

# Pattern formation, Superfluidity and Coherence of Polariton Condensates.

Jonathan Keeling



UMass Amherst, January 2012



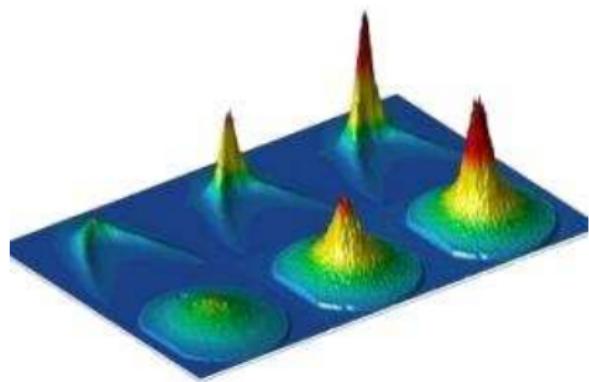
Funding:

**EPSRC**

Engineering and Physical Sciences  
Research Council

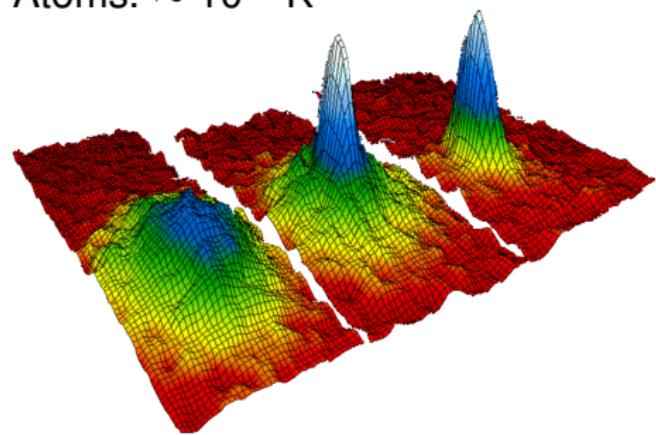
# Bose-Einstein condensation: macroscopic occupation

Polaritons.  $\sim 20\text{K}$



[Kasprzak *et al.* Nature, '06]

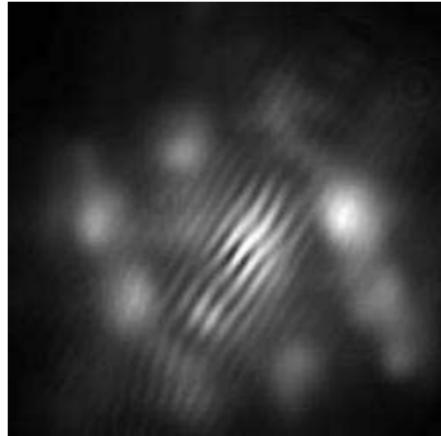
Atoms.  $\sim 10^{-7}\text{K}$



[Anderson *et al.* Science '95]

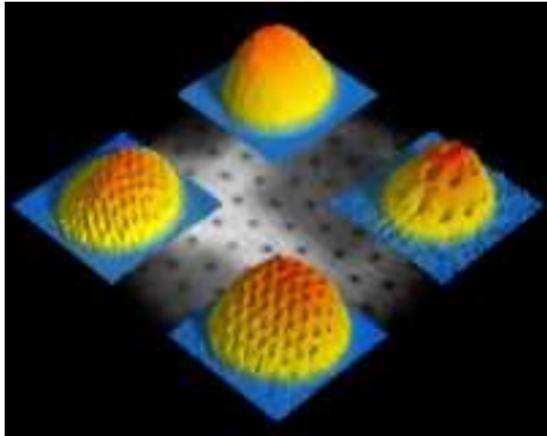
# Macroscopic coherence: vortices

Polaritons:



[Lagoudakis *et al.* Nat. Phys. '08]

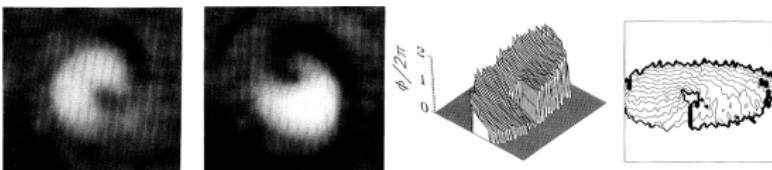
Atoms:



[Abo-Shaeer *et al.* Science '01]

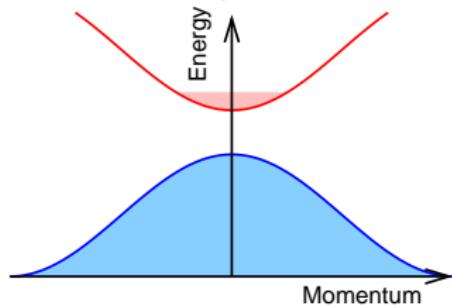
But also, nonlinear optics:

[Arecchi *et al.* PRL '91]



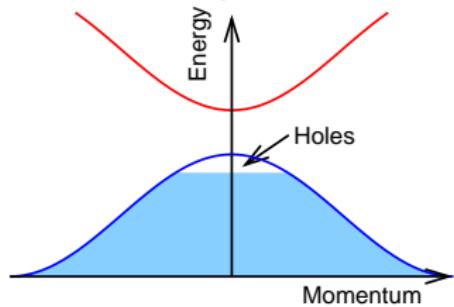
# Excitons

Electronic spectrum:



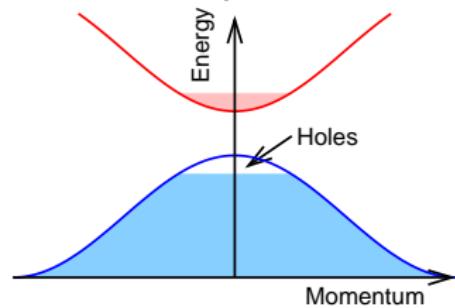
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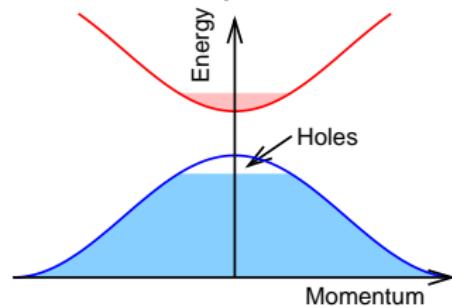


$$H = \sum_i T_i^e + T_i^h + \sum_{ij} V_{ij}^{ee} + V_{ij}^{hh} - V_{ij}^{eh}$$

$$T_i = \frac{p_i^2}{2m_i} \quad V_{ij} = \frac{e^2}{\epsilon_r |r_i - r_j|}$$

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Bound state: Exciton,

$$M \sim m_e + m_h$$

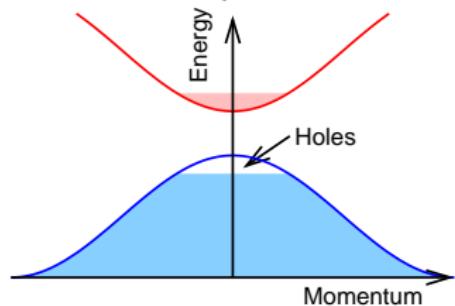
Approximate Bose statistics:

$$[c_{\text{exciton},k}, c_{\text{exciton},k'}^\dagger] \simeq \delta_{k,k'}$$

$$\text{If } \rho(a_{B,\text{exciton}})^D \ll 1$$

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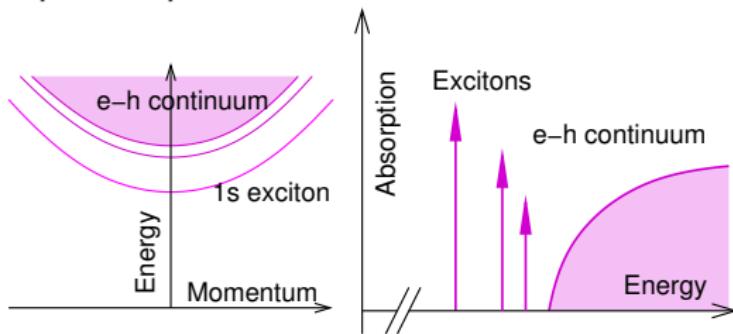
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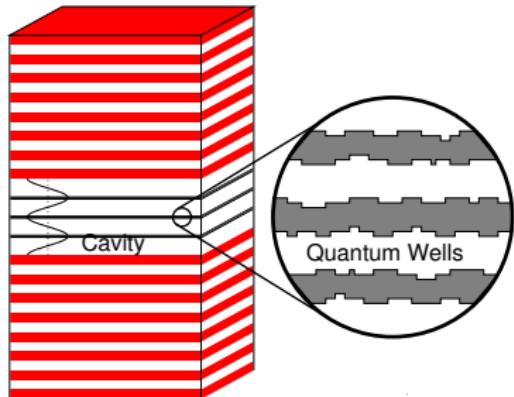
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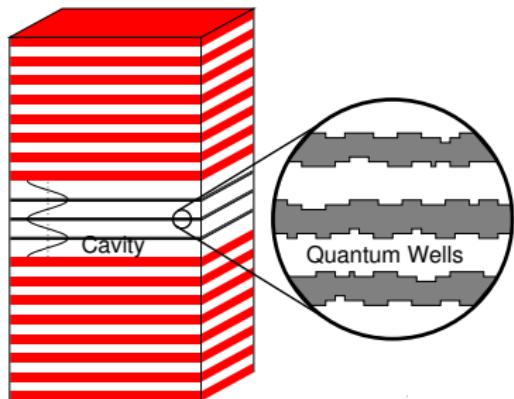
Optical spectrum



# Microcavity polaritons

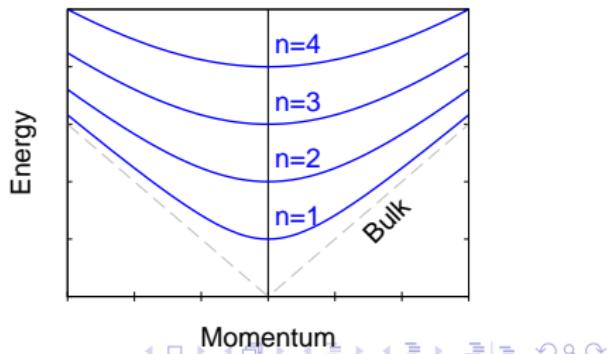


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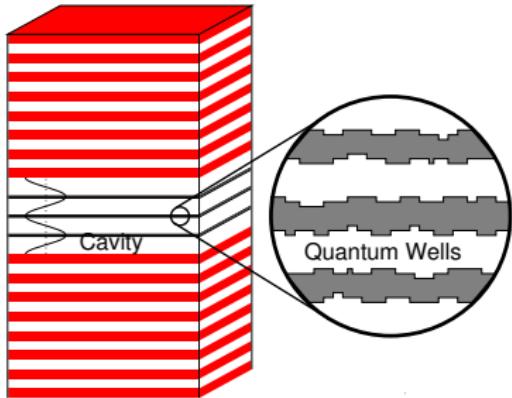


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2 / 2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$



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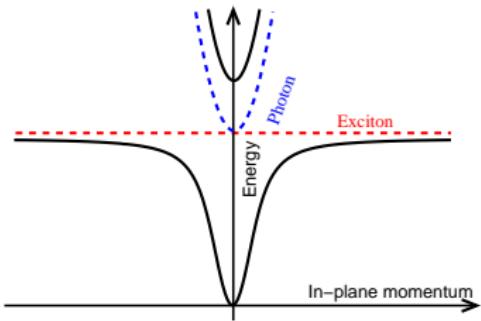


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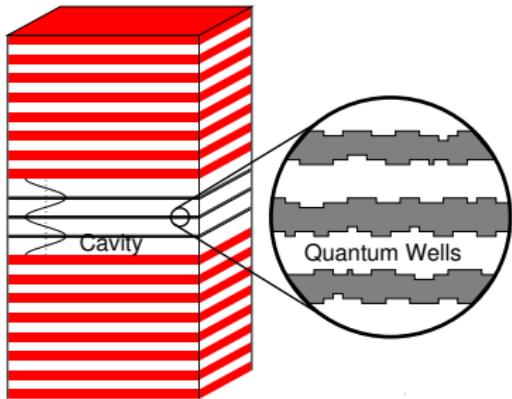
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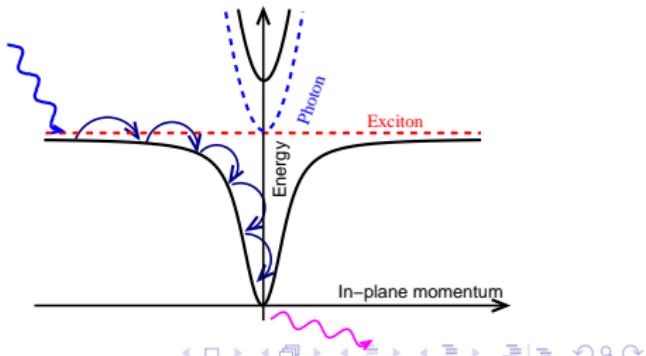


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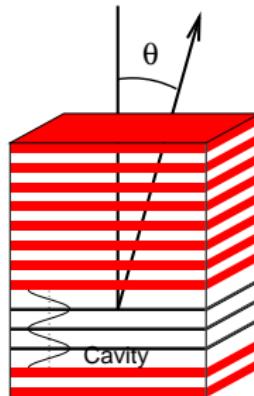
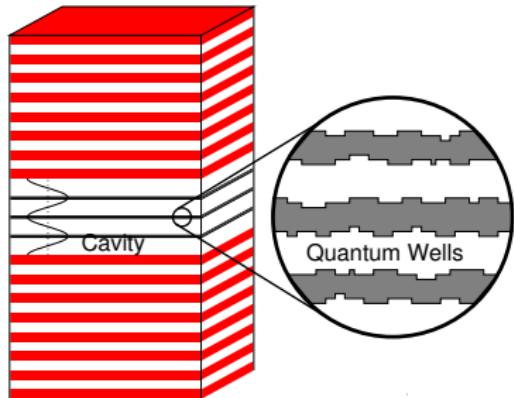


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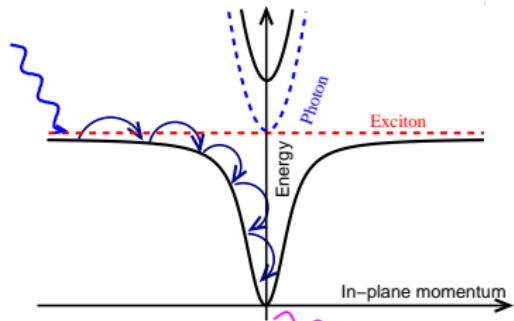


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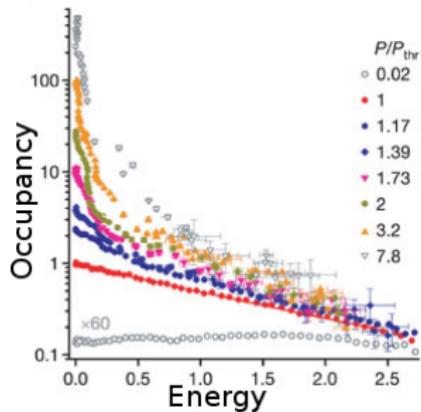
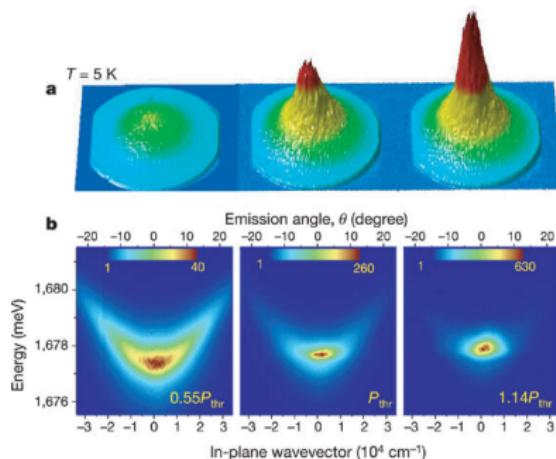
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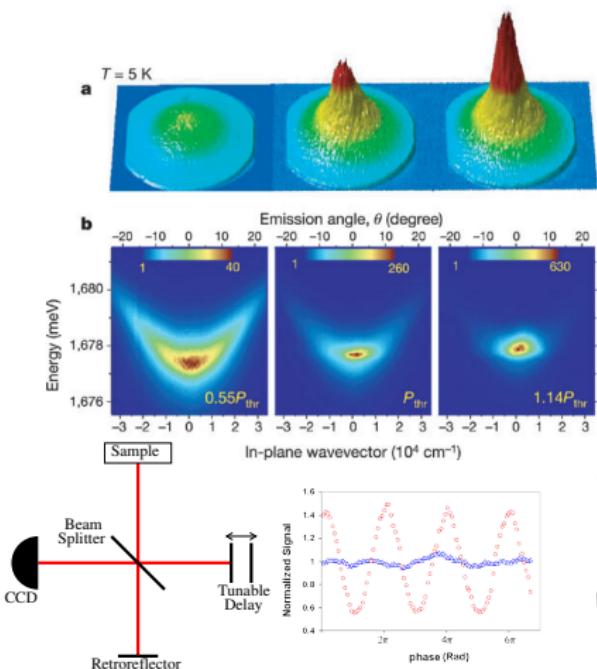


# Polariton experiments: occupation and coherence

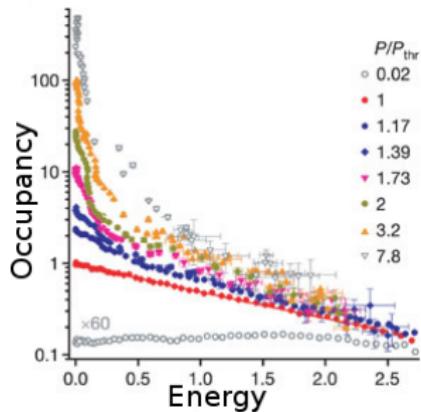


[Kasprzak, *et al.* Nature, '06]

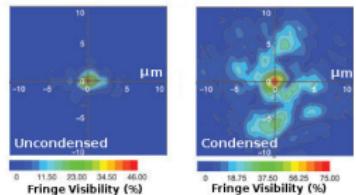
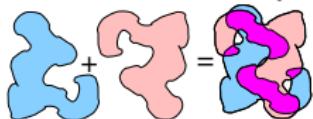
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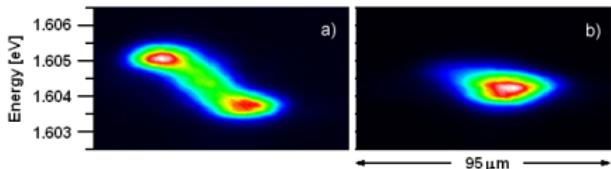


Coherence map:

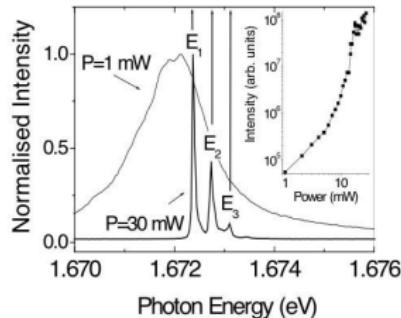
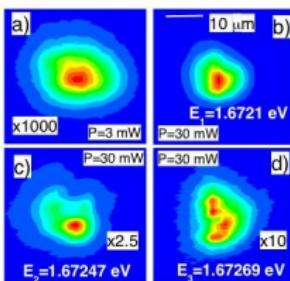


# (Some) other polariton condensation experiments

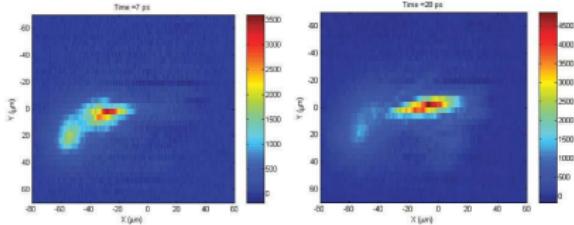
- Polariton traps  
[Balili *et al.* Science '07])



- Multimode condensate and sharp lines  
[Love *et al.* PRL '08]



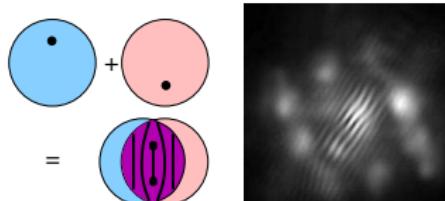
- Wavepacket propagation  
[Amo *et al.* Nature '09]



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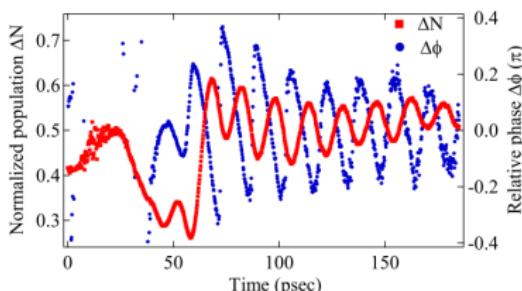
- Quantised vortices

[Lagoudakis *et al.* Nat. Phys. '08. Science '09, PRL '10; Sanvitto *et al.* Nat. Phys. '10; Roumpos *et al.* Nat. Phys. '10 ]



- Josephson oscillations

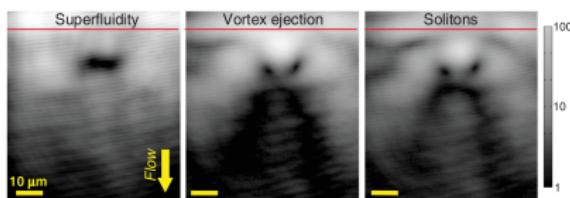
[Lagoudakis *et al.* PRL '10]



- Pattern formation/Hydrodynamics

[Amo *et al.* Science '11, Nature '09;

Wertz *et al.* Nat. Phys '10]



## 1 Introduction to polariton condensation

- What are excitons and polaritons
- Experimental features
- Approaches to modelling

## 2 Pattern formation

- Experiments
- Modelling pattern formation

## 3 Superfluidity

- Non-equilibrium condensate spectrum
- Experiments and aspects of superfluidity
- Current-current response function

## 4 Coherence

- Experiments
- Power law decay of coherence

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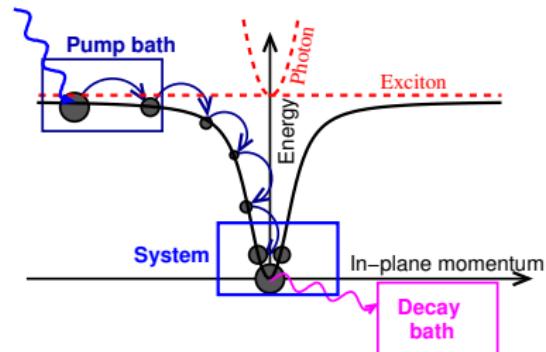
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# Non-equilibrium approach: Steady state, and fluctuations

$$H = H_{\text{sys}} + H_{\text{sys,bath}} + H_{\text{bath}},$$

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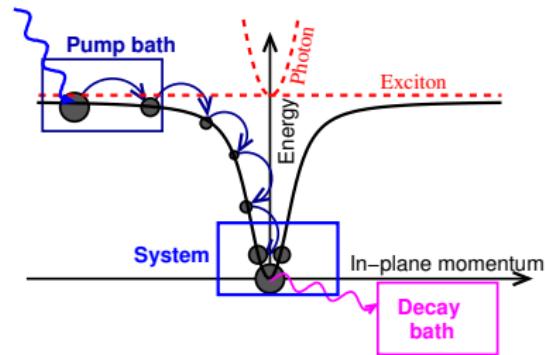


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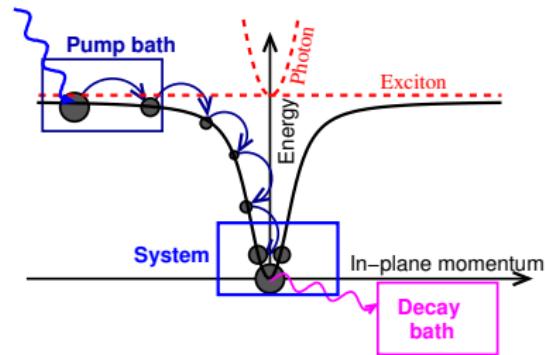
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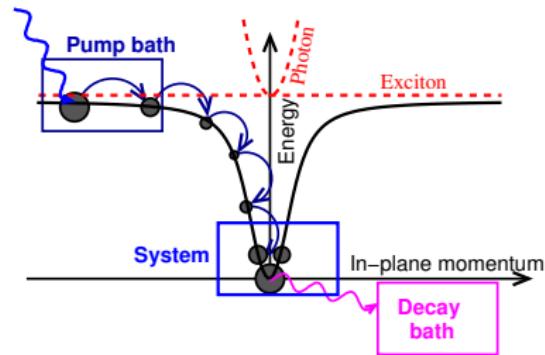
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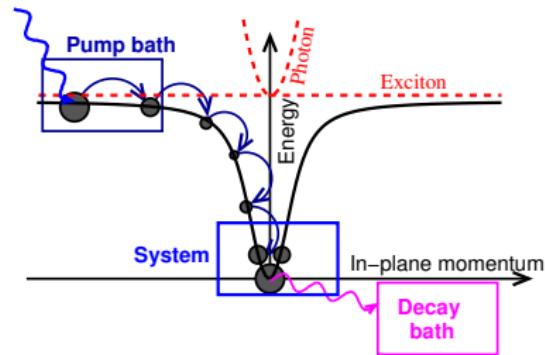
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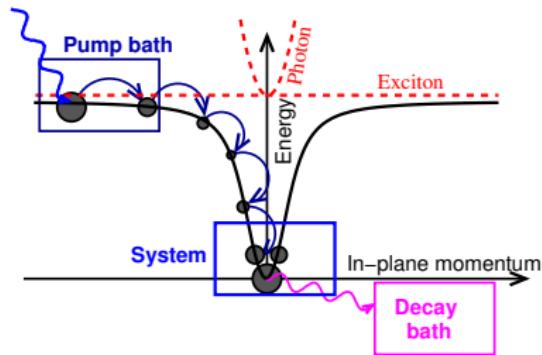
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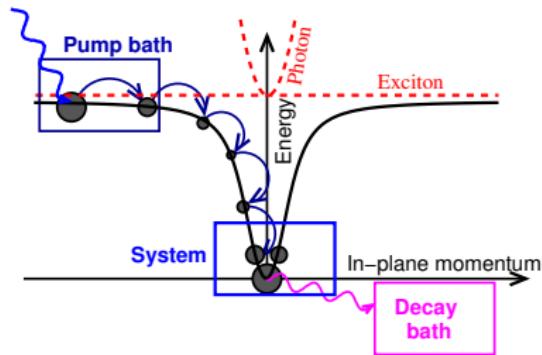
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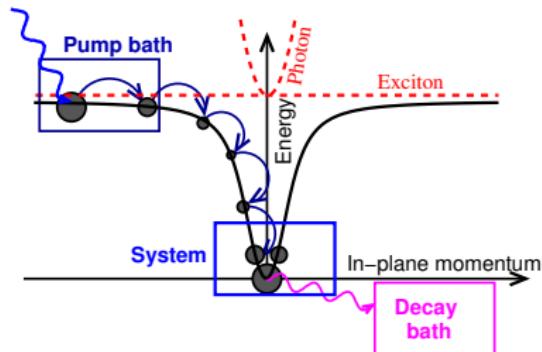
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$$D^K(\omega) = (2n(\omega) + 1) \text{DoS}(\omega)$$





# Pattern formation:

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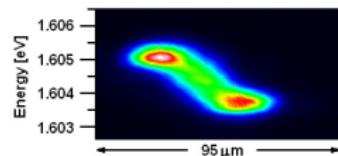
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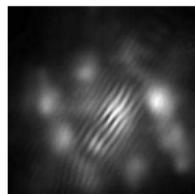
# Pattern formation in experiments

## Polariton Traps



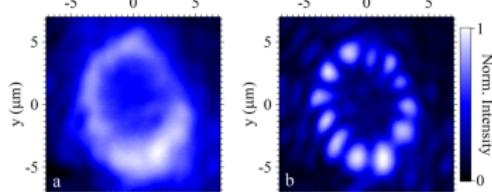
[Balili *et al.* Science '07]

## Vortex formation



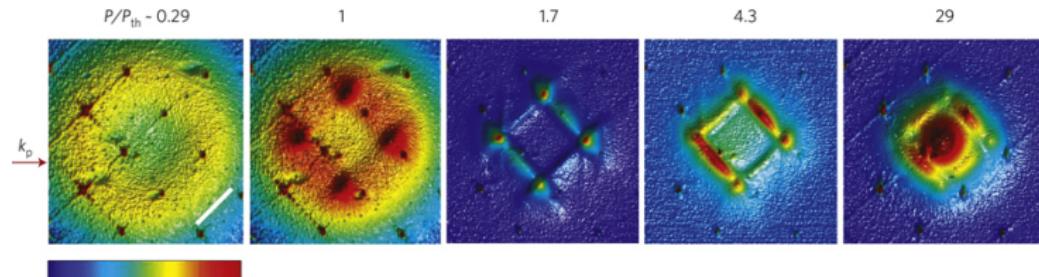
[Lagoudakis *et al.* Nat. Phys '08]

## Elliptical ring pump



[Manni *et al.* PRL '11]

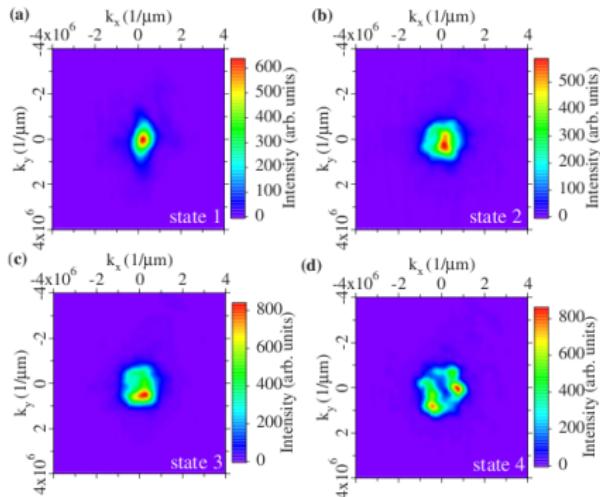
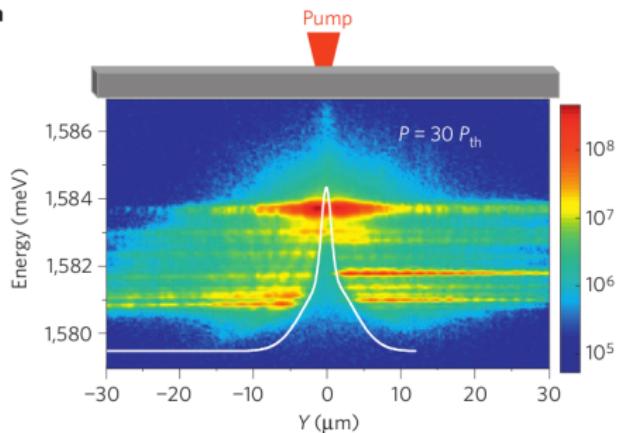
## Patterned lattice: Momentum space image



[Kim *et al.* Nat. Phys '11]

# Non-equilibrium features in experiment

a



Flow from pumping spot  
[Wertz *et al.* Nat. Phys. '10]

$|\psi(\mathbf{k})|^2 \neq |\psi(-\mathbf{k})|^2$ :  
Broken time-reversal symmetry.  
[Krizhanovskii *et al.* PRB '09]

# Complex Gross-Pitaevskii equation

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$$\left( i\partial_t + i\kappa - \left[ V(r) - \frac{\nabla^2}{2m} \right] \right) \psi(r) = \chi(\psi(r, t)) \psi(r, t)$$

Nonlinear, complex susceptibility

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$$i\partial_t \psi|_{\text{nlin}} = U|\psi|^2 \psi$$

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$$\left( i\partial_t + i\kappa - \left[ V(r) - \frac{\nabla^2}{2m} \right] \right) \psi(r) = \chi(\psi(r, t)) \psi(r, t)$$

Nonlinear, complex susceptibility

$$i\partial_t \psi|_{\text{nl}} = U|\psi|^2 \psi$$

$$i\partial_t \psi|_{\text{loss}} = -i\kappa \psi \quad i\partial_t \psi|_{\text{gain}} = i\gamma_{\text{eff}} \psi$$

# Complex Gross-Pitaevskii equation

Steady state equation:

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Nonlinear, complex susceptibility

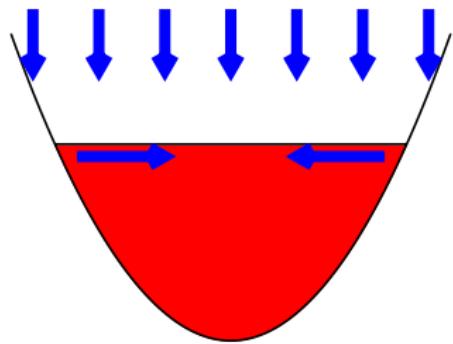
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$$i\partial_t \psi = \left[ -\frac{\nabla^2}{2m} + V(r) + U|\psi|^2 + i \left( \gamma_{\text{eff}} - \kappa - \Gamma |\psi|^2 \right) \right] \psi$$

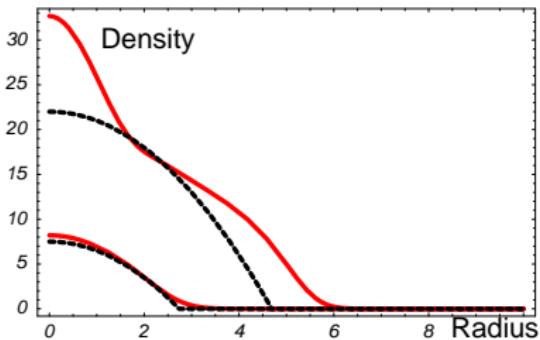
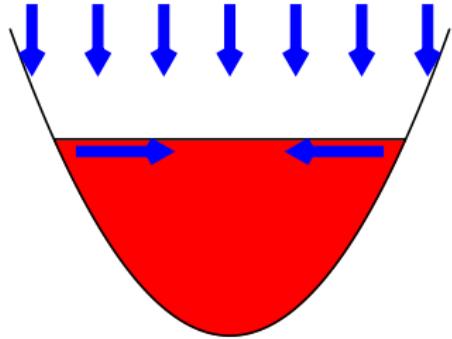
# Gross-Pitaevskii equation: Harmonic trap

$$i\partial_t \psi = \left[ -\frac{\nabla^2}{2m} + \frac{m\omega^2}{2} r^2 + U|\psi|^2 + i(\gamma_{\text{eff}} - \kappa - \Gamma|\psi|^2) \right] \psi$$



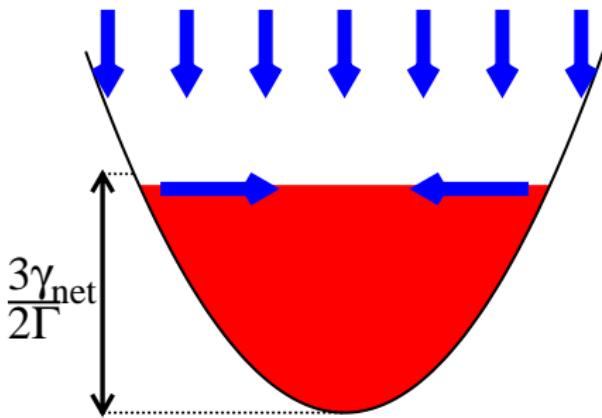
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# Stability of Thomas-Fermi solution

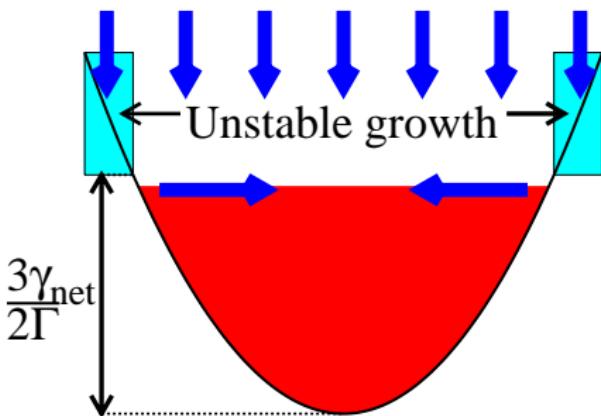
$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}} - \Gamma \rho) \rho$$



# Stability of Thomas-Fermi solution

High  $m$  modes:  $\delta\rho_{n,m} \simeq e^{im\theta} r^m \dots$

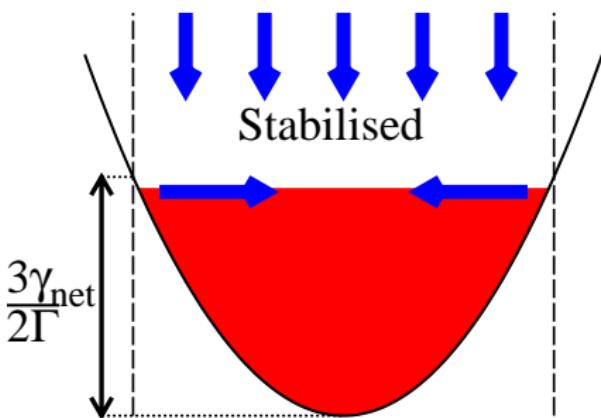
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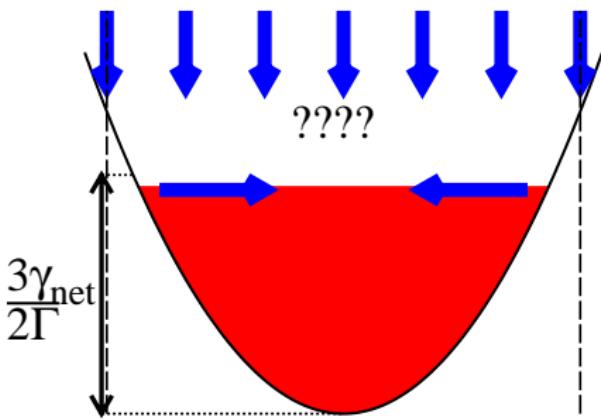
$$\frac{1}{2}\partial_t\rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}}\Theta(r_0 - r) - \Gamma\rho)\rho$$



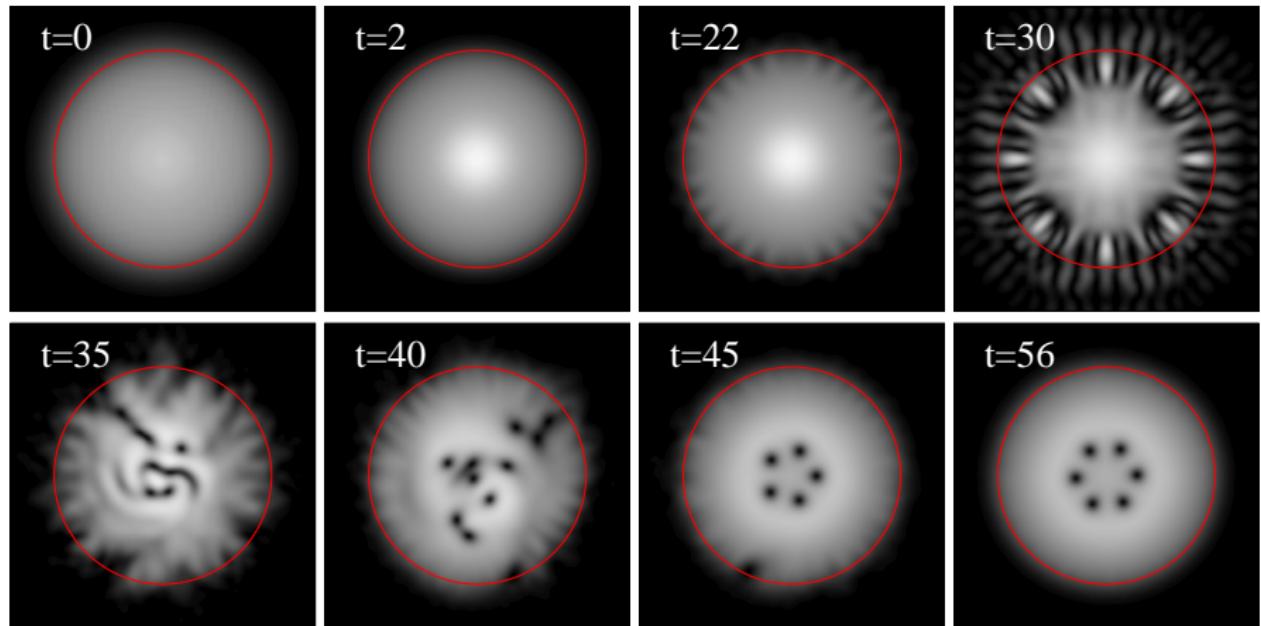
# Stability of Thomas-Fermi solution

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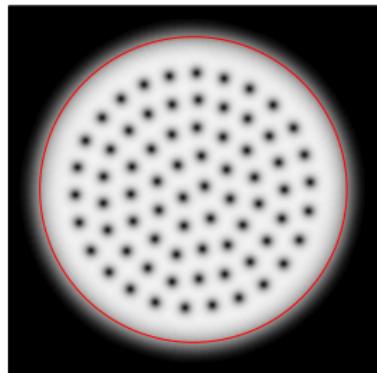
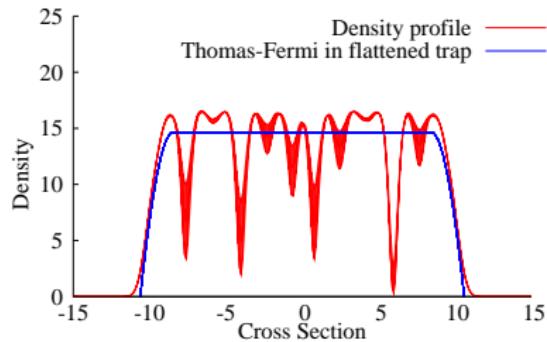


## Time evolution:



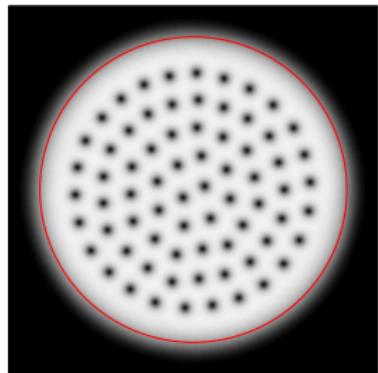
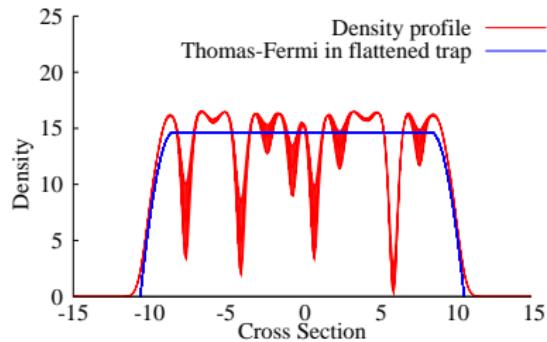
[Keeling & Berloff PRL '08]

# Why vortices



$$\nabla \cdot [p(\mathbf{v} - \mathbf{q} \times \mathbf{r})] = \delta_{\text{tot}} \Theta(r_0 - r) - U_p$$
$$p = \frac{\hbar^2}{2m} [\mathbf{v} - \mathbf{q} \times \mathbf{r}]^2 + \frac{\hbar^2}{2} (\mathbf{v}^2 - \mathbf{r}^2) + U_p - \frac{\nabla^2 \sqrt{p}}{2m\sqrt{p}}$$
$$\mathbf{q} = \mathbf{q} \times \mathbf{r}, \quad \mathbf{q} = \mathbf{q}, \quad r = \frac{\hbar^2}{p} \Theta(r_0 - r) = \frac{p}{\hbar^2}$$

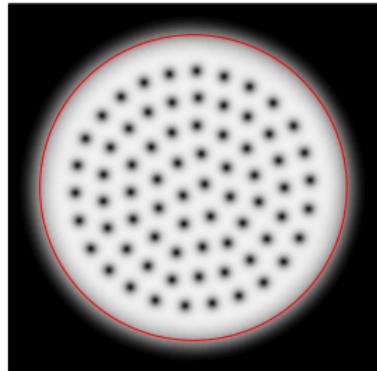
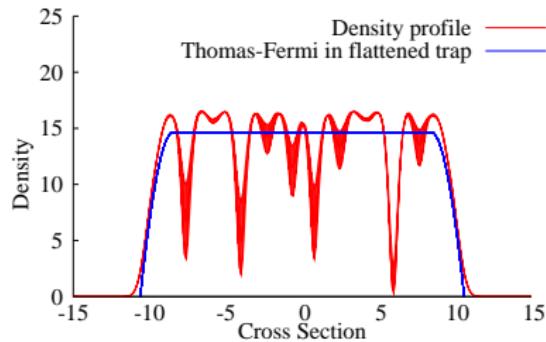
# Why vortices



Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\nabla \cdot [j(\mathbf{v} - \mathbf{g} \times \boldsymbol{\tau})] = \hbar m \Theta(n - n_f) j,$$
$$j = \frac{m}{2} |\mathbf{v} - \mathbf{g} \times \boldsymbol{\tau}|^2 + \frac{m}{2} r^2 (\omega^2 - R^2) + U_p - \frac{\nabla^2 \sqrt{j}}{2m\sqrt{j}}$$
$$\mathbf{g} = \mathbf{g} \times \boldsymbol{\tau}, \quad \mathbf{g} = \mathbf{g}, \quad \boldsymbol{\tau} = \frac{m\omega}{R} \Theta(n - n_f) = \frac{\boldsymbol{\tau}}{R}$$

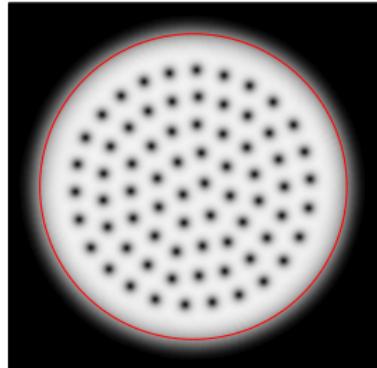
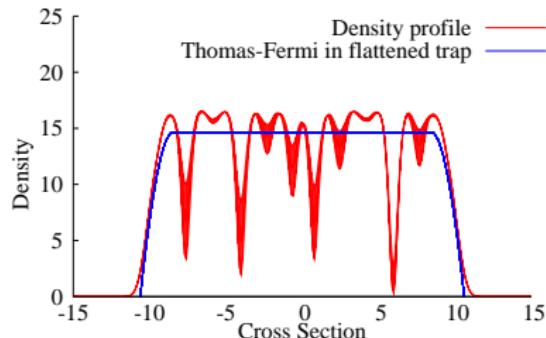
# Why vortices



Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\nabla \cdot [\rho(\mathbf{v} - \boldsymbol{\Omega} \times \mathbf{r})] = (\gamma_{\text{net}}\Theta(r_0 - r) - \Gamma\rho)\rho,$$
$$\mu = \frac{m}{2}|\mathbf{v} - \boldsymbol{\Omega} \times \mathbf{r}|^2 + \frac{m}{2}r^2(\omega^2 - \Omega^2) + U\rho - \frac{\nabla^2\sqrt{\rho}}{2m\sqrt{\rho}}$$

# Why vortices

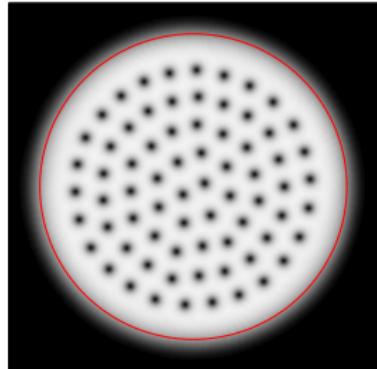
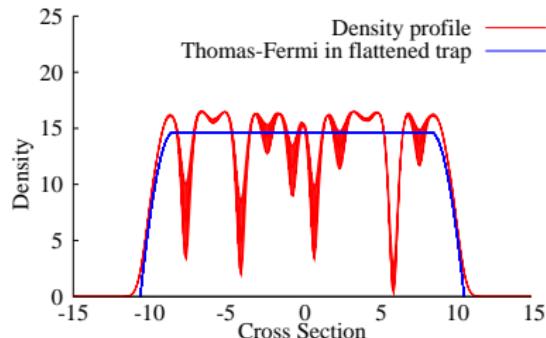


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# Why vortices



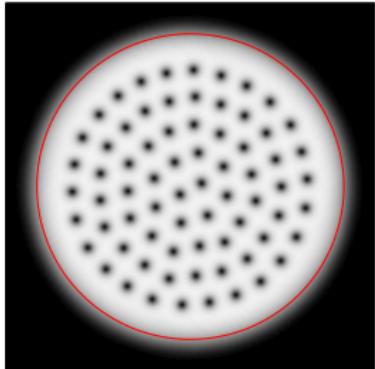
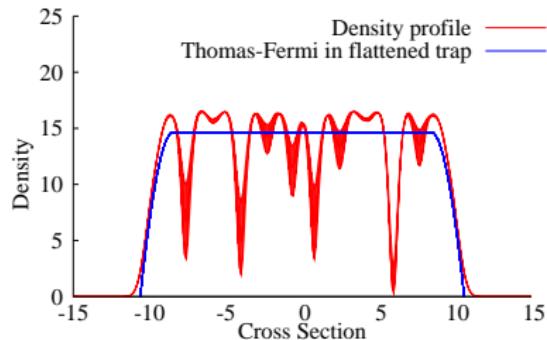
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# Why vortices



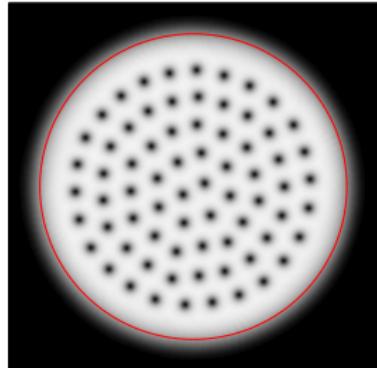
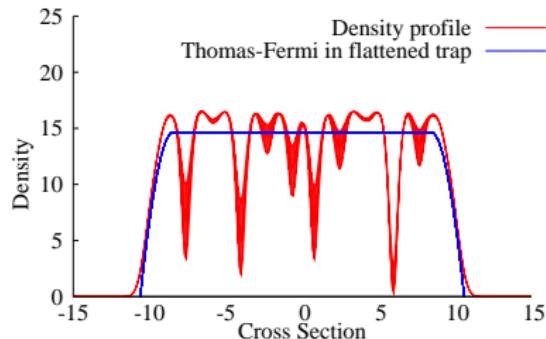
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# Why vortices



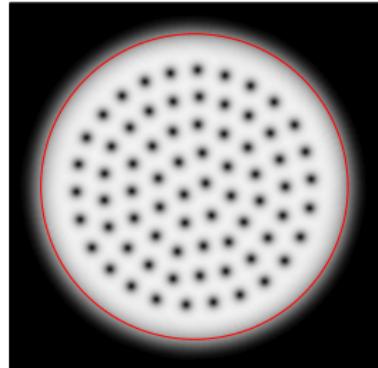
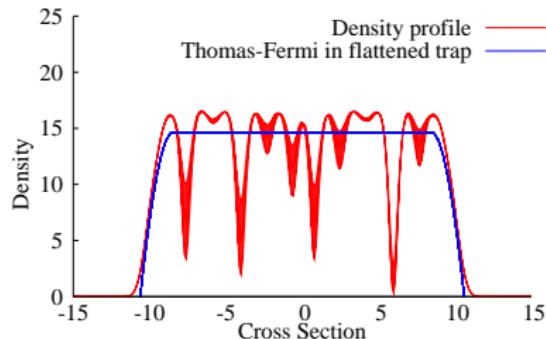
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# Why vortices



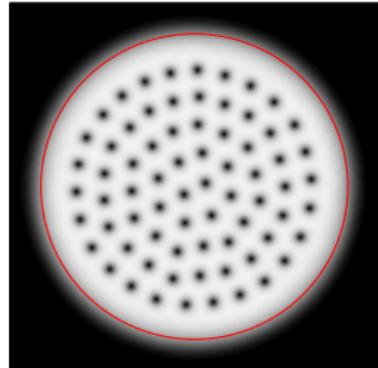
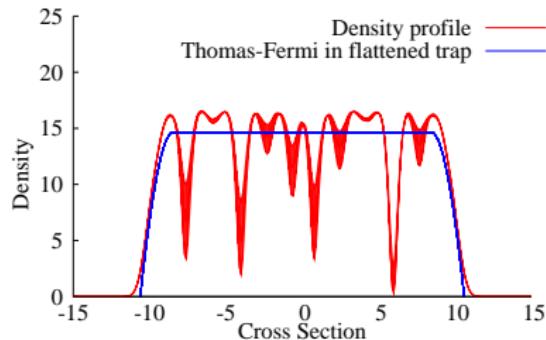
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$$\mathbf{v} = \boldsymbol{\Omega} \times \mathbf{r}, \quad \boldsymbol{\Omega} = \boldsymbol{\omega}, \quad \rho = \frac{\gamma_{\text{net}}}{\Gamma}\Theta(r_0 - r)$$

# Why vortices



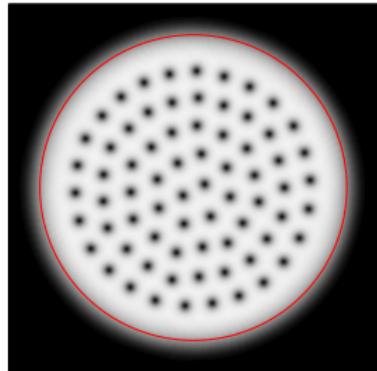
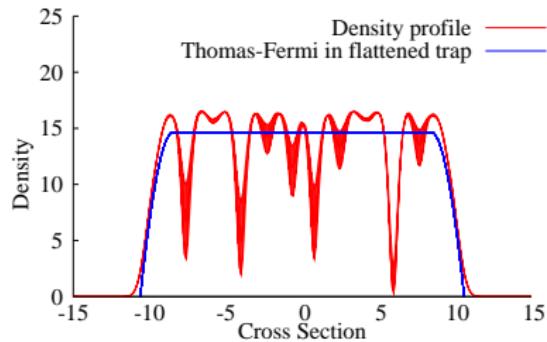
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# Why vortices



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# Superfluidity

## 1 Introduction to polariton condensation

- What are excitons and polaritons
- Experimental features
- Approaches to modelling

## 2 Pattern formation

- Experiments
- Modelling pattern formation

## 3 Superfluidity

- Non-equilibrium condensate spectrum
- Experiments and aspects of superfluidity
- Current-current response function

## 4 Coherence

- Experiments
- Power law decay of coherence

# Fluctuations above transition

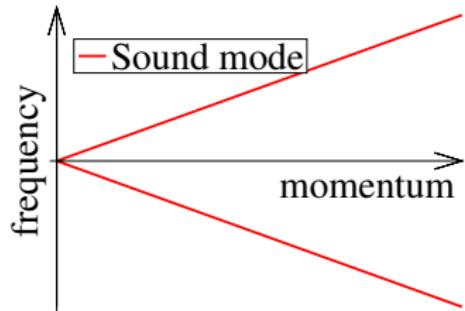
When condensed

$$\text{Det} \left[ D^R(\omega, k) \right]^{-1} = \omega^2 - \xi_k^2$$

With  $\xi_k \simeq ck$

Poles:

$$\omega^* = \xi_k$$



• Generic structure of Green's function:

$$[D^R]^{-1} = \begin{pmatrix} \omega + i\gamma_{\text{res}} - \epsilon_k - \mu & \beta_{\text{res}} - \mu \\ -\beta_{\text{res}} - \mu & -\omega - i\gamma_{\text{res}} - \epsilon_k - \mu \end{pmatrix}$$

# Fluctuations above transition

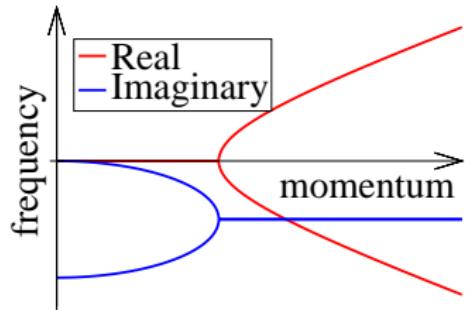
When condensed

$$\text{Det} [D^R(\omega, k)]^{-1} = (\omega + i\gamma_{\text{net}})^2 + \gamma_{\text{net}}^2 - \xi_k^2$$

With  $\xi_k \simeq ck$

Poles:

$$\omega^* = -i\gamma_{\text{net}} \pm \sqrt{\xi_k^2 - \gamma_{\text{net}}^2}$$



• Generic structure of Green's functions:

$$[D^R]^{-1} = \begin{pmatrix} \omega + i\gamma_{\text{net}} - \epsilon_k - \mu & \gamma_{\text{net}} - \mu \\ -\gamma_{\text{net}} - \mu & -\omega - i\gamma_{\text{net}} - \epsilon_k - \mu \end{pmatrix}$$

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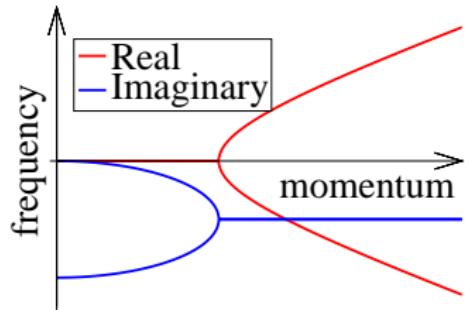
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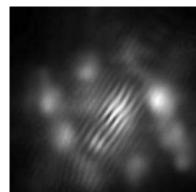


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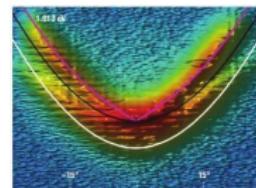
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# Polariton “superfluidity” experiments

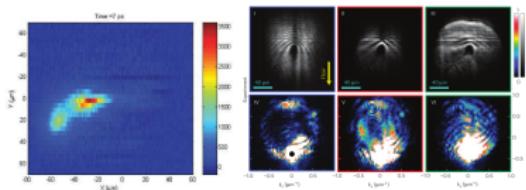
- Quantised vortices in disorder potential  
[Lagoudakis *et al.* Nature Phys. '08]



- Changes to excitation spectrum  
[Utsunomiya *et al.* Nature Phys. '08]

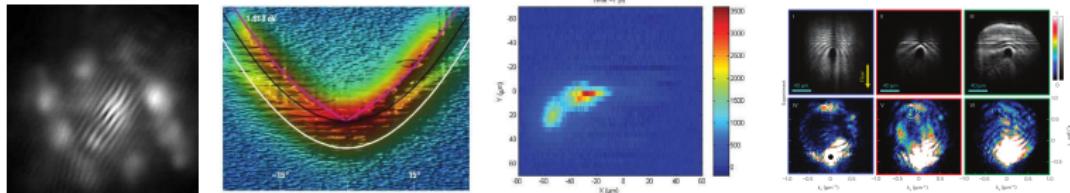


- Wavepacket propagation  
[Amo *et al.* Nature '09]
- Driven superfluidity  
[Amo *et al.* Nature Phys. ('09)]



# Aspects of superfluidity

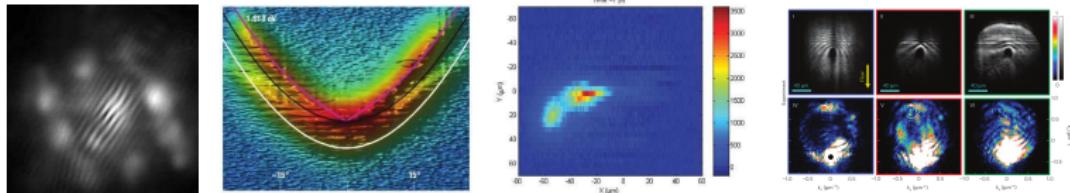
	Quantised vortices	Landau critical velocity	Metastable persistent hydro-flow	Two-fluid hydrodynamics	Local thermal equilibrium	Solitary waves
Superfluid $^4\text{He}$ /cold atom Bose-Einstein condensate	✓	✓	✓	✓	✓	✓
Non-interacting Bose-Einstein condensate	✓	✗	✗	✓	✓	✗
Classical irrotational fluid	✗	✓	✗	✓	✓	✓
Incoherently pumped polariton condensates	✓	✗	?	?	✗	?



Lagoudakis *et al.* Nat. Phys. '08. Utsunomiya *et al.* Nat. Phys. '08. Amo *et al.* Nature '09; Nat. Phys. '09

# Aspects of superfluidity

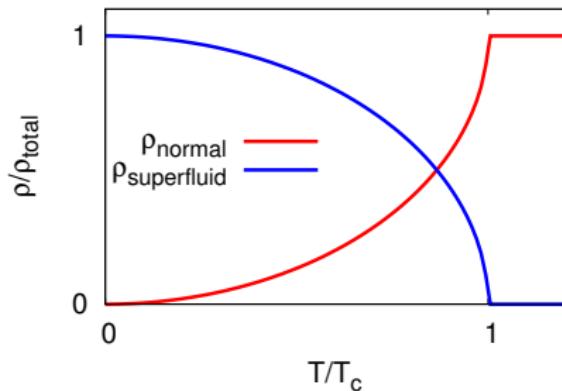
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Classical irrotational fluid	✗	✓	✗	✓	✓	✓
Incoherently pumped polariton condensates	✓	✗	?	?	✗	?



Lagoudakis *et al.* Nat. Phys. '08. Utsunomiya *et al.* Nat. Phys. '08. Amo *et al.* Nature '09; Nat. Phys. '09

# Superfluid density

- Two-fluid hydrodynamics



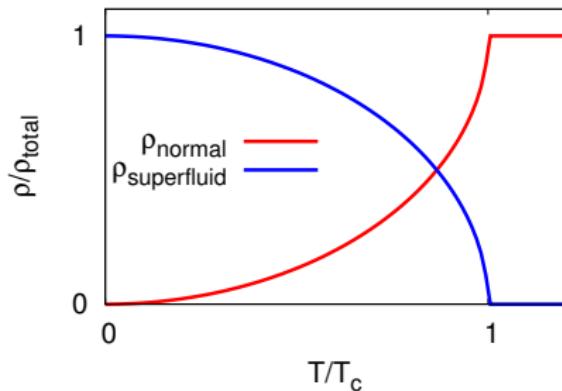
- $\rho_s, \rho_n$  distinguished by slow rotation

Experimentally, rotation:

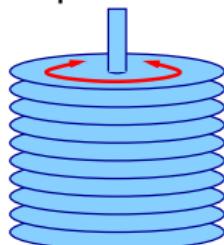
To calculate,  
transverse/longitudinal:

# Superfluid density

- Two-fluid hydrodynamics



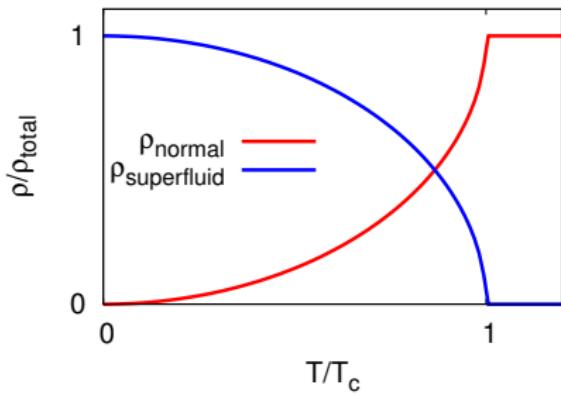
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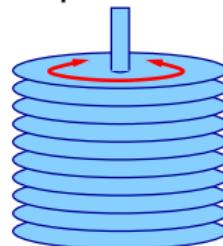
# Superfluid density

- Two-fluid hydrodynamics

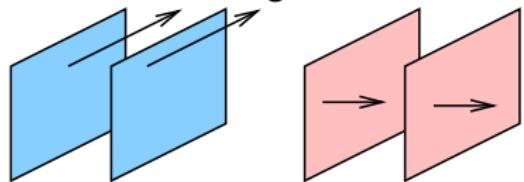


- $\rho_s, \rho_n$  distinguished by slow rotation

- Experimentally, rotation:



- To calculate, transverse/longitudinal:



# Superfluid density

- Current:

$$\mathbf{J} = \rho \mathbf{v} = \psi^\dagger i \nabla \psi = |\psi|^2 \nabla \phi$$

$$J_i(\mathbf{q}) = \psi_{\mathbf{k}+\mathbf{q}}^\dagger \frac{2k_i + q_i}{2m} \psi_{\mathbf{k}}$$

- Response functions:

$$H \rightarrow H - \sum_{\mathbf{q}} \chi(\mathbf{q}) \cdot \mathbf{J}(\mathbf{q}) \quad J(\mathbf{q}) = \chi_J(\mathbf{q}) / (\mathbf{q})$$

- Vertex corrections essential for superfluid part.

# Superfluid density

- Current:

$$\mathbf{J} = \rho \mathbf{v} = \psi^\dagger i \nabla \psi = |\psi|^2 \nabla \phi$$

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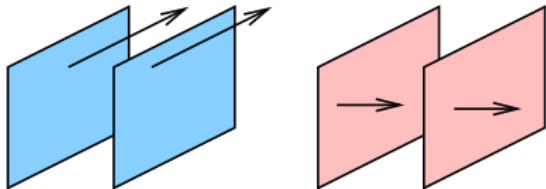
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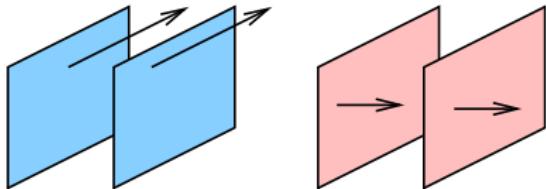
$$\chi_{ij}(\omega = 0, \mathbf{q} \rightarrow 0) = \langle [J_i(\mathbf{q}), J_j(-\mathbf{q})] \rangle = \frac{\rho_s}{m} \frac{q_i q_j}{q^2} + \frac{\rho_N}{m} \delta_{ij}$$

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# Calculating superfluid response function

- Using Keldysh generating functional

$$\chi_{ij}(q) = -\frac{i}{2} \frac{d^2 \mathcal{Z}[f, \theta]}{df_i(q)d\theta_j(-q)}, \quad \mathcal{Z}[f, \theta] = \int \mathcal{D}\psi \exp(iS[f, \theta])$$

• Example: Superfluid density

$$S[f, \theta] = S + \sum_{k,q} (\tilde{\rho}_k - \tilde{\rho}_q)_{k+q} \begin{pmatrix} \theta_k & \theta_{k+q} \\ f_k - \theta_k & -f_{k+q} \end{pmatrix}_q \frac{2k+q}{2\pi} \begin{pmatrix} \psi_k \\ \psi_q \end{pmatrix}_k$$

- Saddle point + fluctuations:

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→ Saddle point + fluctuations

# Calculating superfluid response function

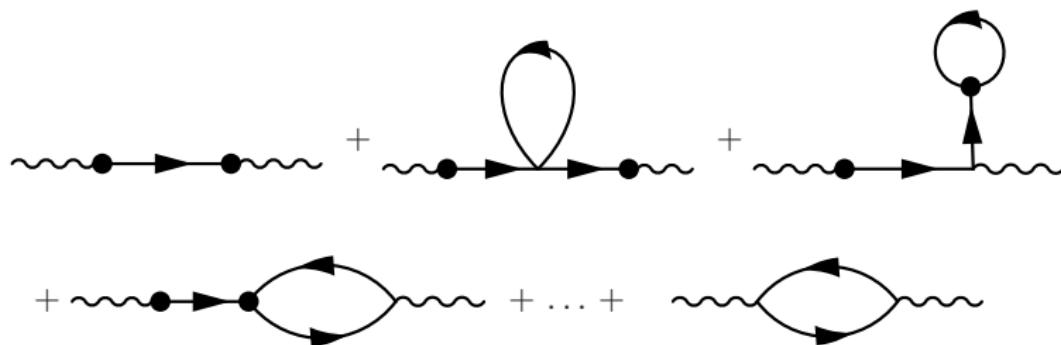
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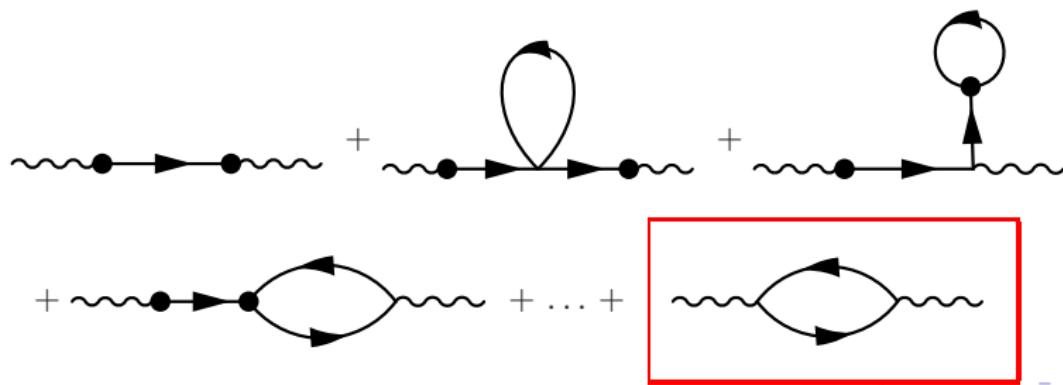
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- Saddle point + fluctuations: Only one diagram for  $\chi_N$



# Non-equilibrium superfluid response

- Superfluid response exists because:

$$\text{Diagram: } \text{Two wavy lines with dots at vertices, connected by a horizontal line with an arrow pointing right.} = \frac{i\psi_0 q_i}{2m} (1, -1) D^R(q, \omega = 0) \begin{pmatrix} 1 \\ -1 \end{pmatrix} \frac{i\psi_0 q_j}{2m}$$

•  $D^R(\omega = 0) \propto 1/\epsilon_0$ , despite pumping/decay — superfluid response exists.

- Normal density:

$$n = \int d^d k \epsilon_k \int \frac{d\omega}{2\pi} \text{Tr} \left[ \sigma_z D^K \sigma_z (D^R + D^A) \right]$$

- Is affected by pump/decay:  
Does not vanish at  $T \rightarrow 0$ .

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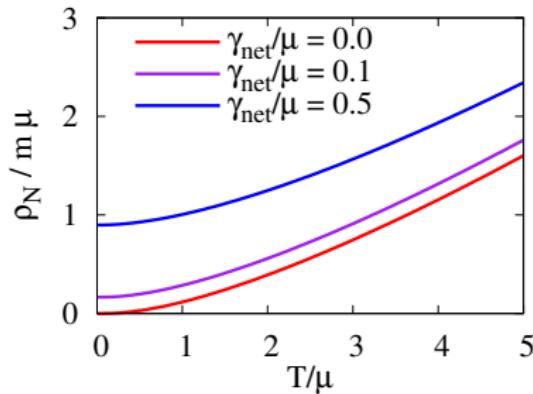
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[JK PRL '11]



# Coherence:

## 1 Introduction to polariton condensation

- What are excitons and polaritons
- Experimental features
- Approaches to modelling

## 2 Pattern formation

- Experiments
- Modelling pattern formation

## 3 Superfluidity

- Non-equilibrium condensate spectrum
- Experiments and aspects of superfluidity
- Current-current response function

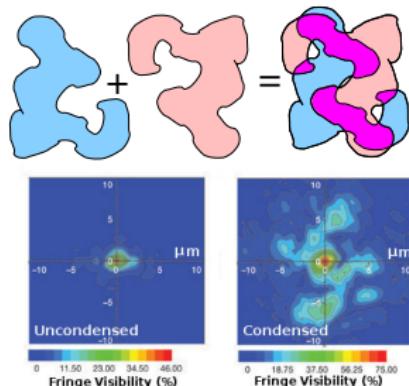
## 4 Coherence

- Experiments
- Power law decay of coherence

# Correlations in a 2D Gas

Correlations:

$$g_1(\mathbf{r}, \mathbf{r}', t) = \langle \psi^\dagger(\mathbf{r}, t) \psi(\mathbf{r}', 0) \rangle$$



$$\rightarrow D^L = D^R = D^B + D^A$$

→ Generally get

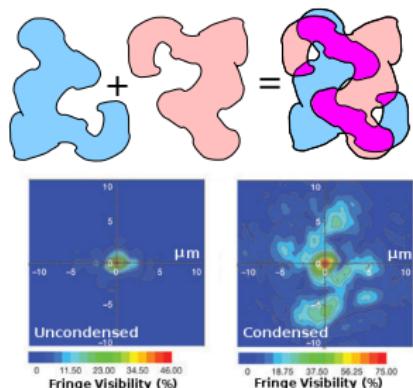
$$\langle \psi^\dagger(\mathbf{r}, t) \psi(0, 0) \rangle = |k_0|^2 \exp \left[ -2\pi \sqrt{\frac{\ln(t/t_0)}{2 \ln(e^2/k_B T_0)}} \right] \quad t > 0$$

[Szymańska *et al.* PRL '06; PRB '07]

# Correlations in a 2D Gas

Correlations: (in 2D)

$$g_1(\mathbf{r}, \mathbf{r}', t) = \langle \psi^\dagger(\mathbf{r}, t) \psi(\mathbf{r}', 0) \rangle \\ \simeq |\psi_0|^2 \exp \left[ -D_{\phi\phi}^<(\mathbf{r}, \mathbf{r}', t) \right]$$



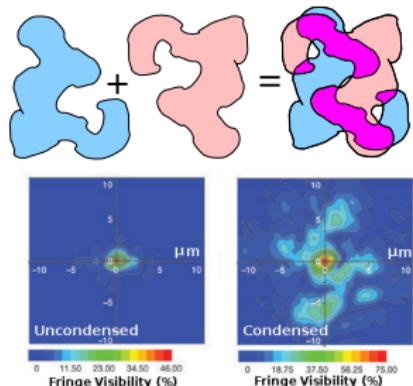
- $D^< = D^K - D^R + D^A$

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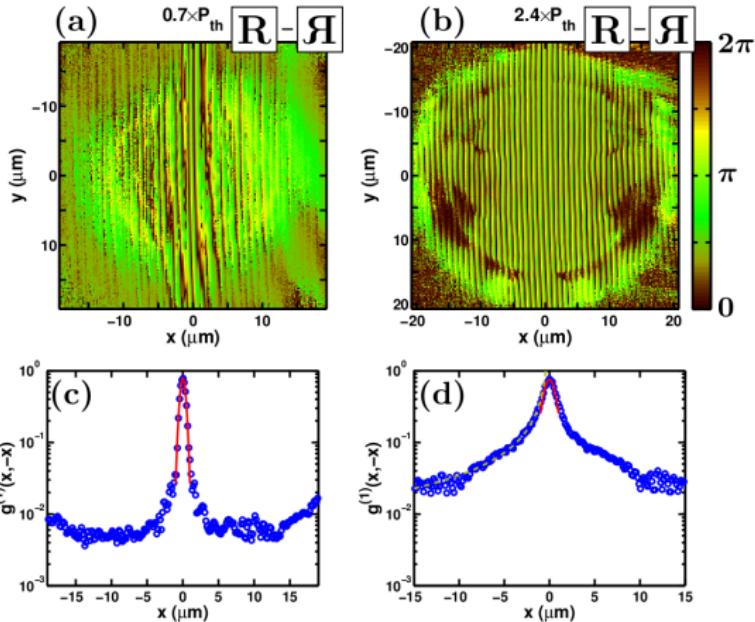


- $D^< = D^K - D^R + D^A$
- Generally, get:

$$\langle \psi^\dagger(\mathbf{r}, t) \psi(0, 0) \rangle \simeq |\psi_0|^2 \exp \begin{cases} \ln(r/r_0) & t \simeq 0 \\ \frac{1}{2} \ln(c^2 t / \gamma_{\text{net}} r_0^2) & r \simeq 0 \end{cases}$$

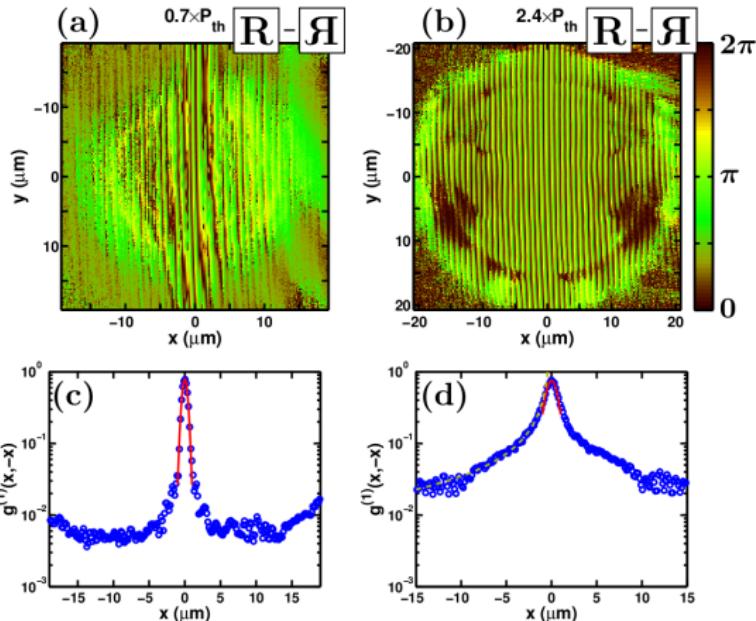
[Szymańska *et al.* PRL '06; PRB '07]

# Experimental observation of power-law decay

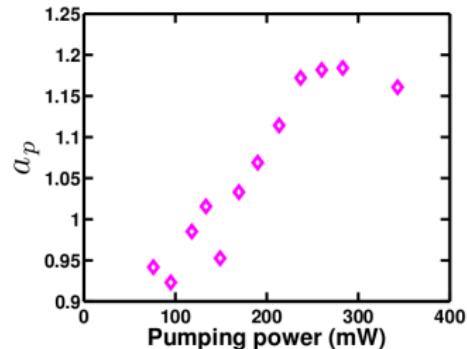


G. Rompos, Y. Yamamoto *et al.* submitted

# Experimental observation of power-law decay



$$g_1(\mathbf{r}, -\mathbf{r}) \propto \left(\frac{r}{r_0}\right)^{-a_p}$$



G. Rompos, Y. Yamamoto *et al.* submitted

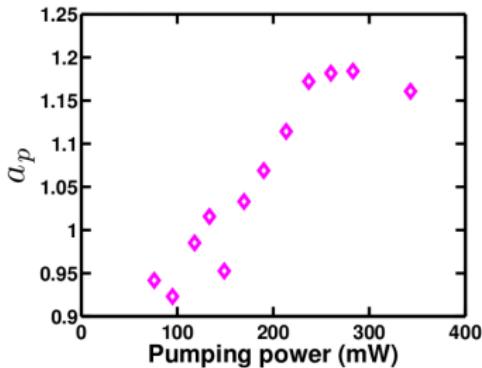
# Exponent in a non-equilibrium 2D gas

$$\lim_{r \rightarrow \infty} \langle \psi^\dagger(\mathbf{r}, 0) \psi(-\mathbf{r}, 0) \rangle = |\psi_0|^2 \exp \left[ -D_{\phi\phi}^<(\mathbf{r}, -\mathbf{r}) \right] \propto \exp \left[ -a_p \ln \left( \frac{2r}{r_0} \right) \right]$$

- Experimentally,  $a_P \simeq 1.2$

• In equilibrium  $a_p = \frac{m k_B T}{2\pi \hbar^2 n_s} < \frac{1}{4}$  (BKT transition)

• Non-equilibrium theory depends on thermalisation.

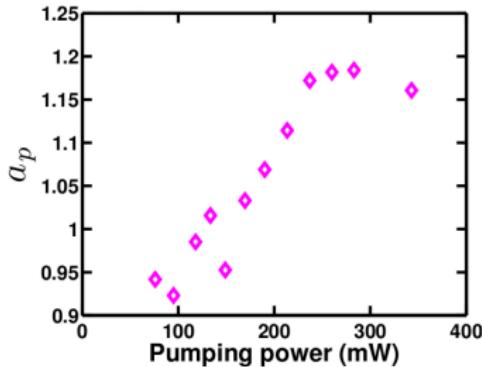


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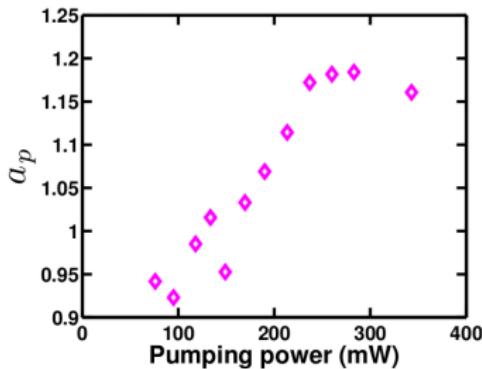
$m k_B T$

$B_p = \frac{\partial \mu}{\partial n_s}$

– Non-thermalised,

Pumping noise

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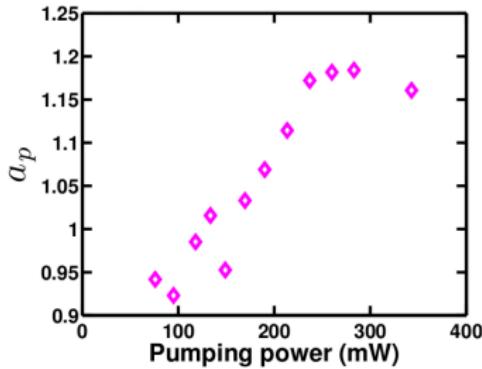


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Nanostructured  
polariton noise  
and  
non-equilibrium  
condensation

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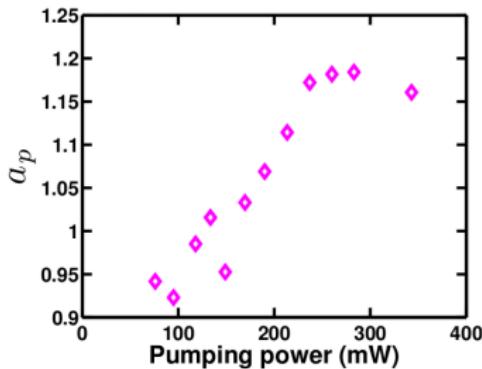
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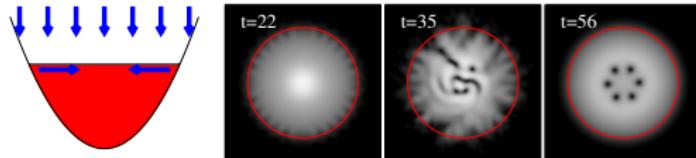
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Pumping noise  
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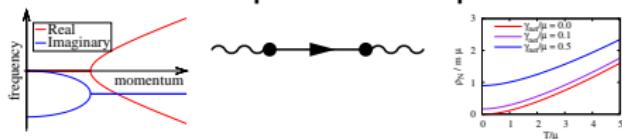


# Conclusion

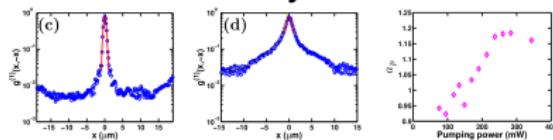
- Instability of Thomas-Fermi and spontaneous rotation



- Survival of superfluid response



- Power law decay of correlations





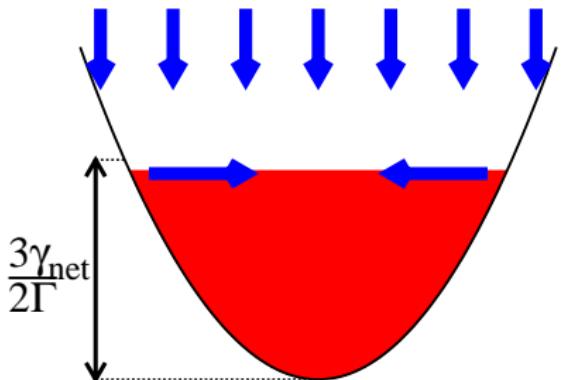
# Extra slides

- 5 GPE stability
- 6 Detecting vortex lattice
- 7 Measuring superfluid density
- 8 Coherence Finite size and Schawlow-Townes

# Instability of Thomas-Fermi: details

$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}} - \Gamma \rho) \rho$$

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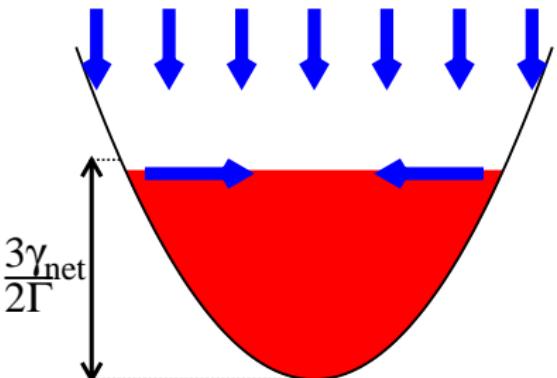


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Normal modes for  $\gamma_{\text{net}}, \Gamma \rightarrow 0$ :



$$\delta \rho_{n,m}(r, \theta, t) = e^{im\theta} h_{n,m}(r) e^{i\omega_{n,m}t}$$

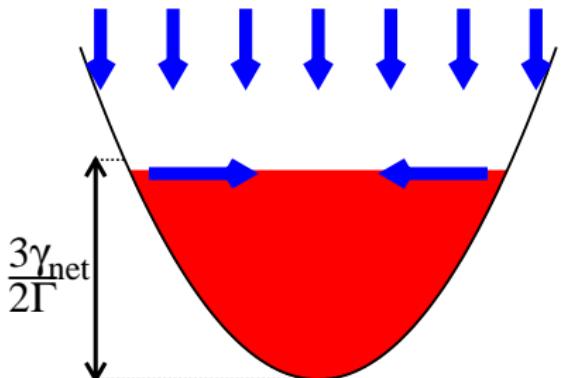
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Add weak pumping/decay:

$$\omega_{n,n} \rightarrow \omega_{n,m} + i\gamma_{\text{net}} \left[ \frac{m(1+2n) + 2n(n+1) - m^2}{2m(1+2n) + 4n(n+1) + m^2} \right]$$

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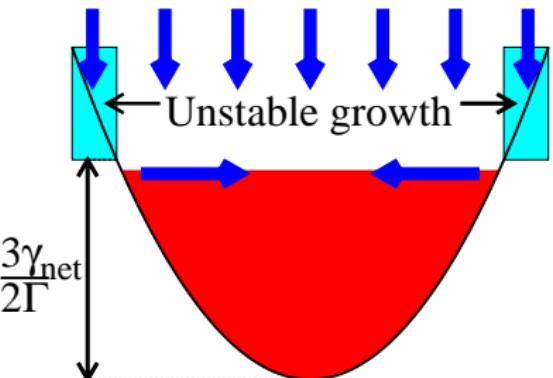
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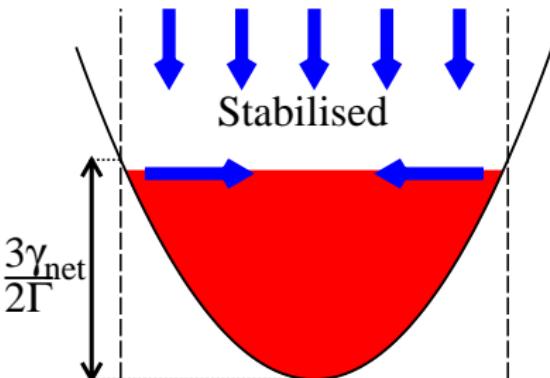
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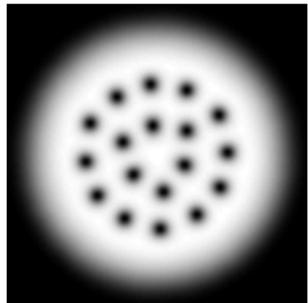
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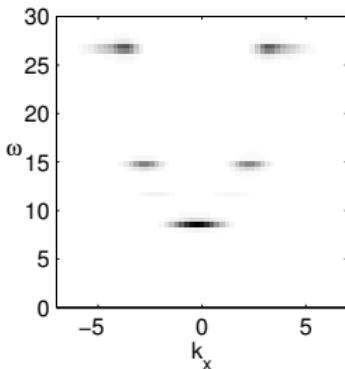
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# Detecting vortex lattices

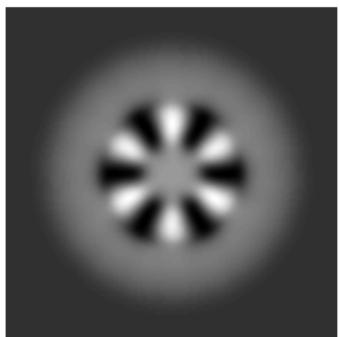
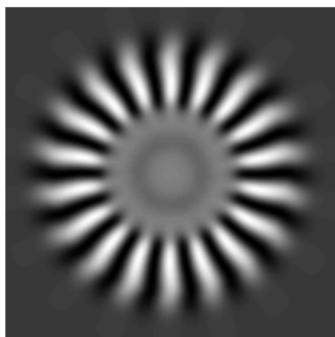
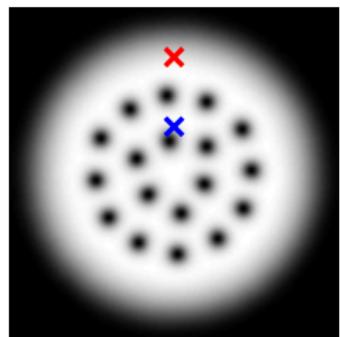
Snapshot



Spectrum:



Defocussed homodyne interference:

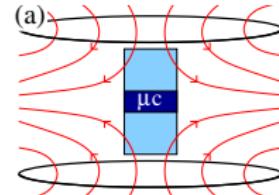


# Measuring superfluid density

## 1. Effect rotating frame

Polariton polarization:  $(\psi_{\circlearrowleft}, \psi_{\circlearrowright})$

$$H = \lambda \begin{pmatrix} \ell^2 & r^2 e^{2i\phi} \\ r^2 e^{-2i\phi} & -\ell^2 \end{pmatrix}$$



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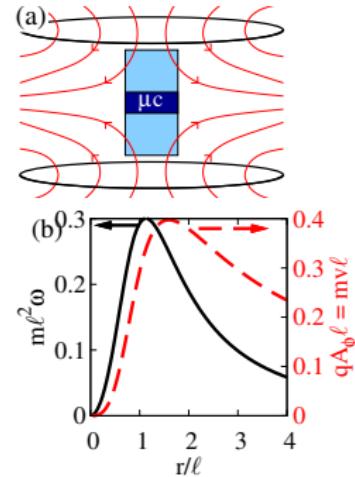
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Ground state Berry phase:

$$q\mathbf{A}_{\text{eff}} = m\omega \times \mathbf{r} = \frac{\hat{\phi}}{r} \left[ 1 - \frac{\ell^2}{\sqrt{r^4 + \ell^4}} \right]$$



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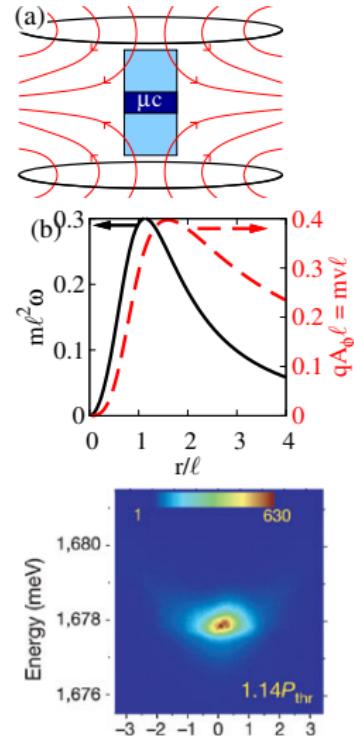
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## 2. Measure resulting current

Energy shift of normal state:

$$\Delta E = (1/2)mv^2 = 0.08/m\ell^2 \simeq 0.1 \text{ meV}$$

## Finite size effects: Single mode vs many mode

$$\langle \psi^\dagger(\mathbf{r}, t) \psi(\mathbf{r}', 0) \rangle \simeq |\psi_0|^2 \exp \left[ -D_{\phi\phi}^<(\mathbf{r}, \mathbf{r}', t) \right]$$

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$D_{\phi\phi}^<(\mathbf{r}, \mathbf{r}', t)$  from sum of phase modes. Study  $ct \gg r$  limit:

$$D_{\phi\phi}^<(\mathbf{r}, \mathbf{r}, t) \propto \sum_n^{n_{max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(\mathbf{r})|^2 (1 - e^{i\omega t})}{|(\omega + i\gamma_{\text{net}})^2 + \gamma_{\text{net}}^2 - \xi_n^2|^2}$$

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$$\Delta\xi \ll \sqrt{\frac{\gamma_{\text{net}}}{t}} \ll E_{\text{max}}$$



$$D_{\phi\phi}^< \sim 1 + \ln(E_{\text{max}} \sqrt{\frac{t}{\gamma_{\text{net}}}})$$

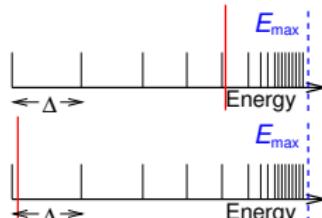
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(Recovers Schawlow-Townes laser linewidth)

$$D_{\phi\phi}^< \sim 1 + \ln(E_{\max}) \sqrt{\frac{t}{\gamma_{\text{net}}}}$$

$$D_{\phi\phi}^< \sim \left( \frac{\pi C}{2\gamma_{\text{net}}} \right) \left( \frac{t}{2\gamma_{\text{net}}} \right)$$