

Non-equilibrium phases of coupled matter-light systems

Jonathan Keeling



University of
St Andrews

600
YEARS



Los Alamos, June 2013

Coupling many atoms to light

Old question: *What happens to radiation when many atoms interact “collectively” with light.*

Superradiance — dynamical and steady state.

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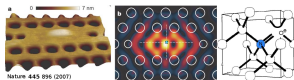
Superradiance — dynamical and steady state.

New relevance

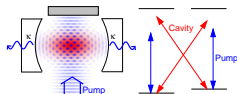
- Superconducting qubits



- Quantum dots & NV centres

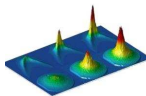


- Ultra-cold atoms



- Rydberg atoms/polaritons

- Microcavity Polaritons



Dicke effect: Enhanced emission

PHYSICAL REVIEW

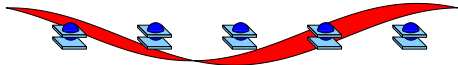
VOLUME 93, NUMBER 1

JANUARY 1, 1954

Coherence in Spontaneous Radiation Processes

R. H. DICKE

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey



$$H_{\text{int}} = \sum_{k,i} g_k \left(\psi_k^\dagger S_i^- e^{-i\mathbf{k}\cdot\mathbf{r}_i} + \text{H.c.} \right)$$

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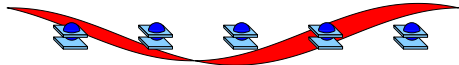
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Collective decay: $\frac{d\rho}{dt} = -\frac{\Gamma}{2} [S^+ S^- \rho - S^- \rho S^+ + \rho S^+ S^-]$

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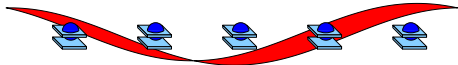
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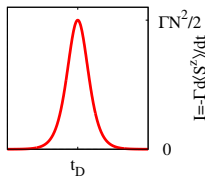
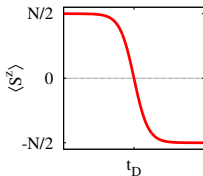
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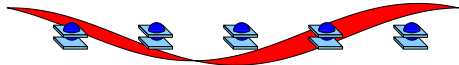
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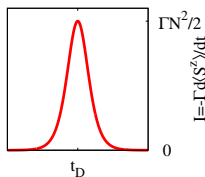
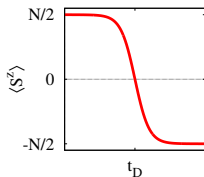
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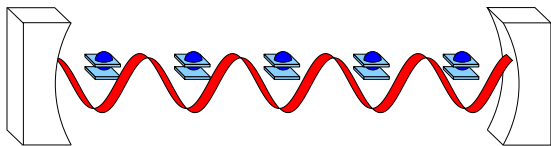
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Problem: dipole interactions dephase. [Friedberg et al, Phys. Lett. 1972]

Collective radiation **with a cavity**: Dynamics

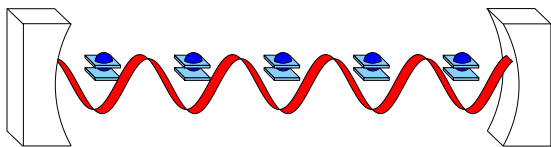


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Single cavity mode: oscillations

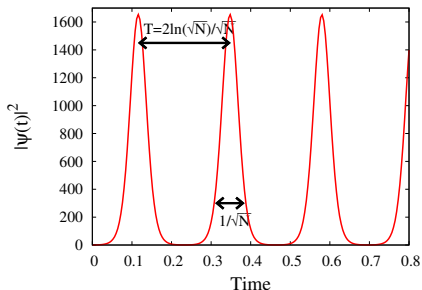
[Bonifacio and Preparata PRA '70]

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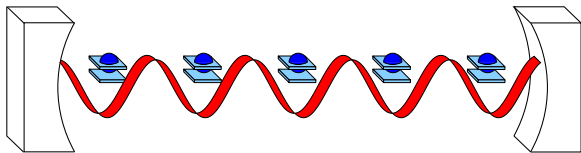
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Dicke model: Equilibrium superradiance transition



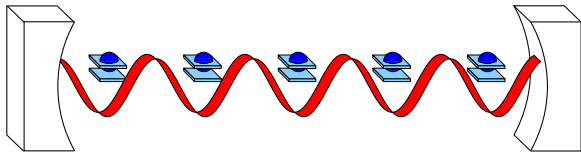
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[Hepp, Lieb, Ann. Phys. '73]

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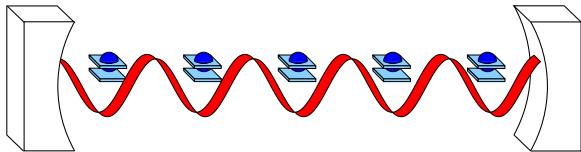
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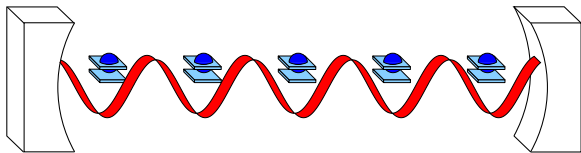
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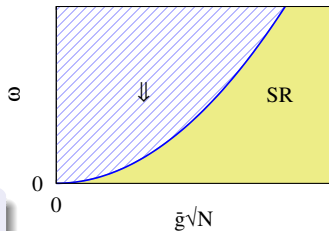
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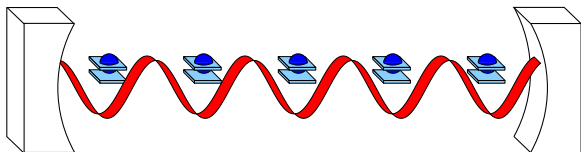
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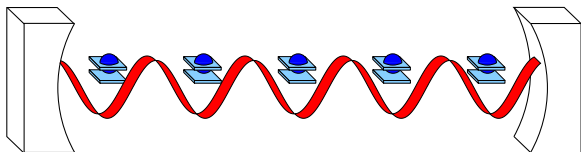
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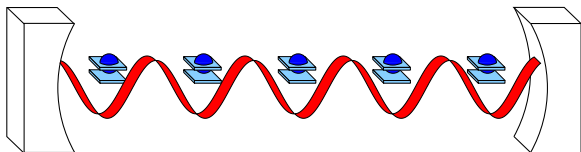
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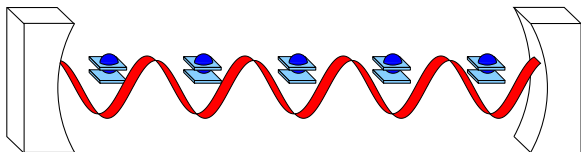
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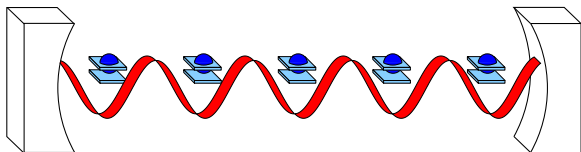
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But Thomas-Reiche-Kuhn sum rule states: $g^2/\omega_0 < 2\zeta$. **No transition**

[Rzazewski *et al* PRL '75]

Dicke phase transition: ways out

Problem: $g^2/\omega_0 < 2\zeta$ for intrinsic parameters. **Solutions:**

- Interpretation
 - Ferroelectric transition in $\mathbf{D} \cdot \mathbf{r}$ gauge.
[JK JPCM '07, Vukics & Domokos PRA 2012]
 - Circuit QED [Nataf and Cluzet, Nat. Comm. '10; Viehmann *et al.* PRL '11]
- Grand canonical ensemble:
 - If $H \rightarrow H - \mu(S^z + \psi^\dagger\psi)$, need only:
 $g^2 N > (\omega - \mu)(\omega_0 - \mu)$
 - Incoherent pumping \rightarrow polariton condensation.
- Dissociate g, ω_0 ,
e.g. Raman scheme: $\omega_0 \ll \omega$.
[Dimer *et al.* PRA '07; Baumann *et al.* Nature '10. Also, Black *et al.* PRL '03]

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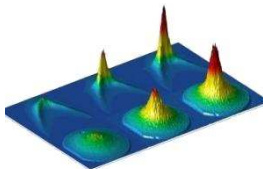
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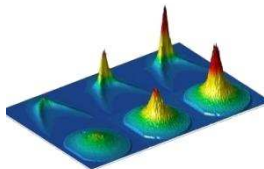
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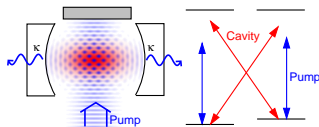
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- 1 Introduction: Dicke model and superradiance
- 2 Dynamics of generalized Dicke model
 - Summary of experiment and classical dynamics
 - Fixed points and dynamical phases
 - Timescales and consequences for experiment
 - Persistent oscillating phases
- 3 Jaynes Cummings Hubbard model
 - JCHM vs Dicke
 - Coherently driven array
 - Disorder

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GROUP:



COLLABORATORS: Simons, Bhaseen, Schmidt, Blatter, Türeci, Krüger
EXPERIMENT: Houck, Wallraff, Fink, Mylnek

FUNDING:



Dynamics of generalized Dicke model



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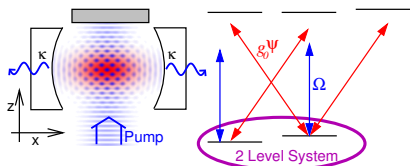
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Cold-atom extended Dicke model



2 Level system, $|\downarrow\rangle, |\uparrow\rangle$:

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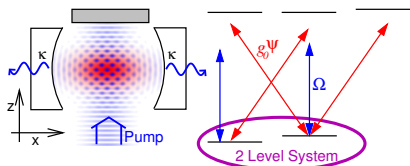
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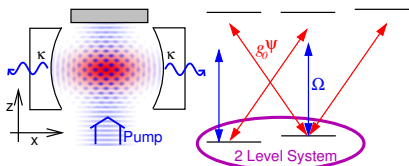
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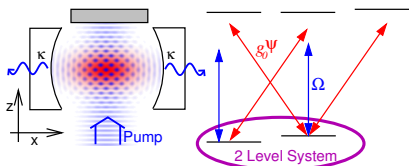
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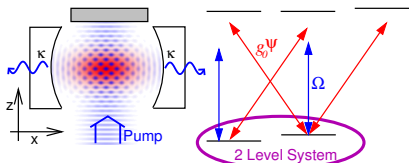
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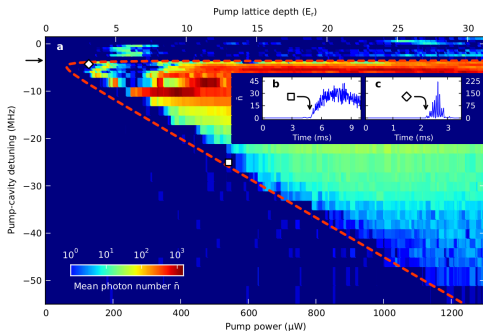
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Classical dynamics of the extended Dicke model

Open dynamical system:

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- Neglects quantum fluctuations — restore via Wigner distributed initial conditions.
- Linearisation about fixed point:
 - Recover Retarded Green's function (spectrum)
 - Cannot recover occupations

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Classical EOM
($|\mathbf{S}| = N/2 \gg 1$)

$$\dot{\mathbf{S}}^- = -i(\omega_0 + U|\psi|^2)\mathbf{S}^- + 2ig(\psi + \psi^*)\mathbf{S}^z$$
$$\dot{\mathbf{S}}^z = ig(\psi + \psi^*)(\mathbf{S}^- - \mathbf{S}^+)$$
$$\dot{\psi} = -[\kappa + i(\omega + U\mathbf{S}^z)]\psi - ig(\mathbf{S}^- + \mathbf{S}^+)$$

- Neglects quantum fluctuations — restore via Wigner distributed initial conditions.
- Linearisation about fixed point:
 - Recover Retarded Green's function (spectrum)
 - Cannot recover occupations

Classical dynamics of the extended Dicke model

Open dynamical system:

$$H = \omega\psi^\dagger\psi + \omega_0\mathbf{S}^z + g(\psi + \psi^\dagger)(\mathbf{S}^- + \mathbf{S}^+) + U\mathbf{S}_z\psi^\dagger\psi.$$
$$\partial_t\rho = -i[H, \rho] - \kappa(\psi^\dagger\psi\rho - 2\psi\rho\psi^\dagger + \rho\psi^\dagger\psi)$$

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- Neglects quantum fluctuations — restore via Wigner distributed initial conditions.
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Fixed points (steady states)

$$0 = i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z$$

$$0 = ig(\psi + \psi^*)(S^- - S^+)$$

$$0 = -[\kappa + i(\omega + US^z)]\psi - ig(S^- + S^+)$$

• $\psi = 0, S = (0, 0, \pm N/2)$
always a solution.

• If $g > g_c, \psi \neq 0$ too

A. $S^z = -S[S^-] = 0$

B. $\psi = \Re[\psi] = 0$

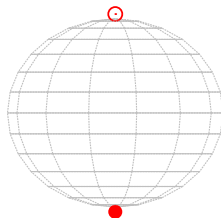
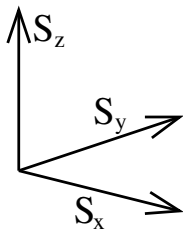
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Small g : \uparrow, \downarrow only.
 $(\omega = 30\text{MHz}, UN = -40\text{MHz})$

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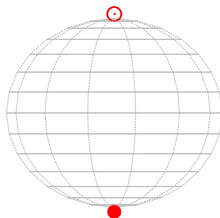
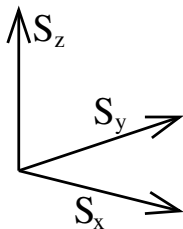
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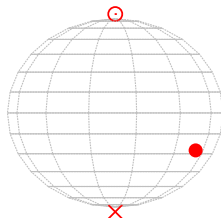
A $S^y = -\Im[S^-] = 0$

B $\psi' = \Re[\psi] = 0$



Small g : \uparrow, \downarrow only.

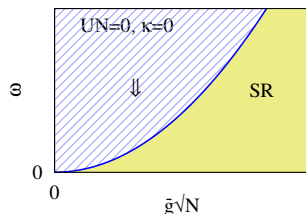
($\omega = 30\text{MHz}, UN = -40\text{MHz}$)



Larger g : SR too.

Steady state phase diagram

$$\begin{aligned}0 &= i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z \\0 &= ig(\psi + \psi^*)(S^- - S^+) \\0 &= -[\kappa + i(\omega + US^z)]\psi - ig(S^- + S^+)\end{aligned}$$



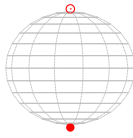
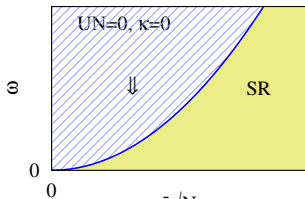
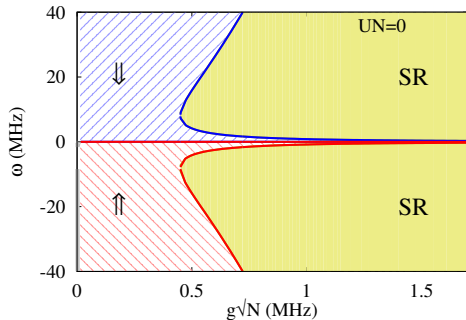
See also Domokos and Ritsch PRL '02, Domokos *et al.* PRL '10

Steady state phase diagram

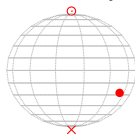
$$0 = i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z$$

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$SR(A): S_y = 0$



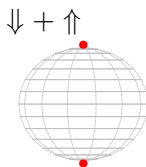
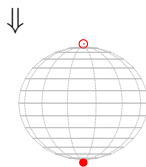
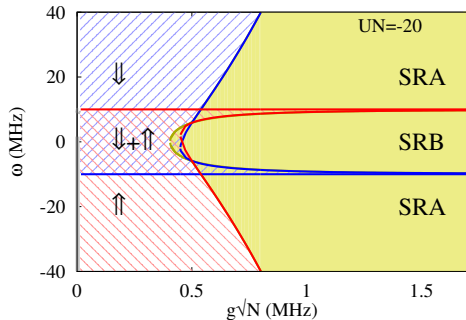
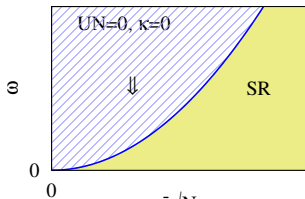
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Steady state phase diagram

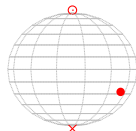
$$0 = i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z$$

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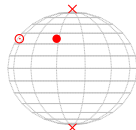
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$\bar{g}\sqrt{N}$
SR(A): $S_y = 0$



SR(B): $\psi' = 0$



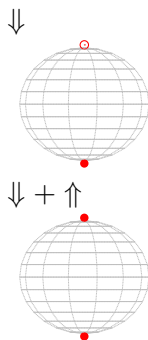
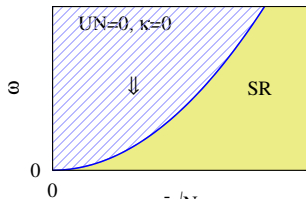
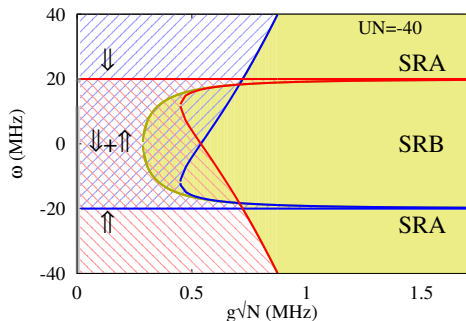
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Steady state phase diagram

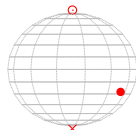
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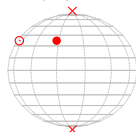
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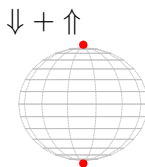
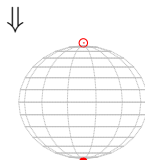
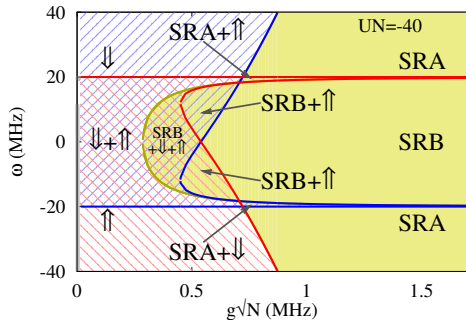
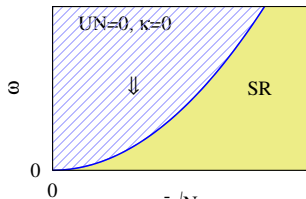
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Steady state phase diagram

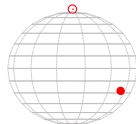
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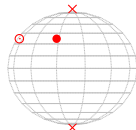
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$\bar{g}\sqrt{N}$
SR(A): $S_y = 0$

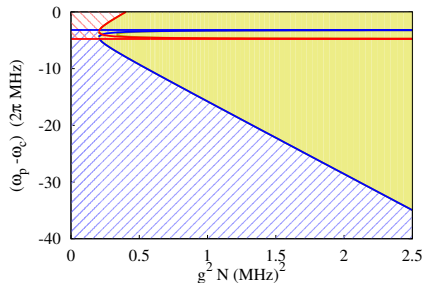


SR(B): $\psi' = 0$



See also Domokos and Ritsch PRL '02, Domokos *et al.* PRL '10

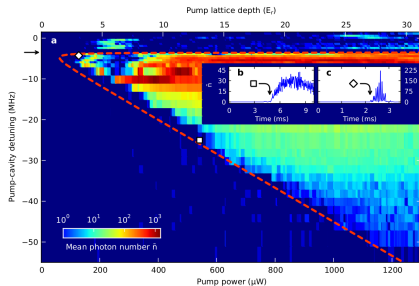
Comparison to experiment



$$UN = -10\text{MHz}$$

Adapted from: [Bhaseen *et al.* PRA '12]

$$\omega = \omega_c - \omega_p + \frac{5}{2}UN,$$



[Baumann *et al* Nature '10]

$$UN = -\frac{g_0^2}{4(\omega_a - \omega_c)}$$

Dynamics of generalized Dicke model



1 Introduction: Dicke model and superradiance

2 Dynamics of generalized Dicke model

- Summary of experiment and classical dynamics
- Fixed points and dynamical phases
- **Timescales and consequences for experiment**
- Persistent oscillating phases

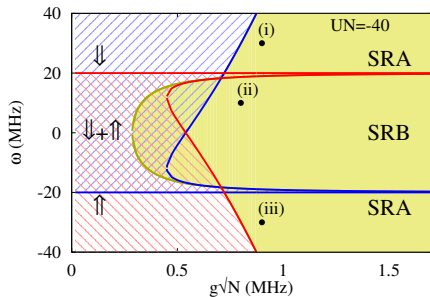
3 Jaynes Cummings Hubbard model

- JCHM vs Dicke
- Coherently driven array
- Disorder

Dynamics: Evolution from normal state

Gray: $\mathbf{S} = (\sqrt{N}, \sqrt{N}, -N/2)$

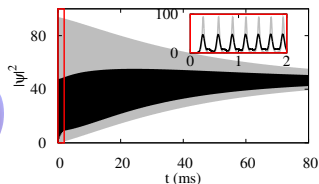
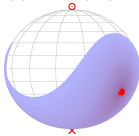
Black: Wigner distribution of \mathbf{S}, ψ



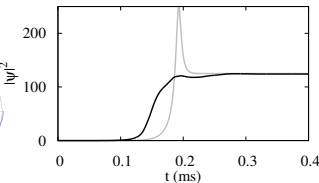
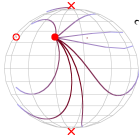
Oscillations: ~ 0.1 ms

Decay: 20ms, 0.1ms, 20ms

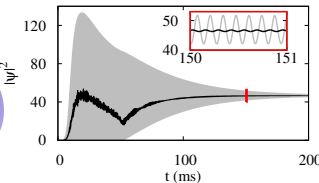
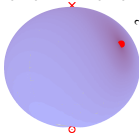
(i) SR(A)



(ii) SR(B)



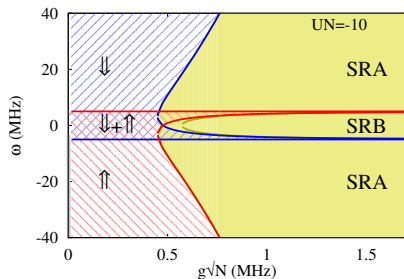
(iii) SR(A)



Asymptotic state: Evolution from normal state

(Near to experimental $UN = -13\text{MHz}$).

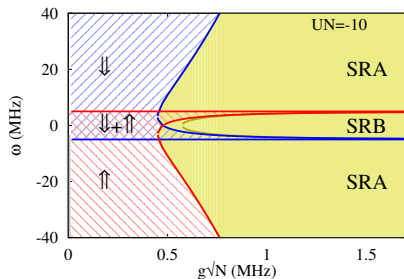
All stable attractors:



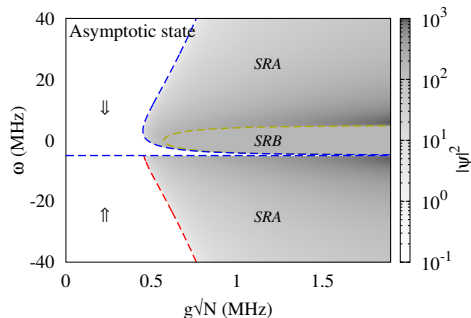
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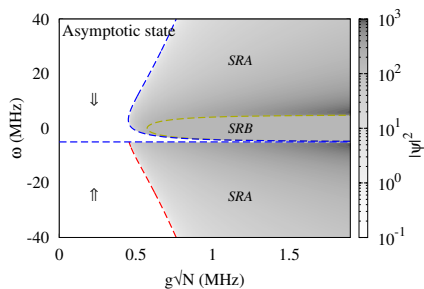
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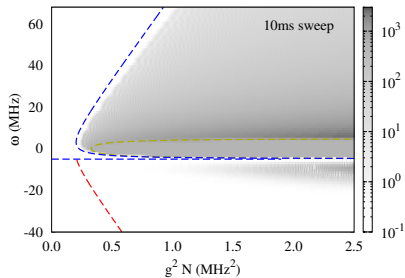
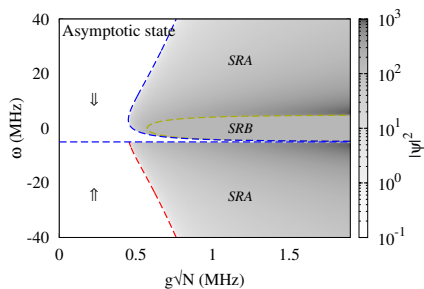
Starting from \Downarrow



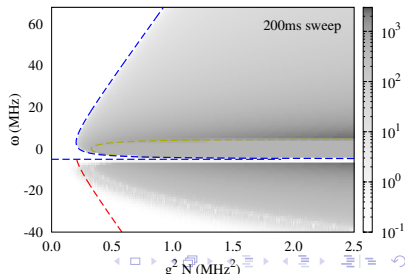
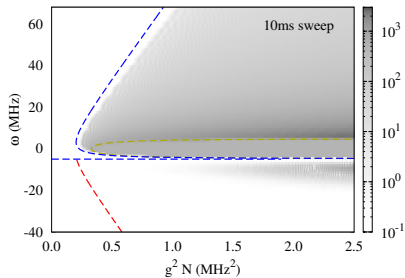
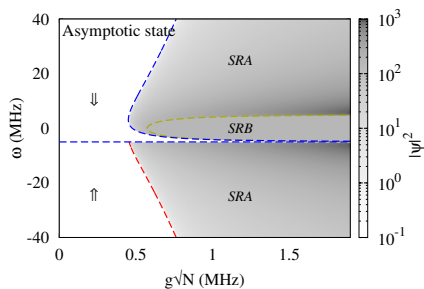
Timescales for dynamics: Consequences for experiment



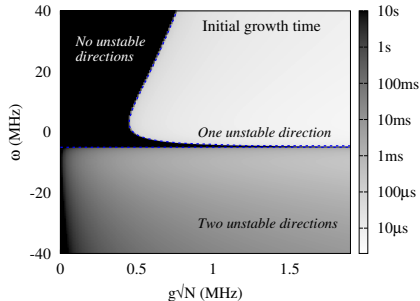
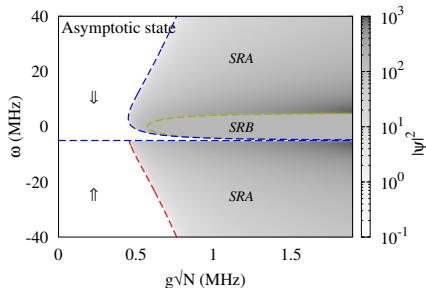
Timescales for dynamics: Consequences for experiment



Timescales for dynamics: Consequences for experiment



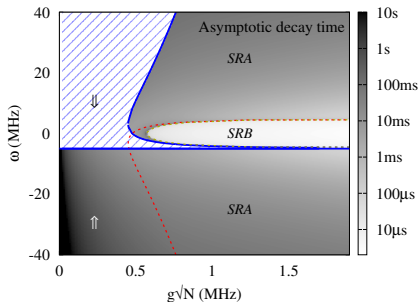
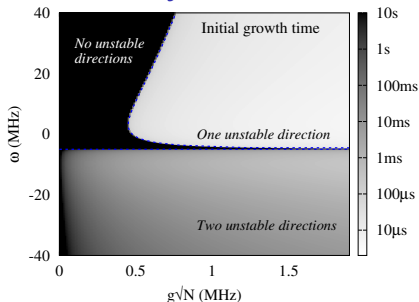
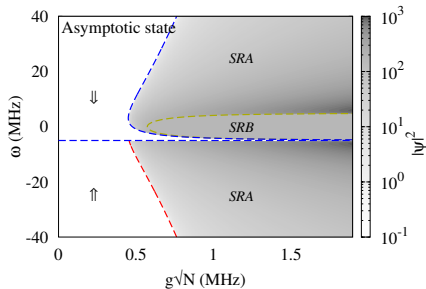
Timescales for dynamics: What are they?



Growth Most unstable eigenvalues near $\mathbf{S} = (0, 0, -N/2)$

Decay Slowest stable eigenvalues near final state

Timescales for dynamics: What are they?

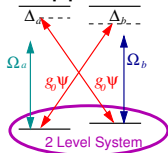


Growth Most unstable eigenvalues near $\mathbf{S} = (0, 0, -N/2)$

Decay Slowest stable eigenvalues near final state

Timescales for dynamics: Why so slow and varied?

Suppose co- and counter-rotating terms differ

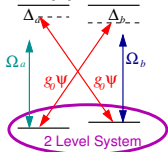


$$H = \dots + g(\psi^\dagger S^- + \psi S^+) + g'(\psi^\dagger S^+ + \psi S^-) + \dots$$

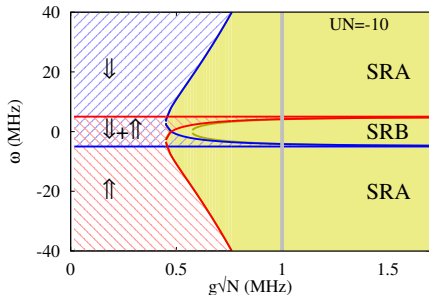
- SR(A) near phase boundary at small $\delta g \rightarrow$ Critical slowing down
- SR(A), SR(B) continuously connect

Timescales for dynamics: Why so slow and varied?

Suppose co- and counter-rotating terms differ



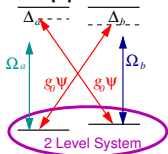
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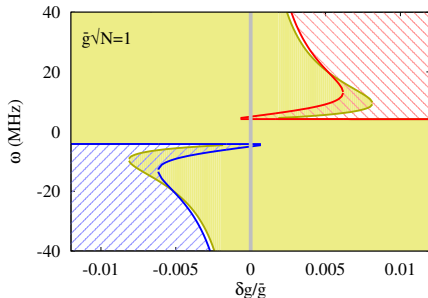
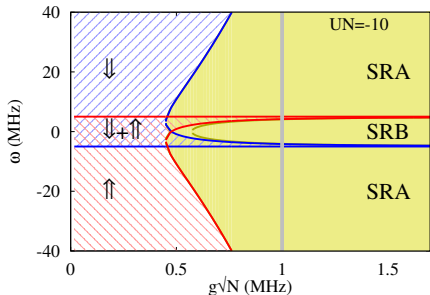
Timescales for dynamics: Why so slow and varied?

Suppose co- and counter-rotating terms differ



$$H = \dots + g(\psi^\dagger S^- + \psi S^+) + g'(\psi^\dagger S^+ + \psi S^-) + \dots$$

$$\delta g = g' - g, \quad 2\bar{g} = g' + g$$



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- SR(A), SR(B) continuously connect

Dynamics of generalized Dicke model



1 Introduction: Dicke model and superradiance

2 Dynamics of generalized Dicke model

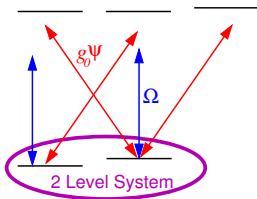
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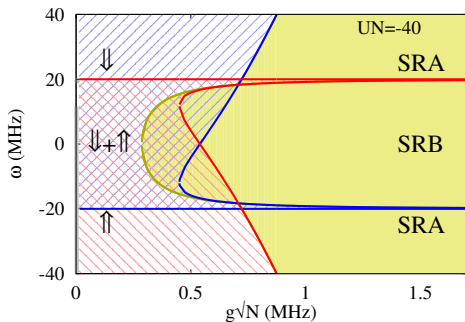
- JCHM vv Dicke
- Coherently driven array
- Disorder

Regions without fixed points

Changing U :

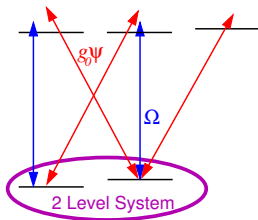


$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$

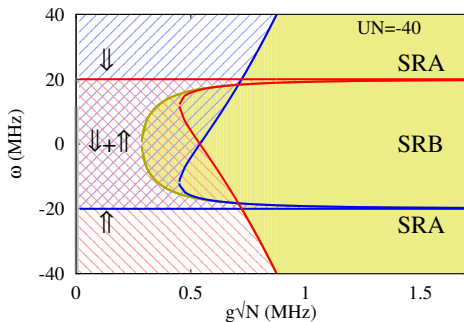


Regions without fixed points

Changing U :

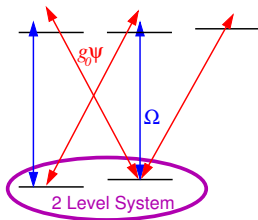


$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$

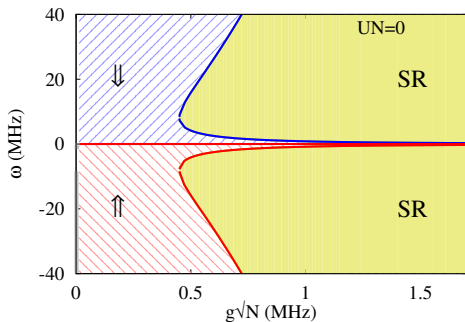


Regions without fixed points

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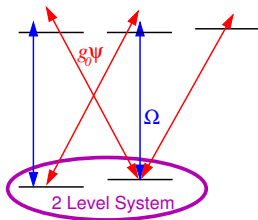


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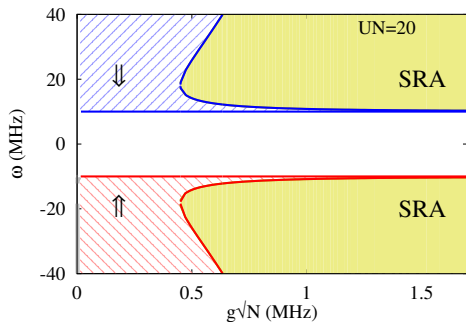


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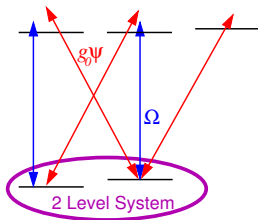


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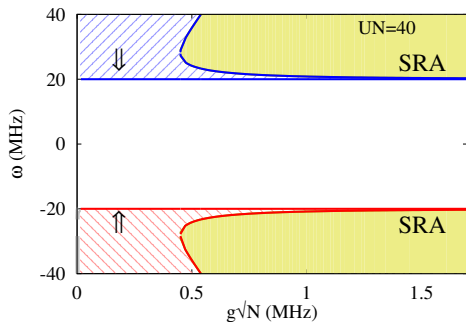


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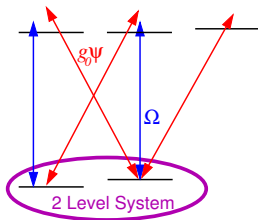


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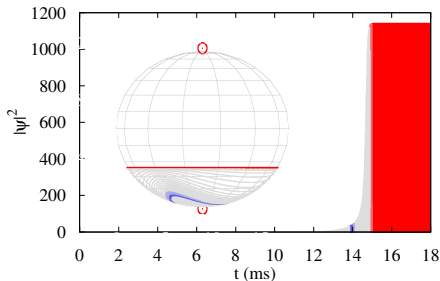
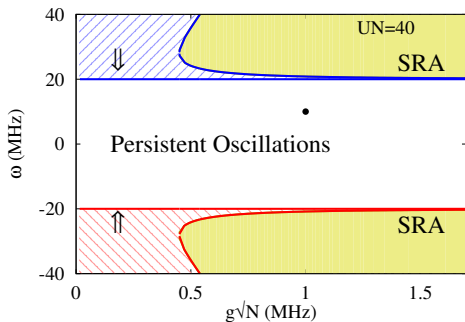


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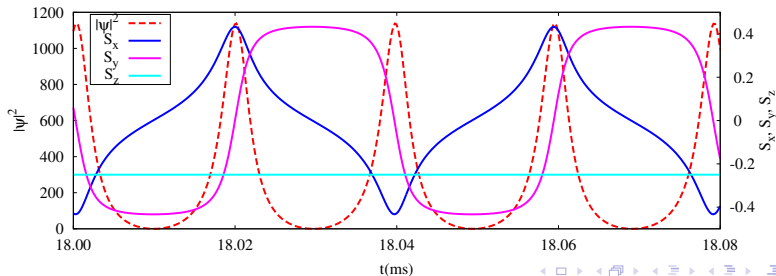
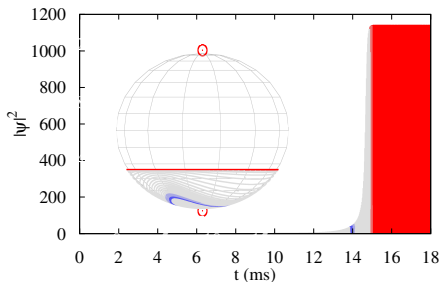
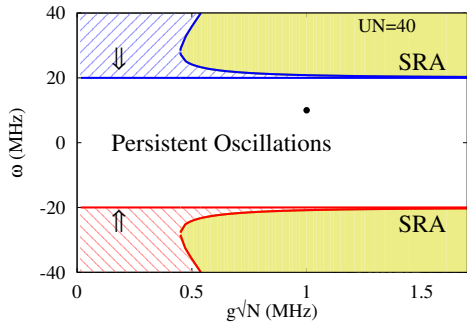
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Persistent (optomechanical) oscillations



Jaynes Cummings Hubbard model



1 Introduction: Dicke model and superradiance

2 Dynamics of generalized Dicke model

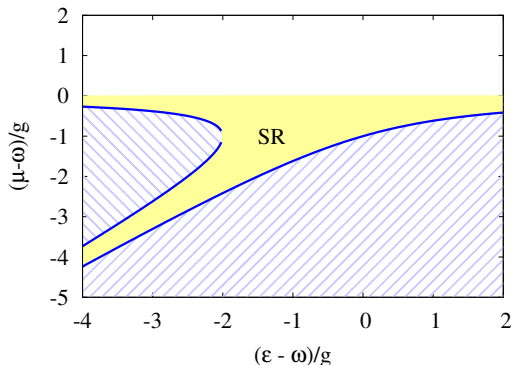
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3 Jaynes Cummings Hubbard model

- JCHM vv Dicke
- Coherently driven array
- Disorder

Equilibrium: Dicke model with chemical potential

$$H - \mu N = (\omega - \mu)\psi^\dagger\psi + (\omega_0 - \mu)S^z + g(\psi^\dagger S^- + \psi S^+)$$



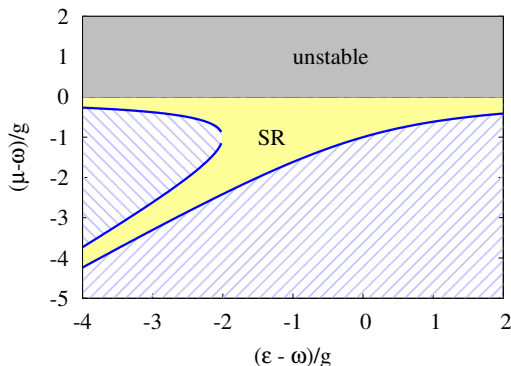
[Eastham and Littlewood, PRB '01]

- Transition at:
 $g^2 N > (\omega - \mu)|\omega_0 - \mu|$
- Reduce critical g

- Unstable if $\mu > \omega$
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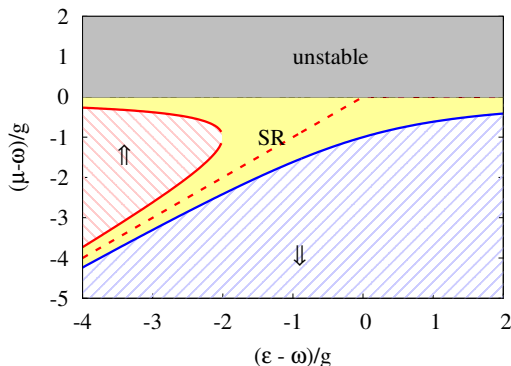


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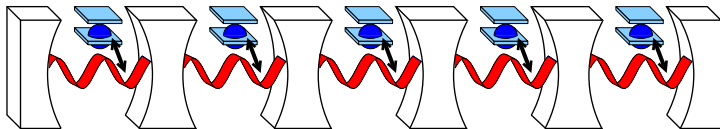
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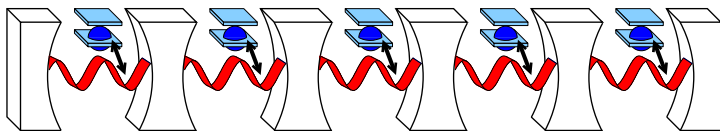
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Jaynes-Cummings Hubbard model

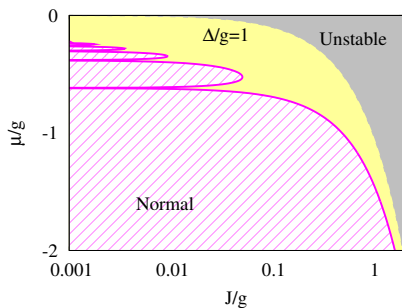


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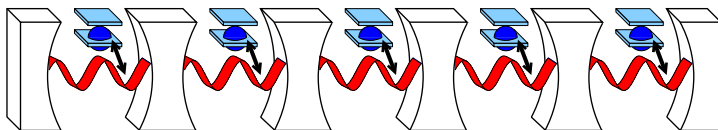
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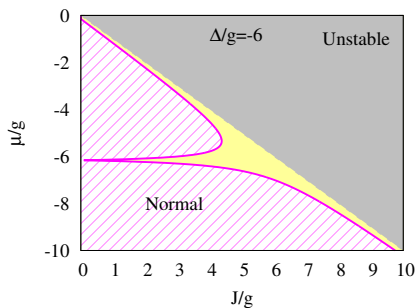
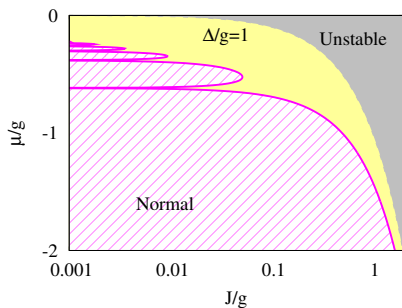
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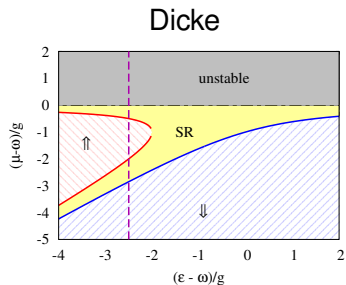
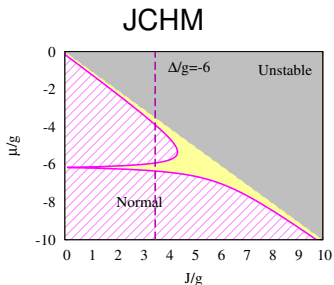
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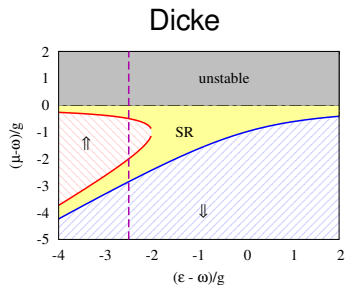
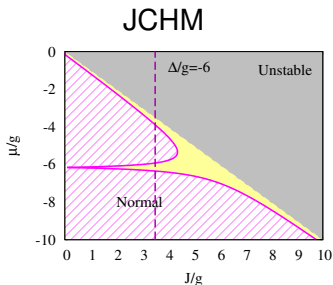
Dicke vs JCHM



- $k = 0$ mode of JCHM \leftrightarrow Dicke photon mode
- $\uparrow \leftrightarrow n = 1$ Mott lobe

[Schmidt, Blatter, JK arXiv:1306.????]

Dicke vs JCHM

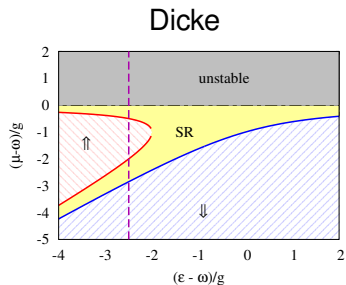
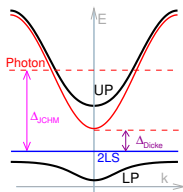
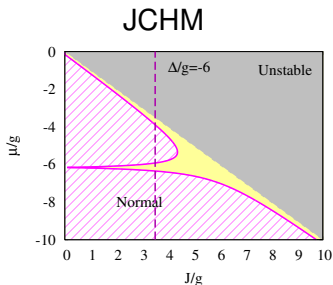


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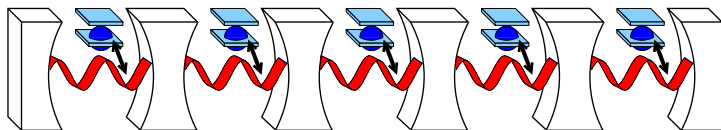
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- Disorder

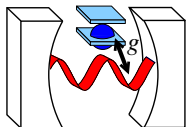
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$$\partial_t \rho = -i[H, \rho] - \frac{\kappa}{2} L_\psi[\rho] - \frac{\gamma}{2} L_{\sigma^-}[\rho]$$

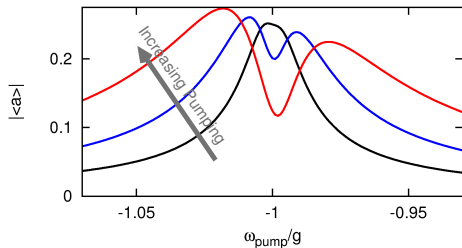
Coherently pumped single cavity [Bishop *et al.* Nat. Phys '09]



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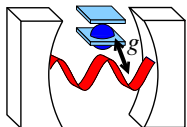
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- Anti-resonance in $|\langle \psi \rangle|$
- Effective 2LS: (Empty) (1 polariton)
- Melow triplet fluorescence



[Lang *et al.* PRL '11]

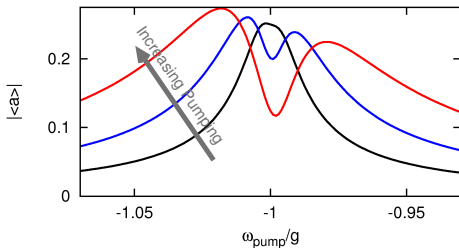
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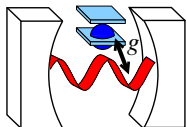
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• Many low triplet fluorescence

[Lang *et al.* PRL '11]

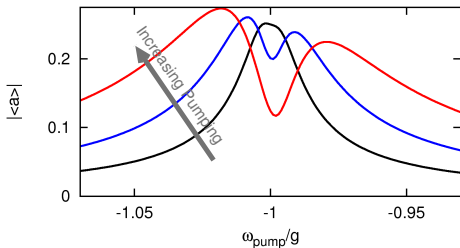
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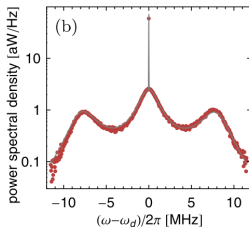
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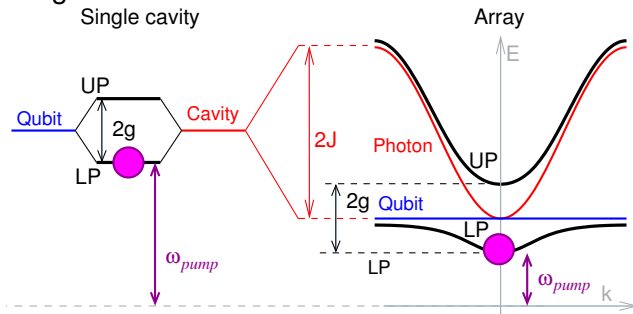
- Mollow triplet fluorescence



[Lang *et al.* PRL '11]

Coherently pumped dimer & array

Chose detuning *a la* Dicke model

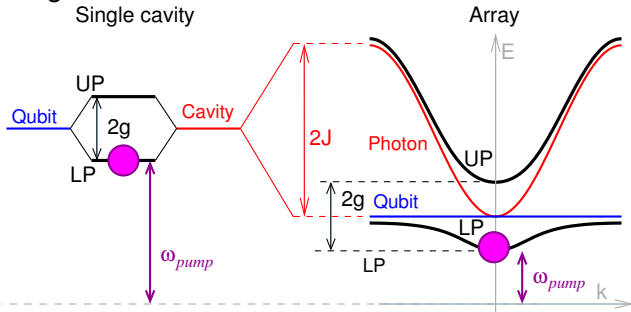


- Bistability at intermediate J
- More/less localised states
- Connects to Dicke limit

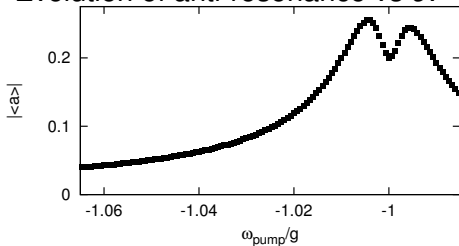
[Nissen *et al.* PRL '12]

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Evolution of anti-resonance vs J .

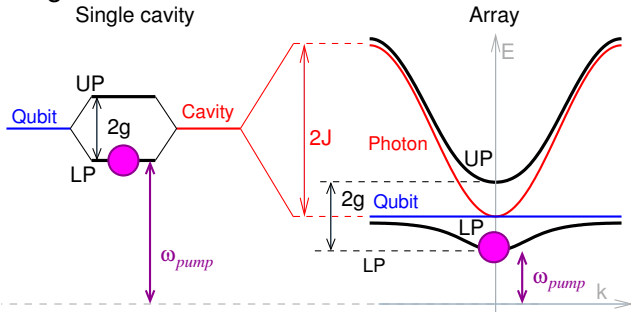


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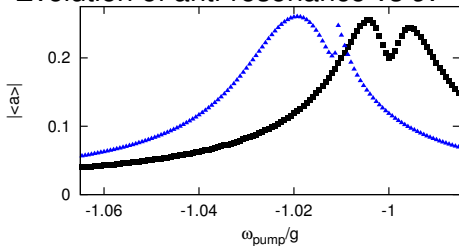
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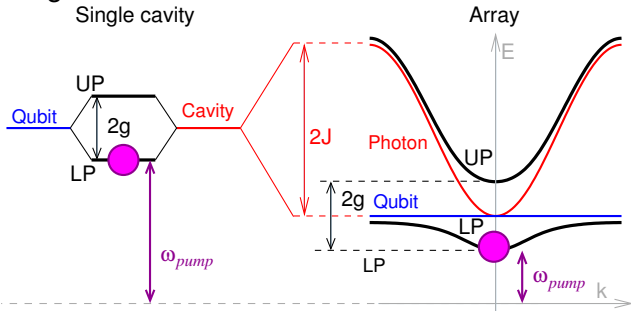


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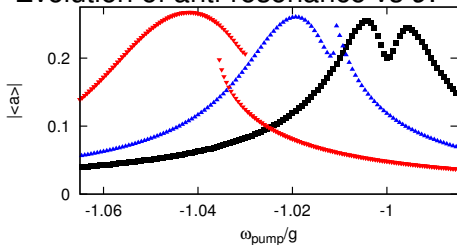
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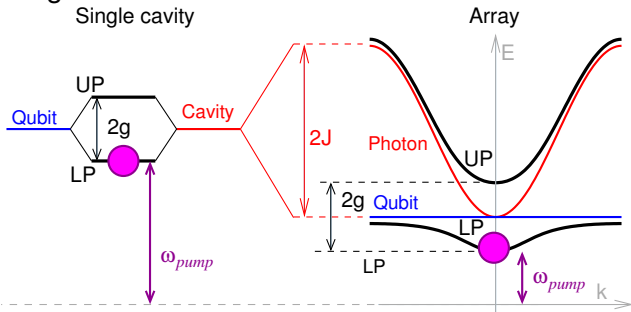


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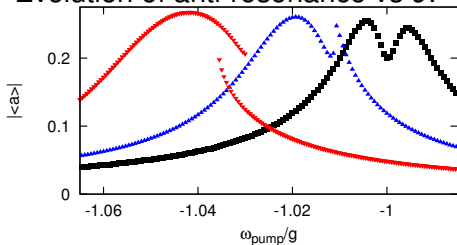
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[Nissen *et al.* PRL '12]

Photon blockade picture $J \lesssim g$

- Polariton basis
- Nonlinearity $|\epsilon_2 - 2\epsilon_1| \propto g$.

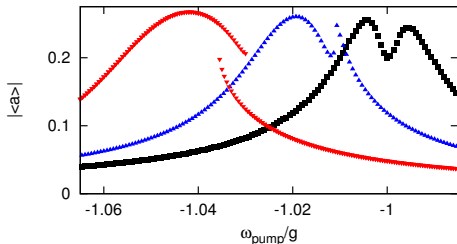
$$H = \sum_i \left(\frac{\epsilon}{2} \tau_i^z + \tilde{f} \tau_i^x \right)$$

- Decouple hopping:

$$\tau_i^+ \tau_j^- \rightarrow \psi^+ \tau_i^+ + \psi^- \tau_j^-$$

- Bistability for

$$J > J_c = \frac{4}{\bar{J}^2} \left(\frac{2\bar{J}^2 + (\bar{\kappa}/2)^2}{3} \right)^{3/2}$$

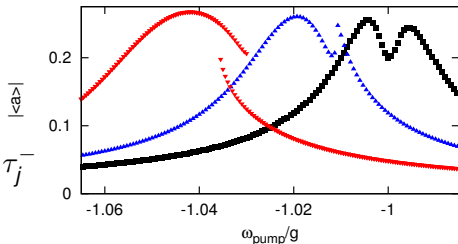


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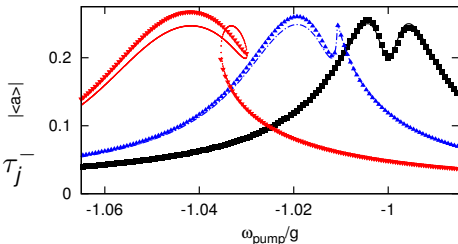
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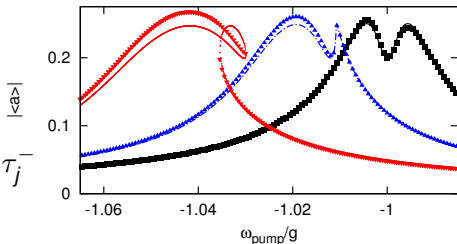
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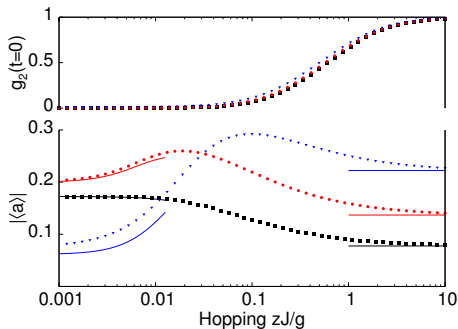


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Coherently pumped array: correlations & fluorescence

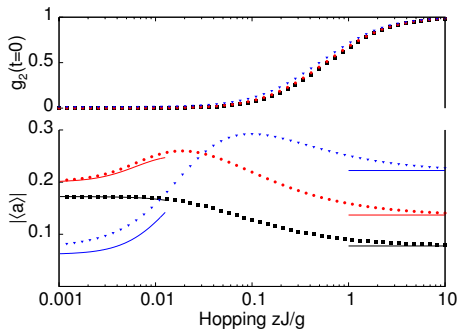


Correlations

● $g_2 : 0 \rightarrow 1$ crossover.

- Small J : Mollow triplet
- Large J : Off resonance fluorescence
- Pump at collective resonance
- Mismatch if $J \neq 0$.

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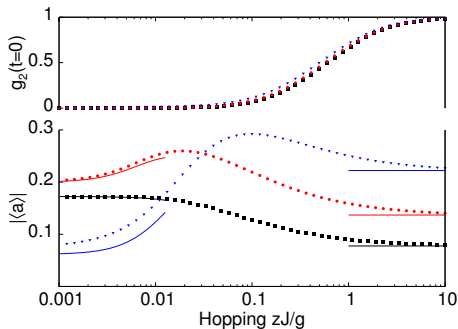


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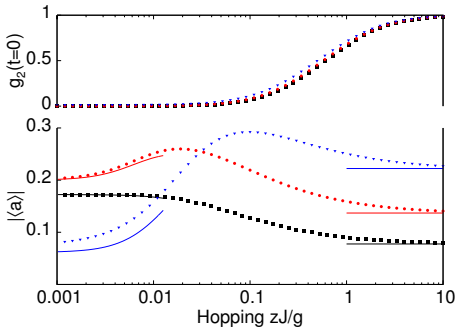
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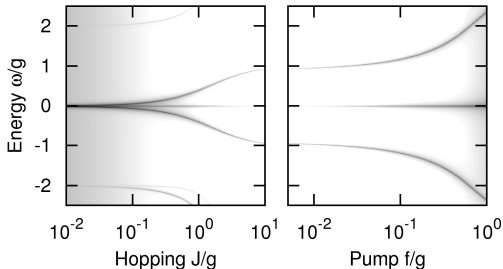


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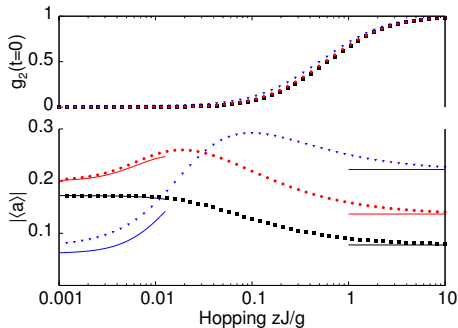
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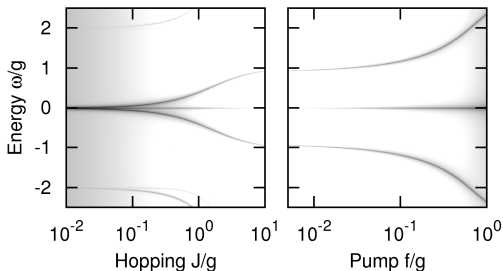
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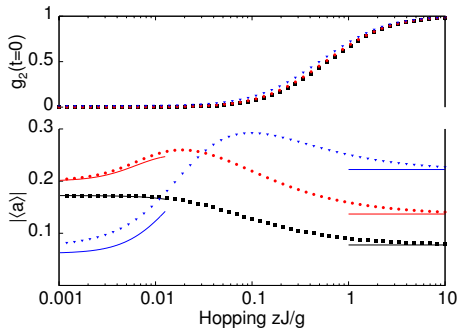
Fluorescence

- Small J : Mollow triplet

- Large J : Off resonance fluorescence
- Pump at collective resonance
- Mismatch if $J \neq 0$.



Coherently pumped array: correlations & fluorescence

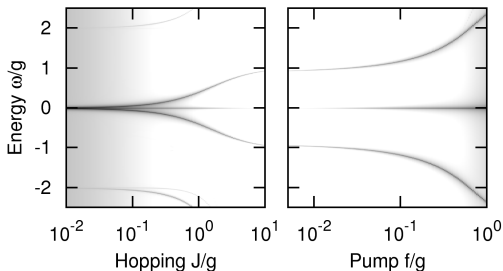


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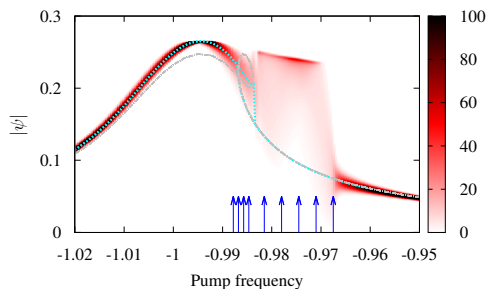
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Coherent pumped array – disorder

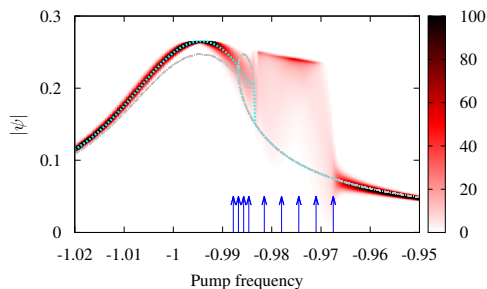
- Effect of disorder, $\Delta \rightarrow \Delta_j$
 - ▶ Distribution of ψ – Washes out bistable jump
 - ◀ Bistability near resonance — phase of ψ depends on Δ_j
 - ◀ Complex ψ distribution
 - ◀ Superfluid phases in driven system?



[Kulaitis *et al.* PRA, '13]

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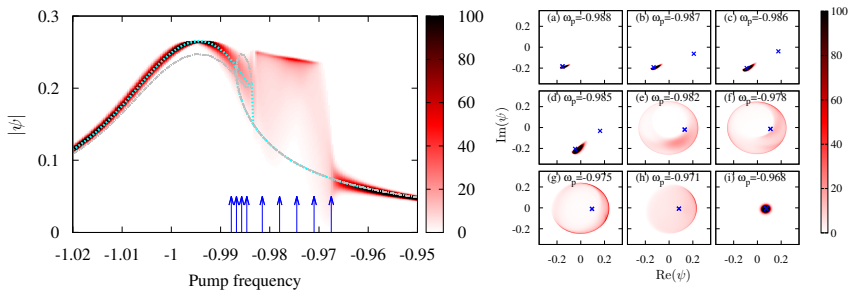


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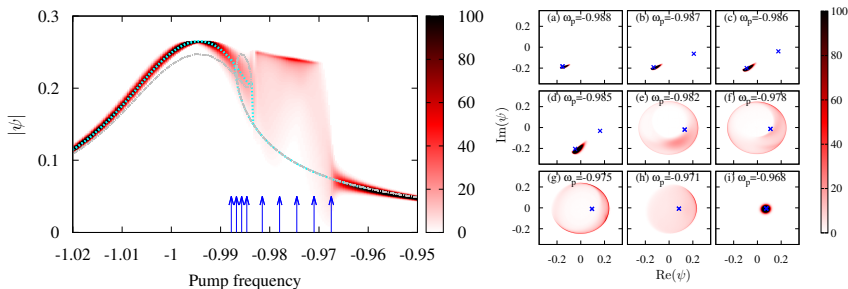
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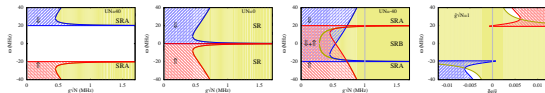
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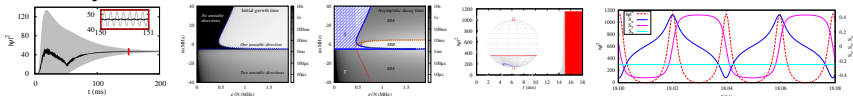
[Kulaitis *et al.* PRA, '13]

Summary

- Wide variety of dynamical phases

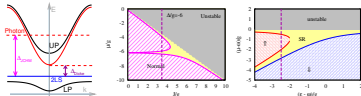


- Slow dynamics for $U < 0$ & Persistent oscillations for $U > 0$

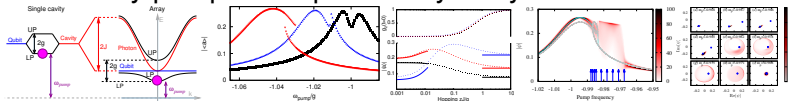


JK *et al.* PRL '10, Bhaseen *et al.* PRA '12

- Dicke model and JCHM: connection at $J \rightarrow \infty$



- Coherently pumped coupled cavity array



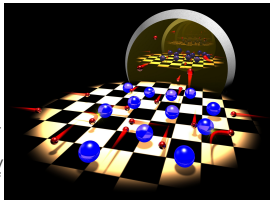
Nissen *et al.* PRL '12, Kulaitis *et al.* PRA '13

Many body quantum optics and correlated states of light

9:00 am on Monday 28 October 2013 – 5:00 pm on Tuesday 29 October 2013

at: **The Royal Society at Chicheley Hall, home of the Kavli Royal Society International Centre, Buckinghamshire**

Theo Murphy international scientific meeting organised by Dr Jonathan Keeling, Professor Steven Girvin, Dr Michael Hartmann and Professor Peter Littlewood FRS.



List of speakers and chairs

Professor Iacopo Carusotto, Professor Andrew Cleland, Professor Hui Deng, Professor Tilman Esslinger, Professor Rosario Fazio, Professor Ed Hinds, Professor Andrew Houck, Professor Ataç İmamoğlu, Professor Jens Koch, Professor Misha Lukin, Professor Martin Plenio, Professor Arno Rauschenbeutel, Professor Timothy Spiller, Professor Jacob Taylor, Professor Hakan Türeci, Professor Andreas Wallraff

Attending this event

This is a residential conference which allows for increased discussion and networking. It is free to attend, however participants need to cover their accommodation and catering costs if required. Places are limited and therefore pre-registration is essential.

- 4 Ferroelectric transition
- 5 Dicke vs JCHM
- 6 Pumping without symmetry breaking
- 7 Collective dephasing

Ferroelectric transition

Atoms in Coulomb gauge

$$H = \sum_k \omega_k a_k^\dagger a_k + \sum_i [p_i - eA(r_i)]^2 + V_{\text{coul}}$$

Ferroelectric transition

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Two-level systems — dipole-dipole coupling

$$H = \omega_0 S^z + \omega \psi^\dagger \psi + g(S^+ + S^-)(\psi + \psi^\dagger) + N\zeta(\psi + \psi^\dagger)^2 - \eta(S^+ - S^-)^2$$

(nb $g^2, \zeta, \eta \propto 1/V$).

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Gauge transform to dipole gauge $\mathbf{D} \cdot \mathbf{r}$

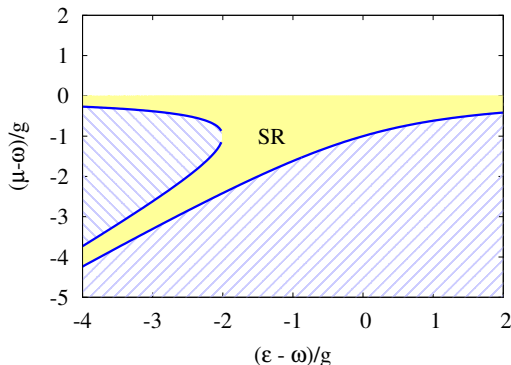
$$H = \omega_0 S^z + \omega \psi^\dagger \psi + \bar{g}(S^+ - S^-)(\psi - \psi^\dagger)$$

“Dicke” transition at $\omega_0 < N\bar{g}^2/\omega \equiv 2\eta N$

But, ψ describes **electric displacement**

Equilibrium: Dicke model with chemical potential

$$H - \mu N = (\omega - \mu)\psi^\dagger\psi + (\omega_0 - \mu)S^z + g(\psi^\dagger S^- + \psi S^+)$$



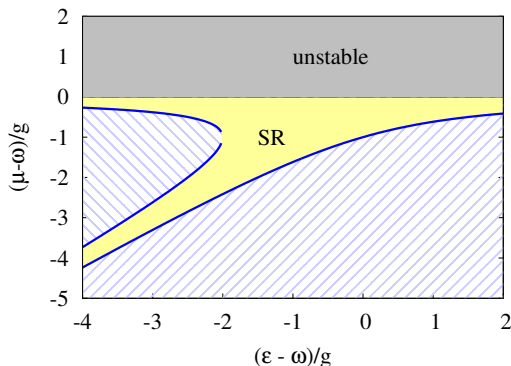
[Eastham and Littlewood, PRB '01]

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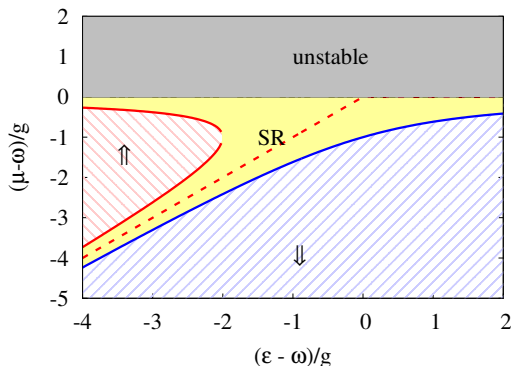


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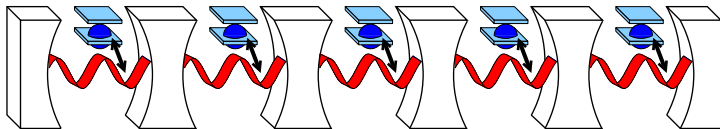
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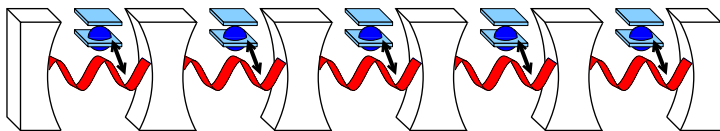
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Jaynes-Cummings Hubbard model

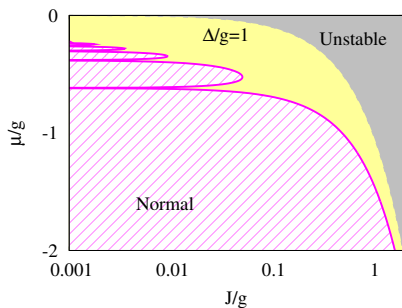


$$H = -\frac{J}{Z} \sum_{ij} \psi_i^\dagger \psi_j + \sum_i \frac{\Delta}{2} \sigma_i^z + g(\psi_i^\dagger \sigma_i^- + \text{H.c.})$$

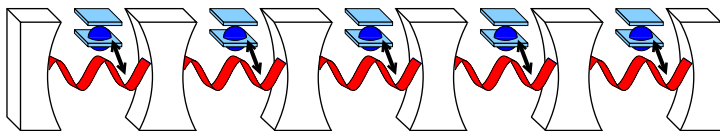
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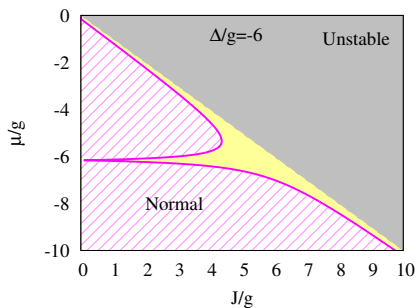
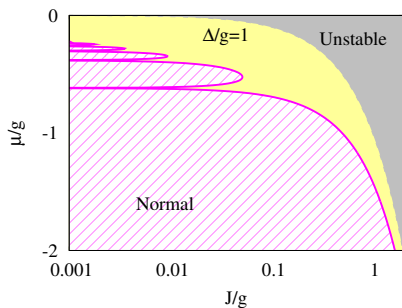
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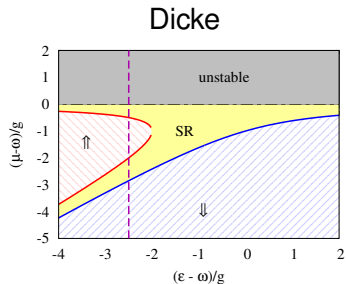
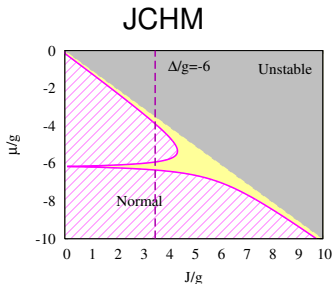
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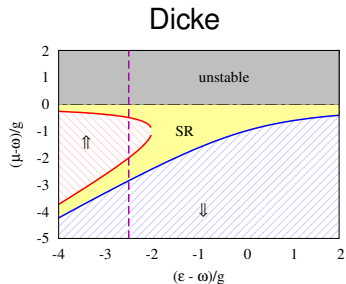
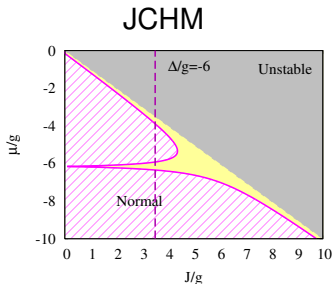
Dicke vs JCHM



- $k = 0$ mode of JCHM \leftrightarrow Dicke photon mode
- $\uparrow\uparrow \leftrightarrow n = 1$ Mott lobe

[Schmidt, Blatter, JK arXiv:1306.????]

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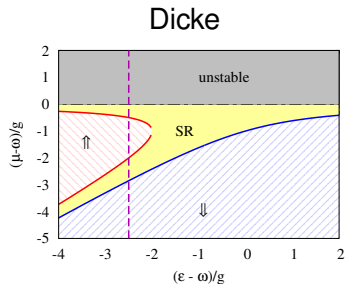
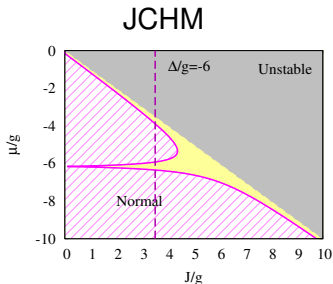


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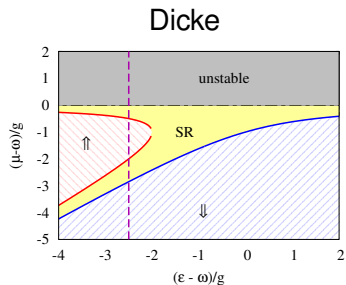
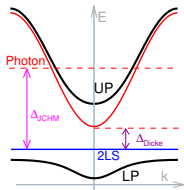
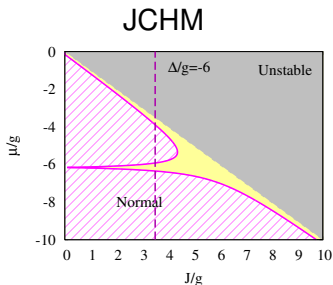
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- How to pump without breaking symmetry
- Counter-rotating terms — Raman pumping
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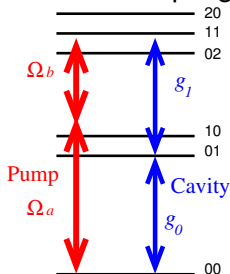
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Tunable-coupling-qubit



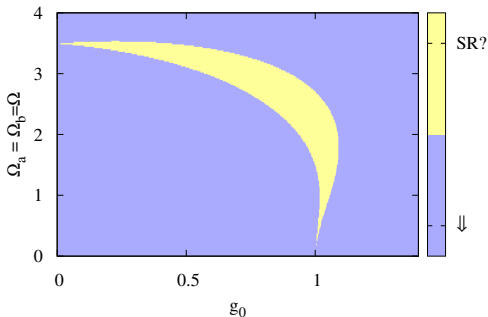
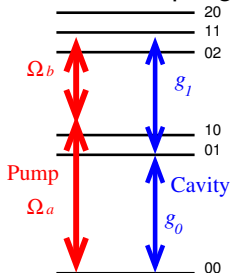
JK, Türeci, Houck in progress

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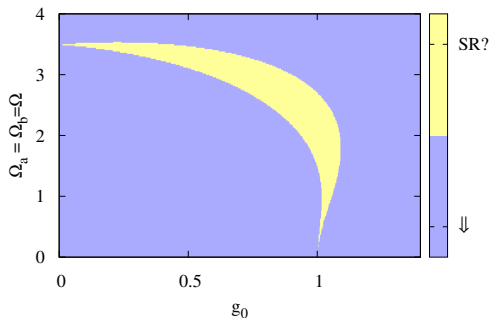
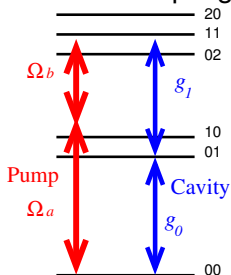


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JK, Türeci, Houck in progress

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