

Condensation, superfluidity, and lasing of coupled light-matter systems.

Jonathan Keeling



IOP TCM Meeting, June 2011



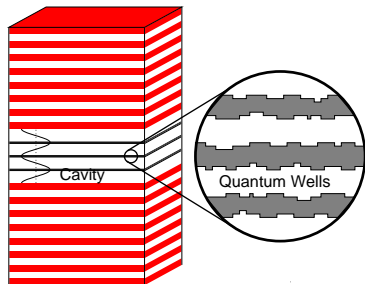
Funding:

EPSRC

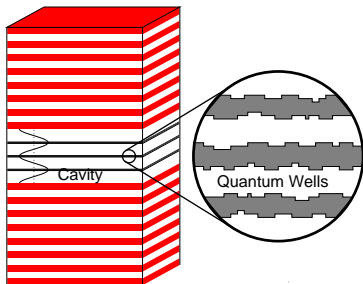
Engineering and Physical Sciences
Research Council



Microcavity polaritons

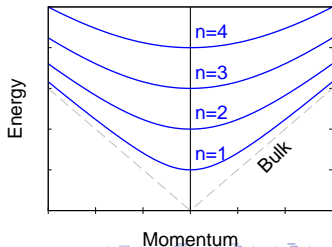


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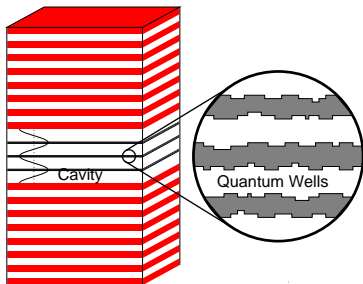


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2/2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$

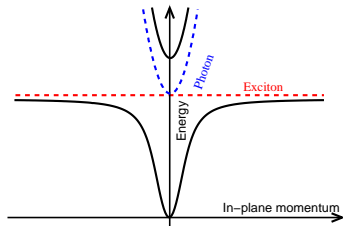


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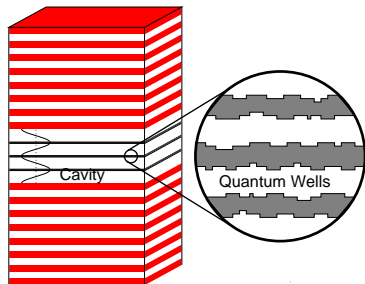


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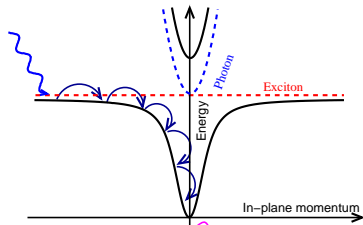


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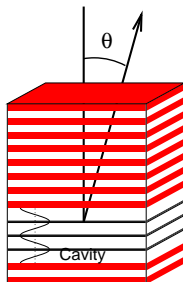
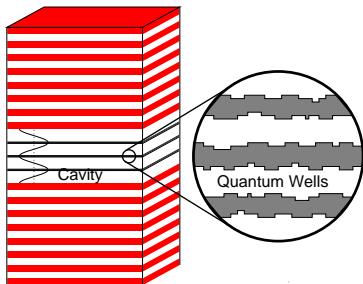


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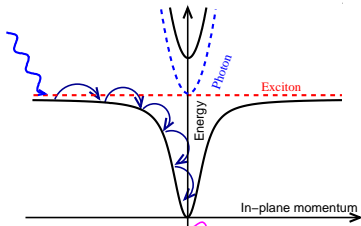


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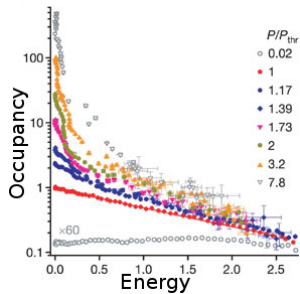
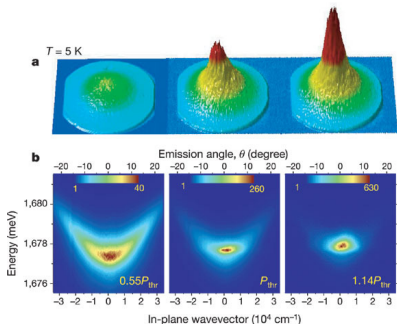


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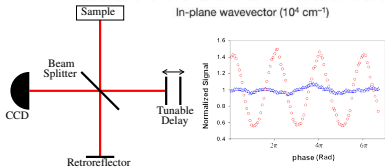
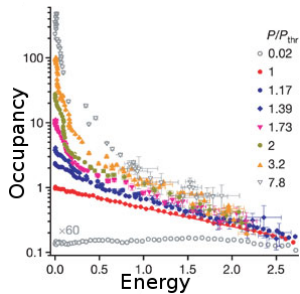
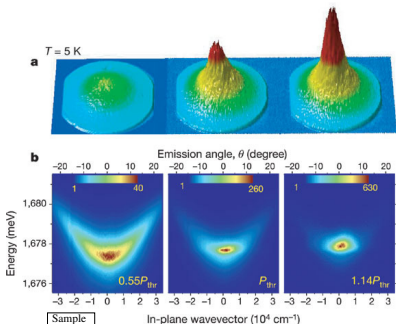


Polariton experiments: occupation and coherence

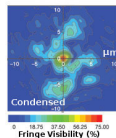
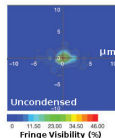
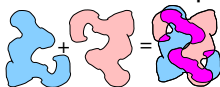


[Kasprzak, *et al.* Nature, 2006]

Polariton experiments: occupation and coherence



Coherence map:

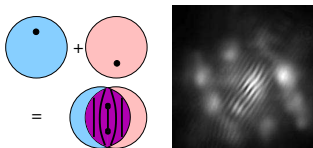


[Kasprzak, *et al.* Nature, 2006]

(Some) other polariton condensation experiments

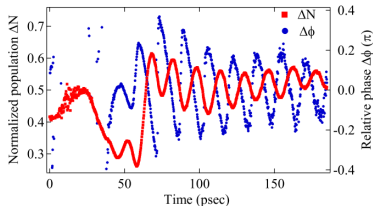
- Quantised vortices

[Lagoudakis *et al.* *Nat. Phys.* '08. *Science* '09, *PRL* '10; Sanvitto *et al.* *Nat. Phys.* '10; Roumpos *et al.* *Nat. Phys.* '10]



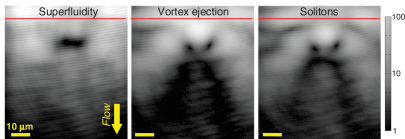
- Josephson oscillations

[Lagoudakis *et al.* *PRL* '10]



- Pattern formation/Hydrodynamics

[Amo *et al.* *Science* '11, *Nature* '09; Wertz *et al.* *Nat. Phys* '10]



1 Microscopic model and lasing vs condensation

- Introducing model
- Stability of normal state
- Comparison to standard laser

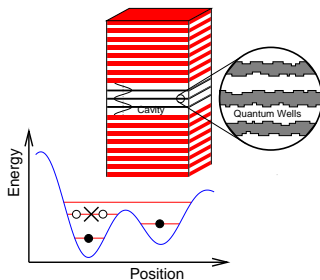
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- Condensed spectrum
- Aspects of superfluidity
- Current-current response function
- Power law decay of coherence

Microscopic model

Polariton model

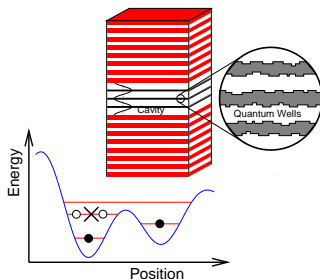
- Disorder-localised excitons
- Treat sites as 2-level systems (exciton/no-exciton)
- Propagating (2D) photons



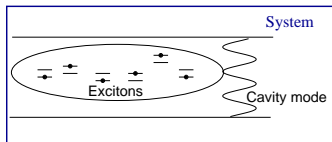
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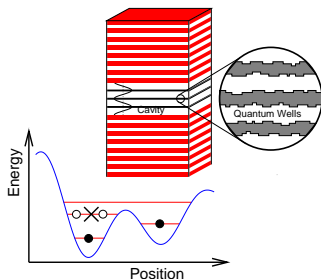
$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} \left[\epsilon_{\alpha} \left(b_{\alpha}^{\dagger} b_{\alpha} - a_{\alpha}^{\dagger} a_{\alpha} \right) + \frac{1}{\sqrt{A}} g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} b_{\alpha}^{\dagger} a_{\alpha} + \text{H.c.} \right]$$



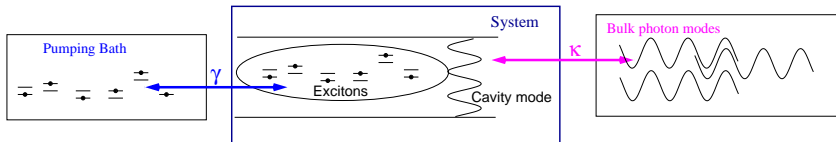
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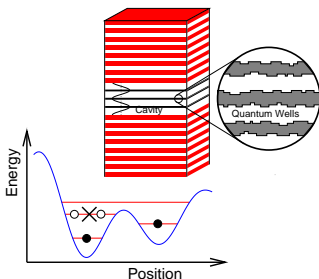
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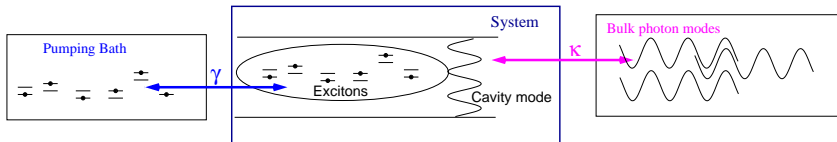
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Pumping bath thermal, μ_B, T_B , Photon bath empty

Luminescence spectrum and Green's functions

$$[D^R - D^A](t, t') = -i \left\langle \left[\psi(t), \psi^\dagger(t') \right]_- \right\rangle$$

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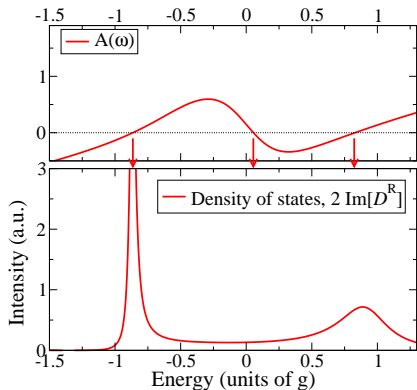
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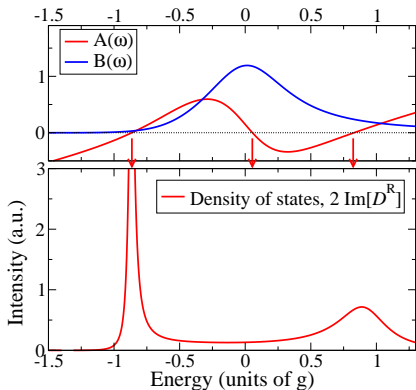
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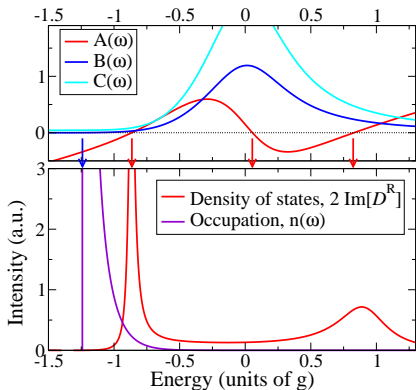
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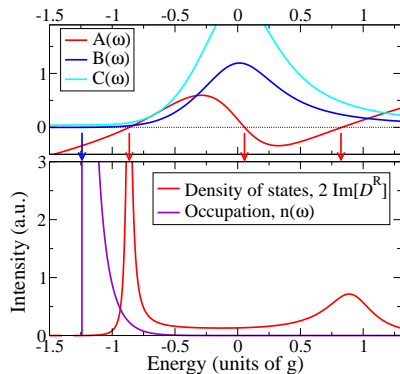
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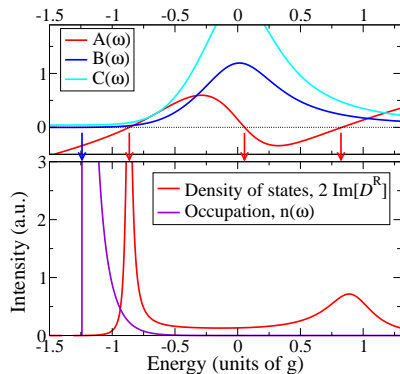
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Stability and evolution with pumping

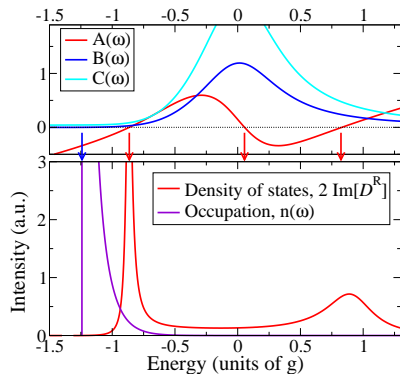


Stability and evolution with pumping



$$\left[D^R(\omega) \right]^{-1} = (\omega - \xi_k) + i\alpha(\omega - \mu_{\text{eff}})$$

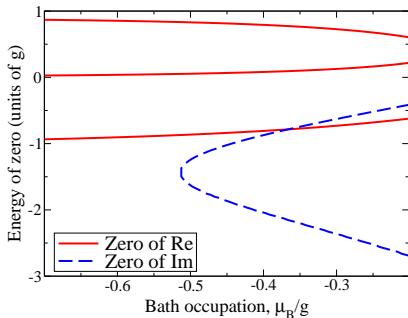
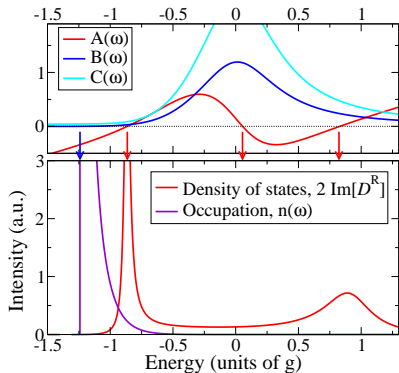
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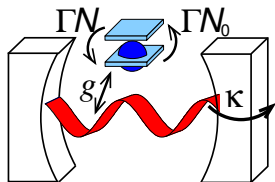
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Green's function and stability for a laser



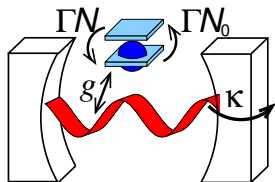
Maxwell-Bloch equations:

$$\partial_t \psi = -i\omega_k \psi - \kappa \psi + gP$$

$$\partial_t P = -2i\epsilon P - 2\gamma P + g\psi N$$

$$\partial_t N = 2\gamma(N_0 - N) - 2g(\psi^* P + P^* \psi)$$

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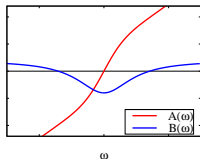
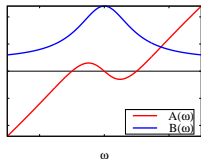
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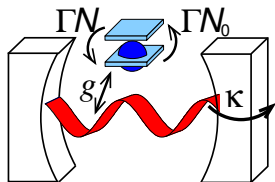
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$$[D^R(\omega)]^{-1} = \omega - \omega_k + i\kappa + \frac{g^2 N_0}{\omega - 2\epsilon + i2\gamma}$$



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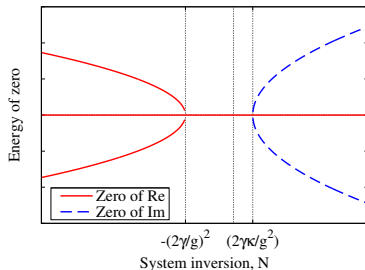
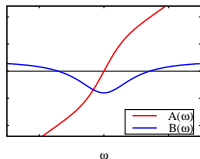
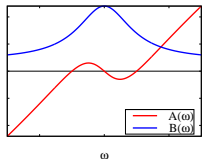
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Spectrum above transition

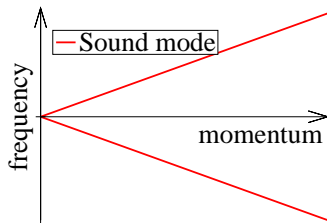
When condensed

$$\text{Det} \left[D^R(\omega, k) \right]^{-1} = \omega^2 - \xi_k^2$$

With $\xi_k \simeq ck$

Poles:

$$\omega^* = \xi_k$$



Spectrum above transition

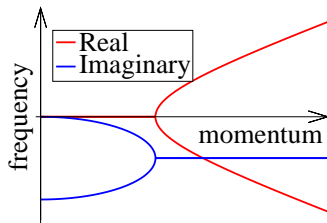
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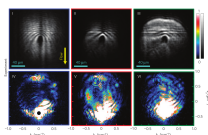
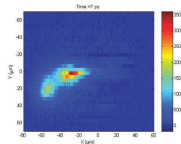
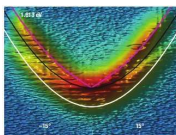
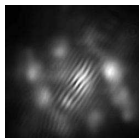
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$$\omega^* = -i\gamma_{\text{net}} \pm \sqrt{\xi_k^2 - \gamma_{\text{net}}^2}$$



Aspects of superfluidity

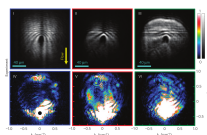
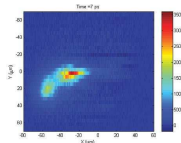
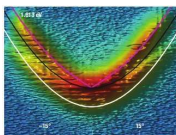
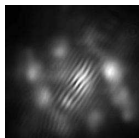
	Quantised vortices	Landau critical velocity	Metastable persistent flow	Two-fluid hydrodynamics	Local thermal equilibrium	Solitary waves
Superfluid ^4He /cold atom Bose-Einstein condensate	✓	✓	✓	✓	✓	✓
Non-interacting Bose-Einstein condensate	✓	✗	✗	✓	✓	✗
Classical irrotational fluid	✗	✓	✗	✓	✓	✓
Incoherently pumped polariton condensates	✓	✗	?	?	✗	?



Lagoudakis *et al.* Nat. Phys. '08. Utsunomiya *et al.* Nat. Phys. '08. Amo *et al.* Nature '09; Nat. Phys. '09

Aspects of superfluidity

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Superfluid ^4He /cold atom Bose-Einstein condensate	✓	✓	✓	✓	✓	✓
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Classical irrotational fluid	✗	✓	✗	✓	✓	✓
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Superfluid density

- Current:

$$\mathbf{J} = \rho \mathbf{v} = \psi^\dagger i \nabla \psi = |\psi|^2 \nabla \phi$$

$$J_i(\mathbf{q}) = \psi_{\mathbf{k}+\mathbf{q}}^\dagger \frac{2k_i + q_i}{2m} \psi_{\mathbf{k}}$$

- Response function:

$$H \rightarrow H - \sum_{\mathbf{q}} \mathbf{f}(\mathbf{q}) \cdot \mathbf{J}_i(\mathbf{q}) \quad J_i(\mathbf{q}) = \chi_{ij}(\mathbf{q}) f_j(\mathbf{q})$$

- Vertex corrections essential for superfluid part.

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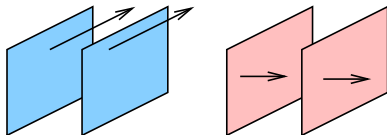
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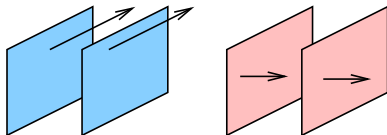
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Non-equilibrium current response functions

- Superfluid response exists because:

$$\text{---}\bullet\text{---}\rightarrow\text{---}\bullet\text{---} = \left(\frac{i\psi_0 q_i}{2m}\right) D^R(q, \omega = 0) \left(\frac{i\psi_0 q_j}{2m}\right)$$

- $D^R(\omega = 0) \propto 1/q^2$ despite pumping/decay — superfluid response exists.
- Normal density:

$$\rho_N = \int d^d k \epsilon_k \int \frac{d\omega}{2\pi} \text{Tr} \left[\sigma_z D^R \sigma_z (D^R + D^A) \right]$$

- Is affected by pump/decay:
Does not vanish at $T \rightarrow 0$.

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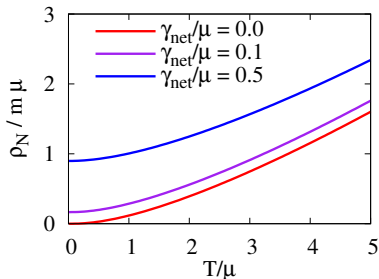
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1 Microscopic model and lasing vs condensation

- Introducing model
- Stability of normal state
- Comparison to standard laser

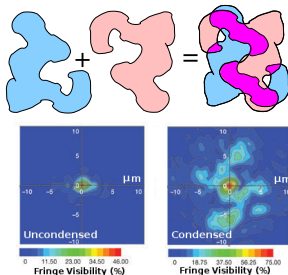
2 Coherence and superfluidity of condensate

- Condensed spectrum
- Aspects of superfluidity
- Current-current response function
- **Power law decay of coherence**

Correlations in a 2D Gas

Correlations:

$$g_1(\mathbf{r}, \mathbf{r}', t) = \langle \psi^\dagger(\mathbf{r}, t) \psi(0, \mathbf{r}') \rangle$$



$$\bullet D^< = D^K - D^R + D^A$$

• Generally, get: $\langle \psi^\dagger(\mathbf{r}, t) \psi(0, 0) \rangle \simeq$

$$|\psi_0|^2 \exp \left[-a_p \begin{cases} \ln(r/r_0) & r \rightarrow \infty, t \simeq 0 \\ \frac{1}{2} \ln(c^2 t / \lambda_{\text{net}}^2) & r \simeq, t \rightarrow \infty \end{cases} \right]$$

[Szymańska *et al.* PRL '06; PRB '07]

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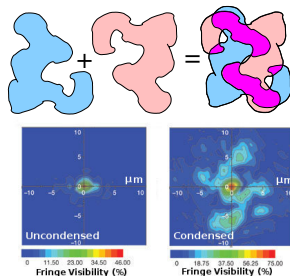
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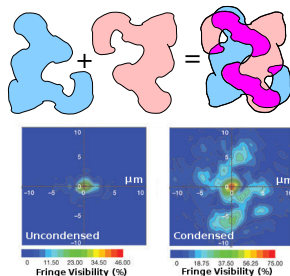
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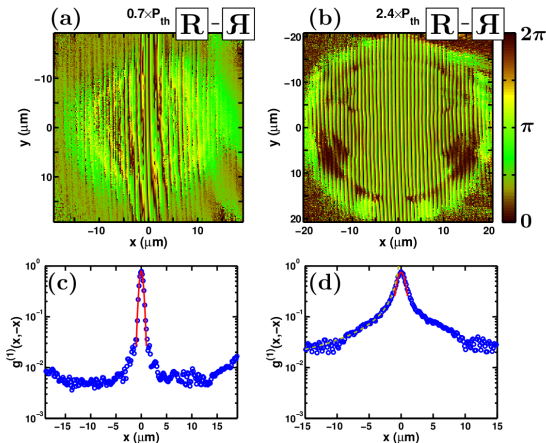
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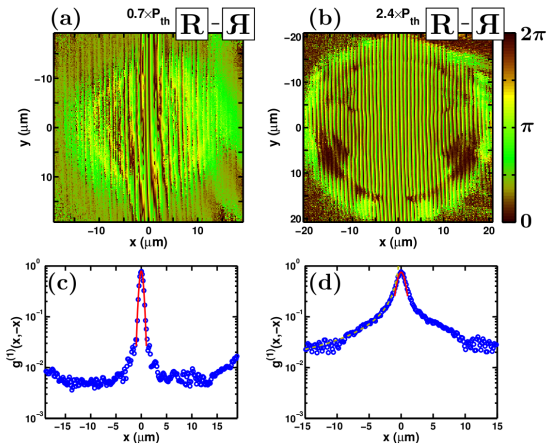


Experimental observation of power-law decay

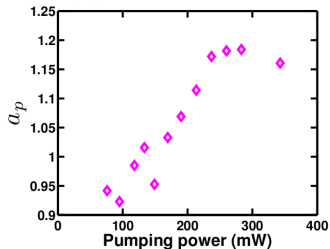


G. Rompos, Y. Yamamoto *et al.* submitted

Experimental observation of power-law decay



$$g_1(r, -r) \propto \left(\frac{r}{r_0} \right)^{-a_p}$$



G. Rompos, Y. Yamamoto *et al.* submitted

Exponent in a non-equilibrium 2D gas

$$\lim_{r \rightarrow \infty} \langle \psi^\dagger(\mathbf{r}, 0) \psi(-\mathbf{r}, 0) \rangle = |\psi_0|^2 \exp \left[-D_{\phi\phi}^<(\mathbf{r}, -\mathbf{r}) \right] \propto \exp \left[-a_p \ln \left(\frac{2r}{r_0} \right) \right]$$

- Experimentally, $a_p \simeq 1.2$

- In equilibrium $a_p = mk_B T / 2\pi \hbar^2 n_s < 1/4$ (BKT transition)
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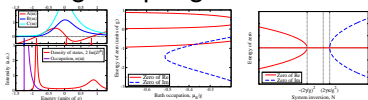
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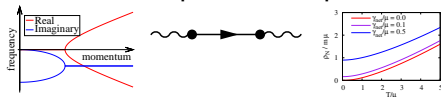
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Conclusion

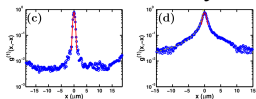
- Strong-coupling & condensation vs lasing.



- Survival of superfluid response



- Power law decay of correlations



Extra slides

- 3 Mean field theory
- 4 Green's functions and stability
- 5 Measuring superfluid density

Non-equilibrium theory; mean-field

Look for mean-field solution, $\psi(\mathbf{r}, t) = \psi_0 e^{-i\mu st}$.

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$$(i\partial_t - \omega_0 + i\kappa)\psi = \sum_{\alpha} g \langle S_{\alpha}^{-} \rangle$$

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Susceptibility:

$$\chi(\psi_0, \mu_s) = -g^2 \gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon - \frac{1}{2}\mu_s)}{[(\nu - E)^2 + \gamma^2][(\nu + E)^2 + \gamma^2]}$$

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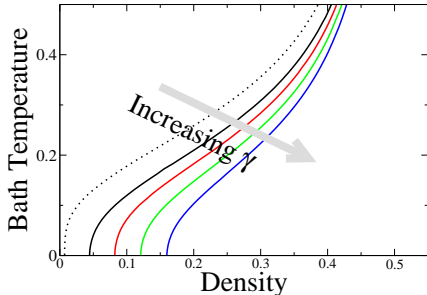
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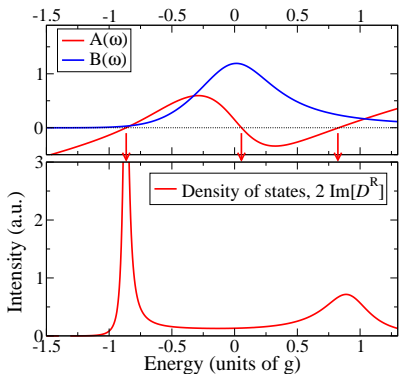
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Poles of Green's function and stability

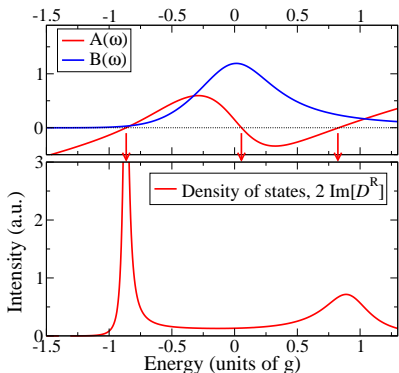
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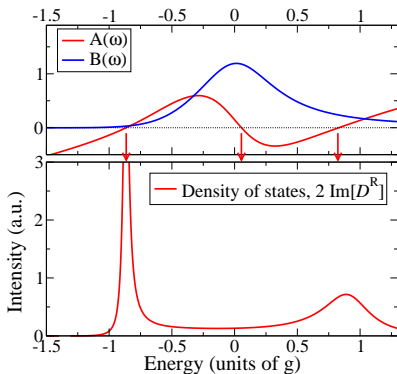
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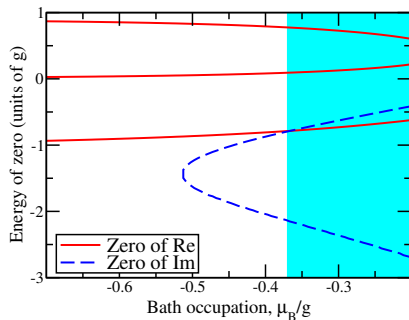
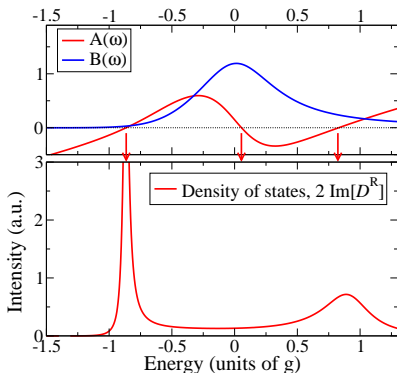
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Strong coupling and lasing — low temperature phenomenon

- Laser result: $B(\omega) = \text{Im} \left\{ \left[D^R(\omega) \right]^{-1} \right\} = \kappa - 2\gamma \frac{g^2 N_0}{(\omega - 2\epsilon)^2 + 4\gamma^2}$

Uniformly invert TLS

• Non-equilibrium model

$$B(\omega) = \kappa + \int \frac{d\nu}{2\pi} \sum_{\alpha} \frac{\gamma^2 g^2 (F_b(\nu + \omega) - F_a(\omega))}{[(\nu + \omega - \epsilon)^2 + \gamma^2][(\nu + \epsilon)^2 + \gamma^2]}$$

Inverts low energy part of homogeneous spectrum

• Non-equilibrium model, $T \gg \gamma$

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Inhomogeneous broadening required for strong-coupling lasing

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Uniformly invert TLS

- Non-equilibrium model

$$B(\omega) = \kappa + \int \frac{d\nu}{2\pi} \sum_{\alpha} \frac{\gamma^2 g^2 (F_b(\nu + \omega) - F_a(\omega))}{\left[(\nu + \omega - \epsilon)^2 + \gamma^2 \right] \left[(\nu + \epsilon)^2 + \gamma^2 \right]}$$

Inverts low energy part of homogeneous spectrum

- Non-equilibrium model, $T \gg \gamma$

$$B(\omega) \simeq \kappa + \sum_{\alpha} g_{\alpha}^2 \gamma \frac{[F_b(\epsilon_{\alpha}) - F_a(\epsilon_{\alpha} - \omega)]}{(\omega - 2\epsilon_{\alpha})^2 + 4\gamma^2}$$

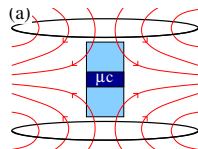
Inhomogeneous broadening required for strong-coupling lasing

Measuring superfluid density

1. Effect rotating frame

Polariton polarization: $(\psi_{\circ}, \psi_{\circ})$

$$H = \lambda \begin{pmatrix} \ell^2 & r^2 e^{2i\phi} \\ r^2 e^{-2i\phi} & -\ell^2 \end{pmatrix}$$



Measuring superfluid density

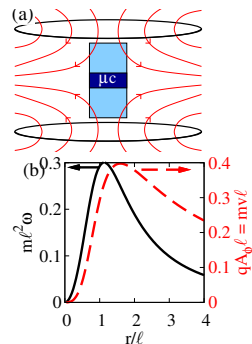
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Ground state Berry phase:

$$q\mathbf{A}_{\text{eff}} = m\omega \times \mathbf{r} = \frac{\hat{\phi}}{r} \left[1 - \frac{\ell^2}{\sqrt{r^4 + \ell^4}} \right]$$



Measuring superfluid density

1. Effect rotating frame

Polariton polarization: (ψ_0, ψ_0)

$$H = \lambda \begin{pmatrix} \ell^2 & r^2 e^{2i\phi} \\ r^2 e^{-2i\phi} & -\ell^2 \end{pmatrix}$$

Ground state Berry phase:

$$q\mathbf{A}_{\text{eff}} = m\omega \times \mathbf{r} = \frac{\hat{\phi}}{r} \left[1 - \frac{\ell^2}{\sqrt{r^4 + \ell^4}} \right]$$

2. Measure resulting current

Energy shift of normal state:

$$\Delta E = (1/2)mv^2 = 0.08/m\ell^2 \simeq 0.1\text{meV}$$

