

Non-equilibrium phases of coupled matter-light systems

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St Andrews

600
YEARS



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Outline

1 Dynamics of generalized Dicke model

- Summary of experiment and classical dynamics
- Fixed points and dynamical phases
- Timescales and consequences for experiment
- Persistent oscillating phases

2 Non-equilibrium states of Jaynes-Cummings-Hubbard Model

- Relating equilibrium JCHM & Dicke model
- Coherently pumped JCHM

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Part 1:



J. Mayoh



M. J. Bhaseen



B. D. Simons

Part 2:



F. Nissen



G. Blatter



M. Biondi



S. Schmidt

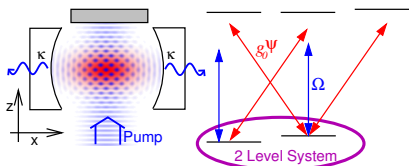


H. Türeci

EPSRC

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Research Council

Reminder of cold-atom extended Dicke model



2 Level system, $|\downarrow\rangle, |\uparrow\rangle$:

$$\downarrow: \Psi(x, z) = 1$$

$$\uparrow: \Psi(x, z) = \sum_{\sigma, \sigma' = \pm} e^{ik(\sigma x + \sigma' z)}$$

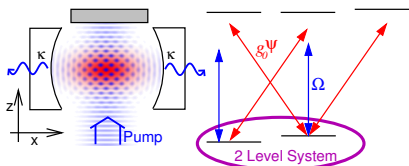
$$H = \omega \psi^\dagger \psi + \omega_0 S^z + g(\psi + \psi^\dagger)(S^- + S^+) + U S_z \psi^\dagger \psi$$

$$\partial_t \rho = -i[H, \rho] - \kappa(\psi^\dagger \psi \rho + 2\rho \psi^\dagger \psi + \rho \psi^\dagger \psi)$$

[Baumann *et al* Nature '10]



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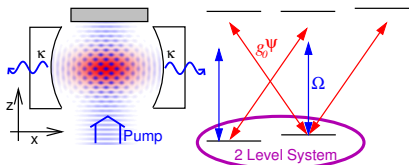
$$H = \omega \psi^\dagger \psi + \omega_0 S^z + g(\psi + \psi^\dagger)(S^- + S^+)$$

$$a_{\sigma\sigma'} = -i[H, a] = -k(\psi^\dagger \psi \sigma - 2\sigma \rho^\dagger + \rho \sigma)$$

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Feedback: $U \propto \frac{g_0^2}{\omega_c - \omega_a}$

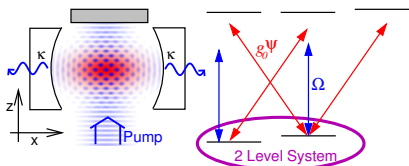
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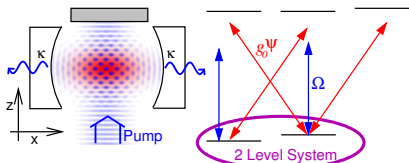
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$$\omega_0 \sim \text{kHz} \ll \omega, \kappa, g\sqrt{N} \sim \text{MHz}.$$

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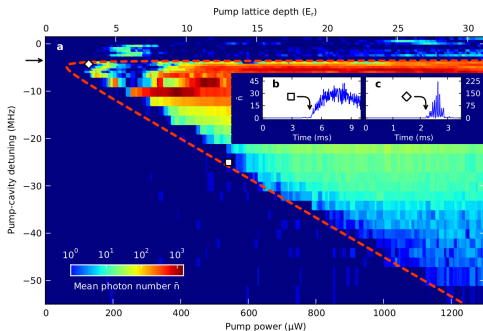
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Classical dynamics of the extended Dicke model

Open dynamical system:

$$H = \omega\psi^\dagger\psi + \omega_0\mathbf{S}^z + g(\psi + \psi^\dagger)(\mathbf{S}^- + \mathbf{S}^+) + US_z\psi^\dagger\psi.$$
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- Neglects quantum fluctuations — restore via Wigner distributed initial conditions.
- Linearisation about fixed point:
 - Recover Retarded Green's function (spectrum)
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Classical EOM
($|\mathbf{S}| = N/2 \gg 1$)

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Fixed points (steady states)

$$0 = i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z$$

$$0 = ig(\psi + \psi^*)(S^- - S^+)$$

$$0 = -[\kappa + i(\omega + US^z)]\psi - ig(S^- + S^+)$$

• $\psi = 0, S = (0, 0, \pm N/2)$
always a solution.

• If $g > g_c, \psi \neq 0$ too

A. $S^z = -S[S^-] = 0$

B. $\psi = \Re[\psi] = 0$

Fixed points (steady states)

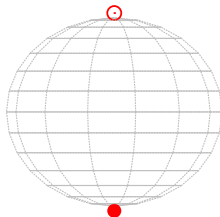
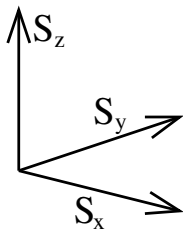
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- $\psi = 0, \mathbf{S} = (0, 0, \pm N/2)$ always a solution.

• If $g > g_c, \psi \neq 0$ too
 $\Delta S^z = -S(S^z) = 0$
 $\Delta \psi = \mathcal{R}[\psi] = 0$



Small g : \uparrow, \downarrow only.
 $(\omega = 30\text{MHz}, UN = -40\text{MHz})$

Fixed points (steady states)

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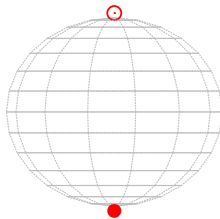
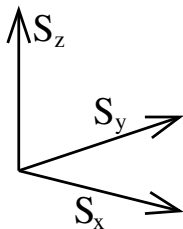
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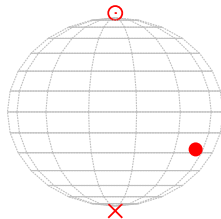
- If $g > g_c, \psi \neq 0$ too

A $S^y = -\Im[S^-] = 0$

B $\psi' = \Re[\psi] = 0$



Small g : \uparrow, \downarrow only.
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Larger g : SR too.

Steady state phase diagram

$$0 = i(\omega_0 + U|\psi|^2)S^- + 2ig(\psi + \psi^*)S^z$$

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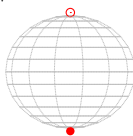
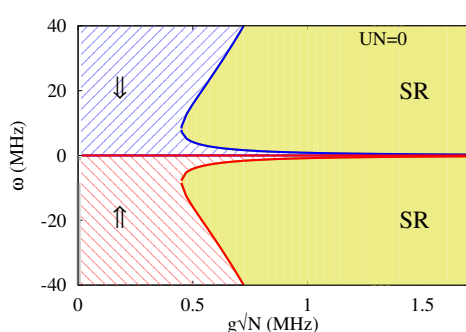
See also Domokos and Ritsch PRL '02, Domokos *et al.* PRL '10

Steady state phase diagram

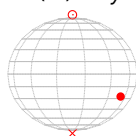
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SR(A): $S_y = 0$



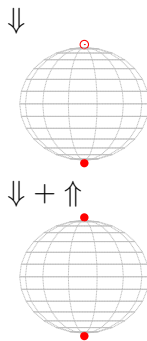
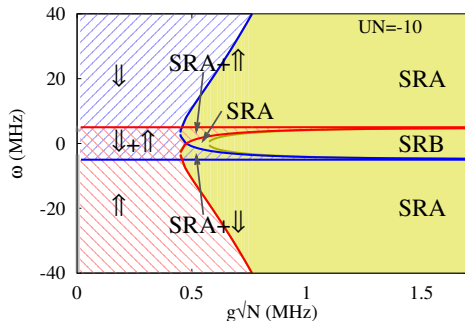
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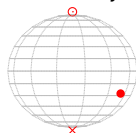
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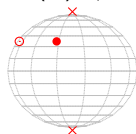
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SR(A): $S_y = 0$



SR(B): $\psi' = 0$



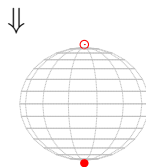
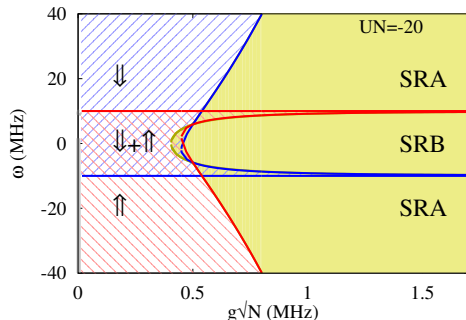
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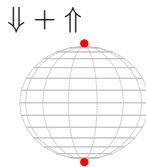
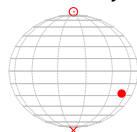
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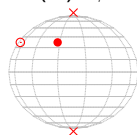
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SR(A): $S_y = 0$



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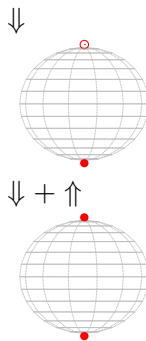
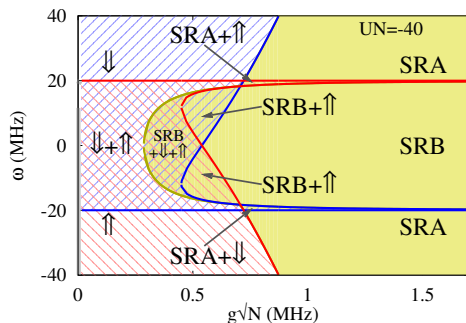
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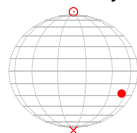
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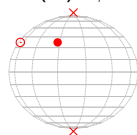
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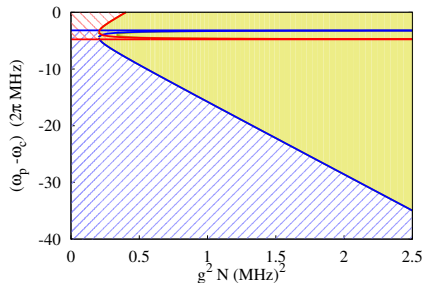


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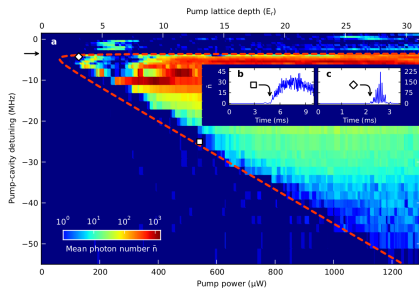
Comparison to experiment



$$UN = -10\text{MHz}$$

Adapted from: [Bhaseen *et al.* PRA '12]

$$\omega = \omega_c - \omega_p + \frac{5}{2}UN,$$



[Baumann *et al* Nature '10]

$$UN = -\frac{g_0^2}{4(\omega_a - \omega_c)}$$

1 Dynamics of generalized Dicke model

- Summary of experiment and classical dynamics
- Fixed points and dynamical phases
- **Timescales and consequences for experiment**
- Persistent oscillating phases

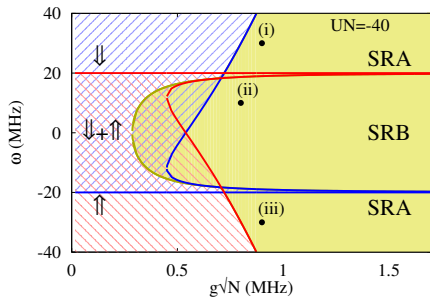
2 Non-equilibrium states of Jaynes-Cummings-Hubbard Model

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- Coherently pumped JCHM

Dynamics: Evolution from normal state

Gray: $\mathbf{S} = (\sqrt{N}, \sqrt{N}, -N/2)$

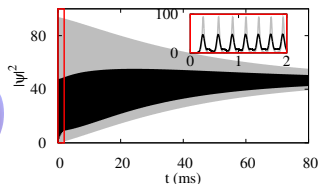
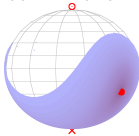
Black: Wigner distribution of \mathbf{S}, ψ



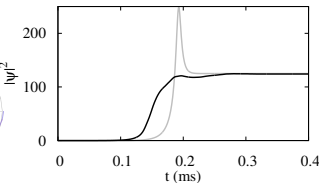
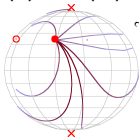
Oscillations: ~ 0.1 ms

Decay: 20ms, 0.1ms, 20ms

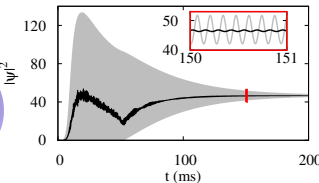
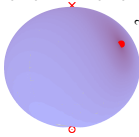
(i) SR(A)



(ii) SR(B)



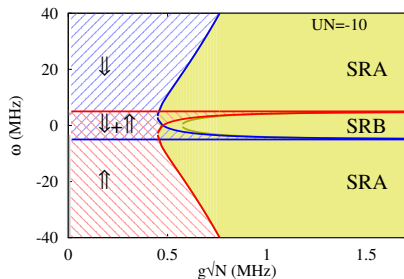
(iii) SR(A)



Asymptotic state: Evolution from normal state

(Near to experimental $UN = -13\text{MHz}$).

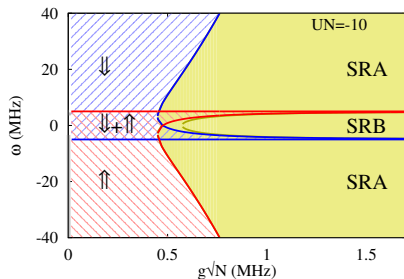
All stable attractors:



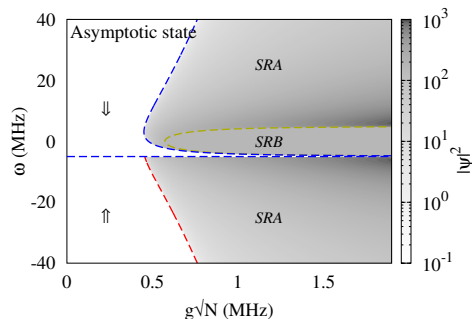
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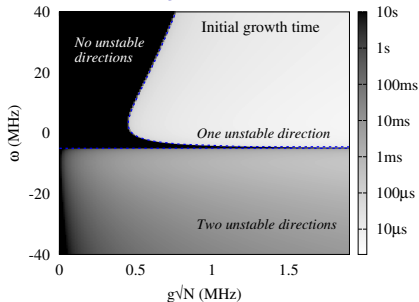
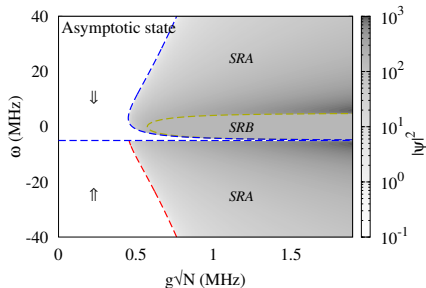
All stable attractors:



Starting from \Downarrow



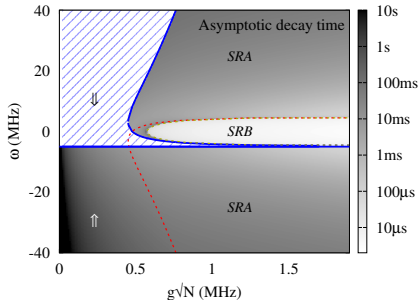
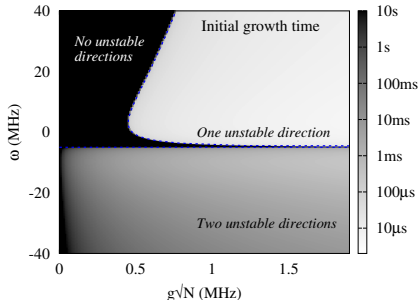
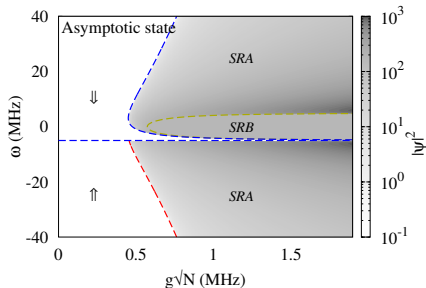
Timescales for dynamics: What are they?



Growth Most unstable eigenvalues near $\mathbf{S} = (0, 0, -N/2)$

Decay Slowest stable eigenvalues near final state

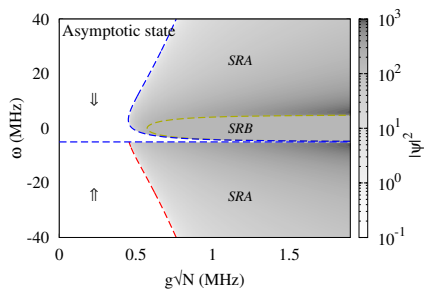
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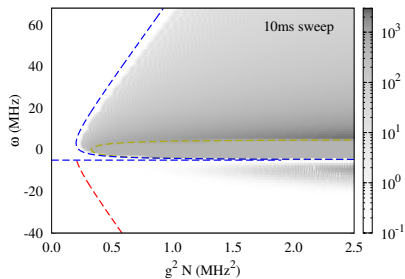
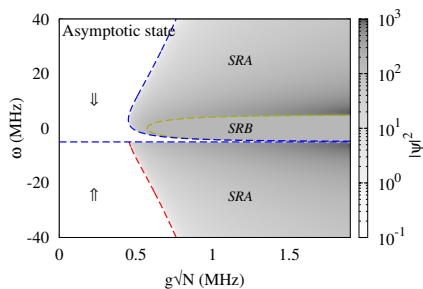
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Decay Slowest stable eigenvalues near final state

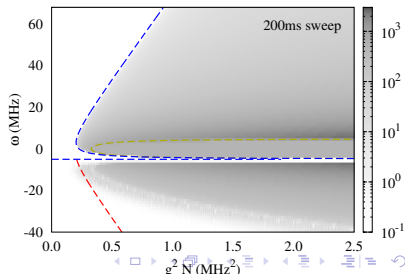
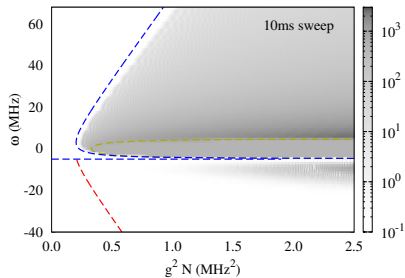
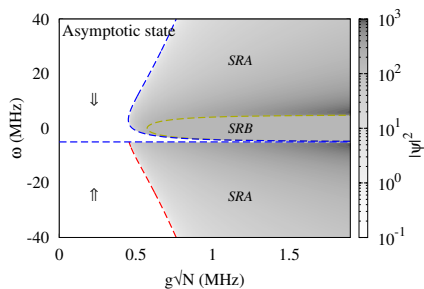
Timescales for dynamics: Consequences for experiment



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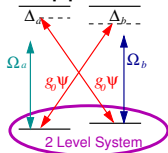


Timescales for dynamics: Consequences for experiment



Timescales for dynamics: Why so slow and varied?

Suppose co- and counter-rotating terms differ

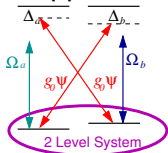


$$H = \dots + g(\psi^\dagger S^- + \psi S^+) + g'(\psi^\dagger S^+ + \psi S^-) + \dots$$

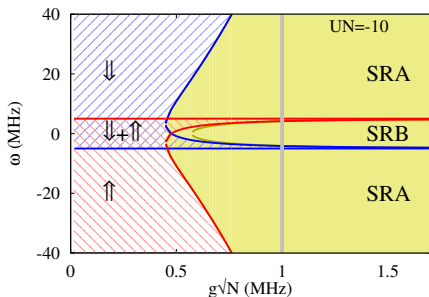
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- SR(A), SR(B) continuously connect

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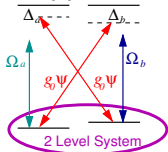
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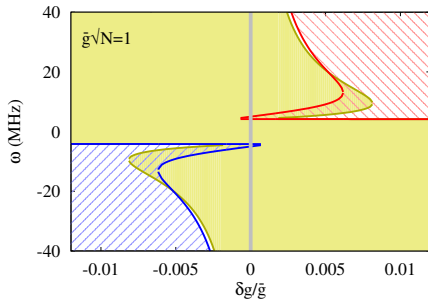
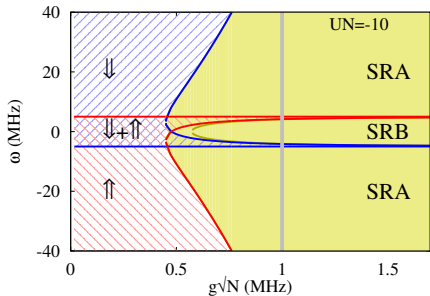
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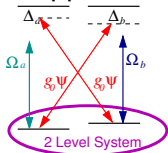
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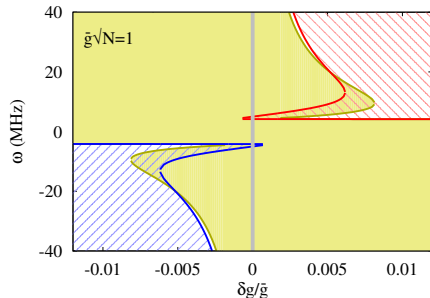
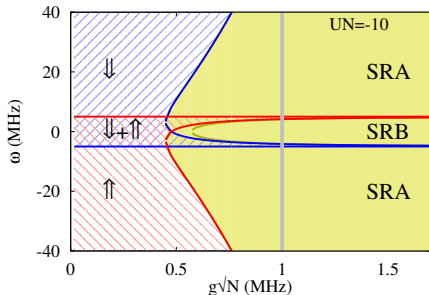
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1 Dynamics of generalized Dicke model

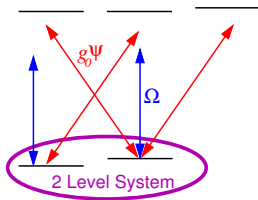
- Summary of experiment and classical dynamics
- Fixed points and dynamical phases
- Timescales and consequences for experiment
- **Persistent oscillating phases**

2 Non-equilibrium states of Jaynes-Cummings-Hubbard Model

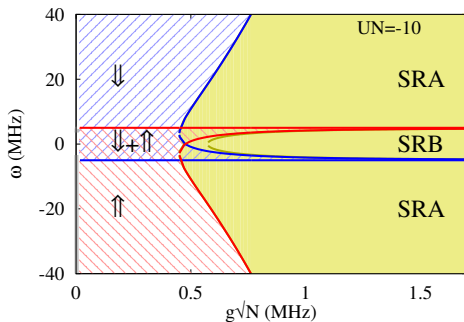
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Regions without fixed points

Changing U :

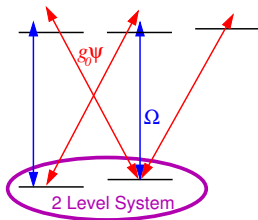


$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$

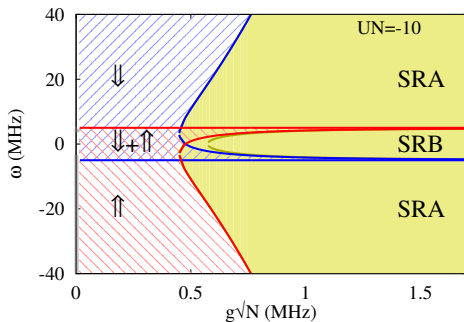


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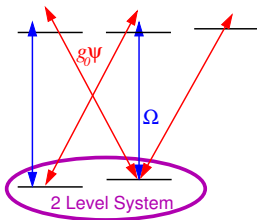


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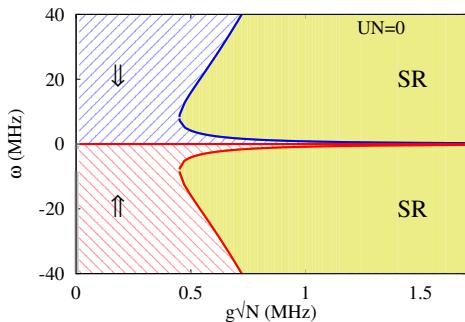


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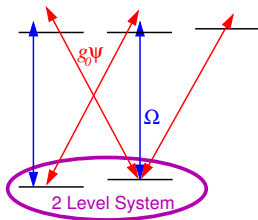


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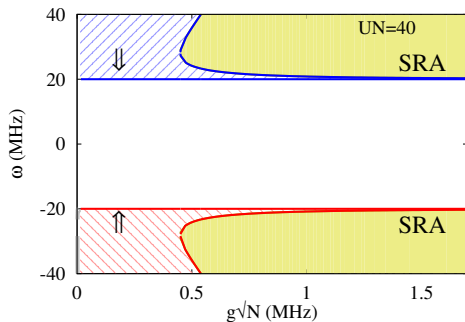


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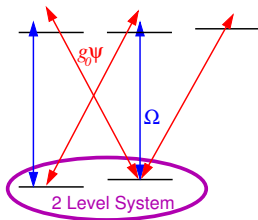


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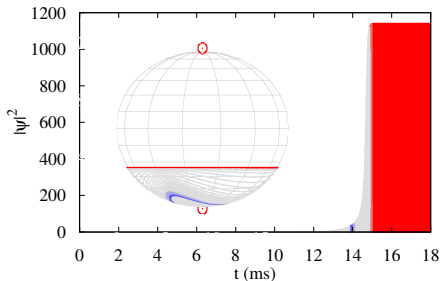
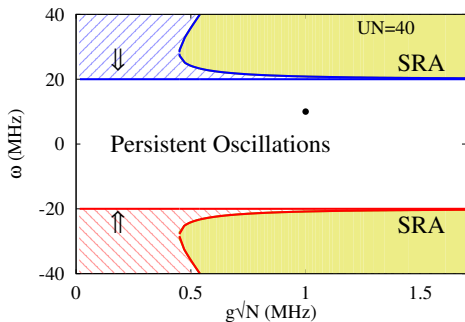


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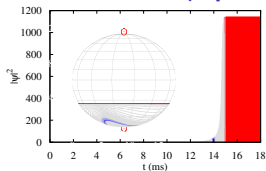
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Persistent (optomechanical) oscillations

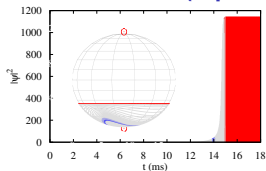


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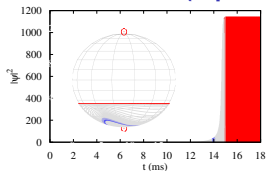
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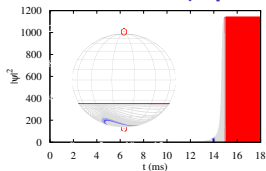
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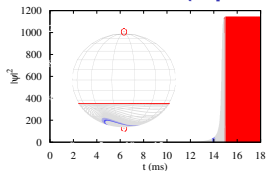
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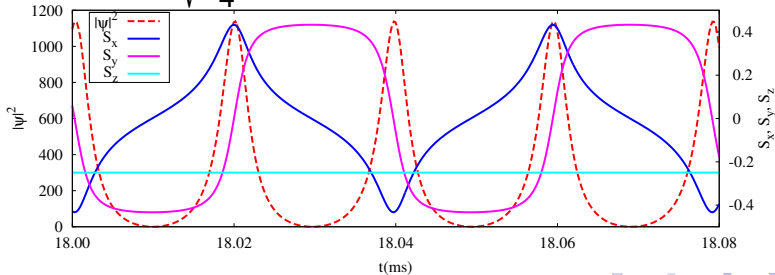
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Outline

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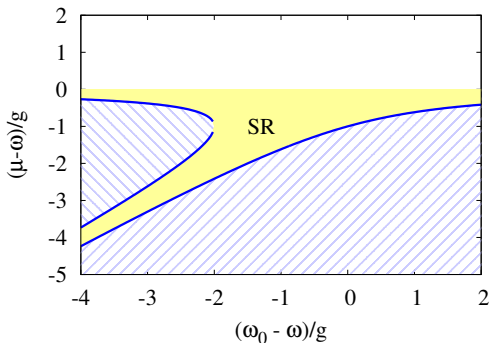
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- Coherently pumped JCHM

Equilibrium: Dicke model with chemical potential

$$H - \mu N = (\omega - \mu)\psi^\dagger\psi + (\omega_0 - \mu)S^z + g(\psi^\dagger S^- + \psi S^+)$$



- Transition at:
 $g^2 N > (\omega - \mu)(\omega_0 - \mu)$
- Reduce critical g

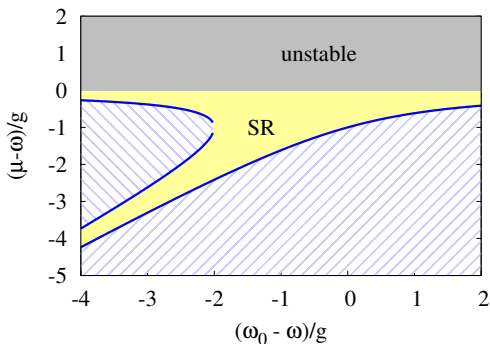
• Unstable if $\mu > \omega$

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[Eastham and Littlewood, PRB '01]

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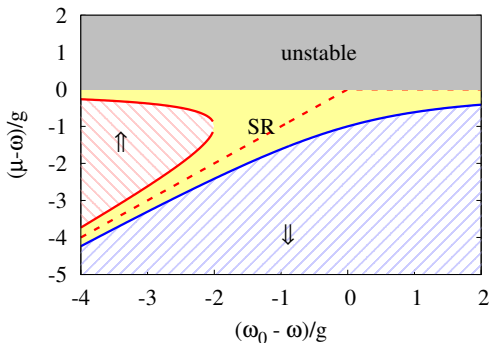


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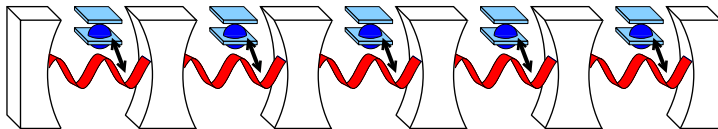
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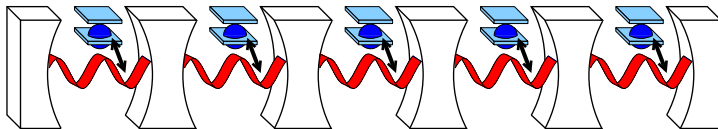
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Jaynes-Cummings Hubbard model

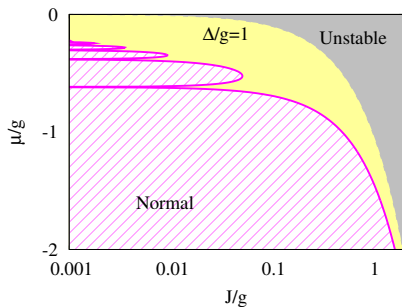


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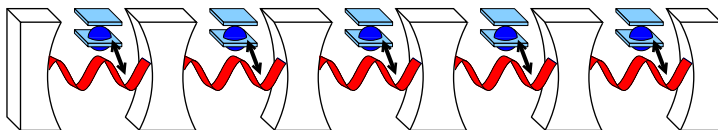
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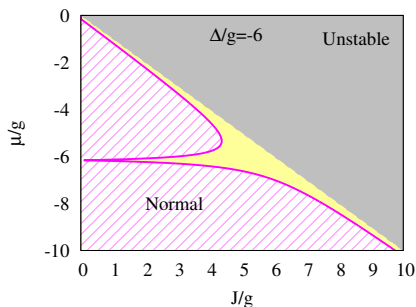
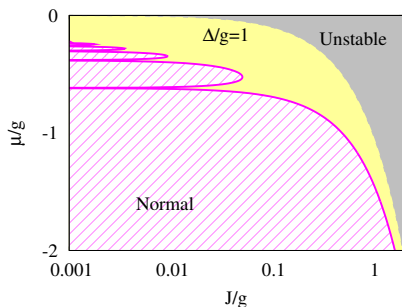
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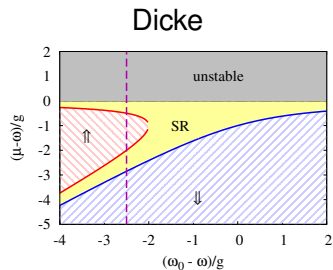
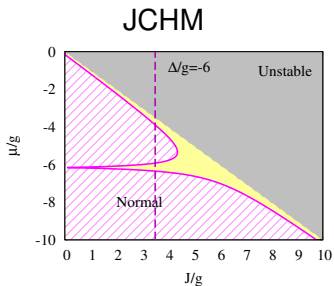
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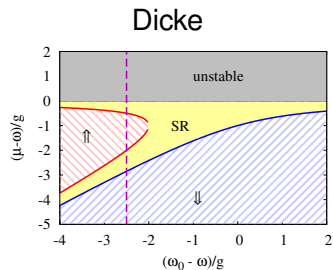
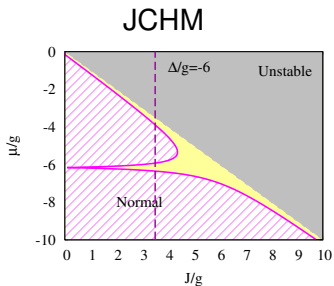


Dicke vs JCHM



- $k = 0$ mode of JCHM \leftrightarrow Dicke photon mode
- $\uparrow \leftrightarrow n = 1$ Mott lobe

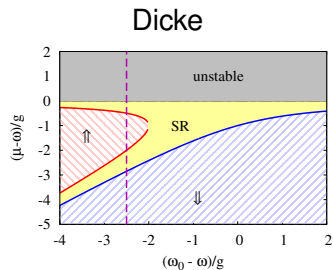
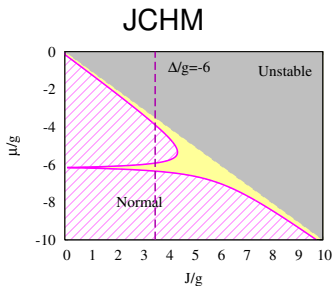
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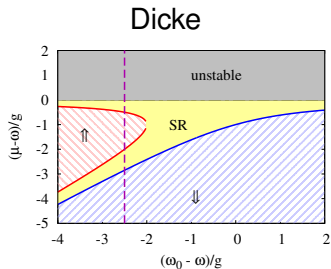
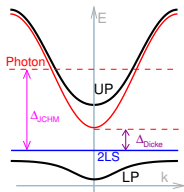
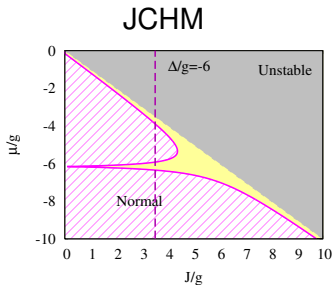
● $\Uparrow \leftrightarrow \Downarrow$ $n = 1$ Mott lobe

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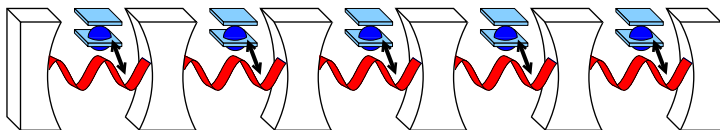
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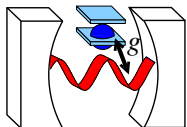
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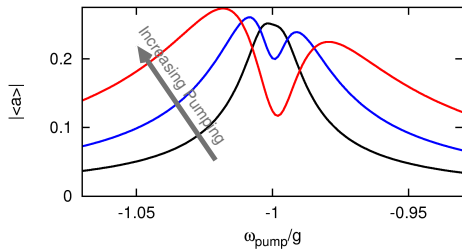
Coherently pumped single cavity [Bishop *et al.* Nat. Phys '09]



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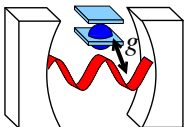
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- Anti-resonance in $|\langle \psi \rangle|$
- Effective 2LS: (Empty) (1 polariton)
- Motz triplet fluorescence



[Lang *et al.* PRL '11]

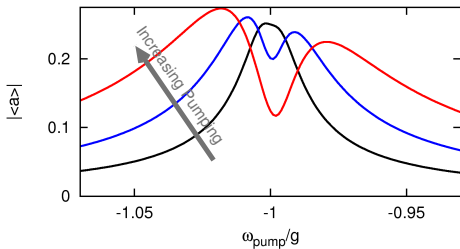
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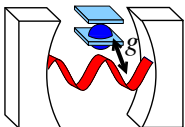
- Anti-resonance in $|\langle \psi \rangle|$.
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• Macro triplet fluorescence

[Lang *et al.* PRL '11]

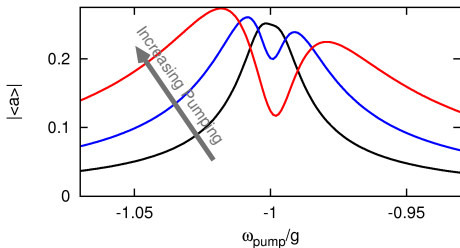
Coherently pumped single cavity [Bishop *et al.* Nat. Phys '09]



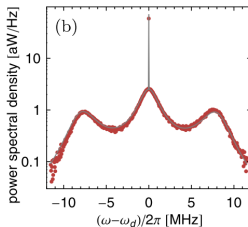
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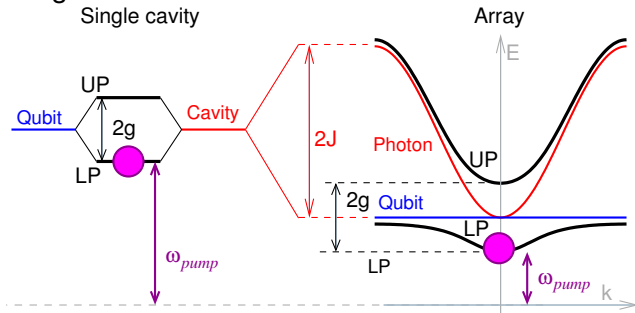
- Mollow triplet fluorescence



[Lang *et al.* PRL '11]

Coherently pumped dimer & array

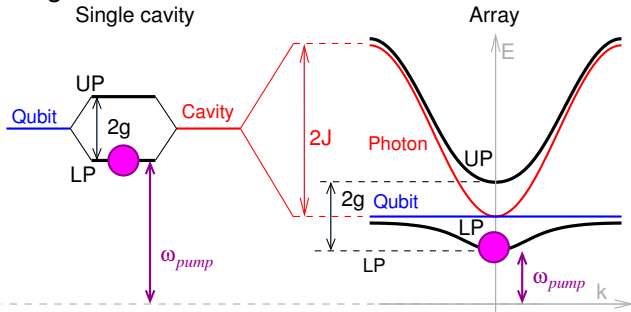
Chose detuning *a la* Dicke model



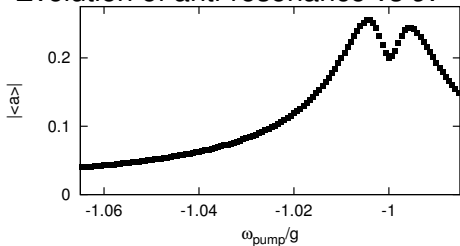
- Bistability at intermediate J
 - More/less localised states
 - Connects to Dicke limit
- Lines → effective TLS

Coherently pumped dimer & array

Chose detuning *a la* Dicke model



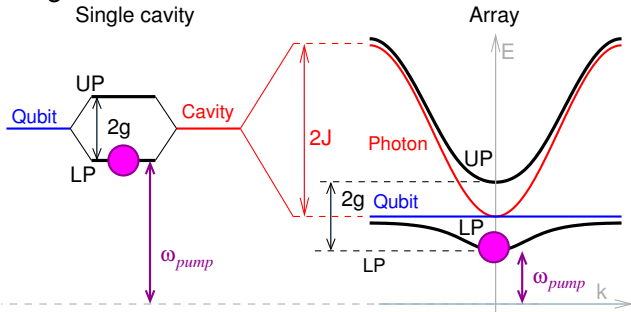
Evolution of anti-resonance vs J .



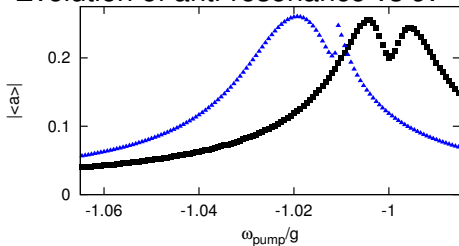
- Bistability at intermediate J
- More/less localised states
- Connects to Dicke limit
- Lines \rightarrow effective TLS

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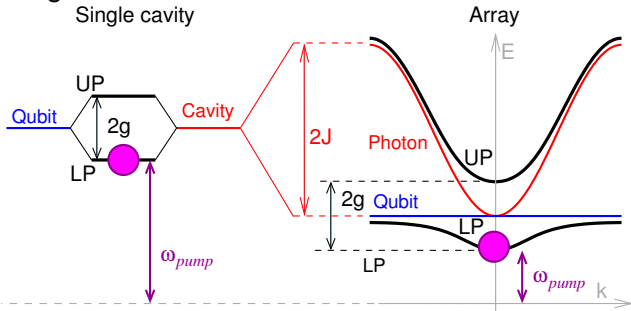
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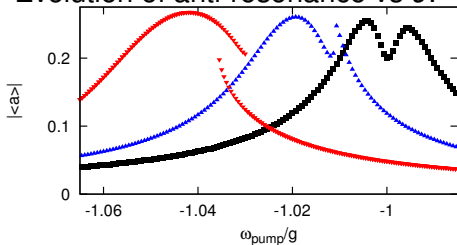
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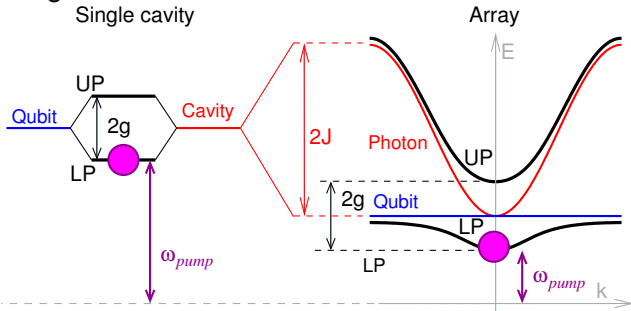
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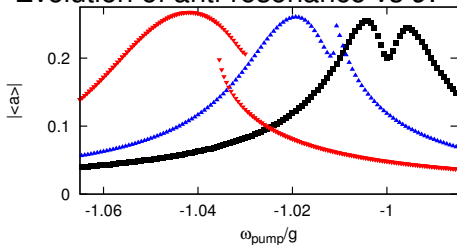
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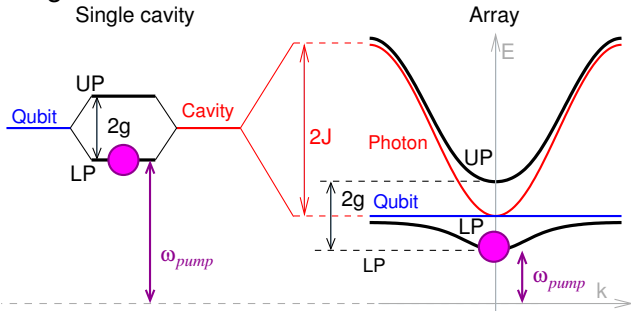


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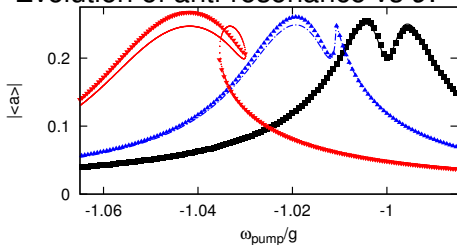
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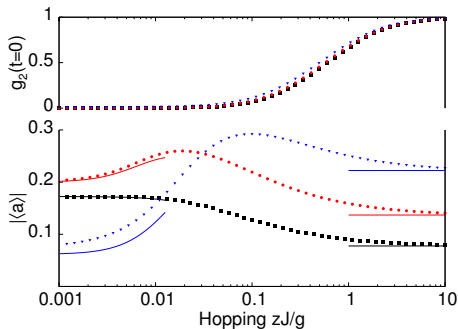


Evolution of anti-resonance vs J .



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Coherently pumped array: correlations & fluorescence

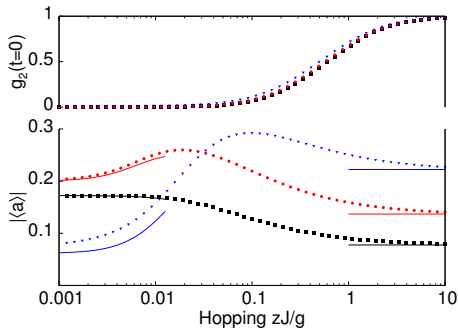


Correlations

* $g_2 : 0 \rightarrow 1$ crossover.

- Small J : Mollow triplet
- Large J : Off resonance fluorescence
 - Pump at collective resonance
 - Mismatch if $J \neq 0$.

Coherently pumped array: correlations & fluorescence

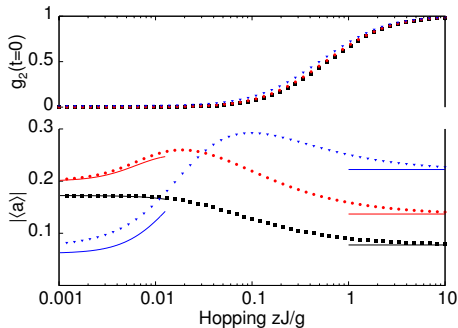


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Coherently pumped array: correlations & fluorescence



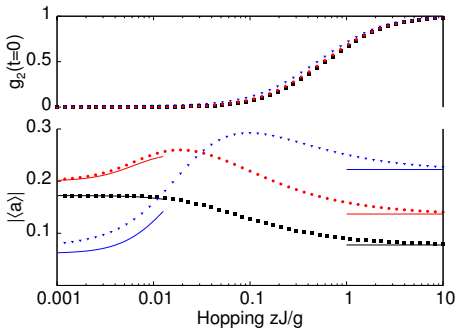
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Coherently pumped array: correlations & fluorescence

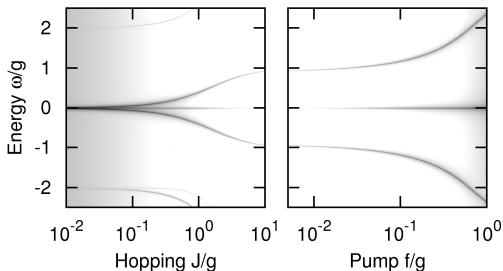


Correlations

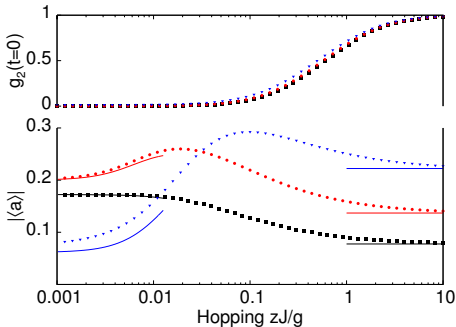
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Coherently pumped array: correlations & fluorescence



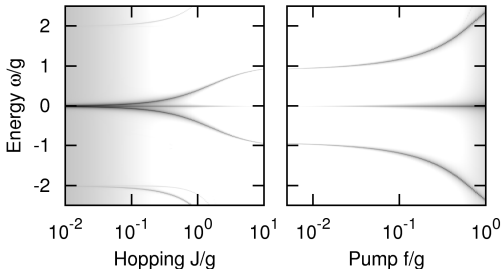
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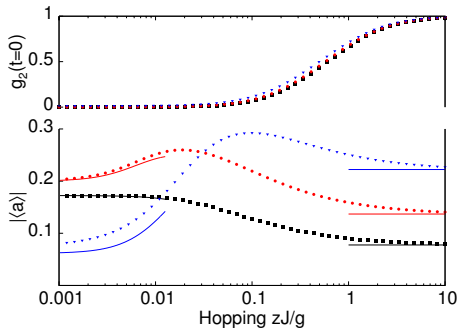
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Coherently pumped array: correlations & fluorescence

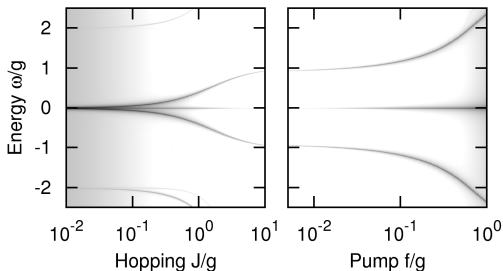


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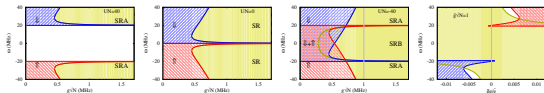
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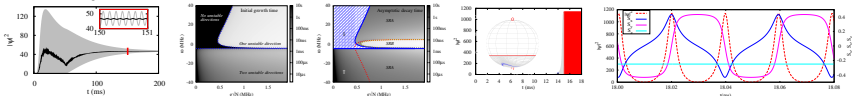


Summary

- Wide variety of dynamical phases

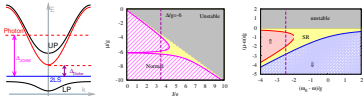


- Slow dynamics for $U < 0$ & Persistent oscillations for $U > 0$

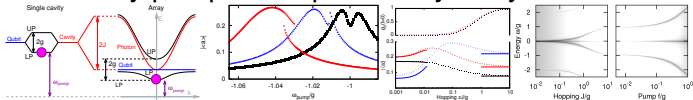


JK *et al.* PRL '10, Bhaseen *et al.* PRA '12

- Dicke model and JCHM: connection at $J \rightarrow \infty$



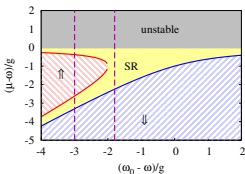
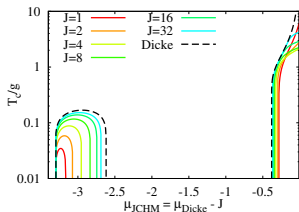
- Coherently pumped coupled cavity array



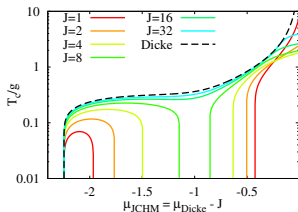
Nissen *et al.* PRL in press '12

Dicke vs JCHM, $T \neq 0$

$$\Delta_{Dicke} = -3$$



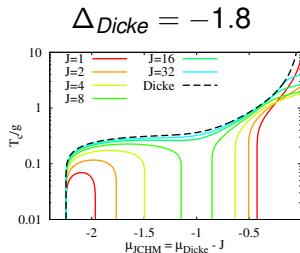
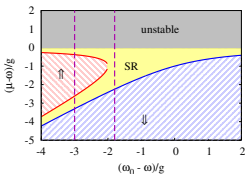
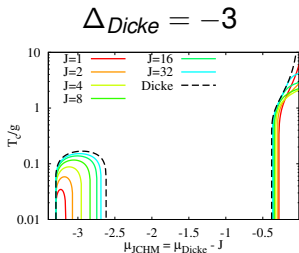
$$\Delta_{Dicke} = -1.8$$



- Match improves as $J \rightarrow \infty$

• Finite bandwidth: fluctuations suppresses T_c .

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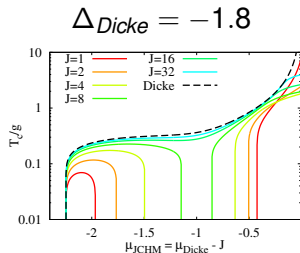
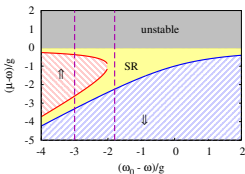
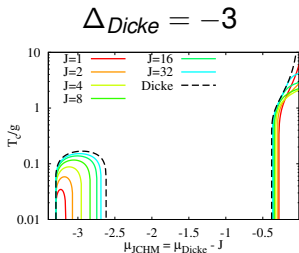


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• Fluctuation mass $m \sim 1/J$, fluctuations suppress $T_c < \mu/m$

• Fluctuations can induce re-entrance [JK et al. PRB '05]

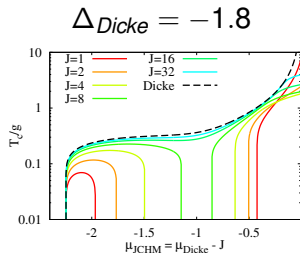
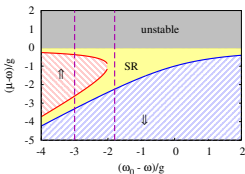
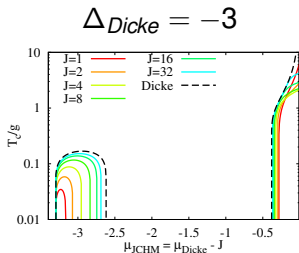
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