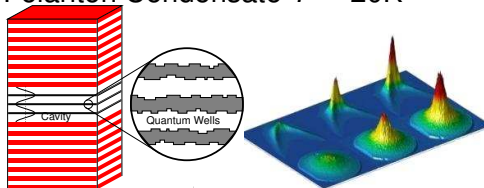
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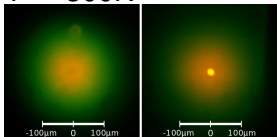
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Photon Condensate

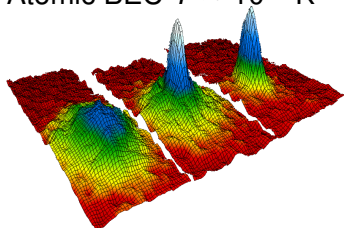
$T \sim 300\text{K}$



[Klaers *et al.* Nature, '10]

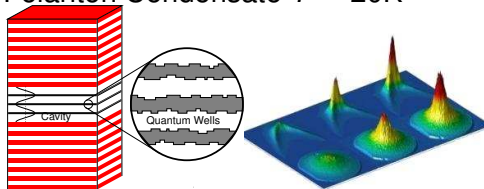
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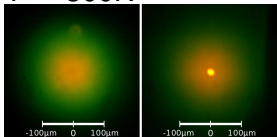
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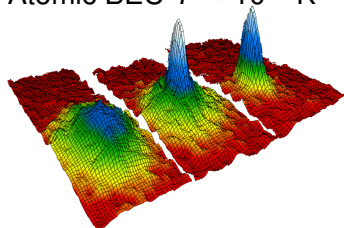
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Laser
 $T \sim ?, < 0, \infty$



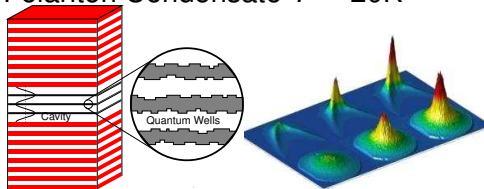
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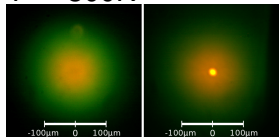
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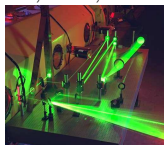
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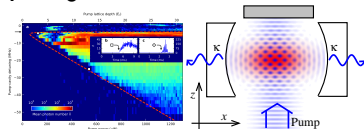


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Superradiance transition
 $T \sim 0$



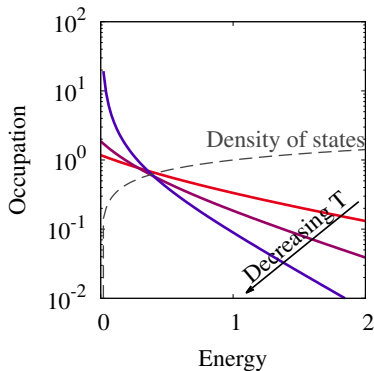
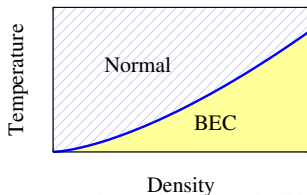
[Baumann *et al.* Nature, '10]

“Textbook” BEC

• Non-interacting viewpoint

▶ BE distribution: $\mu < \omega_0$

▶ $T_c = \frac{2\pi\hbar^2}{m} \left(\frac{n}{\xi_d} \right)^{2/d}$



• Interacting approach (WIDBG)

$$H = \sum_k \omega_k \psi_k^\dagger \psi_k + \frac{g}{2V} \sum_{k,k',q} \psi_{k+q}^\dagger \psi_{k-q}^\dagger \psi_k \psi_{k'}$$

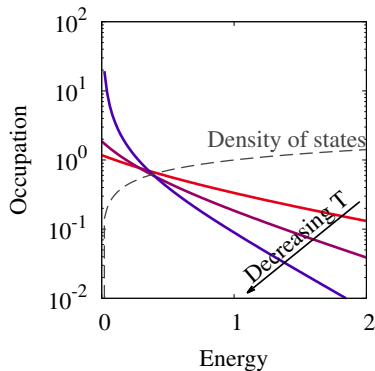
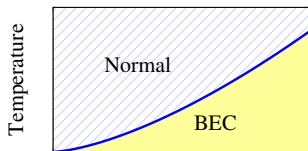
• Mean field: $\langle \psi \rangle^2 = (\mu - \omega_0)/V$

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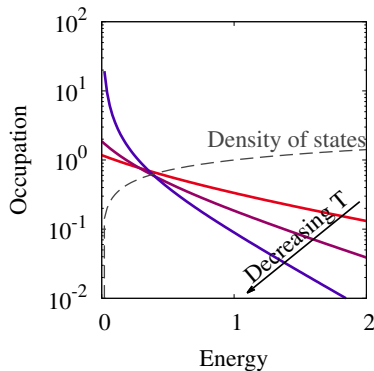
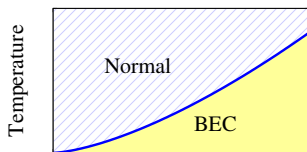
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Fluctuations deplete condensate vanishes at $T > T_c$

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“Textbook” Laser: Maxwell Bloch equations

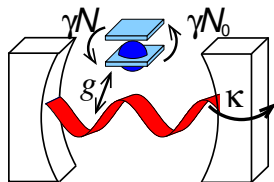
$$H = \omega_0 \psi^\dagger \psi + \sum_{\alpha} \epsilon_{\alpha} S_{\alpha}^Z + g_{\alpha, \mathbf{k}} \left(\psi S_{\alpha}^+ + \psi^\dagger S_{\alpha}^- \right)$$

$$\text{Maxwell-Bloch eqns: } P = -i \langle S^- \rangle, N = 2 \langle S^Z \rangle$$

$$\partial_t \psi = -i \omega_0 \psi - \kappa \psi + \sum_{\alpha} g_{\alpha} P_{\alpha}$$

$$\partial_t P_{\alpha} = -2i \epsilon_{\alpha} P_{\alpha} - 2\gamma P_{\alpha} + g_{\alpha} \psi N_{\alpha}$$

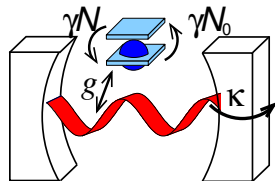
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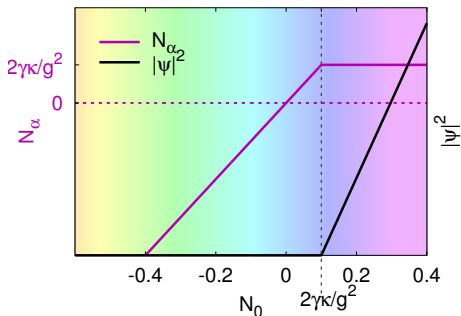
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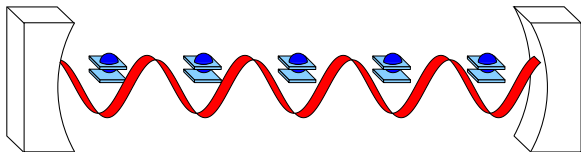
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$$|\psi|^2 > 0 \text{ if } N_0 g^2 > 2\gamma\kappa$$

“Textbook” Dicke-Hepp-Lieb superradiance



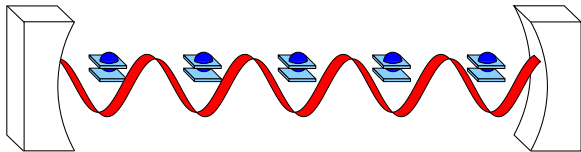
$$H = \omega \psi^\dagger \psi + \epsilon S^z + g (\psi^\dagger S^- + \psi S^+)$$

• Coherent state: $|\psi\rangle \rightarrow e^{\lambda \psi^\dagger + \eta S^+} |\Omega\rangle$

• Small g , min at $\lambda, \eta = 0$

[Hepp, Lieb, Ann. Phys. '73]

“Textbook” Dicke-Hepp-Lieb superradiance



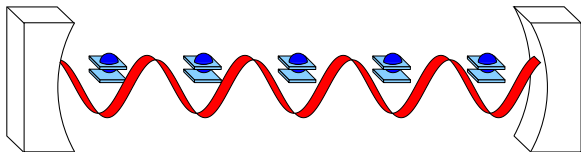
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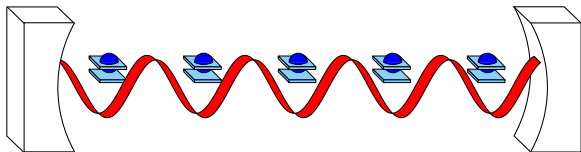
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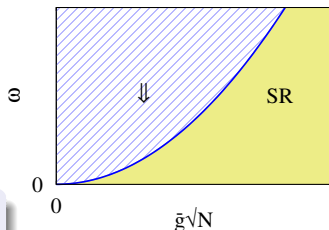
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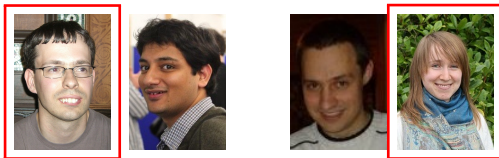
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Outline

- 1 Condensation, superradiance, lasing
- 2 Polariton condensation and Dicke model
 - Dicke model and condensation
 - Non-equilibrium condensation vs lasing
- 3 Room temperature condensates: Photons
 - Lasing model and thermalisation
 - Critical properties
- 4 Room temperature condensates: Organic polaritons
 - Dicke phase diagram with phonons
 - Condensation of phonon replicas?
 - (Ultra-strong phonon coupling?)
- 5 Conclusions

Acknowledgements

GROUP:

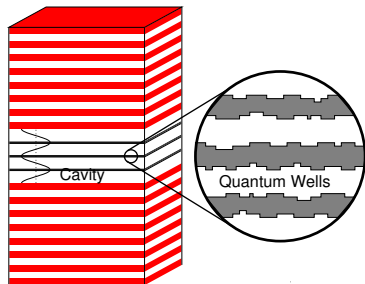


COLLABORATORS: Szymanska (Warwick), Reja (Cam.), Littlewood (ANL)

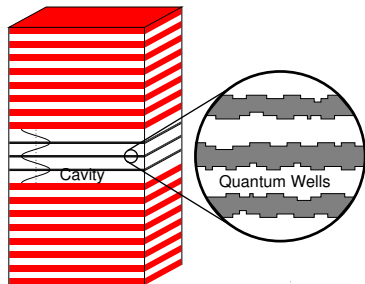
FUNDING:



Microcavity polaritons

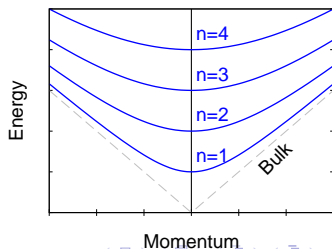


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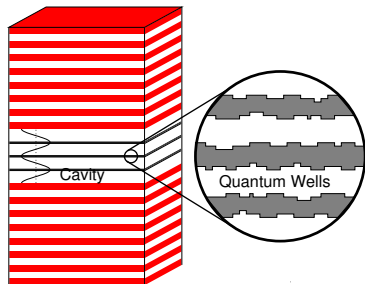


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2/2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$

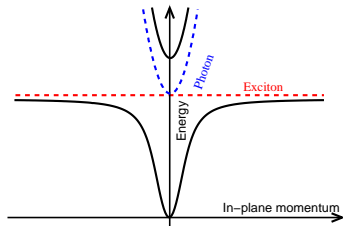


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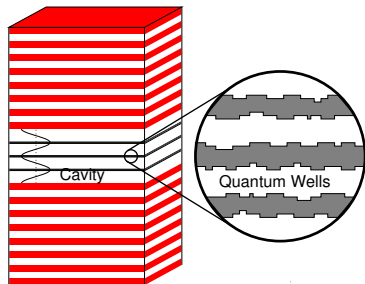


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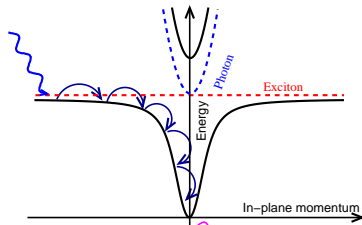


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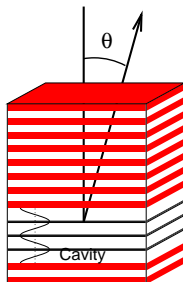
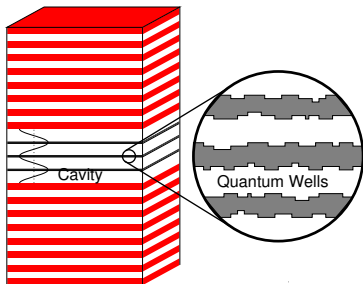


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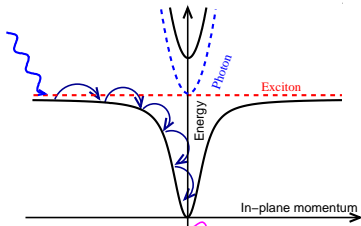


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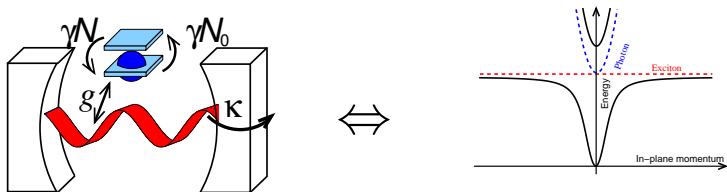
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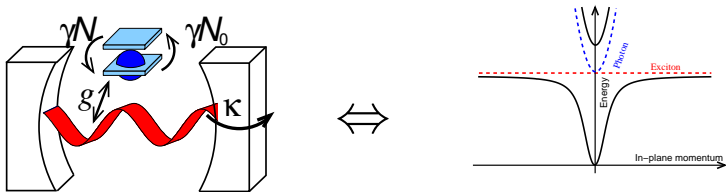
Lasing-condensation crossover model

- Use model that can show lasing and condensation:



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Dicke model:

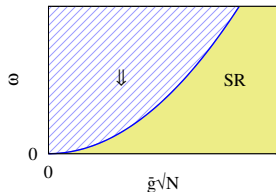
$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} [\epsilon S_{\alpha}^Z + g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} S_{\alpha}^{+} + \text{H.c.}]$$

Dicke-Hepp-Lieb superradiance and modes

$$H = \omega\psi^\dagger\psi + \epsilon S^z + g(\psi^\dagger S^- + \psi S^+)$$

Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

- Normal state, $S^z = -N/2 + B^\dagger B$
 $H = \omega\psi^\dagger\psi + \epsilon B^\dagger B + g\sqrt{N}(\psi^\dagger B + \psi B^\dagger)$
- Excitation cost E :
 $(E - \omega)(E - \epsilon) = g^2 N$
- Transition when $E = 0$



[Hepp, Lieb, Ann. Phys. '73]

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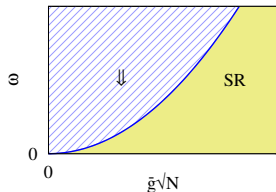
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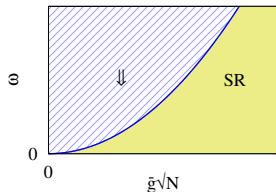
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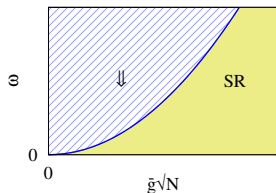
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Grand canonical ensemble

Grand canonical ensemble:

- If $H \rightarrow H - \mu(S^z + \psi^\dagger\psi)$, need only: $g^2 N > (\omega - \mu)|\epsilon - \mu|$

- Fix density / fix $\mu > 0$ — pumping

- Transition at:
 $g^2 N > (\omega - \mu)(\epsilon - \mu)$
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- Unstable if $\mu > \omega$
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[Eastham and Littlewood, PRB '01]

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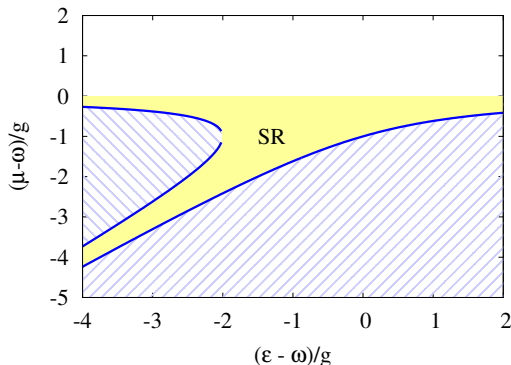
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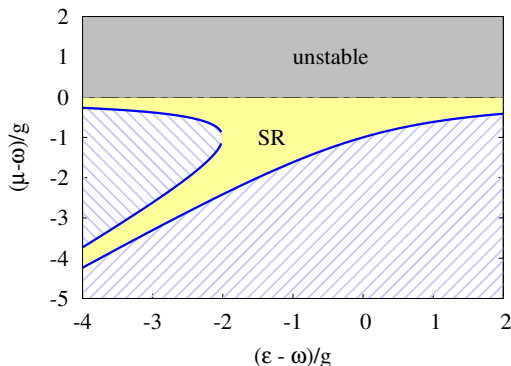
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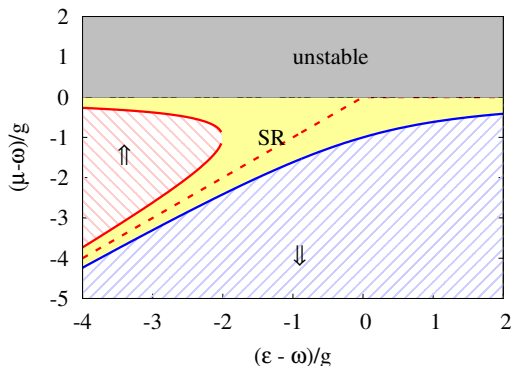
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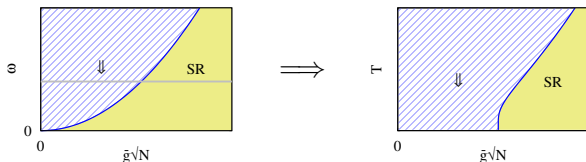
[Eastham and Littlewood, PRB '01]

- Transition at:
 $g^2 N > (\omega - \mu)(\epsilon - \mu)$
- μ hits lowest mode
- Unstable if $\mu > \omega$
- Inverted if $\mu > \epsilon$

Grand canonical Dicke, finite temperature

- Finite temperature:

$$Ng^2 \tanh(\beta\epsilon) > \omega\epsilon$$



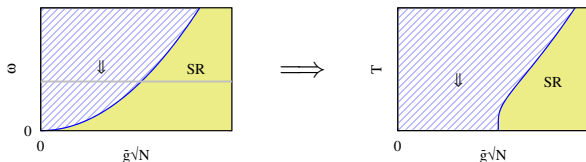
[Hepp, Lieb, Ann. Phys. '73]

With chemical potential $Ng^2 \tanh(\beta(\epsilon - \mu)) > (\omega - \mu)(\epsilon - \mu)$

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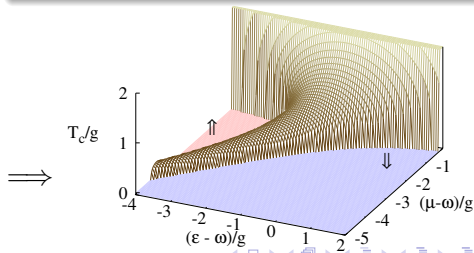
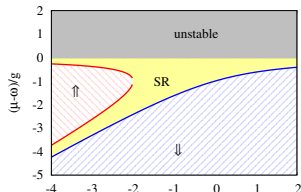
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- With chemical potential

$$Ng^2 \tanh(\beta(\epsilon - \mu)) > (\omega - \mu)(\epsilon - \mu)$$



Polariton model and equilibrium results

Localised excitons, propagating photons

$$H - \mu N = \sum_{\mathbf{k}} (\omega_{\mathbf{k}} - \mu) \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} (\epsilon_{\alpha} - \mu) S_{\alpha}^Z + g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} S_{\alpha}^{+} + \text{H.c.}$$

Self-consistent polarisation and field

$$(\omega - \mu) \psi = \sum_{\alpha} \frac{g_{\alpha}^2 \psi}{2E_{\alpha}} \tanh(\beta E_{\alpha}), \quad E_{\alpha}^2 = \left(\frac{\epsilon_{\alpha} - \mu}{2} \right)^2 + g_{\alpha}^2 |\psi|^2$$

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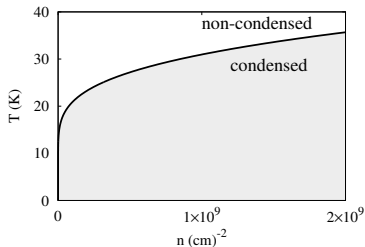
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Phase diagram:



Polariton model and equilibrium results

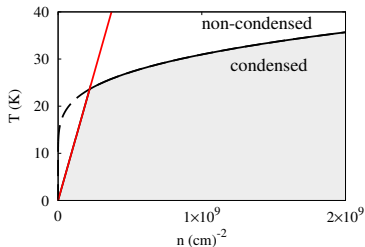
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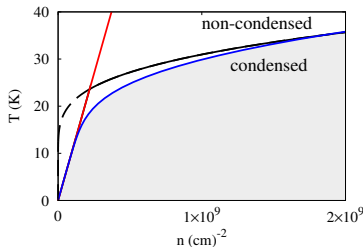
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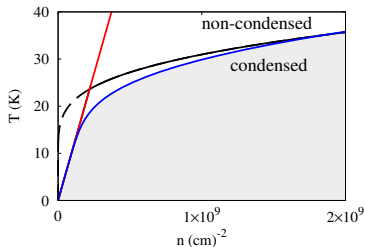
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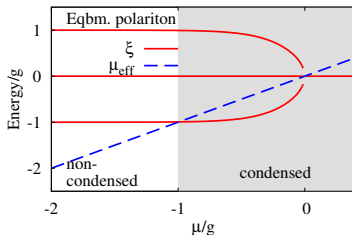
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Phase diagram:



Modes (at $k = 0$)



Non-equilibrium condensation vs lasing

1 Condensation, superradiance, lasing

2 Polariton condensation and Dicke model

- Dicke model and condensation
- **Non-equilibrium condensation vs lasing**

3 Room temperature condensates: Photons

- Lasing model and thermalisation
- Critical properties

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- Dicke phase diagram with phonons
- Condensation of phonon replicas?
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5 Conclusions

Simple Laser: Maxwell Bloch equations

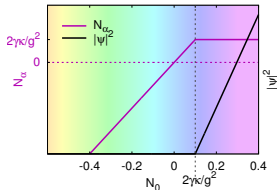
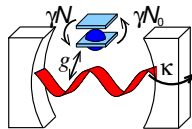
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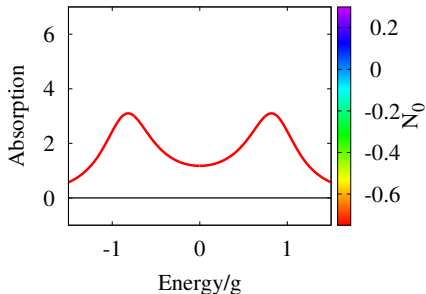
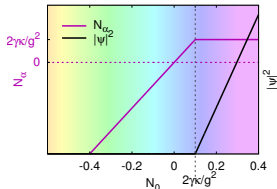
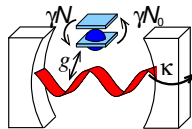
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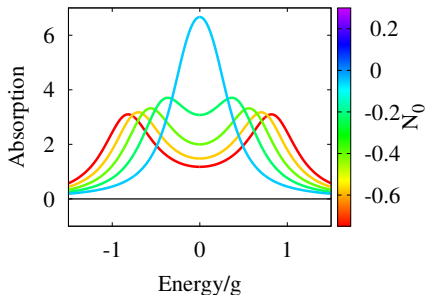
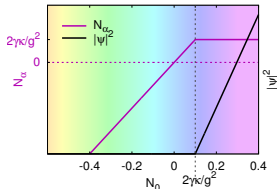
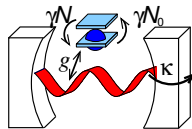
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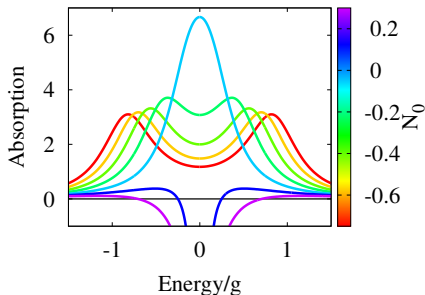
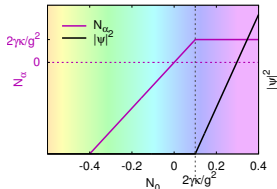
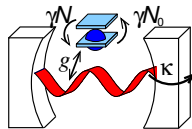
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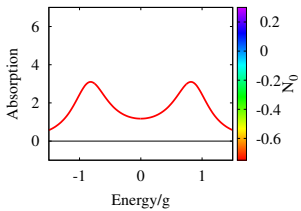
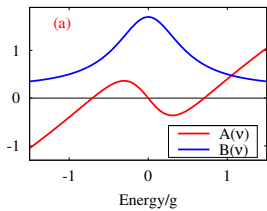
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Poles of Retarded Green's function and gain

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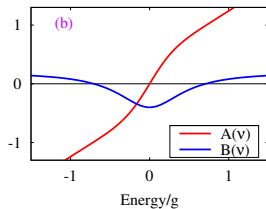
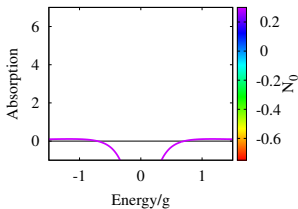
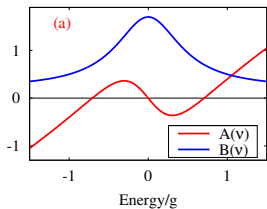
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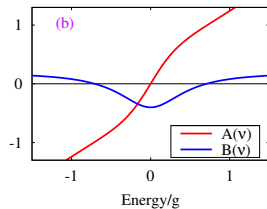
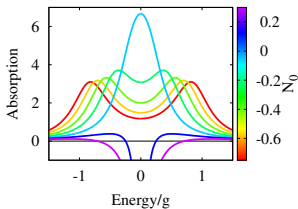
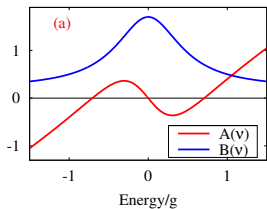
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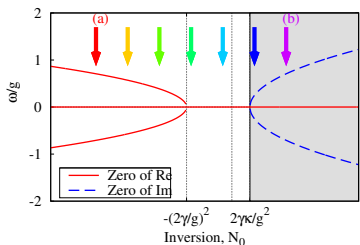


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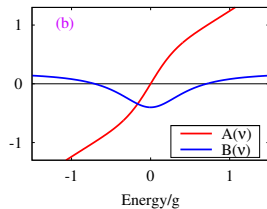
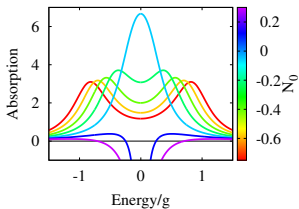
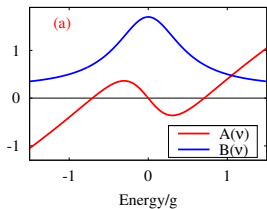


Laser:

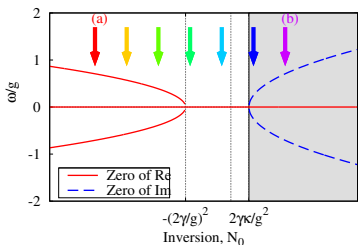


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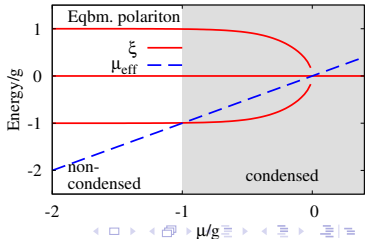
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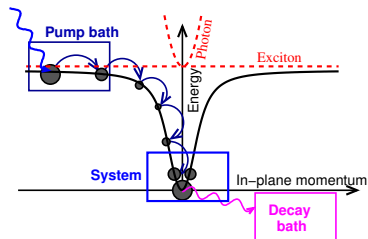
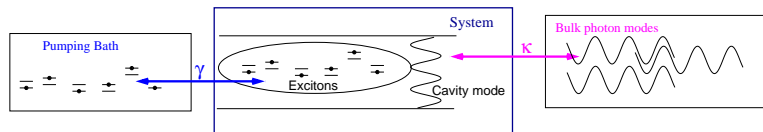
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Equilibrium:



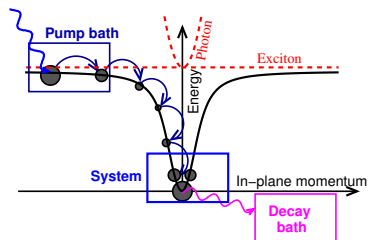
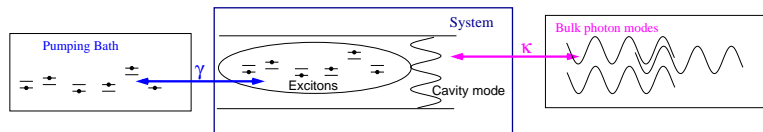
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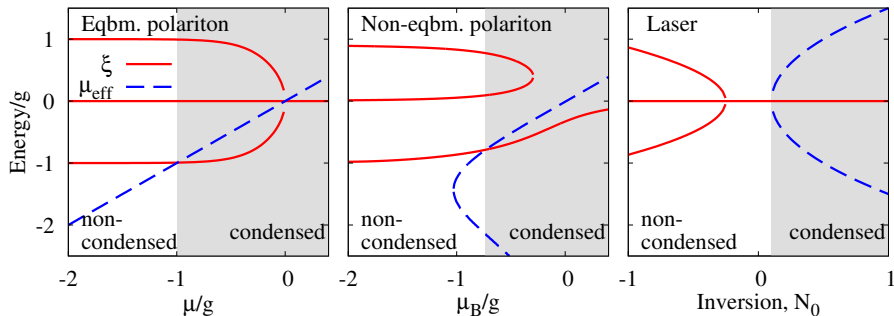
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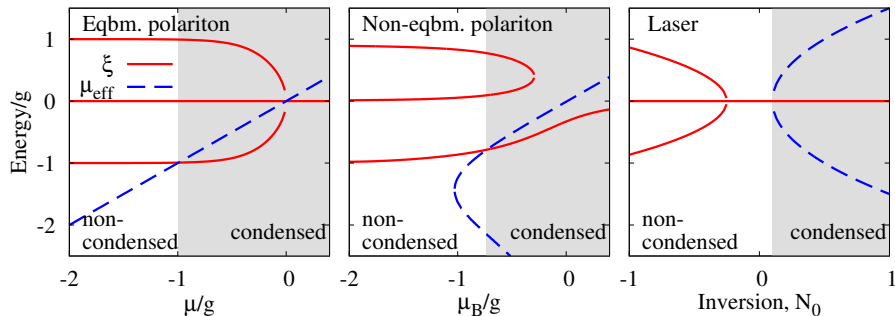
Strong coupling and lasing — low temperature phenomenon



- inversionless
- allows strong coupling
- requires low $T \leftrightarrow$ condensation
- Related weak-coupling inversionless lasing

[Szymanska *et al.* PRL '06; Keeling *et al.* book chapter 1010.3338]

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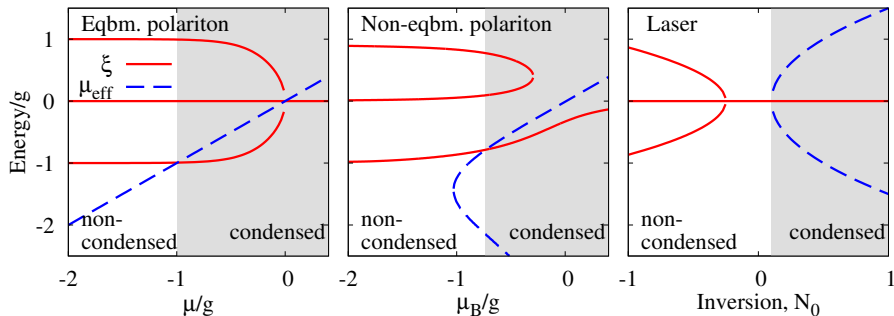


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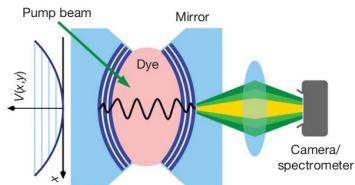
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Polariton and photon Condensation

- 1 Condensation, superradiance, lasing
- 2 Polariton condensation and Dicke model
 - Dicke model and condensation
 - Non-equilibrium condensation vs lasing
- 3 Room temperature condensates: Photons**
 - Lasing model and thermalisation
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- 4 Room temperature condensates: Organic polaritons
 - Dicke phase diagram with phonons
 - Condensation of phonon replicas?
 - (Ultra-strong phonon coupling?)
- 5 Conclusions

Photon BEC experiments

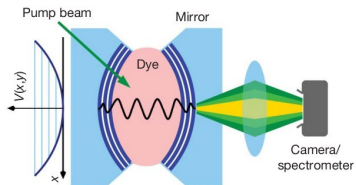


- Dye filled microcavity

➤ No strong coupling

[Klaers et al, Nature, 2010]

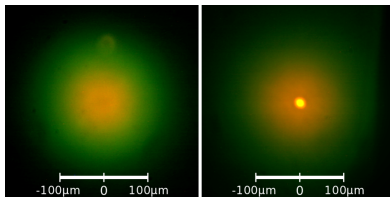
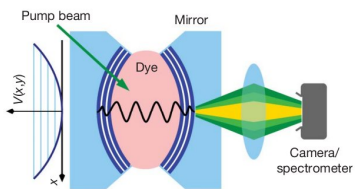
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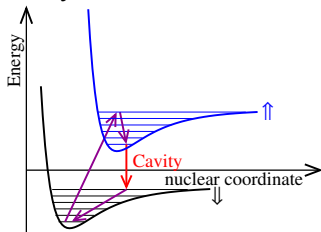
Relation to dye laser

- No electronic inversion
- No strong coupling
- No single cavity mode
 - ▶ Condensate mode is not maximum gain
 - ▶ Gain/Absorption in balance
- Thermalised many-mode system

Relation to dye laser

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4 Level Dye Laser

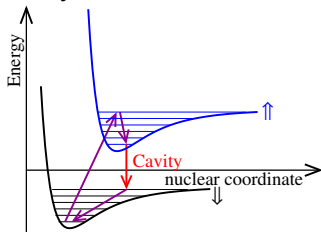


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4 Level Dye Laser



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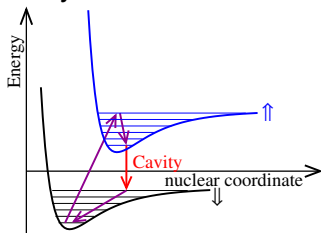
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Modelling

$$H_{\text{sys}} = \sum_m \omega_m \psi_m^\dagger \psi_m + \sum_\alpha [\epsilon S_\alpha^z + g (\psi_m S_\alpha^+ + \text{H.c.})]$$

- 2D harmonic cavity

$$\omega_m = \omega_{\text{cutoff}} + m\omega_{H.O.}$$

$$\text{Degeneracies } g_m = m + 1$$

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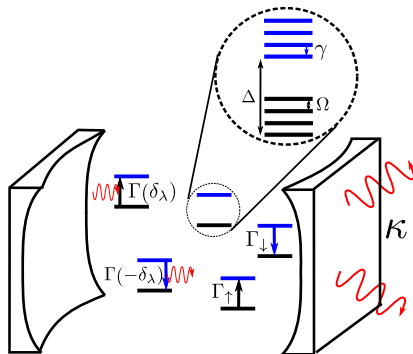
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$$\omega_m = \omega_{\text{cutoff}} + m\omega_{H.O.}$$

$$\text{Degeneracies } g_m = m + 1$$

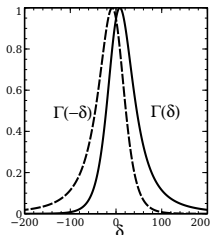
- Local vibrational mode



Modelling

Rate equation

$$\dot{\rho} = -i[H_0, \rho] - \sum_m \frac{\kappa}{2} \mathcal{L}[\psi_m] - \sum_{\alpha} \left[\frac{\Gamma_{\uparrow}}{2} \mathcal{L}[S_{\alpha}^{+}] + \frac{\Gamma_{\downarrow}}{2} \mathcal{L}[S_{\alpha}^{-}] \right] \\ - \sum_{m, \alpha} \left[\frac{\Gamma(\delta_m = \omega_m - \epsilon)}{2} \mathcal{L}[S_{\alpha}^{+} \psi_m] + \frac{\Gamma(-\delta_m = \epsilon - \omega_m)}{2} \mathcal{L}[S_{\alpha}^{-} \psi_m^{\dagger}] \right]$$



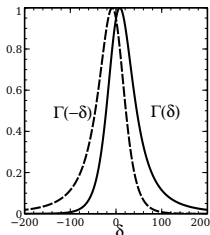
- $\Gamma(+\delta) \simeq \Gamma(-\delta) e^{-2\delta}$
- $\Gamma \rightarrow 0$ at large δ

[Marthaler et al PRL '11, Kirton & JK arXiv:1303.3459]

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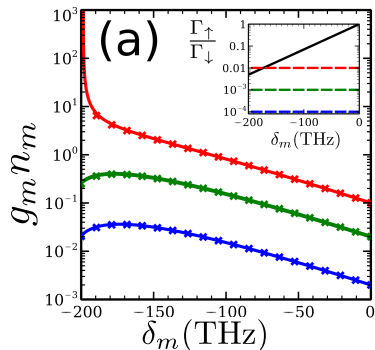
Distribution $g_m n_m$

- Rate equation — include spontaneous emission
- Bose-Einstein distribution without losses

[Kirton & JK arXiv:1303.3459]

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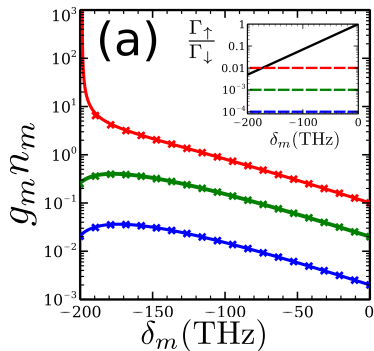


Low loss: Thermal

[Kirton & JK arXiv:1303.3459]

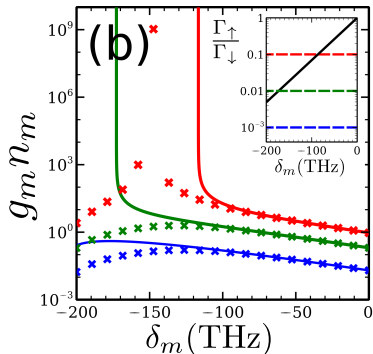
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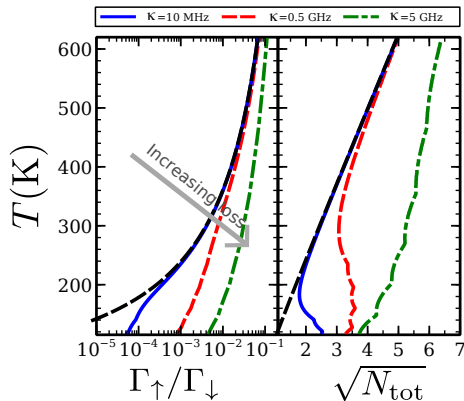
Low loss: Thermal

[Kirton & JK arXiv:1303.3459]



High loss \rightarrow Laser

Threshold condition



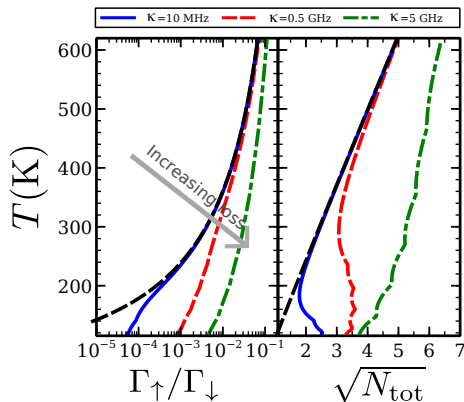
Compare threshold:

- Pump rate (Laser)
- Critical density (condensate)

- Thermal at low κ /high temperature
- High loss, κ competes with $\Gamma(\pm\delta_0)$
- Low temperature, $\Gamma(\pm\delta_0)$ shrinks

[Kirton & JK arXiv:1303.3459]

Threshold condition



Compare threshold:

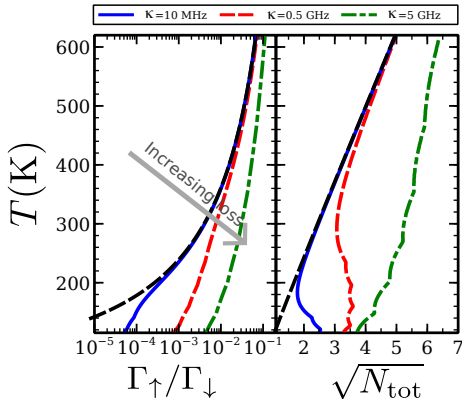
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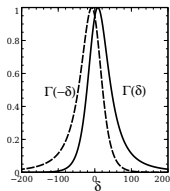
Threshold condition



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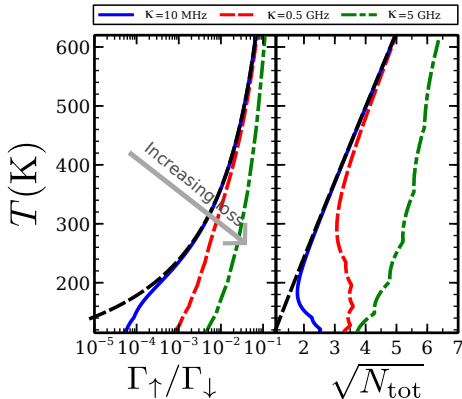
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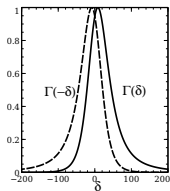


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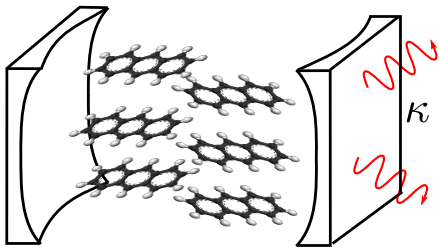


Organic polaritons: photon-exciton-phonon coupling

- 1 Condensation, superradiance, lasing
- 2 Polariton condensation and Dicke model
 - Dicke model and condensation
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- 3 Room temperature condensates: Photons
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 - (Ultra-strong phonon coupling?)
- 5 Conclusions

Organic materials in microcavities

- What?

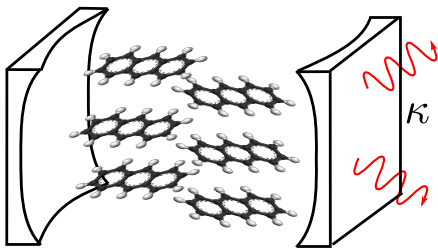


- Why?

- Lasing threshold at room T

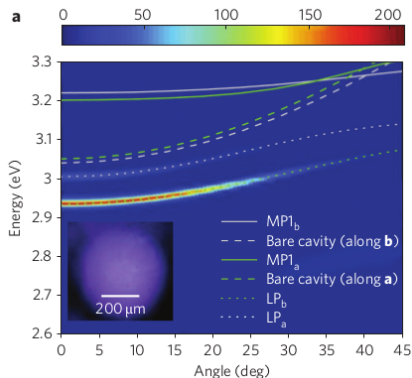
Organic materials in microcavities

• What?



• Lasing threshold at room T

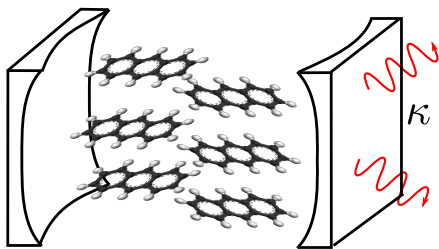
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Polariton splitting: 0.1 eV \leftrightarrow 1000 K.
[Kena Cohen and Forrest, Nat. Photon 2010]

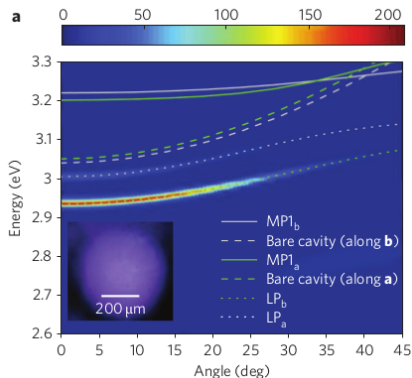
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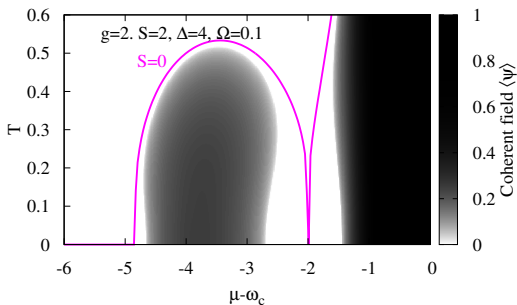
Phase diagram

$$H = \omega \psi^\dagger \psi + \sum_{\alpha} \left[\epsilon S_{\alpha}^z + g \left(\psi S_{\alpha}^+ + \psi^\dagger S_{\alpha}^- \right) + \Omega \left\{ b_{\alpha}^\dagger b_{\alpha} + \sqrt{S} \left(b_{\alpha}^\dagger + b_{\alpha} \right) S_{\alpha}^z \right\} \right]$$

- S suppresses condensation — reduces overlap
- Reentrant behaviour — Min μ at $T \sim 0.2$

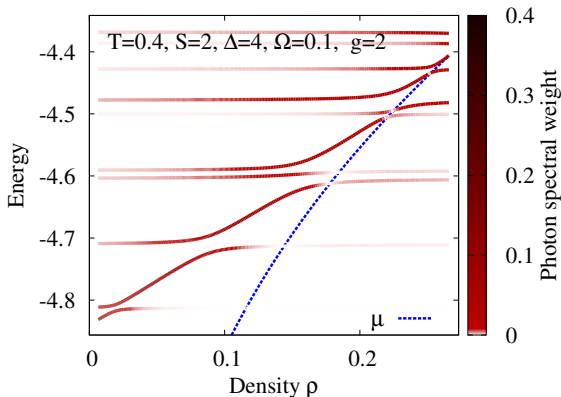
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Polariton spectrum: photon weight



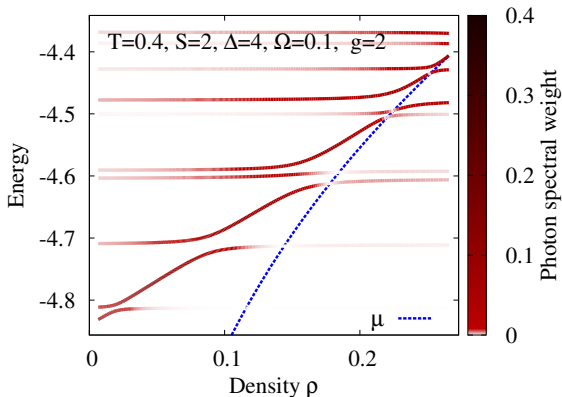
- Saturating 2LS: $g_{\text{eff}}^2 \sim g^2(1 - 2\rho)$

• What is nature of polariton mode?

• $D(t) = -\langle \psi^\dagger(t)\psi(0) \rangle$, $D(\omega) = \sum_n \frac{Z_n}{\omega - \omega_n}$

[Cwik *et al.* arXiv:1303.3702]

Polariton spectrum: photon weight

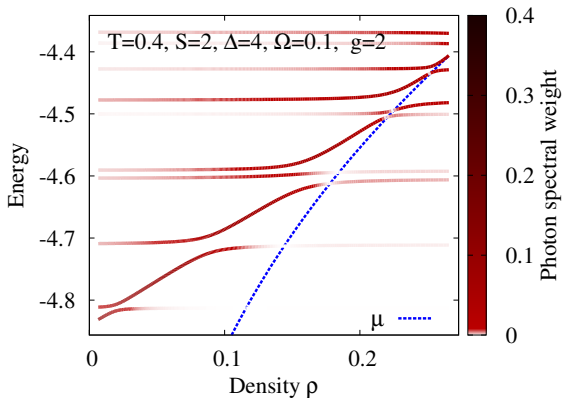


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[Cwik *et al.* arXiv:1303.3702]

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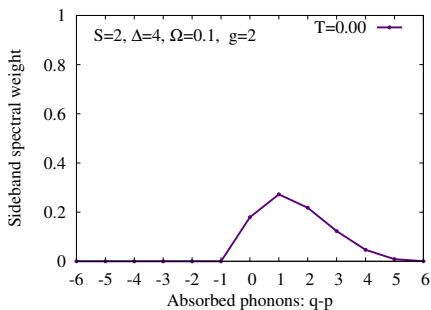
Polariton spectrum: what condensed

- Repeat weight for n -phonon channel
- Eigenvector that is macroscopically occupied
- Optimal $T \sim 2\Omega$

[Cwik *et al.* arXiv:1303.3702]

Polariton spectrum: what condensed

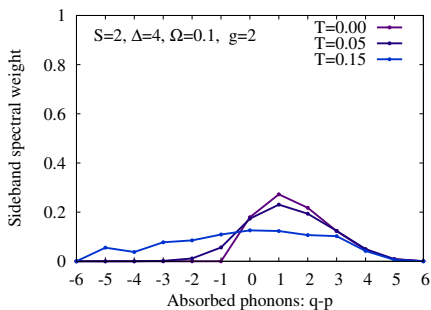
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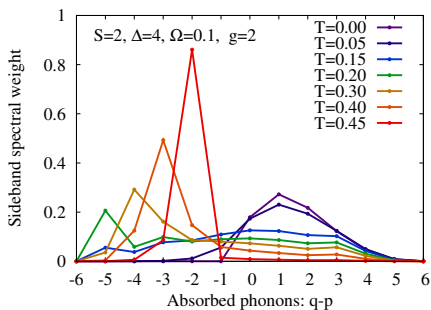
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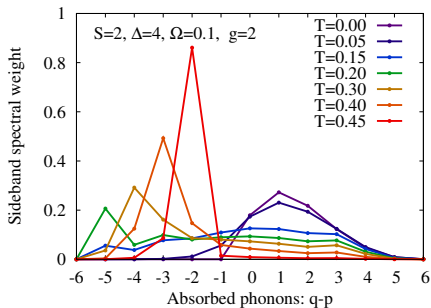


Optimal $T \sim 2\Omega$

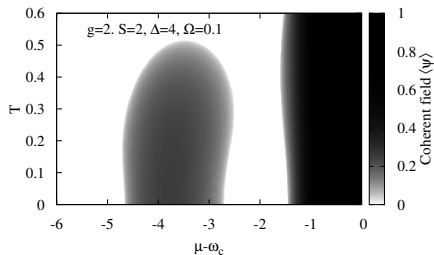
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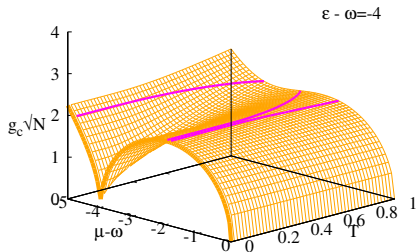
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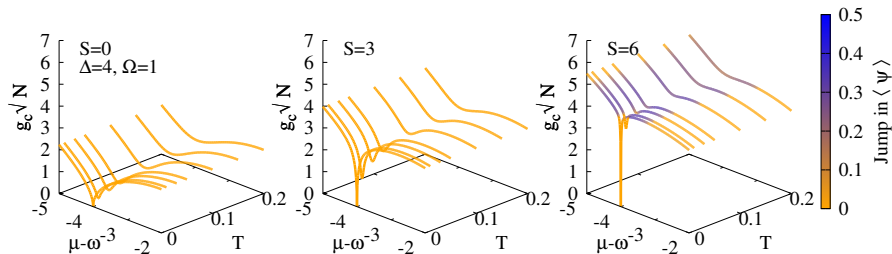
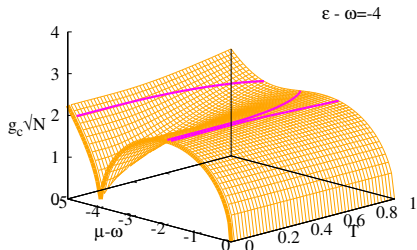
Critical coupling with increasing S

- Re-orient phase diagram
- g vs μ, T
- Colors \rightarrow Jump of $\langle \psi \rangle$



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Explanation: Polaron formation

- Unitary transform

$$H_\alpha \rightarrow \tilde{H}_\alpha = e^{K_\alpha} H_\alpha e^{-K_\alpha} \quad K = \sqrt{S} S_\alpha^z (b_\alpha^\dagger - b_\alpha)$$

- Coupling moves to S^\pm

$$\tilde{H}_\alpha = \text{const.} + \epsilon S_\alpha^z + \Omega b_\alpha^\dagger b_\alpha + g \left[\psi S_\alpha^+ e^{\sqrt{S}(b_\alpha^\dagger - b_\alpha)} + \text{H.c.} \right]$$

- Optimal phonon displacements, $\sim \sqrt{S}$

- Reduced $g_{\text{eff}} \sim g \times \exp(-S/2)$

- For $\psi \neq 0$, competition

$$\text{Variational MFT } |\psi\rangle_\alpha \sim \exp(-\eta K_\alpha - \zeta (b_\alpha^\dagger)) |0, S\rangle_\alpha$$

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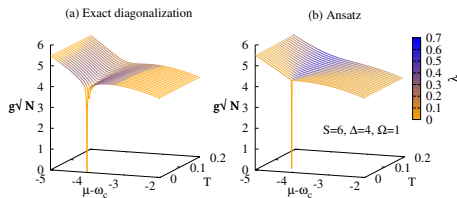
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Collective polaron formation

- Compares well at $S \gg 1$
- Coherent bosonic state



- Feedback: Large/small $g_{\text{eff}} \leftrightarrow \lambda = \langle \psi \rangle$
- Variational free energy

$$F = (\omega_c - \mu)\lambda^2 + N \left\{ \Omega \left[\zeta^2 - S \frac{\eta(2-\eta)}{4} \right] - T \ln \left[2 \cosh \left(\frac{\xi}{T} \right) \right] \right\}$$

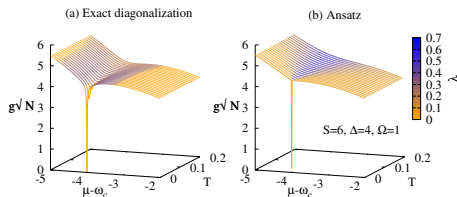
Effective 2LS energy in field:

$$\zeta^2 = \left(\frac{\omega_c - \mu}{2} + \Omega \sqrt{S} (1 - \eta) \zeta \right)^2 + g^2 \lambda^2 e^{-5\eta}$$

[Cwik *et al.* arXiv:1303.3702]

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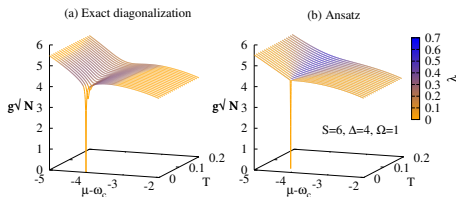
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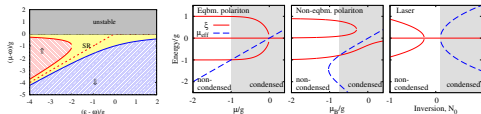
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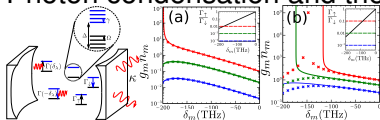
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Summary

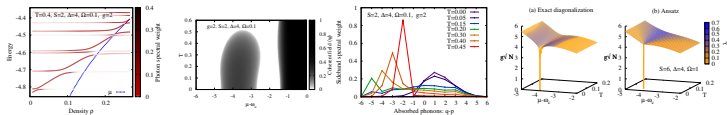
● Polariton condensation vs lasing; superradiance



● Photon condensation and thermalisation



● Reentrance, phonon assisted transition, 1st order at $S \gg 1$

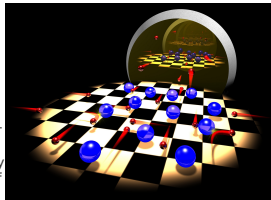


Many body quantum optics and correlated states of light

9:00 am on Monday 28 October 2013 – 5:00 pm on Tuesday 29 October 2013

at: **The Royal Society at Chicheley Hall, home of the Kavli Royal Society International Centre, Buckinghamshire**

Theo Murphy international scientific meeting organised by Dr Jonathan Keeling, Professor Steven Girvin, Dr Michael Hartmann and Professor Peter Littlewood FRS.



List of speakers and chairs

Professor Iacopo Carusotto, Professor Andrew Cleland, Professor Hui Deng, Professor Tilman Esslinger, Professor Rosario Fazio, Professor Ed Hinds, Professor Andrew Houck, Professor Ataç İmamoğlu, Professor Jens Koch, Professor Misha Lukin, Professor Martin Plenio, Professor Arno Rauschenbeutel, Professor Timothy Spiller, Professor Jacob Taylor, Professor Hakan Türeci, Professor Andreas Wallraff

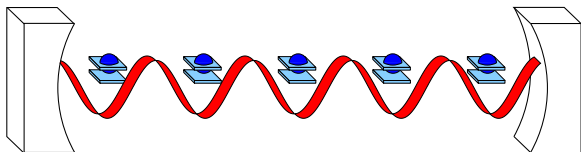
Attending this event

This is a residential conference which allows for increased discussion and networking. It is free to attend, however participants need to cover their accommodation and catering costs if required. Places are limited and therefore pre-registration is essential.

Extra slides

- 6 No go theorem
- 7 Retarded Green's function for laser
- 8 Organic properties
- 9 Anticrossing vs ρ

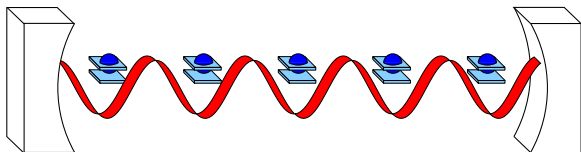
No go theorem and transition



Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

[Rzazewski *et al* PRL '75]

No go theorem and transition



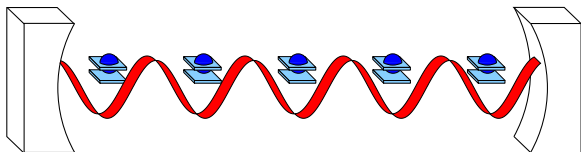
Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

No go theorem: Minimal coupling $(p - eA)^2/2m$

$$-\sum_i \frac{e}{m} A \cdot p_i \Leftrightarrow g(\psi^\dagger S^- + \psi S^+), \quad \sum_i \frac{A^2}{2m} \Leftrightarrow N\zeta(\psi + \psi^\dagger)^2$$

[Rzazewski *et al* PRL '75]

No go theorem and transition



Spontaneous polarisation if: $Ng^2 > \omega\epsilon$

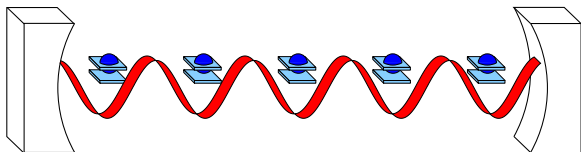
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No go theorem and transition



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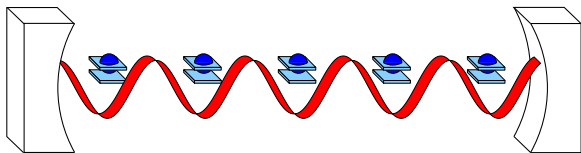
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Dicke phase transition: ways out

Problem: $g^2/\omega_0 < 2\zeta$ for intrinsic parameters. **Solutions:**

- Interpretation
 - Ferroelectric transition in $\mathbf{D} \cdot \mathbf{r}$ gauge.
[JK JPCM '07, Vukics & Domokos PRA 2012]
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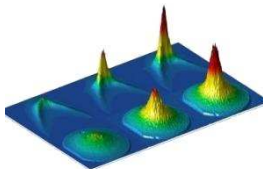
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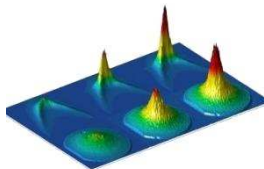
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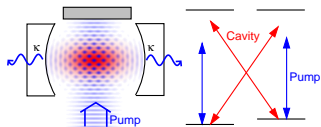
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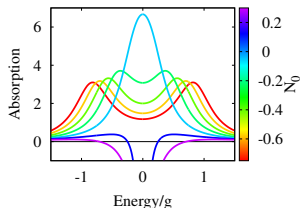
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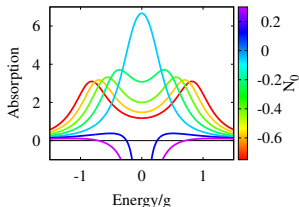


Maxwell-Bloch Equations: Retarded Green's function



- Introduce $D^R(\omega)$:
Response to perturbation
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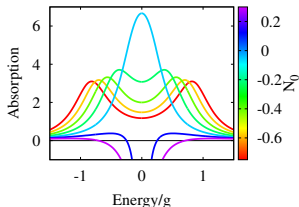
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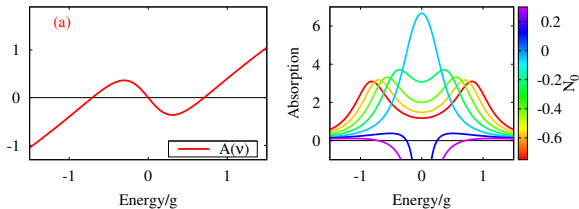
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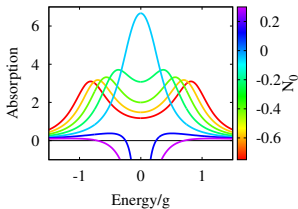
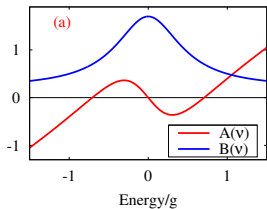
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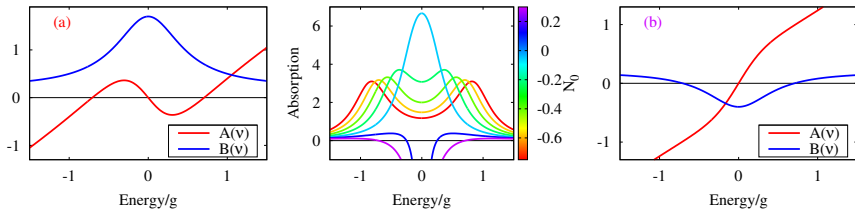
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Organic materials in microcavities

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- ★ J aggregates [Bulovic *et al.*]

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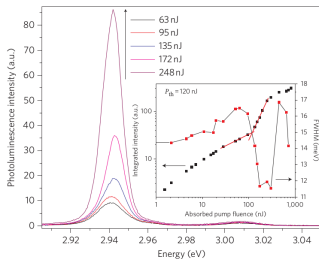
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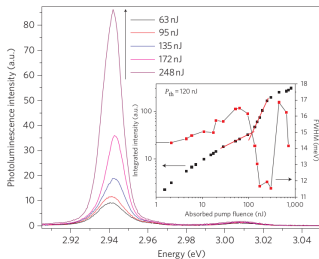
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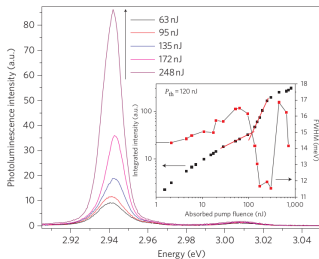
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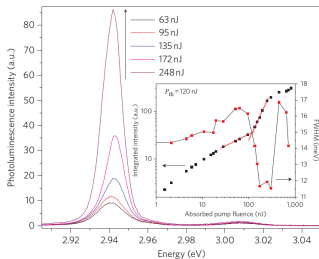
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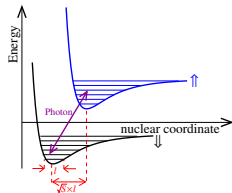
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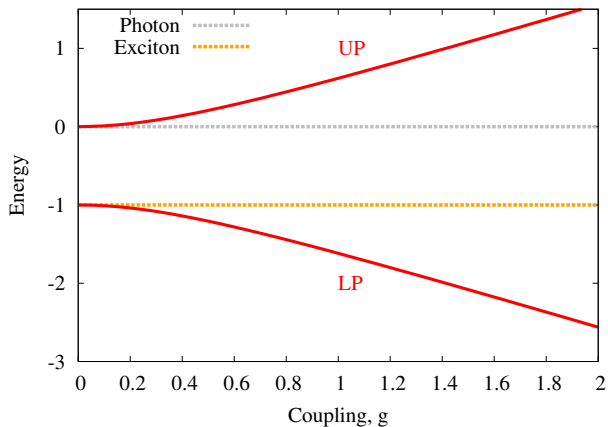
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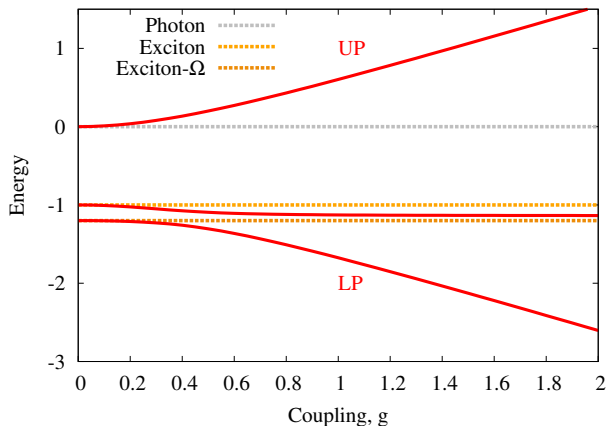


Polariton spectrum — coupled oscillators

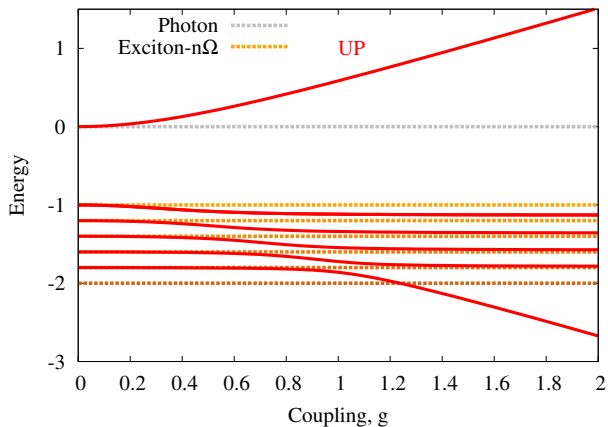
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