

Non-equilibrium phases of coupled matter-light systems

Jonathan Keeling



University of
St Andrews

600
YEARS



Oxford, May 2012

Coupling many atoms to light

Old question: *What happens to radiation when many atoms interact “collectively” with light.*

Superradiance — dynamical and steady state.

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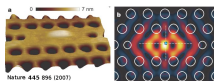
Superradiance — dynamical and steady state.

New relevance

- Superconducting qubits



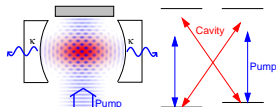
- Quantum dots



- Nitrogen-vacancies in diamond



- Ultra-cold atoms



- Rydberg atoms

Dicke effect: Enhanced emission

PHYSICAL REVIEW

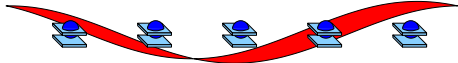
VOLUME 93, NUMBER 1

JANUARY 1, 1954

Coherence in Spontaneous Radiation Processes

R. H. DICKE

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey



$$H_{\text{int}} = \sum_{k,i} g_k \left(\psi_k^\dagger S_i^- e^{-ik \cdot r_i} + \text{H.c.} \right)$$

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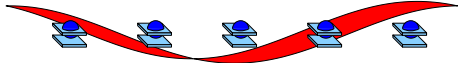
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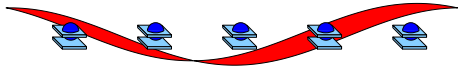
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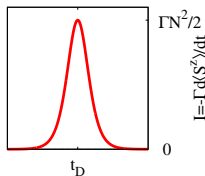
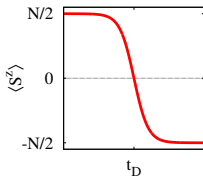
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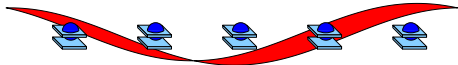
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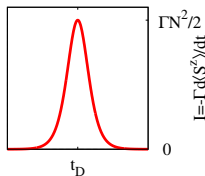
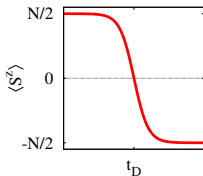
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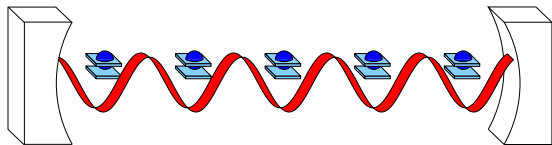
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Problem: dipole interactions dephase. [Friedberg et al, Phys. Lett. 1972]

Collective radiation **with a cavity**: Dynamics

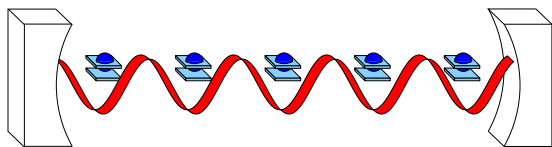


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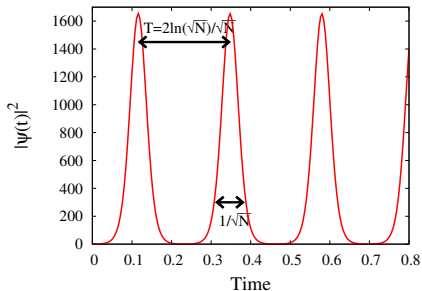
[Bonifacio and Preparata PRA '70; JK PRA '09]

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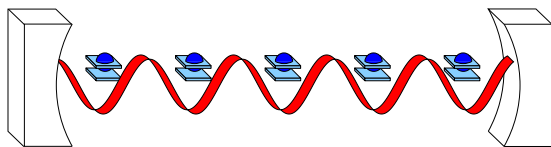
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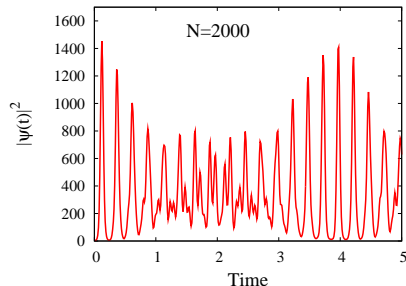
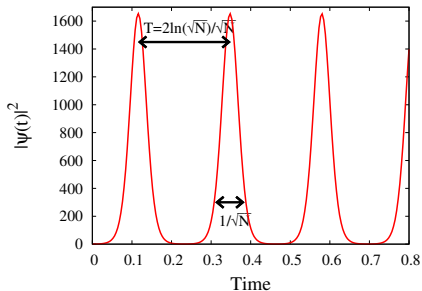
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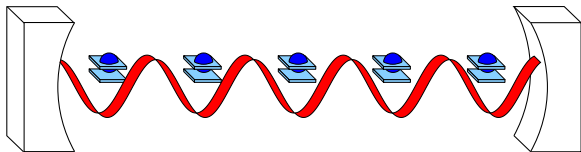
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Dicke model: Equilibrium superradiance transition



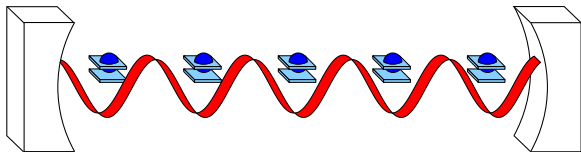
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• Coherent state: $|\psi\rangle \rightarrow e^{\lambda \psi^\dagger + \eta S^+} |\Omega\rangle$

• Small g , min at $\lambda, \eta = 0$

[Hepp, Lieb, Ann. Phys. '73]

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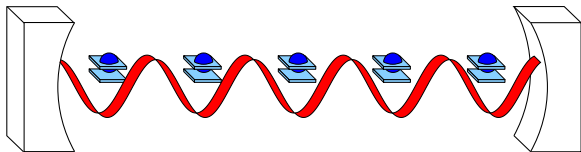
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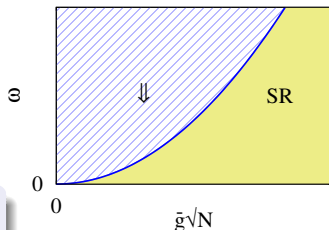
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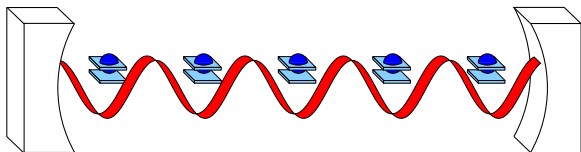
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Spontaneous polarisation if: $Ng^2 > \omega\omega_0$



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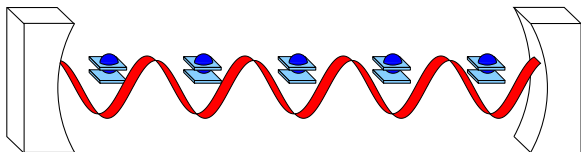
No go theorem and transition



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[Rzazewski *et al* PRL '75]

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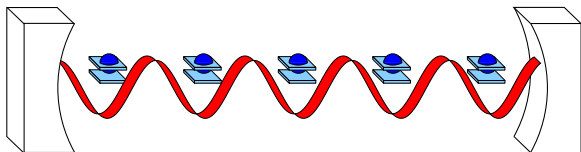
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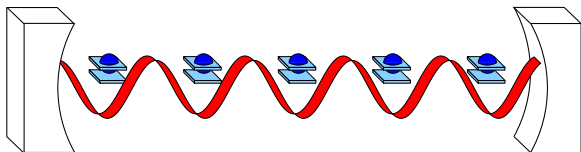
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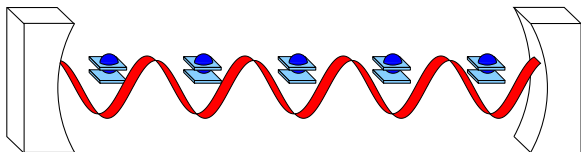
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But Thomas-Reiche-Kuhn sum rule states: $g^2/\omega_0 < 2\zeta$. **No transition**

[Rzazewski *et al* PRL '75]

Dicke phase transition: ways out

Problem: $g^2/\omega_0 < 2\zeta$ for intrinsic parameters. **Solutions:**

- Interpretation
 - Ferroelectric transition in $\mathbf{D} \cdot \mathbf{r}$ gauge. [JK-JPGM '07]
 - Circuit QED [Nataf and Cluzet, Nat. Comm. '10; Viehmann et al. PRL '11]
- Grand canonical ensemble:
 - If $H \rightarrow H - \mu(S^z + \psi^\dagger\psi)$, need only:
 $g^2 N > (\omega - \mu)(\omega_0 - \mu)$
 - Incoherent pumping \rightarrow polariton condensation.
- Dissociate g, ω_0 ,
 - e.g. Raman scheme: $\omega_0 \ll \omega$. [Dimer et al. PRA '07; Baumann et al. Nature '10. Also, Black et al. PRL '03]

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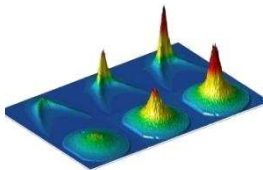
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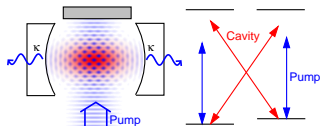
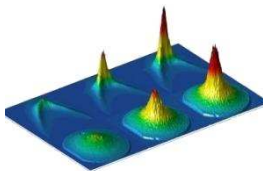
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Outline

1 Introduction: Dicke model and superradiance

2 Dynamics of generalized Dicke model

- Summary of experiment and classical dynamics
- Fixed points and dynamical phases
- Timescales and consequences for experiment
- Persistent oscillating phases

3 Non-equilibrium states of Jaynes-Cummings-Hubbard Model

- Relating equilibrium JCHM & Dicke model
- Coherently pumped JCHM

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Part 1:



J. Mayoh



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Part 2:



F. Nissen



G. Blatter



M. Biondi



S. Schmidt

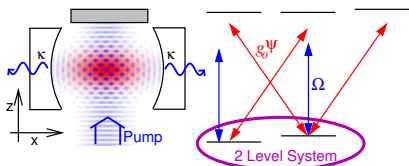


H. Türeci

EPSRC

Funding: Engineering and Physical Sciences
Research Council

Reminder of cold-atom extended Dicke model



2 Level system, $|\downarrow\rangle, |\uparrow\rangle$:

$$\downarrow: \Psi(x, z) = 1$$

$$\uparrow: \Psi(x, z) = \sum_{\sigma, \sigma' = \pm} e^{ik(\sigma x + \sigma' z)}$$

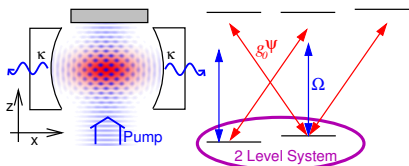
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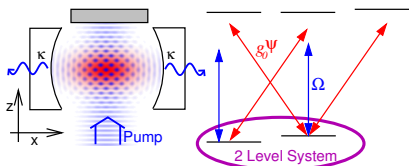
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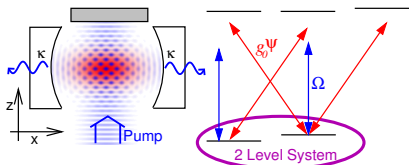
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[Baumann *et al* Nature '10]



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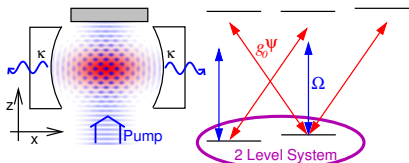
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$$\omega_0 \sim \text{kHz} \ll \omega, \kappa, g\sqrt{N} \sim \text{MHz}.$$

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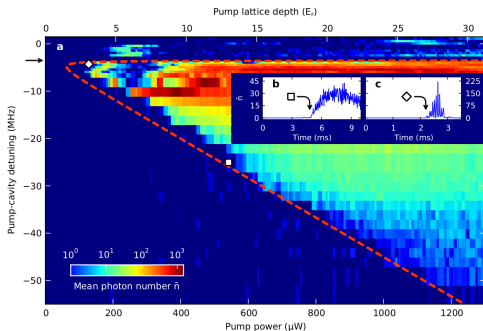
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Classical dynamics of the extended Dicke model

Open dynamical system:

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- Neglects quantum fluctuations — restore via Wigner distributed initial conditions.
- Linearisation about fixed point:
 - Recover Retarded Green's function (spectrum)
 - Cannot recover occupations

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Classical EOM
($|\mathbf{S}| = N/2 \gg 1$)

$$\dot{\mathbf{S}}^- = -i(\omega_0 + U|\psi|^2)\mathbf{S}^- + 2ig(\psi + \psi^*)\mathbf{S}^Z$$
$$\dot{\mathbf{S}}^Z = ig(\psi + \psi^*)(\mathbf{S}^- - \mathbf{S}^+)$$
$$\dot{\psi} = -[\kappa + i(\omega + U\mathbf{S}^Z)]\psi - ig(\mathbf{S}^- + \mathbf{S}^+)$$

- Neglects quantum fluctuations — restore via Wigner distributed initial conditions.
- Linearisation about fixed point:
 - Recover Retarded Green's function (spectrum)
 - Cannot recover occupations

Classical dynamics of the extended Dicke model

Open dynamical system:

$$H = \omega\psi^\dagger\psi + \omega_0\mathbf{S}^z + g(\psi + \psi^\dagger)(\mathbf{S}^- + \mathbf{S}^+) + U\mathbf{S}_z\psi^\dagger\psi.$$
$$\partial_t\rho = -i[H, \rho] - \kappa(\psi^\dagger\psi\rho - 2\psi\rho\psi^\dagger + \rho\psi^\dagger\psi)$$

Classical EOM
($|\mathbf{S}| = N/2 \gg 1$)

$$\dot{\mathbf{S}}^- = -i(\omega_0 + U|\psi|^2)\mathbf{S}^- + 2ig(\psi + \psi^*)\mathbf{S}^z$$
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Fixed points (steady states)

$$0 = i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z$$

$$0 = ig(\psi + \psi^*)(S^- - S^+)$$

$$0 = -[\kappa + i(\omega + US^z)]\psi - ig(S^- + S^+)$$

• $\psi = 0, S = (0, 0, \pm N/2)$
always a solution.

• If $g > g_c, \psi \neq 0$ too

A. $S^z = -S[S^-] = 0$

B. $\psi = \Re[\psi] = 0$

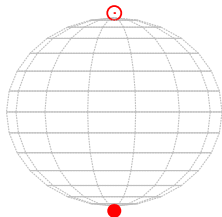
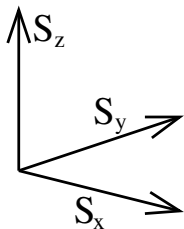
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Small g : \uparrow, \downarrow only.
($\omega = 30\text{MHz}$, $UN = -40\text{MHz}$)

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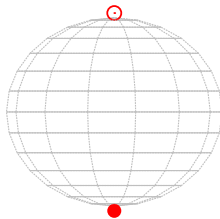
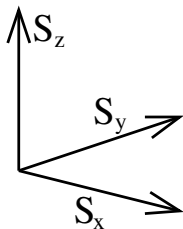
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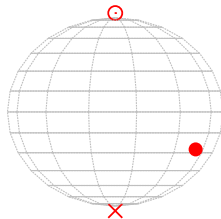
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A $S^y = -\Im[S^-] = 0$

B $\psi' = \Re[\psi] = 0$



Small g : \uparrow, \downarrow only.
 $(\omega = 30\text{MHz}, UN = -40\text{MHz})$



Larger g : SR too.

Steady state phase diagram

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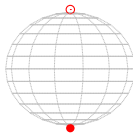
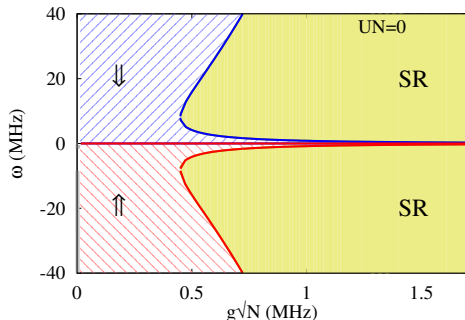
See also Domokos and Ritsch PRL '02, Domokos *et al.* PRL '10

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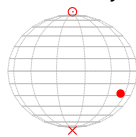
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SR(A): $S_y = 0$



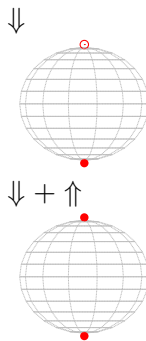
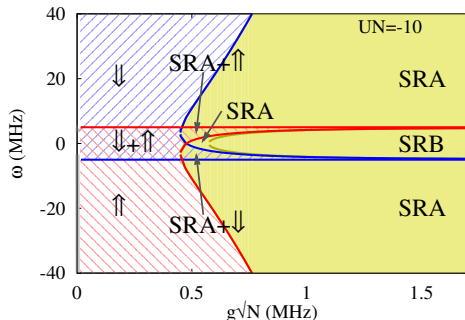
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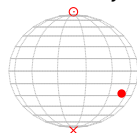
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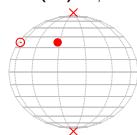
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SR(A): $S_y = 0$



SR(B): $\psi' = 0$



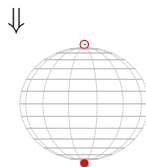
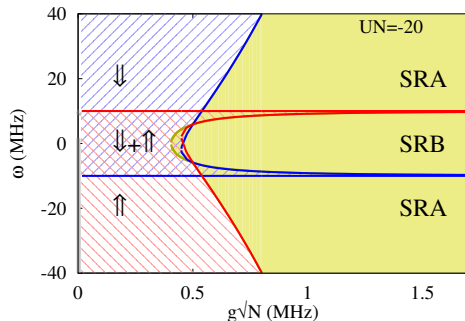
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Steady state phase diagram

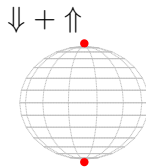
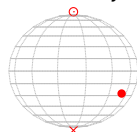
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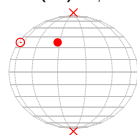
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SR(A): $S_y = 0$



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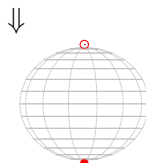
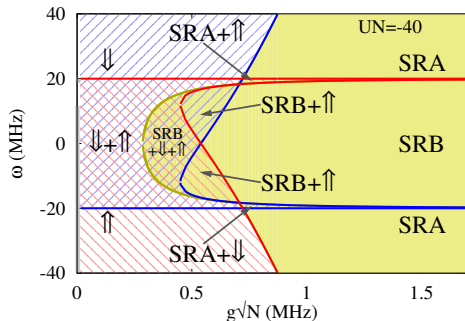
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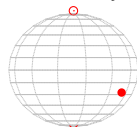
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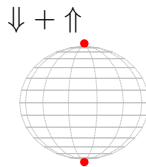
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SR(A): $S_y = 0$

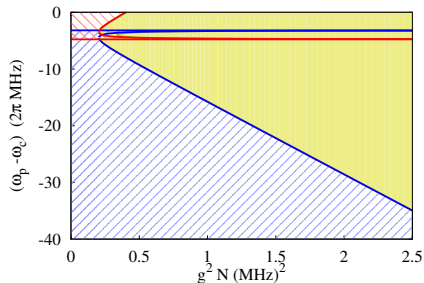


SR(B): $\psi' = 0$



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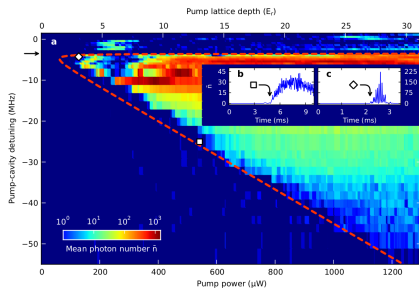
Comparison to experiment



$$UN = -10\text{MHz}$$

Adapted from: [Bhaseen *et al.* PRA '12]

$$\omega = \omega_c - \omega_p + \frac{5}{2}UN,$$



[Baumann *et al* Nature '10]

$$UN = -\frac{g_0^2}{4(\omega_a - \omega_c)}$$

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2 Dynamics of generalized Dicke model

- Summary of experiment and classical dynamics
- Fixed points and dynamical phases
- **Timescales and consequences for experiment**
- Persistent oscillating phases

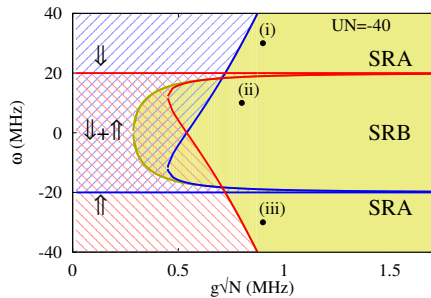
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Dynamics: Evolution from normal state

Gray: $\mathbf{S} = (\sqrt{N}, \sqrt{N}, -N/2)$

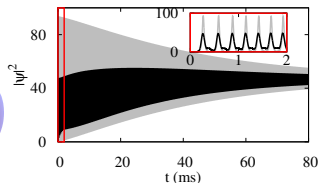
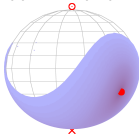
Black: Wigner distribution of \mathbf{S}, ψ



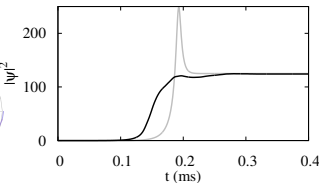
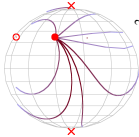
Oscillations: ~ 0.1 ms

Decay: 20ms, 0.1ms, 20ms

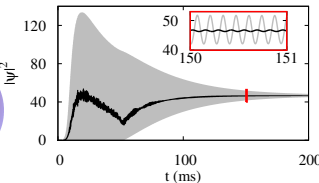
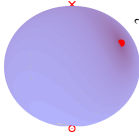
(i) SR(A)



(ii) SR(B)



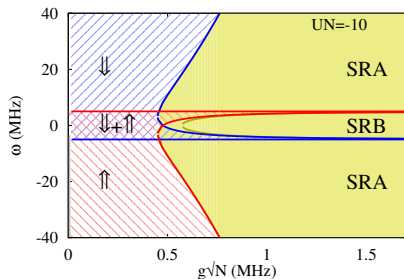
(iii) SR(A)



Asymptotic state: Evolution from normal state

(Near to experimental $UN = -13\text{MHz}$).

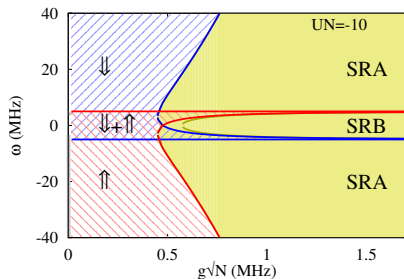
All stable attractors:



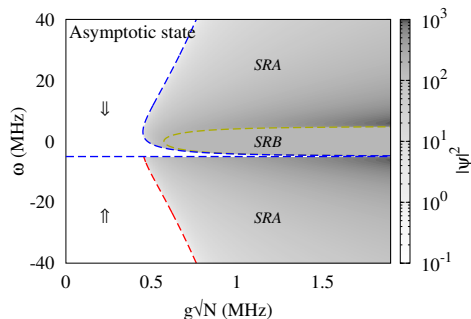
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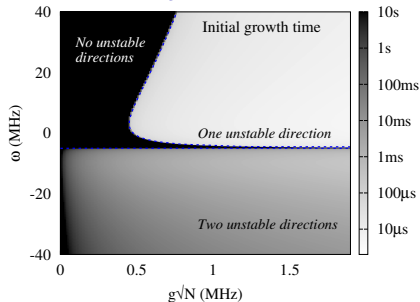
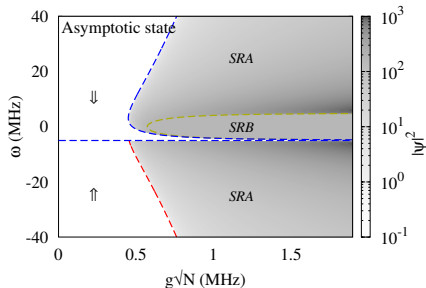
All stable attractors:



Starting from \Downarrow



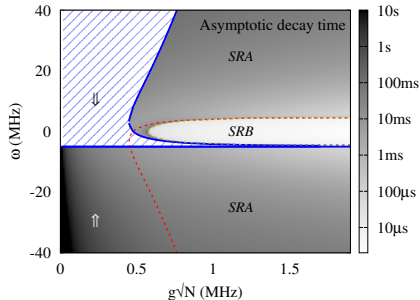
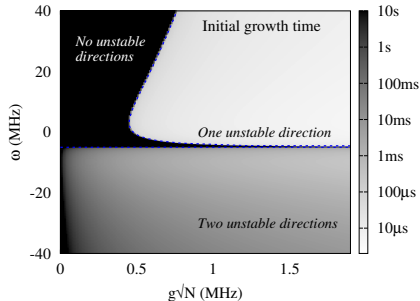
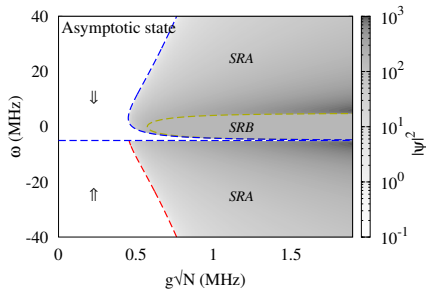
Timescales for dynamics: What are they?



Growth Most unstable eigenvalues near $\mathbf{S} = (0, 0, -N/2)$

Decay Slowest stable eigenvalues near final state

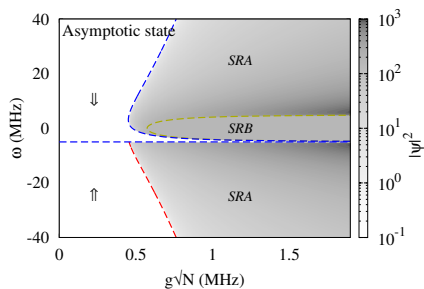
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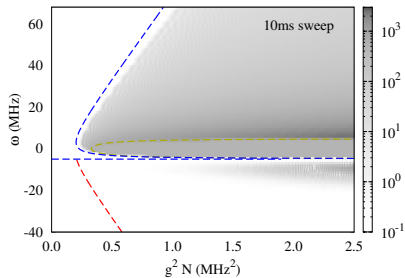
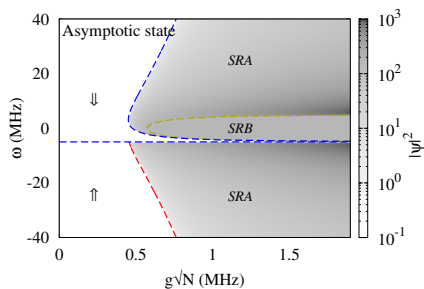
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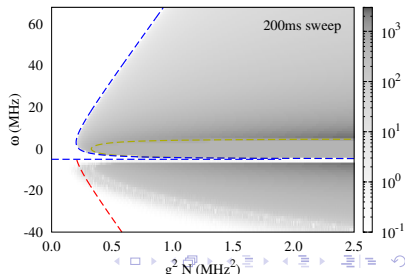
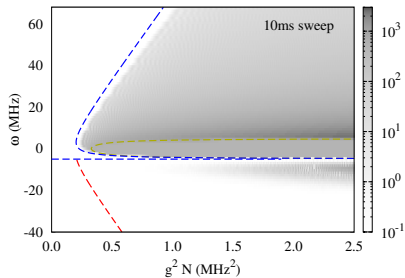
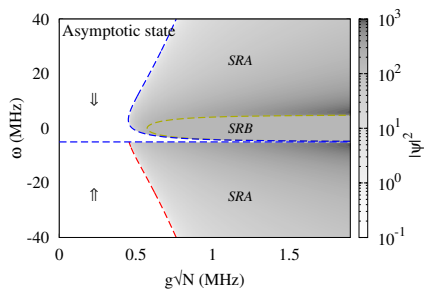
Timescales for dynamics: Consequences for experiment



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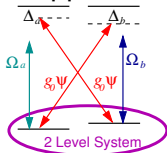


Timescales for dynamics: Consequences for experiment



Timescales for dynamics: Why so slow and varied?

Suppose co- and counter-rotating terms differ

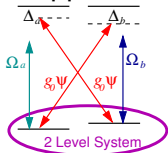


$$H = \dots + g(\psi^\dagger S^- + \psi S^+) + g'(\psi^\dagger S^+ + \psi S^-) + \dots$$

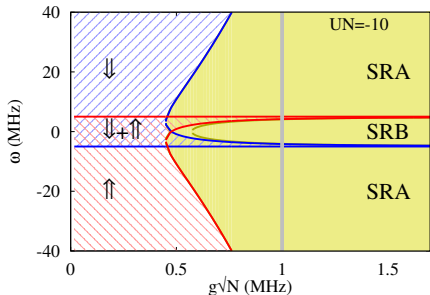
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- SR(A), SR(B) continuously connect

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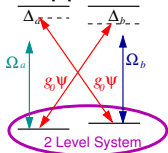
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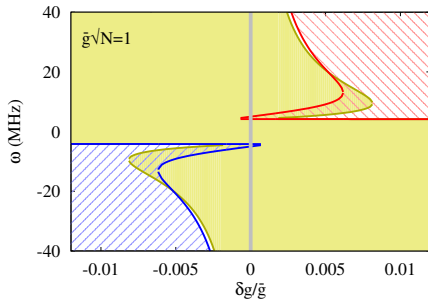
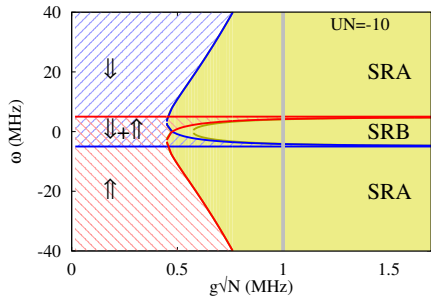
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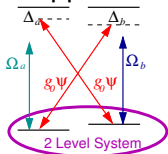
$$\delta g = g' - g, \quad 2\bar{g} = g' + g$$



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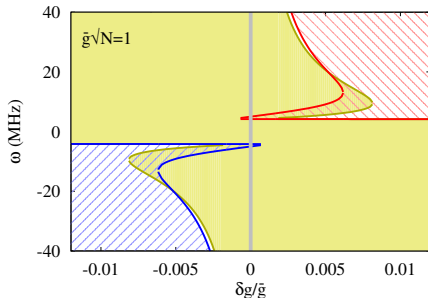
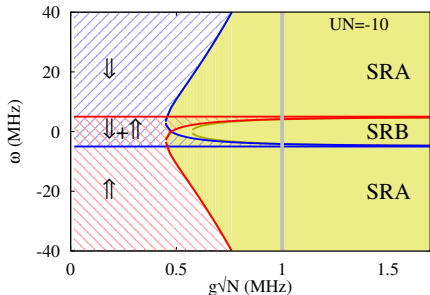
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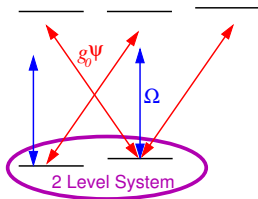
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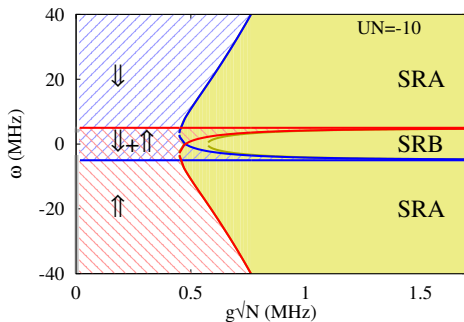
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Regions without fixed points

Changing U :

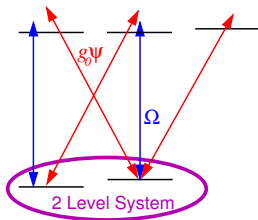


$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$

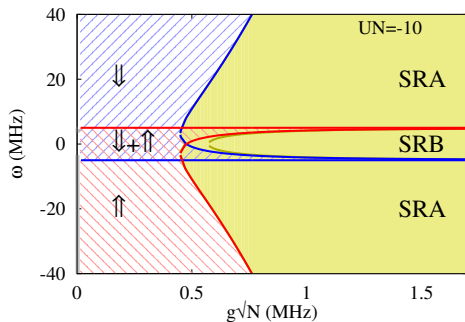


Regions without fixed points

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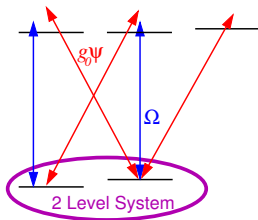


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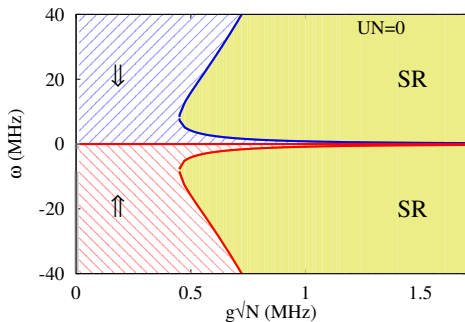


Regions without fixed points

Changing U :

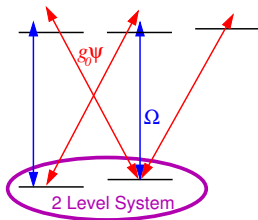


$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$

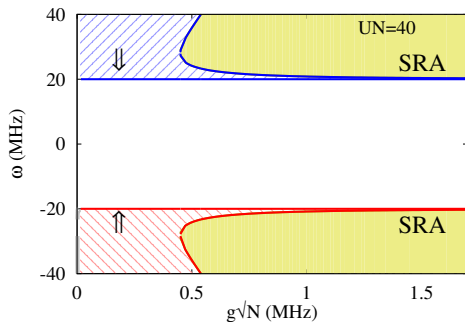


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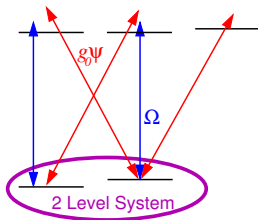


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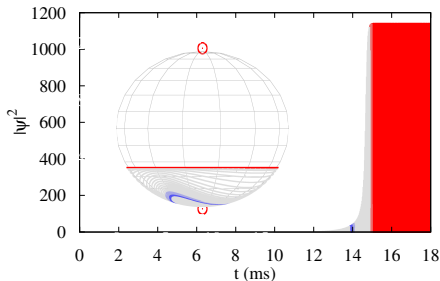
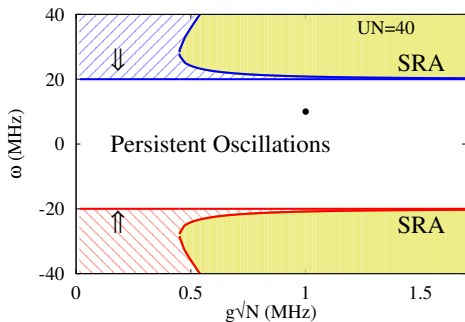


Regions without fixed points

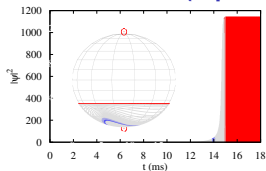
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Persistent (optomechanical) oscillations

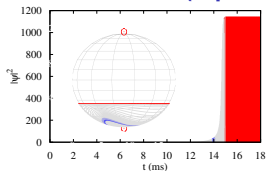


$$\dot{S}^- = -i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^Z$$

$$\dot{S}^Z = ig(\psi + \psi^*)(S^- - S^+)$$

$$\dot{\psi} = -[\kappa + i(\omega + US^Z)]\psi - ig(S^- + S^+)$$

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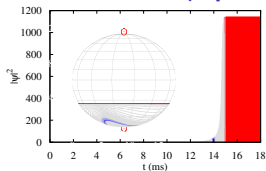
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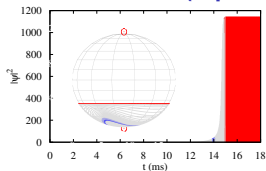
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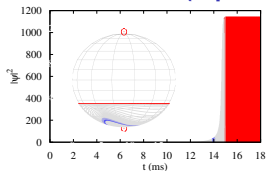
Fix $\omega + US^z = 0$ if $\psi' = 0$.

$$S^- = re^{-i\theta} \quad r = \sqrt{\frac{N^2}{4} - (S^z)^2}$$

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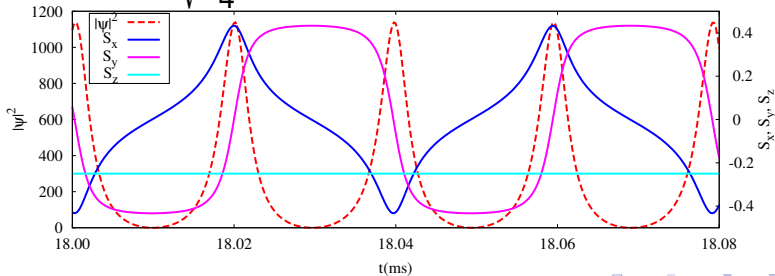
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Outline

1 Introduction: Dicke model and superradiance

2 Dynamics of generalized Dicke model

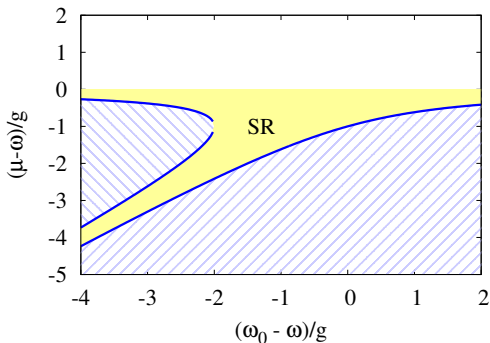
- Summary of experiment and classical dynamics
- Fixed points and dynamical phases
- Timescales and consequences for experiment
- Persistent oscillating phases

3 Non-equilibrium states of Jaynes-Cummings-Hubbard Model

- Relating equilibrium JCHM & Dicke model
- Coherently pumped JCHM

Equilibrium: Dicke model with chemical potential

$$H - \mu N = (\omega - \mu)\psi^\dagger\psi + (\omega_0 - \mu)S^z + g(\psi^\dagger S^- + \psi S^+)$$



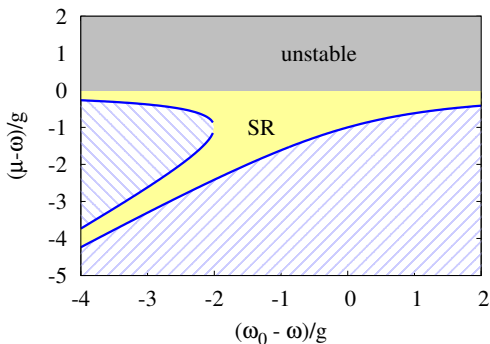
- Transition at:
 $g^2 N > (\omega - \mu)(\omega_0 - \mu)$
- Reduce critical g

- Unstable if $\mu > \omega$
- Inverted if $\mu > \omega_0$

[Eastham and Littlewood, PRB '01]

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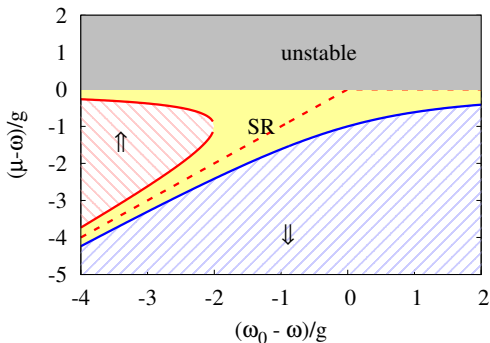


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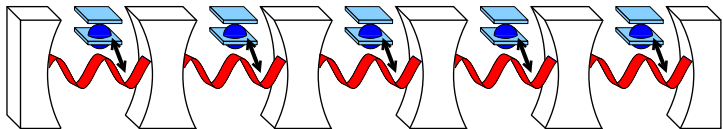
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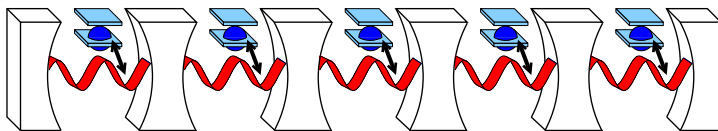
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Jaynes-Cummings Hubbard model

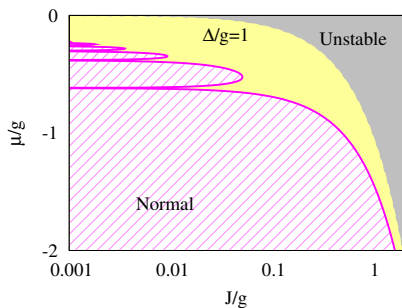


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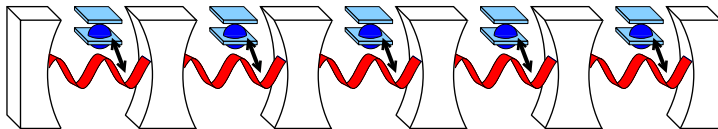
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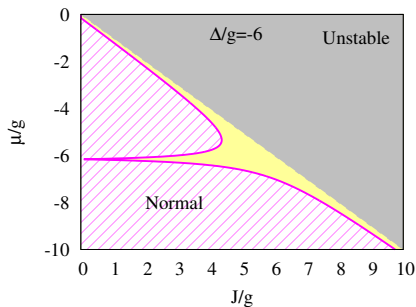
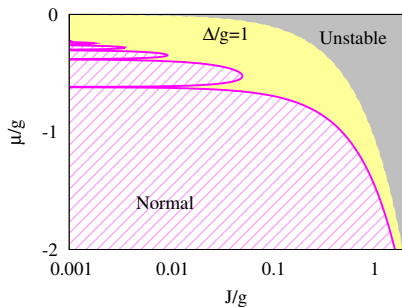
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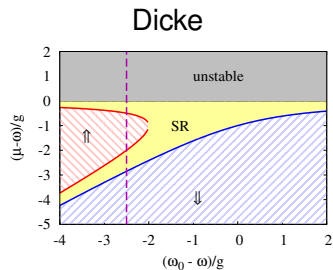
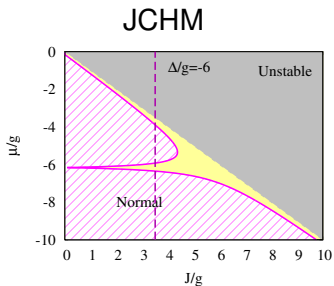
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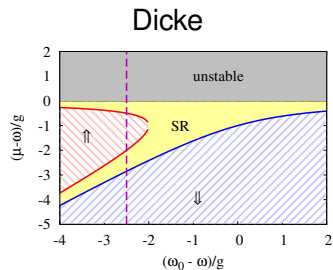
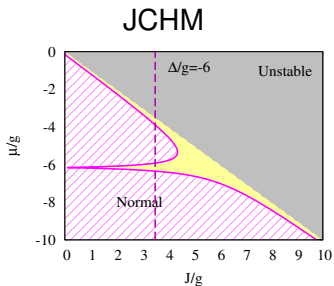


Dicke vs JCHM



- $k = 0$ mode of JCHM \leftrightarrow Dicke photon mode
- $\uparrow \leftrightarrow n = 1$ Mott lobe

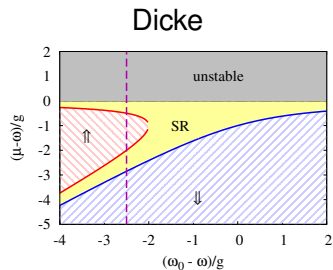
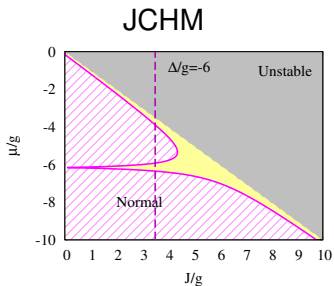
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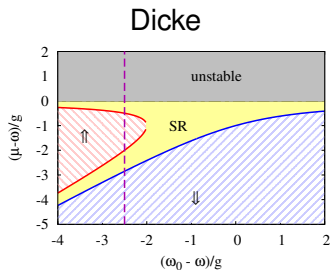
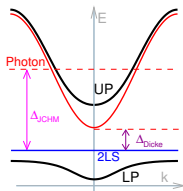
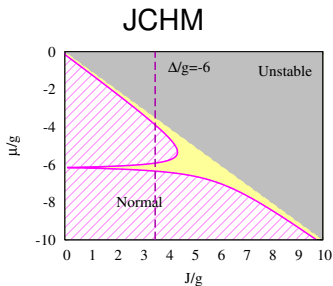
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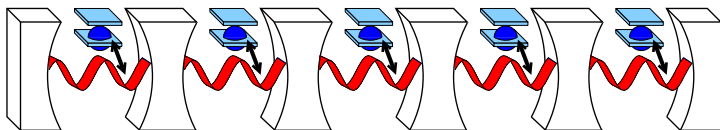
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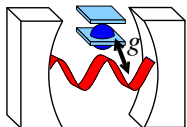
Coherently pumped JCHM



$$H = -\frac{J}{z} \sum_{ij} \psi_i^\dagger \psi_j + \sum_i \frac{\Delta}{2} \sigma_i^z + g(\psi_i^\dagger \sigma_i^- + \text{H.c.}) + f(\psi_i e^{i\omega_L t} + \text{H.c.})$$

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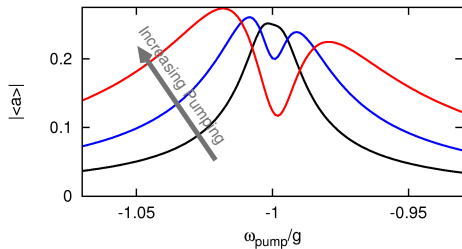
Coherently pumped single cavity [Bishop *et al.* Nat. Phys '09]



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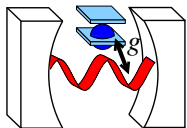
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- Anti-resonance in $|\langle \psi \rangle|$
- Effective 2LS: (Empty) (1 polariton)
- Motzky triplet fluorescence



[Lang *et al.* PRL '11]

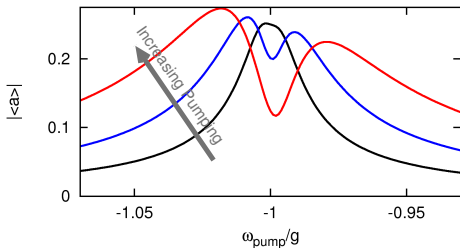
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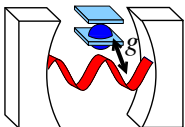
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• Macro triplet fluorescence

[Lang *et al.* PRL '11]

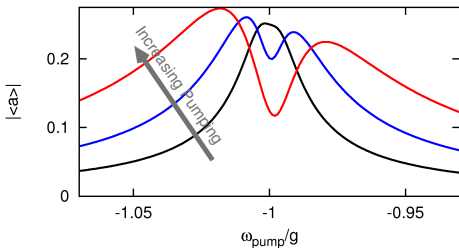
Coherently pumped single cavity [Bishop *et al.* Nat. Phys '09]



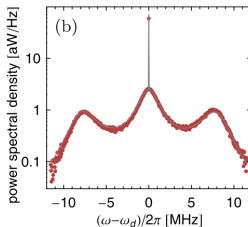
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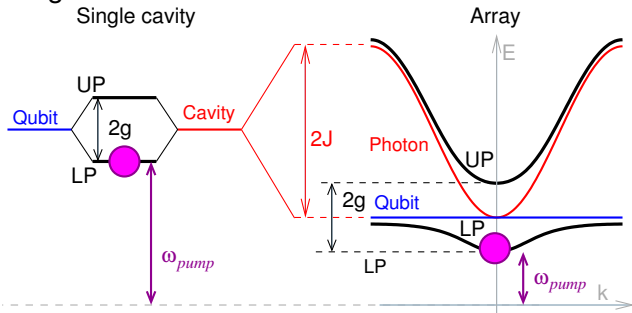
- Mollow triplet fluorescence



[Lang *et al.* PRL '11]

Coherently pumped dimer & array

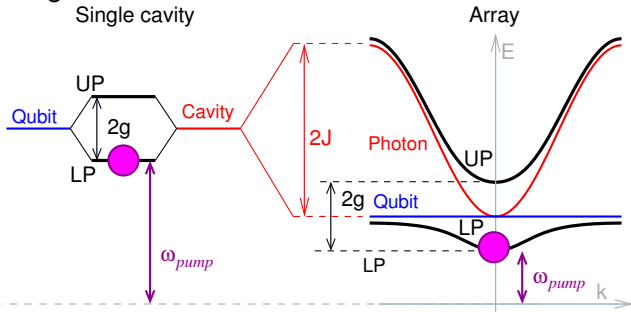
Chose detuning *a la* Dicke model



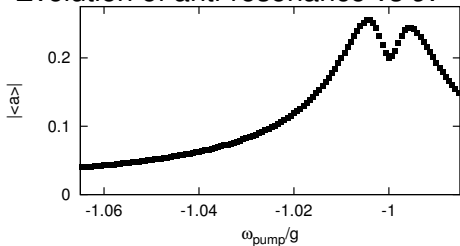
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 - More/less localised states
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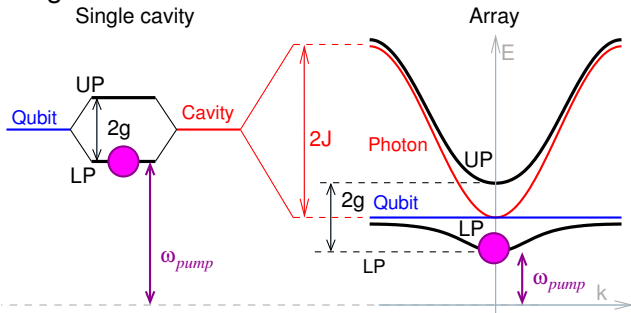
Evolution of anti-resonance vs J .



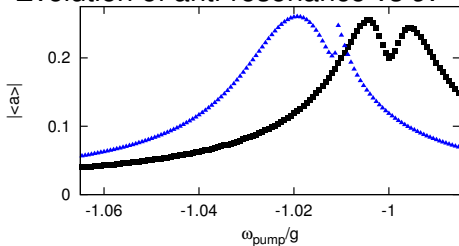
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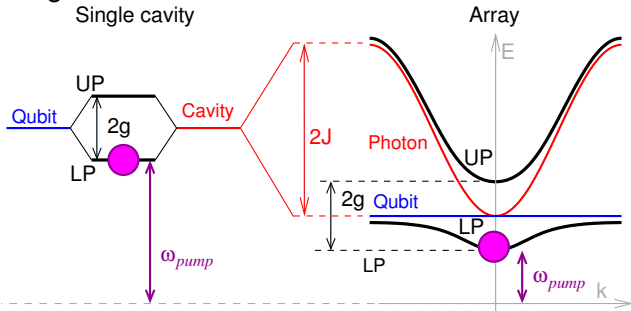
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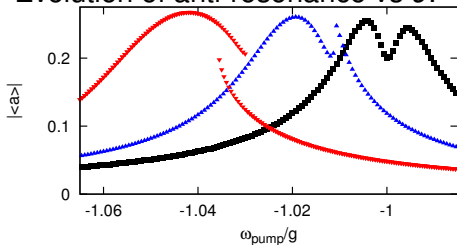
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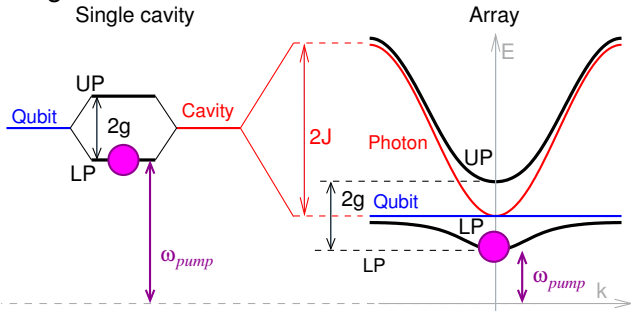
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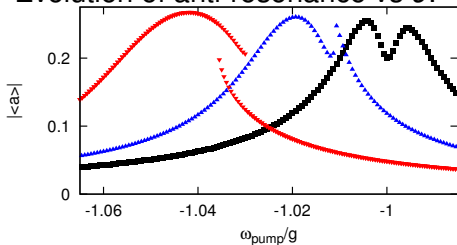
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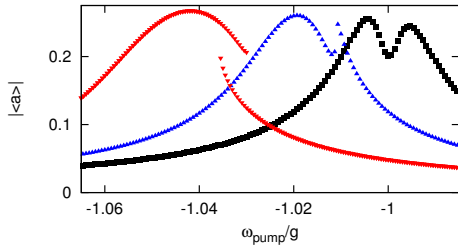
Photon blockade picture $J \lesssim g$

- Polariton basis
- Nonlinearity $|\epsilon_2 - 2\epsilon_1| \propto g$.

$$H = \sum_i \left(\frac{\epsilon}{2} \tau_i^z + \tilde{f} \tau_i^x \right)$$

- Decouple hopping:

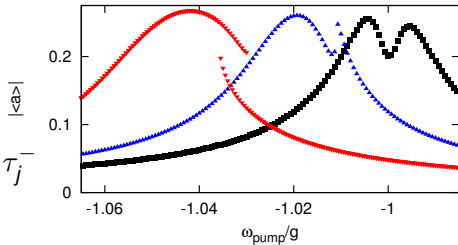
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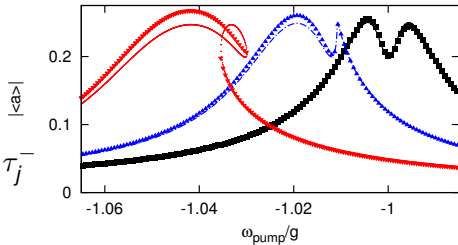
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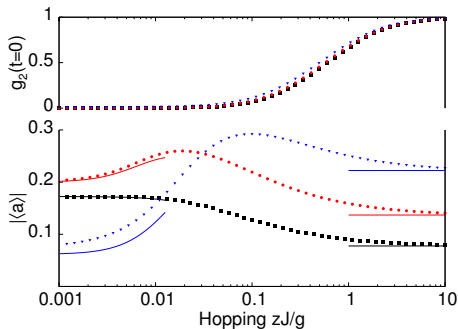
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Coherently pumped array: correlations & fluorescence

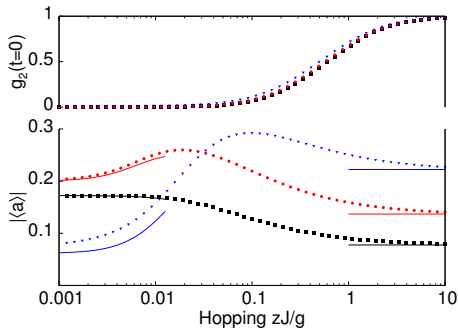


Correlations

* $g_2 : 0 \rightarrow 1$ crossover.

- Small J : Mollow triplet
- Large J : Off resonance fluorescence
- Pump at collective resonance
- Mismatch if $J \neq 0$.

Coherently pumped array: correlations & fluorescence

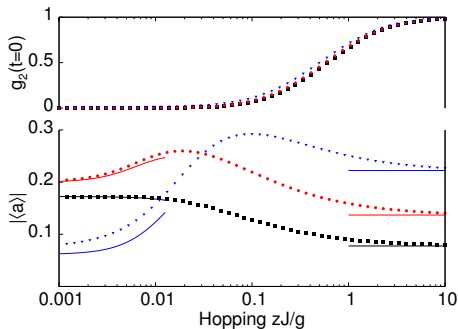


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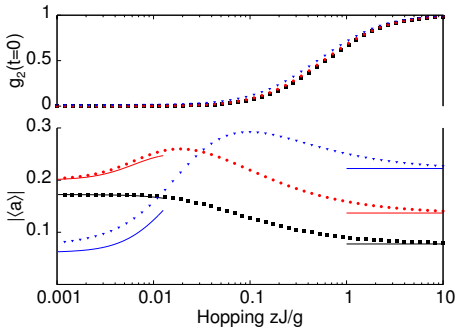
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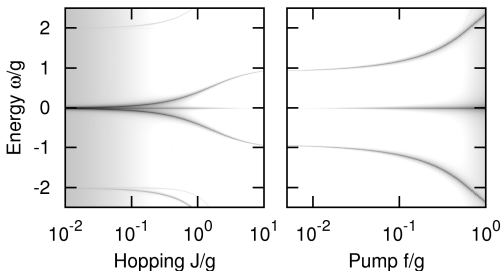


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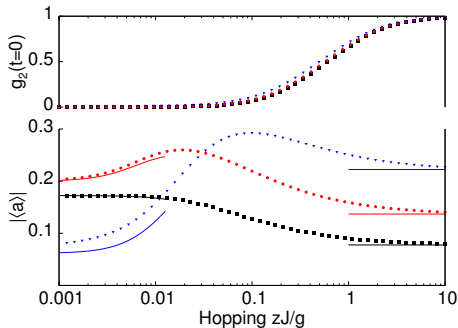
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Coherently pumped array: correlations & fluorescence



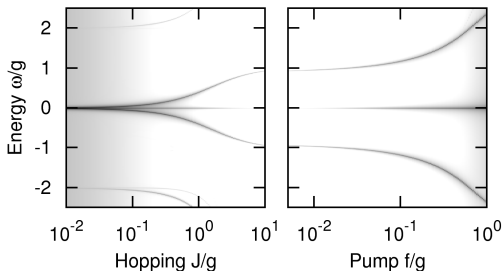
Correlations

- $g_2 : 0 \rightarrow 1$ crossover.

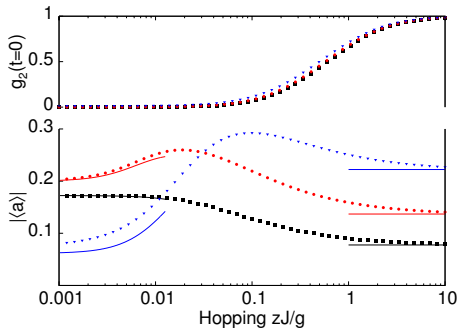
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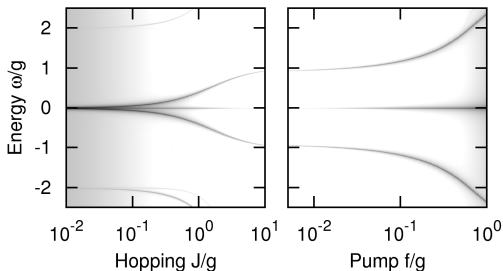


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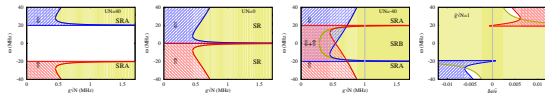
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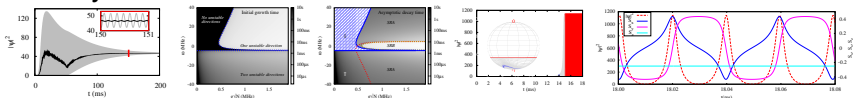


Summary

- Wide variety of dynamical phases

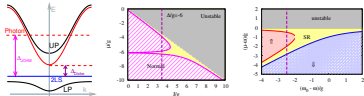


- Slow dynamics for $U < 0$ & Persistent oscillations for $U > 0$

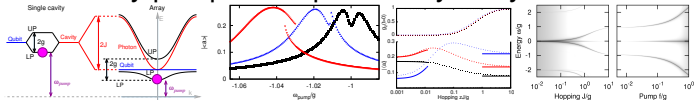


JK *et al.* PRL '10, Bhasen *et al.* PRA '12

- Dicke model and JCHM: connection at $J \rightarrow \infty$



- Coherently pumped coupled cavity array



Nissen *et al.* PRL in press '12

4 Ferroelectric transition

5 Dicke vs JCHM

Ferroelectric transition

Atoms in **Coulomb gauge**

$$H = \sum_k \omega_k a_k^\dagger a_k + \sum_i [p_i - eA(r_i)]^2 + V_{\text{coul}}$$

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Two-level systems — dipole-dipole coupling

$$H = \omega_0 S^z + \omega \psi^\dagger \psi + g(S^+ + S^-)(\psi + \psi^\dagger) + N\zeta(\psi + \psi^\dagger)^2 - \eta(S^+ - S^-)^2$$

(nb $g^2, \zeta, \eta \propto 1/V$).

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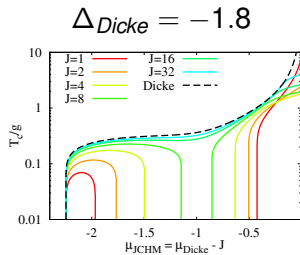
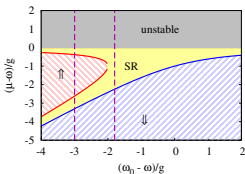
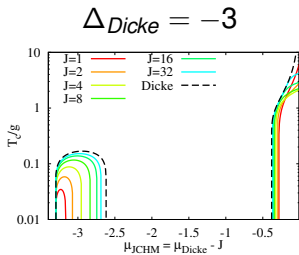
Gauge transform to dipole gauge $\mathbf{D} \cdot \mathbf{r}$

$$H = \omega_0 S^z + \omega \psi^\dagger \psi + \bar{g}(S^+ - S^-)(\psi - \psi^\dagger)$$

“Dicke” transition at $\omega_0 < N\bar{g}^2/\omega \equiv 2\eta N$

But, ψ describes **electric displacement**

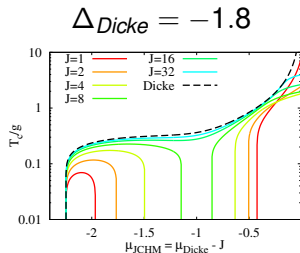
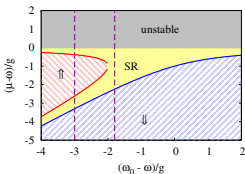
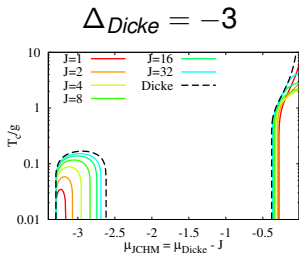
Dicke vs JCHM, $T \neq 0$



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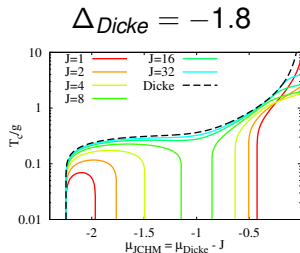
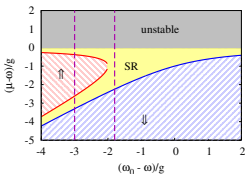
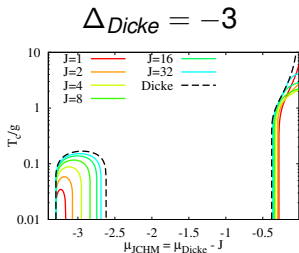


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→ Fluctuation mass $m \sim 1/J$, fluctuations suppress $T_c < \mu/m$

→ Fluctuations can induce re-entrance [JK et al. PRB '05]

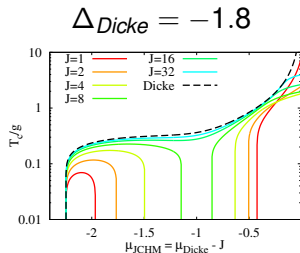
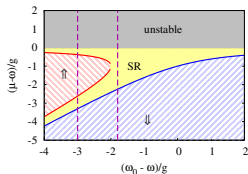
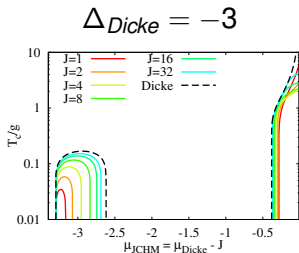
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