

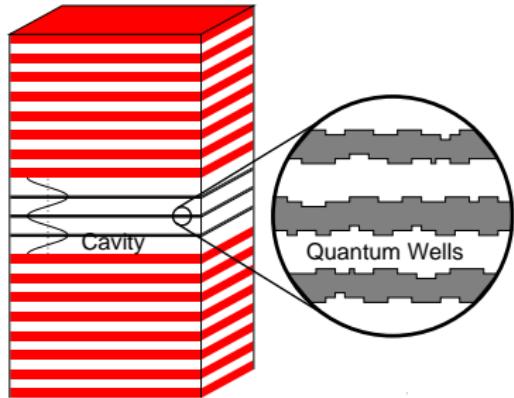
# Superfluidity and pattern formation in non-equilibrium polariton condensates

Jonathan Keeling

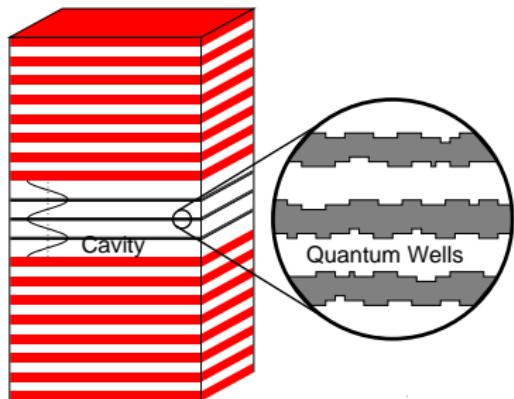
The Burn, April 2011



# Microcavity polaritons

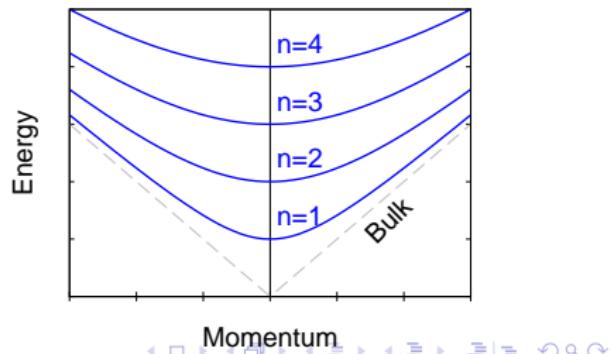


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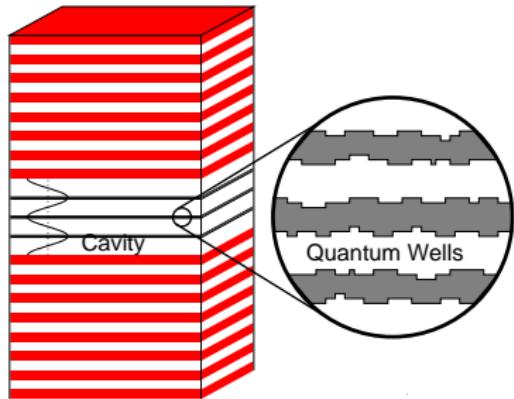


Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2 / 2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$



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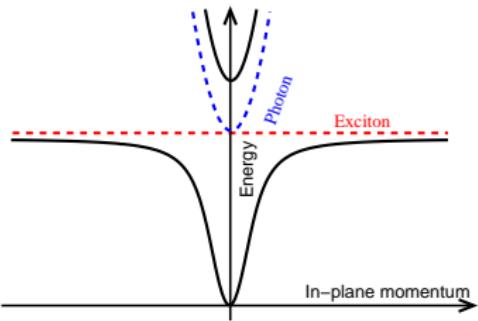


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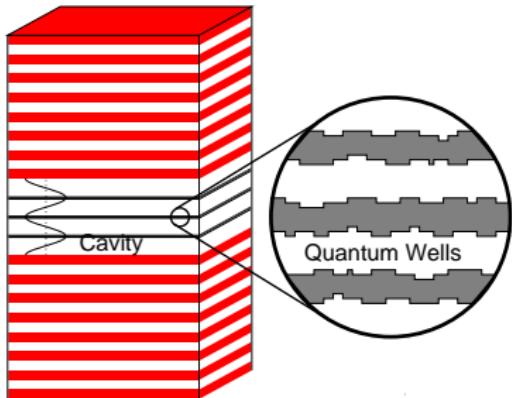
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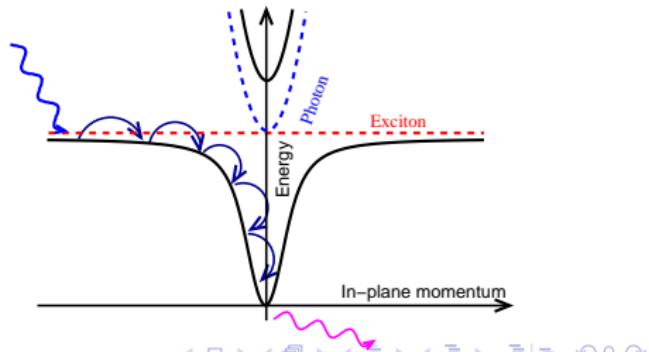


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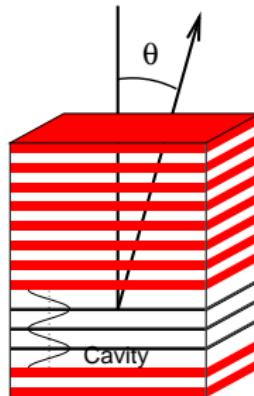
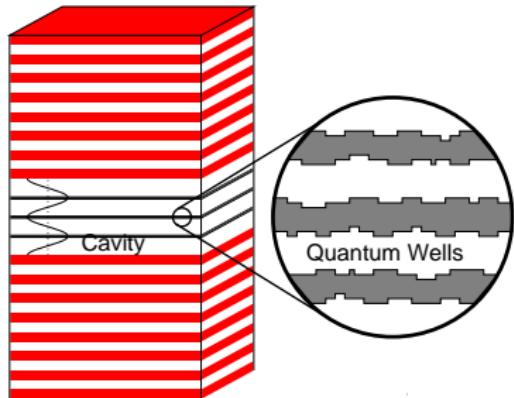
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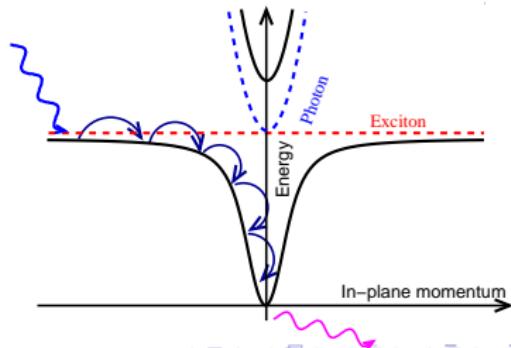


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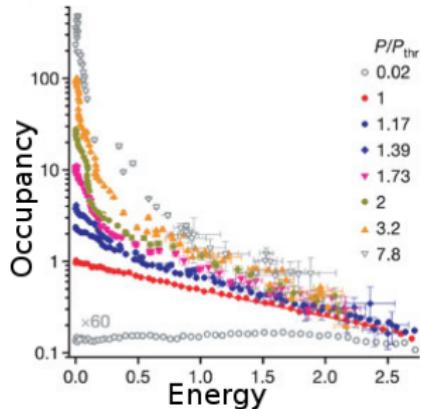
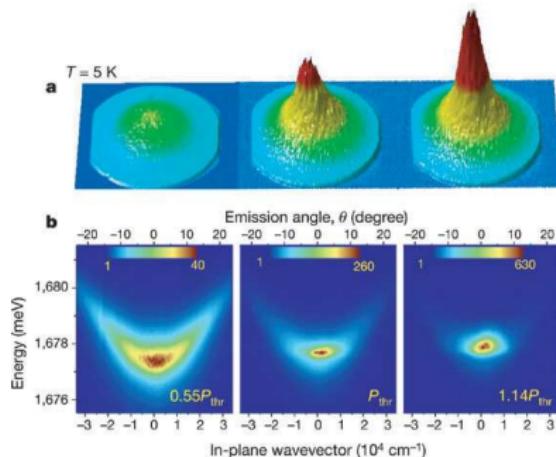
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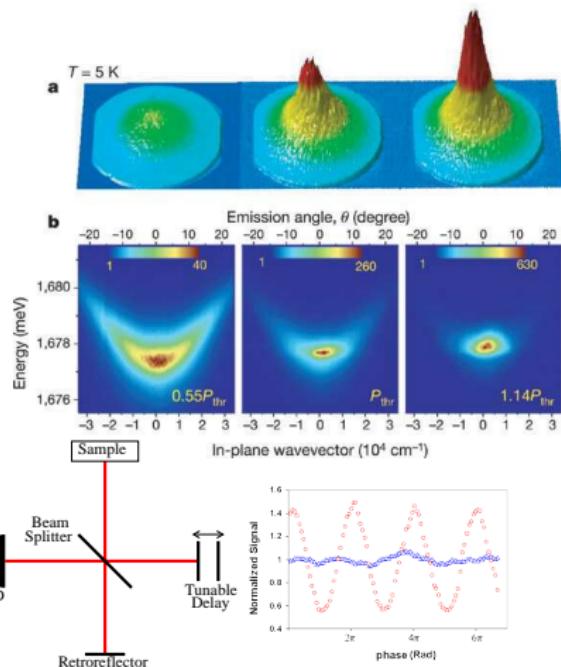


# Polariton experiments: Momentum/Energy distribution

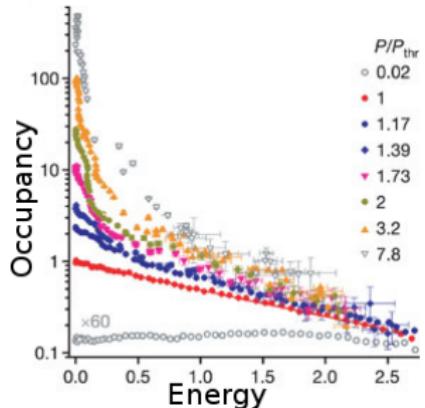


[Kasprzak, et al., Nature, 2006]

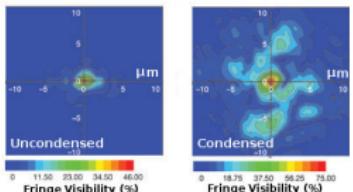
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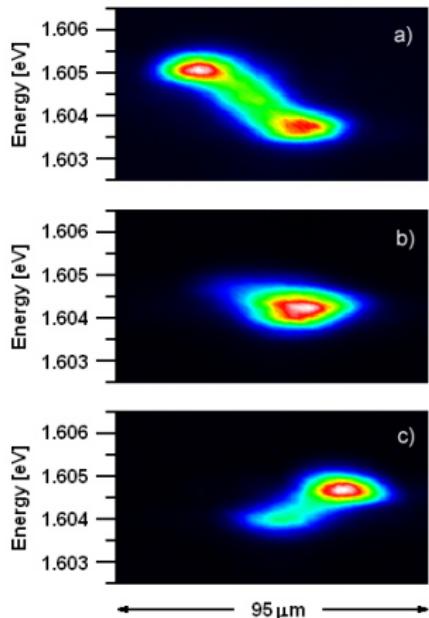
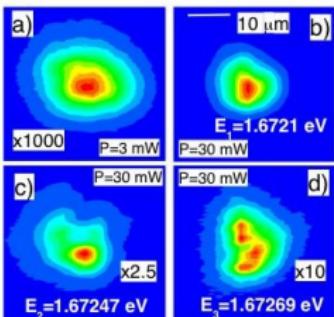
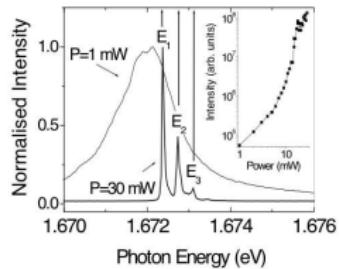


Coherence map:



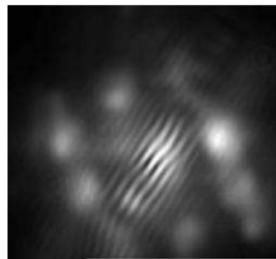
# Other polariton condensation experiments

- Stress traps for polaritons  
[Balili *et al* Science 316 1007 (2007)]
- Temporal coherence and line narrowing  
[Love *et al* Phys. Rev. Lett. 101 067404 (2008)]

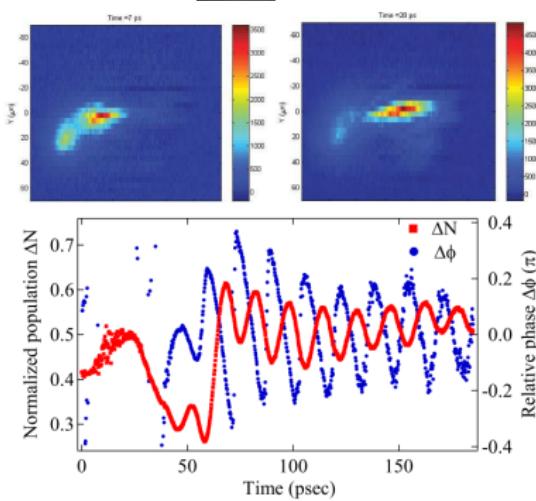


# Other polariton condensation experiments

- Quantised vortices in disorder potential  
[Lagoudakis *et al* Nature Phys. 4, 706 (2008)]



- Soliton propagation  
[Amo *et al* Nature 457 291 (2009)]
- Driven superfluidity  
[Amo *et al* Nature Phys. (2009)]
- Josephson Oscillations  
[Lagoudakis *et al* PRL (2010)]



## 1 Non-equilibrium model — coherence and strong coupling

## 2 Pattern formation

- Instability of Thomas-Fermi profile
- Polarisation degree of freedom
- Steady vortex lattices

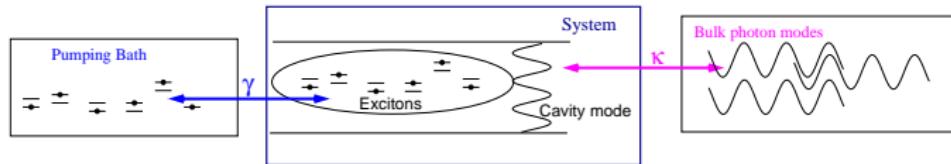
## 3 Condensed spectrum and superfluidity

- Condensed spectrum
- Current-current response function
- Power law decay of coherence

# Non-equilibrium approach: Steady state, and fluctuations

$$H = H_{\text{sys}} + H_{\text{sys,bath}} + H_{\text{bath}},$$

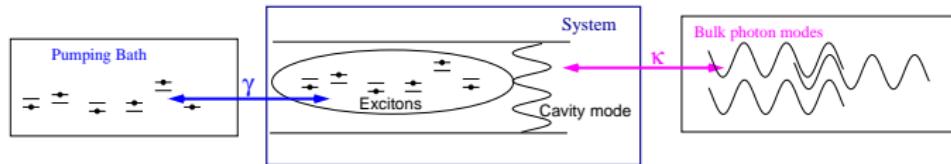
$$H_{\text{sys}} = \sum_k \omega_k \psi_k \psi_k^\dagger + \sum_\alpha g_\alpha (\phi_\alpha^\dagger \psi_k + \text{H.c.}) + H_{\text{ex}}[\phi_\alpha, \phi_\alpha^\dagger]$$



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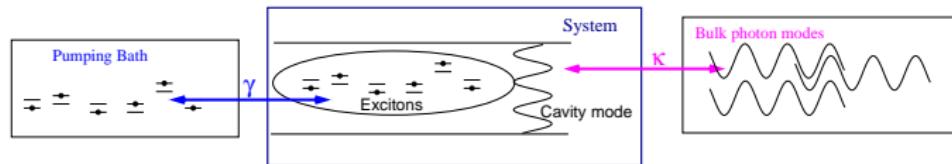


**Steady state**,  $\psi(\mathbf{r}, t) = \psi_0 e^{-i\mu_S t}$ .

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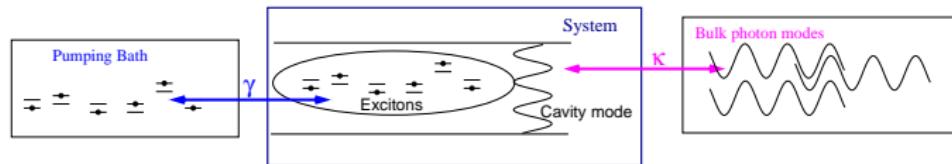
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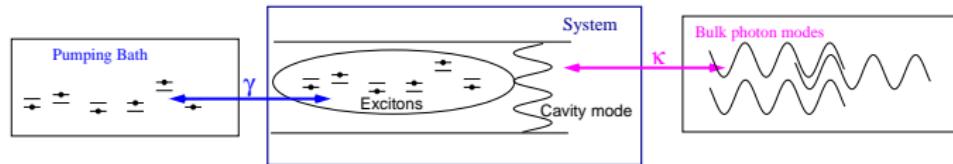
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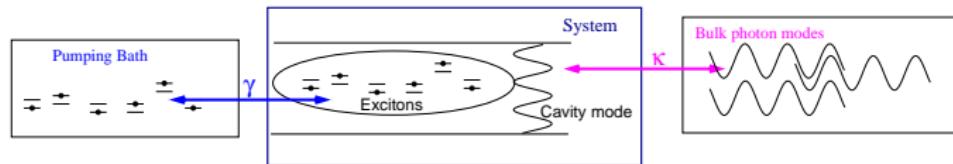
## Fluctuations

$$[D^R - D^A](t, t') = -i \left\langle [\psi(t), \psi^\dagger(t')]_- \right\rangle$$

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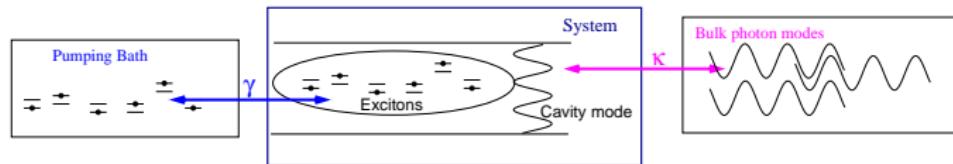
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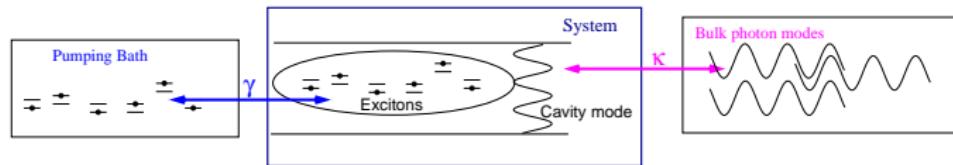
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1 Non-equilibrium model — coherence and strong coupling

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- Instability of Thomas-Fermi profile
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Nonlinear, complex susceptibility

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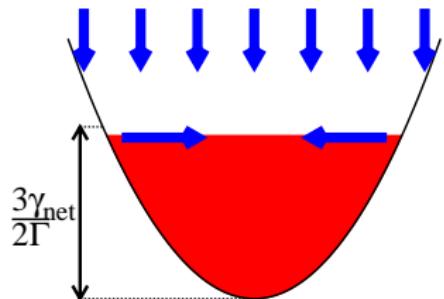
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$$i\partial_t \psi = \left[ -\frac{\nabla^2}{2m} + V(r) + U|\psi|^2 + i \left( \gamma_{\text{eff}}(\mu_B) - \kappa - \Gamma |\psi|^2 \right) \right] \psi$$

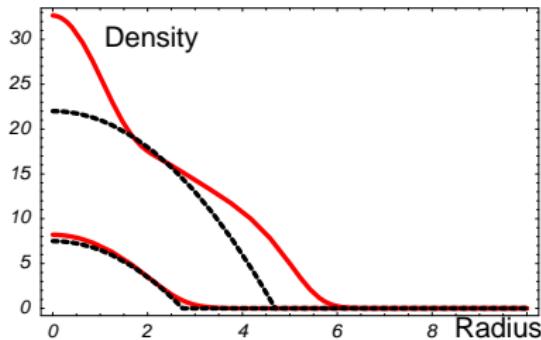
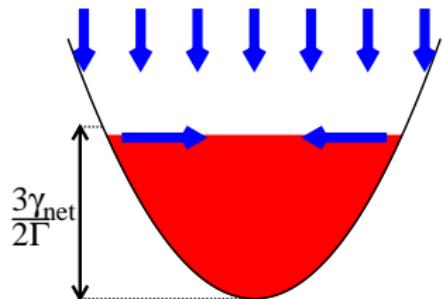
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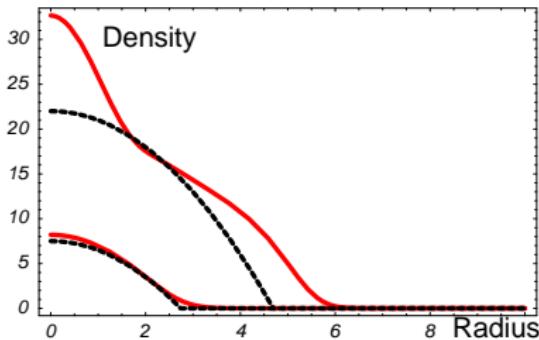
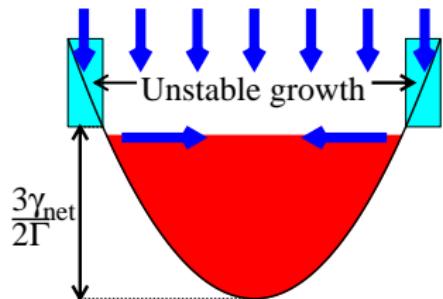
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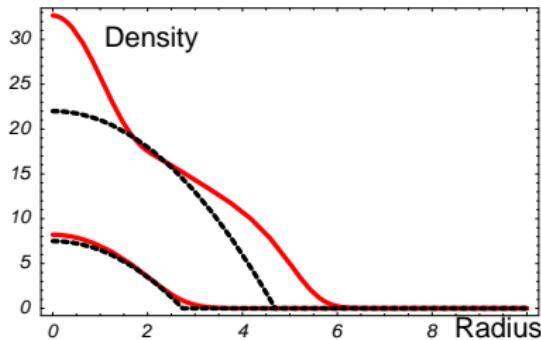
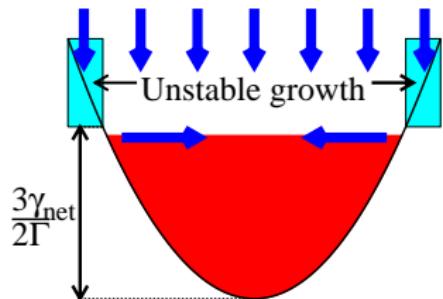
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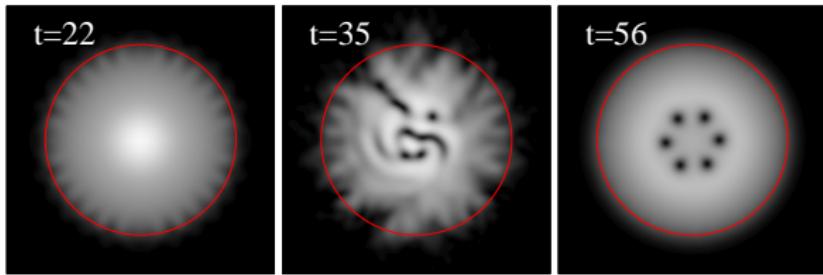


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Pattern formation: Vortex lattices



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- Instability of Thomas-Fermi profile
- **Polarisation degree of freedom**
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# Polariton spin degree of freedom

- Left- and Right-circular polarised states.

- For weakly interacting dilute gas gas model

- Tendency to biexciton formation  $\rightarrow L_1$

- Magnetic field:  $\Delta_1$

- $\Delta_1$  — inequivalent axes.

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$$H = \frac{|\nabla \Psi_L|^2}{2m} + \frac{|\nabla \Psi_R|^2}{2m} + \frac{U_0}{2} \left( |\Psi_L|^2 + |\Psi_R|^2 \right)^2$$

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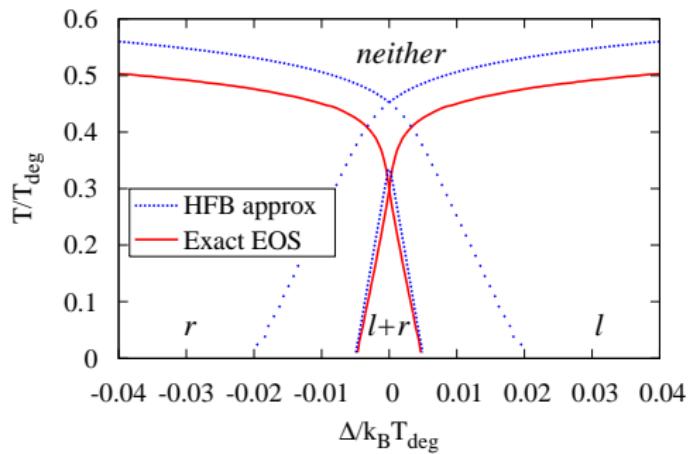
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- Magnetic field:  $\Delta$
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# Equilibrium phase diagrams

$$\Delta_{\parallel} = 0.$$

For  $U_1 = 0.5$ ,  $\Psi_{L,R}$  decouple.



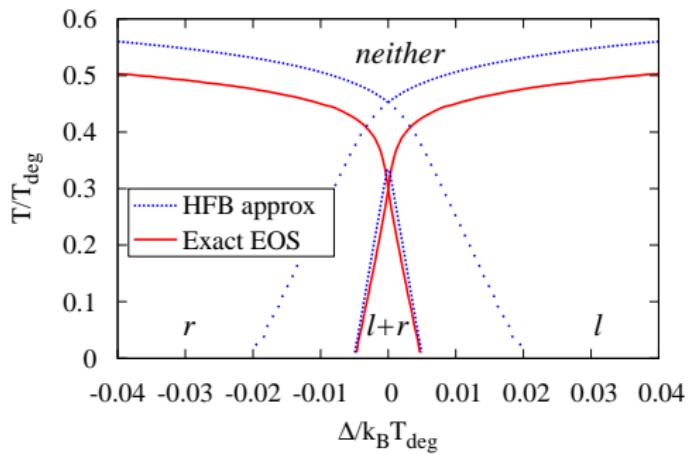
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[Rubo *et al* Phys. Lett. A '06; Keeling, Phys. Rev. B '08]

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$\Delta_{\parallel} \neq 0$ : Eqbm state locked.

[Rubo *et al* Phys. Lett. A '06; Keeling, Phys. Rev. B '08]

# Non-equilibrium spinor condensate

$$H = H_0[\psi_L] + H_0[\psi_R] + (U_0 - 2\textcolor{blue}{U}_1)|\psi_L|^2|\psi_R|^2 + \frac{\Delta_{\perp}}{2} \left( |\psi_L|^2 - |\psi_R|^2 \right) + \Delta_{\parallel} (\psi_L^{\dagger}\psi_R + \text{H.c.})$$

Spinor Gross-Pitaevskii equation

$$i\partial_t\psi_L = \left[ -\frac{\nabla^2}{2m} + V(r) + U_0|\psi_L|^2 + i \left( \gamma_{\text{eff}} - \kappa - \Gamma_0|\psi_L|^2 \right) \right] \psi_L$$

- Left-right coupling:  $U_1$
- Cross-spin loss terms  $\Gamma_0$
- Magnetic field:  $\gamma_{\text{eff}}$
- Energy-dependent gain  $\kappa$
- [Moulton et al. PRL 103]

# Non-equilibrium spinor condensate

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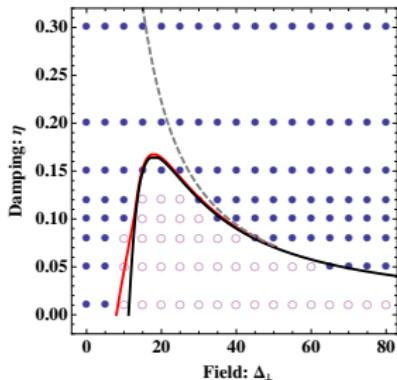
- Left-right coupling:  $\textcolor{blue}{U}_1$
- Magnetic field:  $\Delta_{\perp}$ ,  $\Delta_{\parallel}$
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- Energy-dependent gain  $\eta$   
[Wouters *et al* PRB '10]

# Desynchronised phase in homogeneous system

$$i\partial_t \psi_L = \left[ -\frac{\nabla^2}{2m} + V(r) + U_0 |\psi_L|^2 + (U_0 - 2\textcolor{blue}{U}_1) |\psi_R|^2 + \frac{\Delta_{\perp}}{2} \right. \\ \left. + i \left( \gamma_{\text{eff}} - \kappa - \Gamma_0 |\psi_L|^2 - \Gamma_1 |\psi_R|^2 - \eta i\partial_t \right) \right] \psi_L + \Delta_{||} \psi_R$$

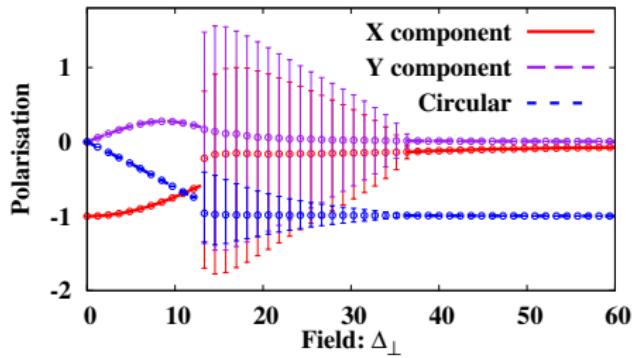
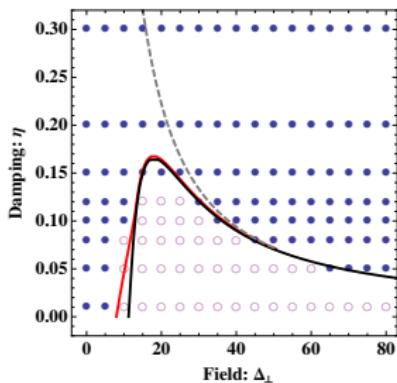
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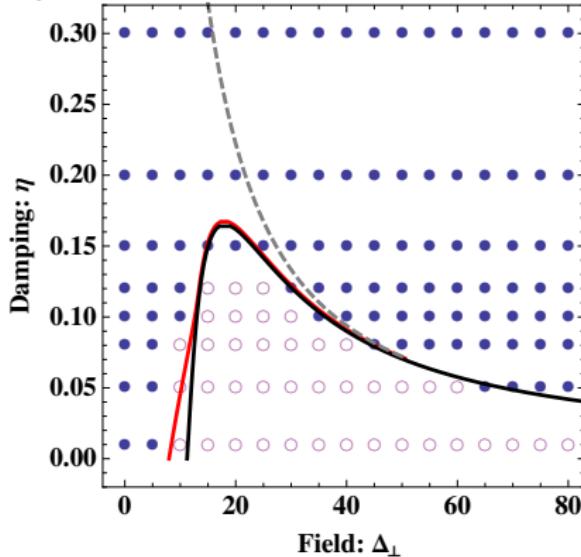
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# Desynchronisation and pattern formation

Trapping:

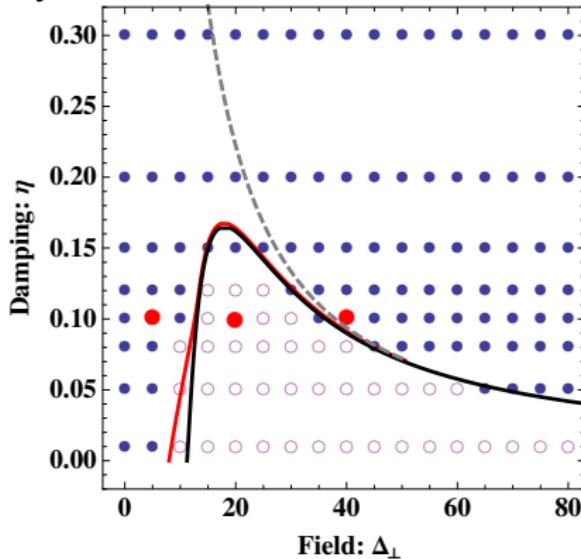
Synchronisation unaffected



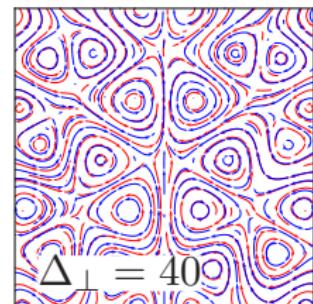
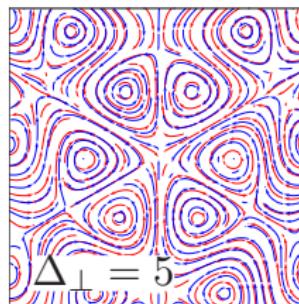
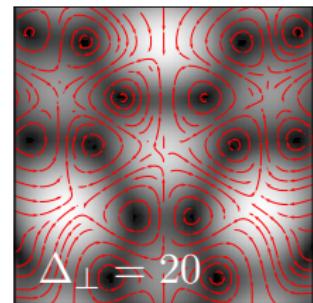
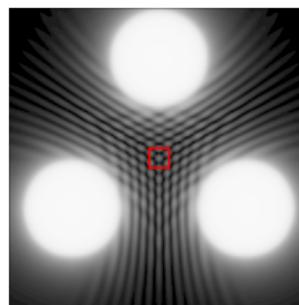
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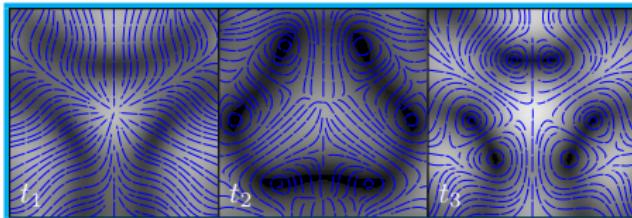
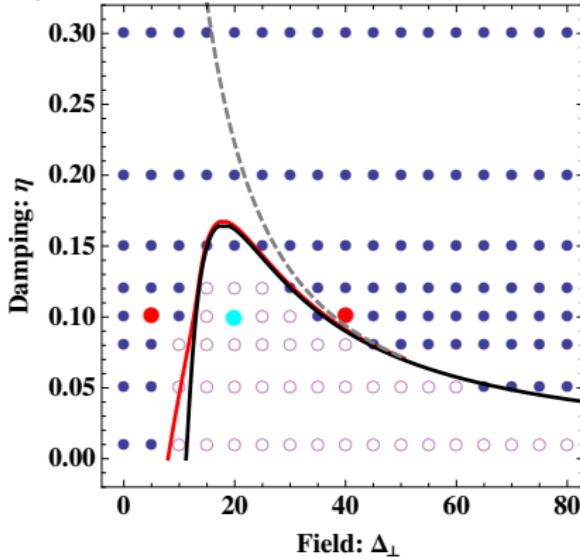
Desynchronisation  $\rightarrow$  half-vortex separation:



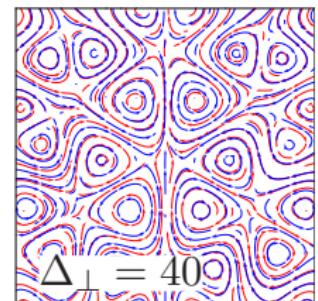
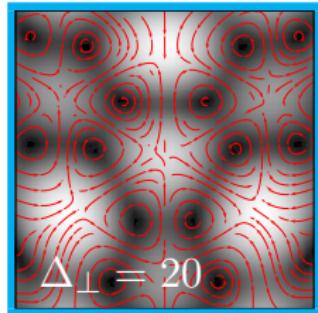
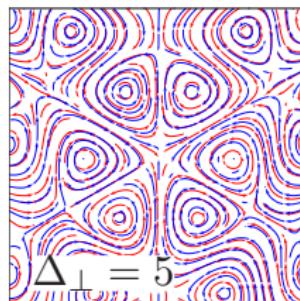
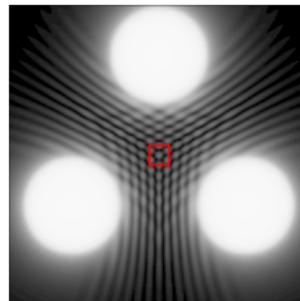
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Desynchronisation  $\rightarrow$  half-vortex separation:



## 1 Non-equilibrium model — coherence and strong coupling

## 2 Pattern formation

- Instability of Thomas-Fermi profile
- Polarisation degree of freedom
- Steady vortex lattices

## 3 Condensed spectrum and superfluidity

- Condensed spectrum
- Current-current response function
- Power law decay of coherence

# Superfluidity: Landau Criterion

Why superfluidity:

- Macroscopic occupation of single wavefunction
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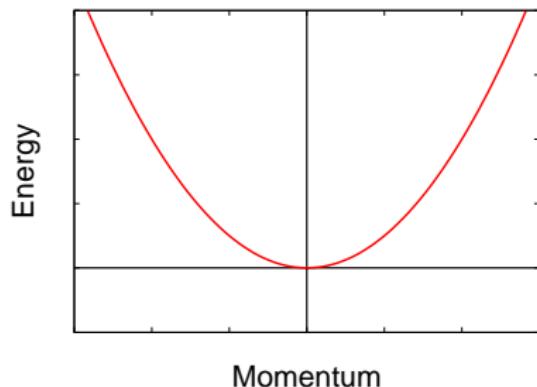
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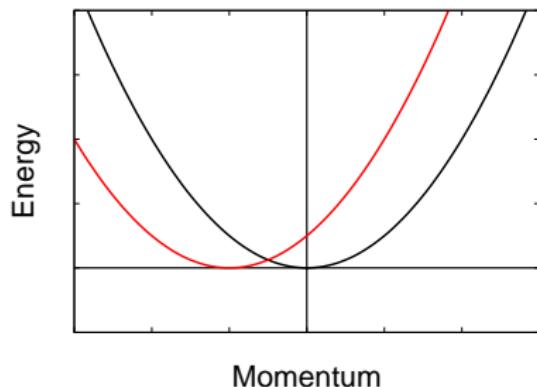
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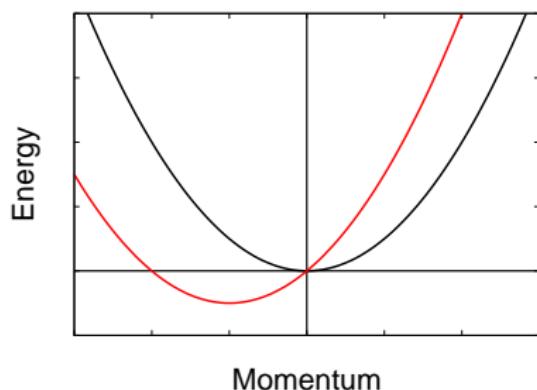
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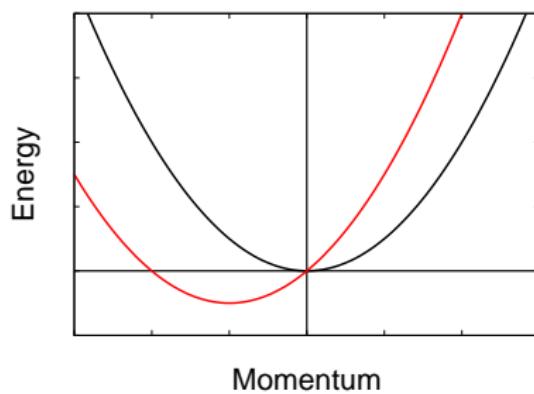
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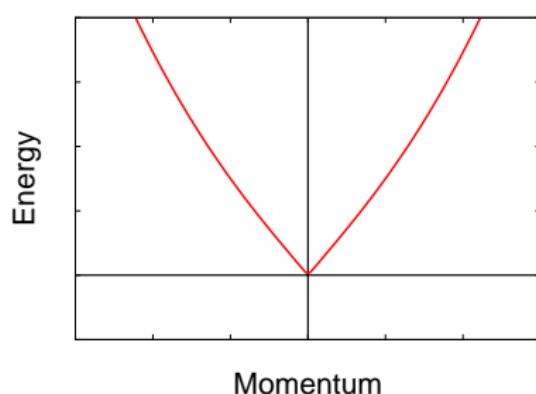
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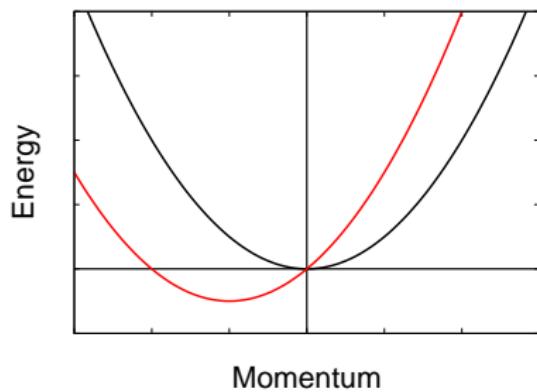
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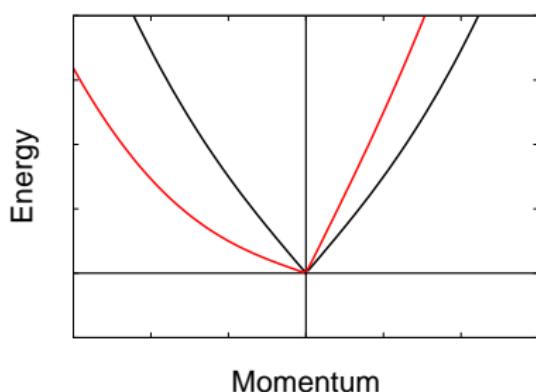
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Superfluid:



# Fluctuations above transition

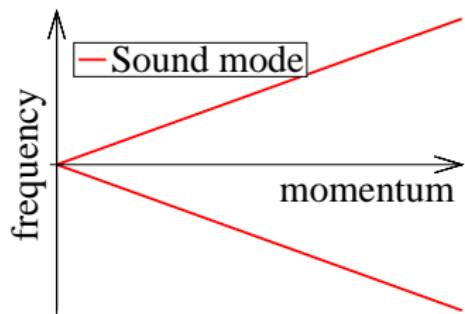
When condensed

$$\text{Det} \left[ D^R(\omega, k) \right]^{-1} = \omega^2 - \xi_k^2$$

With  $\xi_k \simeq ck$

Poles:

$$\omega^* = \xi_k$$



# Fluctuations above transition

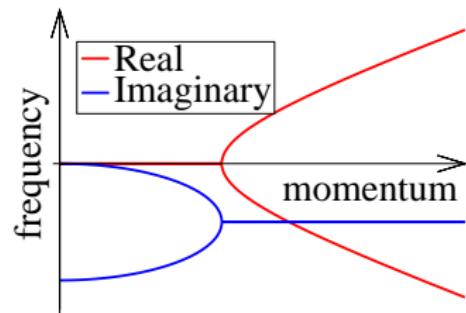
When condensed

$$\text{Det} \left[ D^R(\omega, k) \right]^{-1} = (\omega + i\gamma_{\text{net}})^2 + \gamma_{\text{net}}^2 - \xi_k^2$$

With  $\xi_k \simeq ck$

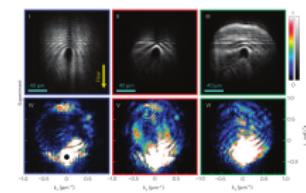
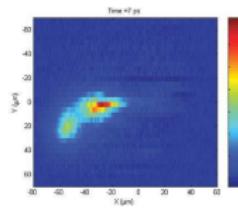
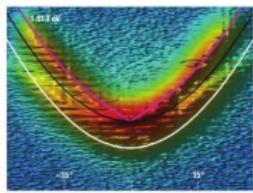
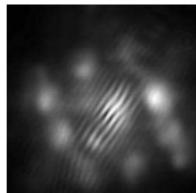
Poles:

$$\omega^* = -i\gamma_{\text{net}} \pm \sqrt{\xi_k^2 - \gamma_{\text{net}}^2}$$



# Non-equilibrium superfluidity checklist

	Quantised vortices	Landau critical velocity	Metastable persistent hydrodynamic flow	Two-fluid hydrodynamics	Local thermal equilibrium	Solitary waves
Superfluid $^4\text{He}$ /cold atom Bose-Einstein condensate	✓	✓	✓	✓	✓	✓
Non-interacting Bose-Einstein condensate	✓	✗	✗	✗	✓	✗
Classical irrotational fluid	✗	✓	✗	✓	✓	✓
Incoherently pumped polariton condensates	✓	✗	?	?	✗	?



Lagoudakis *et al* Nature Phys. 4, 706 (2008). Utsunomiya *et al* Nature Phys. 4 700 (2008). Amo *et al* Nature 457 291 (2009); Nature Phys (2009)

# Asking about non-equilibrium superfluidity

Current:

$$\mathbf{J} = \rho \mathbf{v} = \psi^\dagger i \nabla \psi = |\psi|^2 \nabla \phi$$

• Response functions

$$x_0(\omega=0, q \rightarrow 0) = \langle J(q), J(-q) \rangle = \frac{\rho_S(q)}{m} + \frac{\rho_N}{m} \delta_q$$

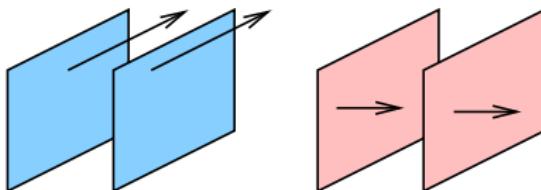
• Given  $D$  and  $J_i = \sum_{k>0} \frac{2k+q}{2m} v_k$

• Vertex corrections essential for superfluid part

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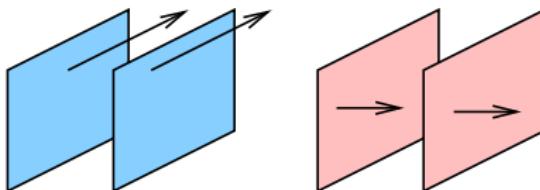
$$\chi_{ij}(\omega = 0, \mathbf{q} \rightarrow 0) = \langle [J_i(\mathbf{q}), J_j(-\mathbf{q})] \rangle = \frac{\rho_S}{m} \frac{q_i q_j}{q^2} + \frac{\rho_N}{m} \delta_{ij}$$

- Given  $D$  and  $J_i = \partial_{\mathbf{q}, i} \frac{\Delta_0(\mathbf{q})}{2m} \partial_{\mathbf{q}, i}$
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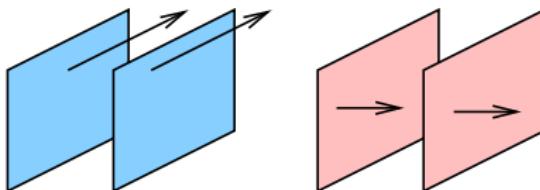
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What does this tell us about the superfluid part?

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# Non-equilibrium superfluid response

- Superfluid response exists because:

$$\text{Diagram: } \textcirclearrowleft \bullet \rightarrow \bullet \textcirclearrowright = \frac{i\psi_0 q_i}{2m} (1, -1) D^R(q, \omega = 0) \begin{pmatrix} 1 \\ -1 \end{pmatrix} \frac{i\psi_0 q_j}{2m}$$

$D^R(\omega = 0) \propto 1/\epsilon_q$ , despite pumping/decay — superfluid response exists.

- Normal density:

$$n_N = \int d^d k \sigma_z \int \frac{d\omega}{2\pi} \text{Tr} \left[ \sigma_z D^K \sigma_z (D^R + D^A) \right]$$

- Is affected by pump/decay:  
Does not vanish at  $T \rightarrow 0$ .

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- Superfluid response exists because:

$$\text{Diagram: } \text{Two wavy lines with dots at vertices, connected by a horizontal line with an arrow pointing right.} = \frac{i\psi_0 q_i}{2m} (1, -1) D^R(q, \omega = 0) \begin{pmatrix} 1 \\ -1 \end{pmatrix} \frac{i\psi_0 q_j}{2m}$$

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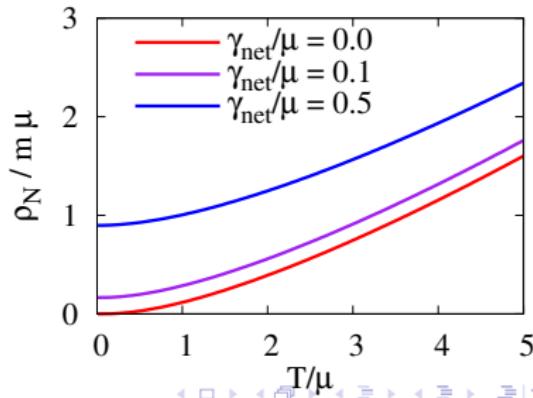
$$\text{Diagram: Two wavy lines meeting at a point with a dot, followed by a horizontal line with a dot, then two wavy lines meeting at a point with a dot.} = \frac{i\psi_0 q_i}{2m} (1, -1) D^R(q, \omega = 0) \begin{pmatrix} 1 \\ -1 \end{pmatrix} \frac{i\psi_0 q_j}{2m}$$

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# Correlations in a 2D Gas

Correlations (in 2D):

$$g_1(\mathbf{r}, \mathbf{r}') = \langle \psi^\dagger(\mathbf{r}, t) \psi(0, \mathbf{r}') \rangle \simeq |\psi_0|^2 \exp \left[ -D_{\phi\phi}^<(t, \mathbf{r}, \mathbf{r}') \right]$$

- $D^< = D^K - D^R + D^A$

• Generally get  $\langle \psi^\dagger(\mathbf{r}, t) \psi(0, 0) \rangle \sim$

$$|\psi_0|^2 \exp \left[ -\beta_p \begin{cases} \ln(t/\tau_0) & t \rightarrow \infty, T \geq 0 \\ \beta \ln(c^2 t / \pi m \hbar^2) & T \simeq 0, t \rightarrow \infty \end{cases} \right]$$

[Szymańska et al., PRL '06; PRB '07]

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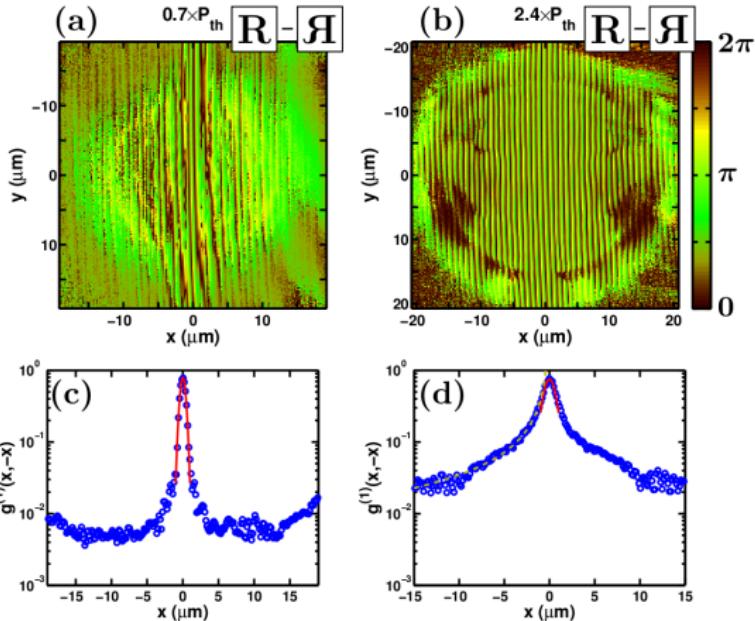
$$g_1(\mathbf{r}, \mathbf{r}') = \langle \psi^\dagger(\mathbf{r}, t) \psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-D_{\phi\phi}^<(t, r, r')]$$

- $D^< = D^K - D^R + D^A$

- Generally, get:  $\langle \psi^\dagger(\mathbf{r}, t) \psi(0, 0) \rangle \simeq |\psi_0|^2 \exp \left[ -a_p \begin{cases} \ln(r/r_0) & r \rightarrow \infty, t \simeq 0 \\ \frac{1}{2} \ln(c^2 t / \gamma_{\text{net}} r_0^2) & r \simeq, t \rightarrow \infty \end{cases} \right]$

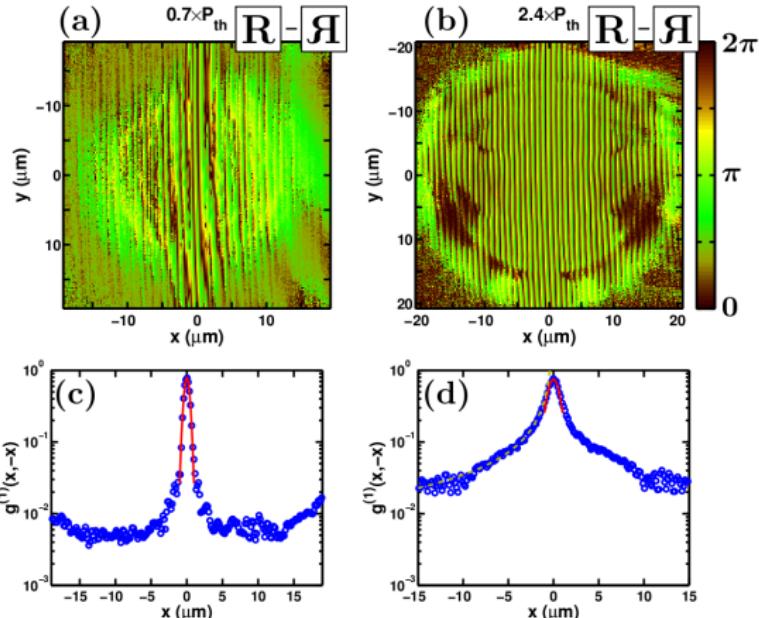
[Szymańska et al., PRL '06; PRB '07]

# Power law experiment



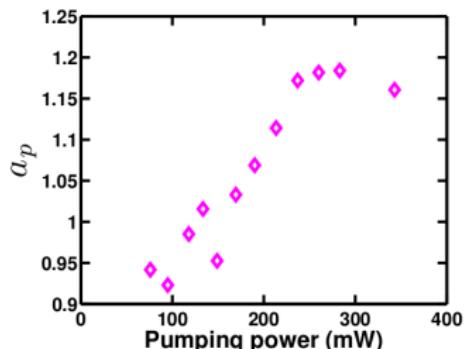
G. Rompos, Y. Yamamoto et al., submitted

# Power law experiment



G. Rompos, Y. Yamamoto et al., submitted

$$g_1(r, -r) \propto \left(\frac{r}{r_0}\right)^{-a_p}$$



# Power law experiment — non-equilibrium theory

$$\lim_{r \rightarrow \infty} \langle \psi^\dagger(\mathbf{r}, 0) \psi(-\mathbf{r}, 0) \rangle = |\psi_0|^2 \exp \left[ -D_{\phi\phi}^<(r, -r) \right] \propto \exp \left[ -a_p \ln \left( \frac{2r}{r_0} \right) \right]$$

- Experimentally,  $a_P \simeq 1.1$

↳  $a_p = \pi^2 k_B T / 2 \hbar^2 n_s < 1/4$  (BKT transition)

↳ Non-equilibrium theory depends on thermalisation.

G. Rompos, Y. Yamamoto et al., submitted

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- In equilibrium  $a_p = m k_B T / 2\pi \hbar^2 n_s < 1/4$  (BKT transition)

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Non-thermalised state or Pumping noise

G. Rompos, Y. Yamamoto et al., submitted

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  - ▶ Thermalised (yet diffusive modes)  $a_p = m k_B T / 2\pi \hbar^2 n_s$

© Jonathan Keeling, University of Cambridge 2011

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# Acknowledgements

## People:



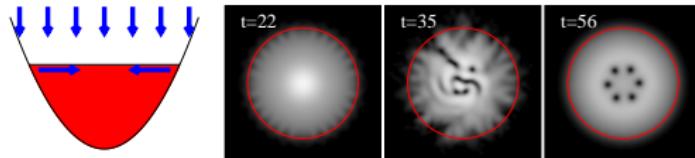
## Funding:

**EPSRC**

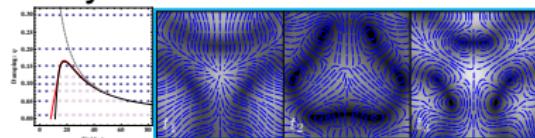
Engineering and Physical Sciences  
Research Council

# Summary

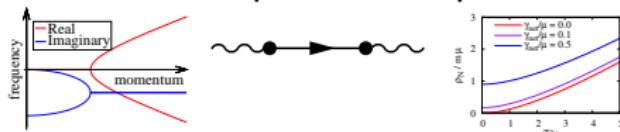
- Instability of Thomas-Fermi and spontaneous rotation



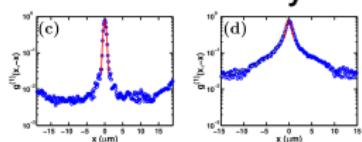
- Desynchronisation and vortex lattices



- Survival of superfluid response vs change to spectrum



- Power law decay of correlations

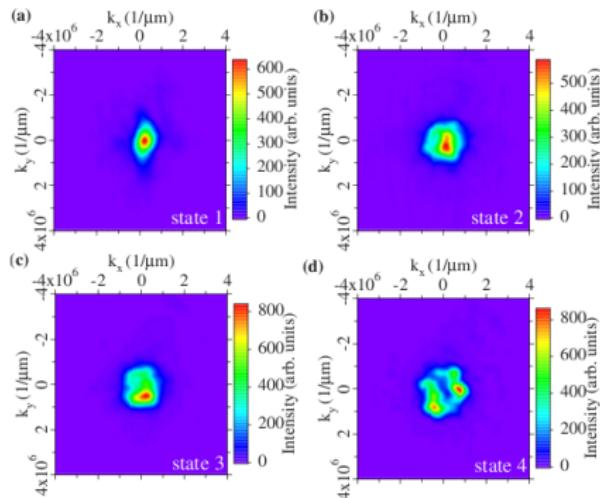




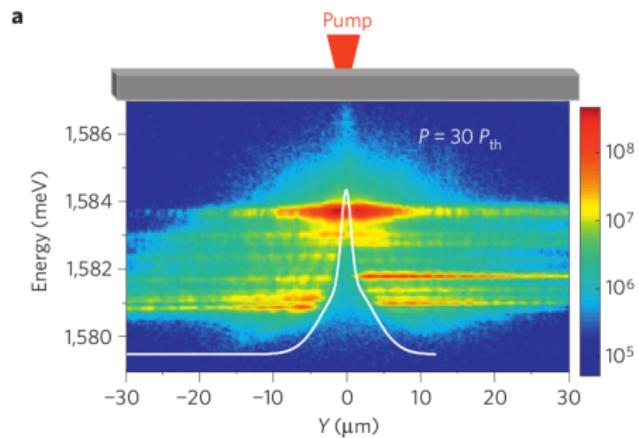
# Extra slides

- ④ Non-equilibrium pattern formation
- ⑤ Other polariton experiments
- ⑥ Green's functions and stability
- ⑦ Spinor problem
- ⑧ T=0 Keldysh results

# Other polariton condensation experiments: Non-equilibrium features



$|\psi(\mathbf{k})|^2 \neq |\psi(-\mathbf{k})|^2$ :  
Broken time-reversal symmetry.  
[Krizhanovskii *et al*, PRB (2009)]



Flow from pumping spot  
[Wertz *et al.*, Nat. Phys. (2010)]

# Complex Gross-Pitaevskii equation

Steady state equation:

$$(\mu_s - \omega_0 + i\kappa) \psi = \chi(\psi, \mu_s) \psi$$

- Local density limit:

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$$\left( i\partial_t + i\kappa - \left[ V(r) - \frac{\nabla^2}{2m} \right] \right) \psi(r) = \chi(\psi(r, t)) \psi(r, t)$$

Nonlinear, complex susceptibility

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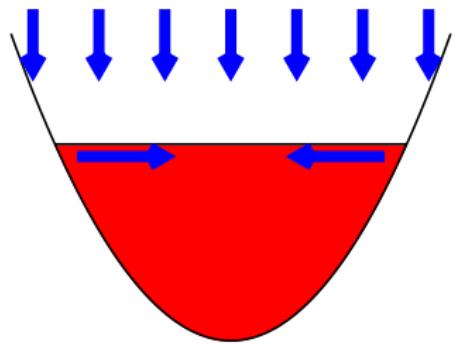
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$$i\partial_t \psi = \left[ -\frac{\nabla^2}{2m} + V(r) + U|\psi|^2 + i \left( \gamma_{\text{eff}}(\mu_B) - \kappa - \Gamma |\psi|^2 \right) \right] \psi$$

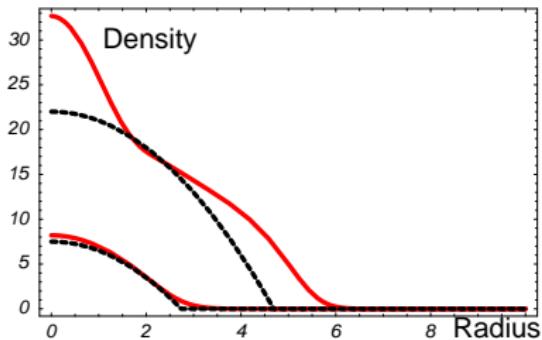
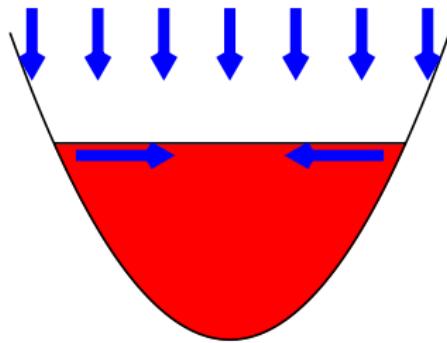
# Gross-Pitaevskii equation: Harmonic trap

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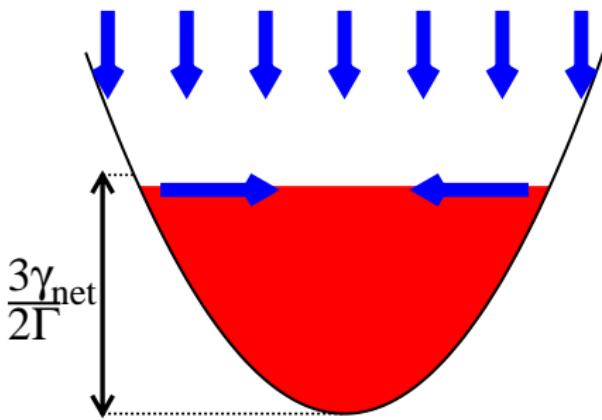
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# Stability of Thomas-Fermi solution

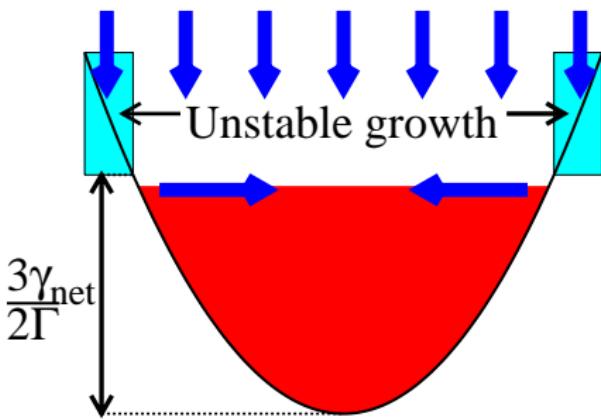
$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}} - \Gamma \rho) \rho$$



# Stability of Thomas-Fermi solution

High  $m$  modes:  $\delta\rho_{n,m} \simeq e^{im\theta} r^m \dots$

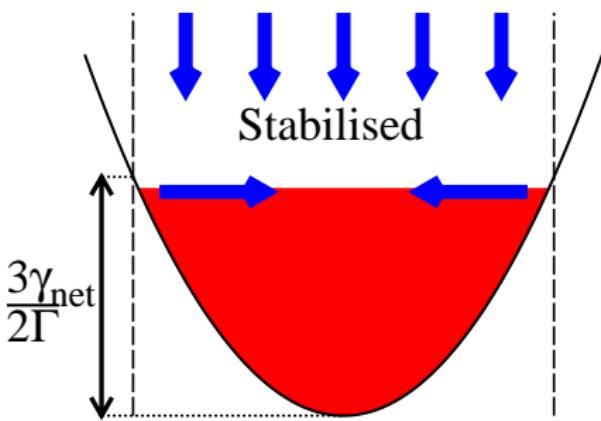
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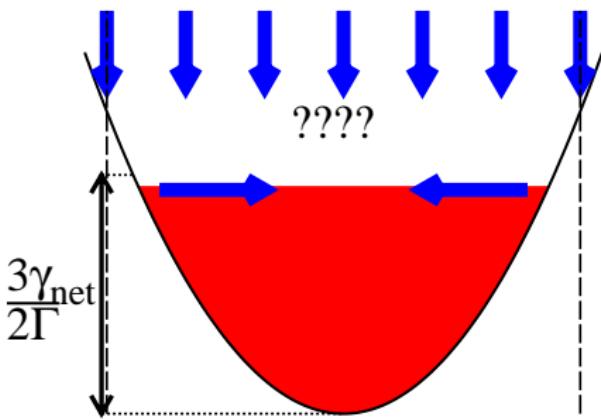
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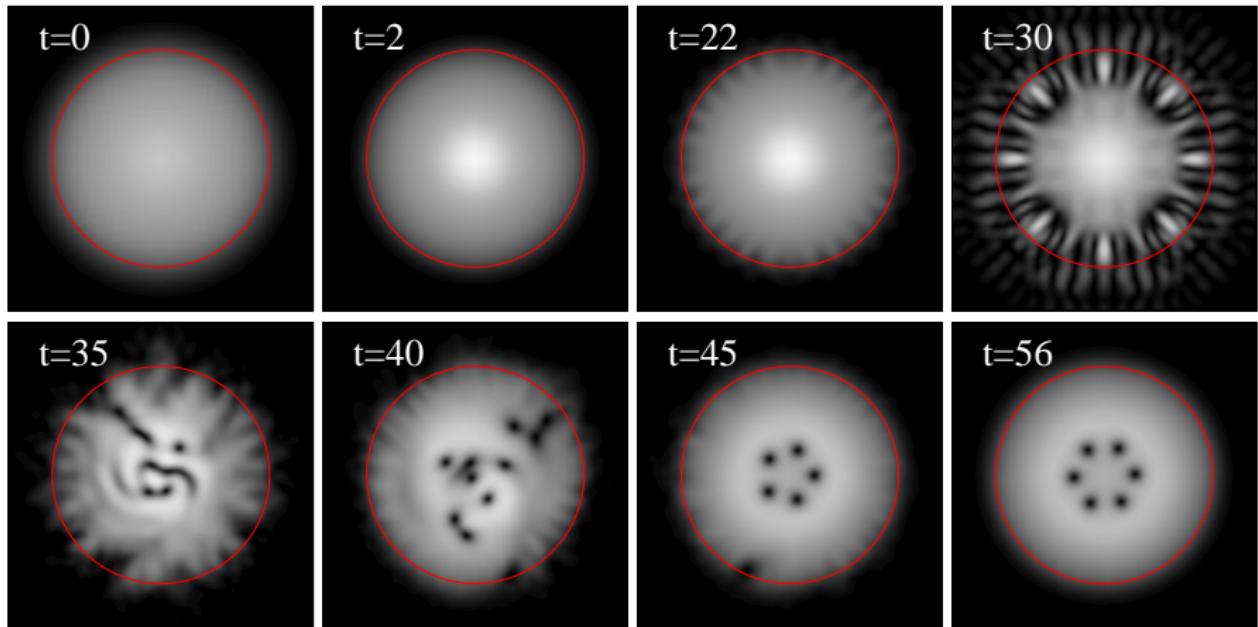
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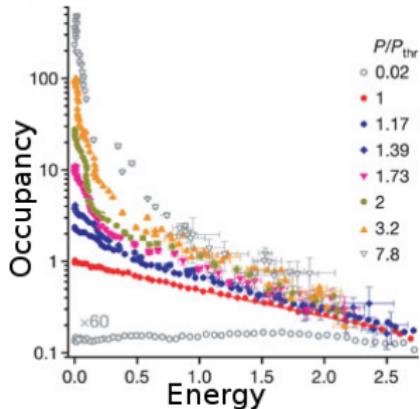
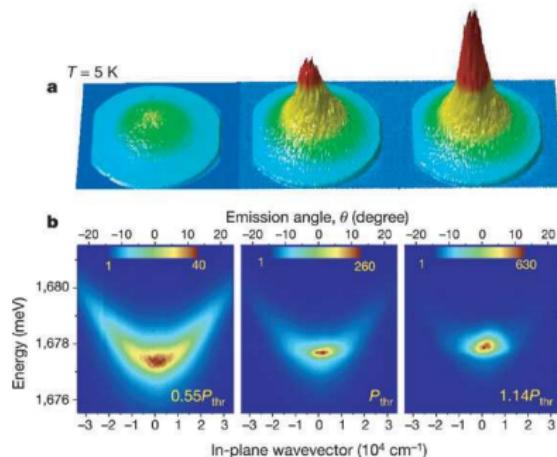
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## Time evolution:

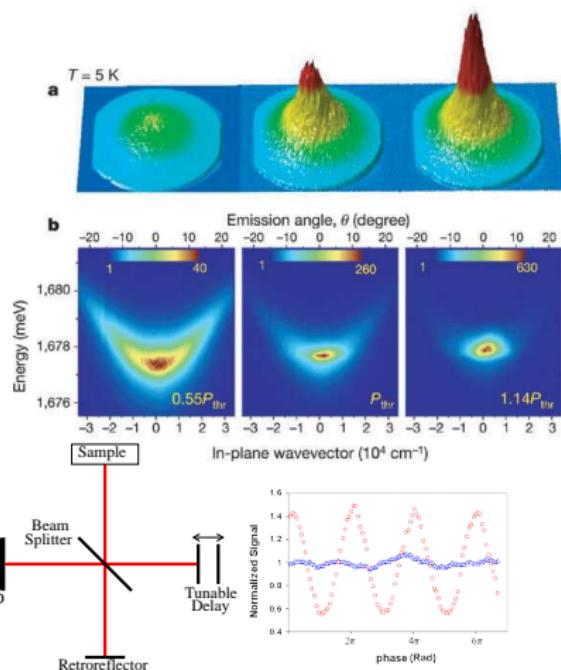


# Polariton experiments: Momentum/Energy distribution

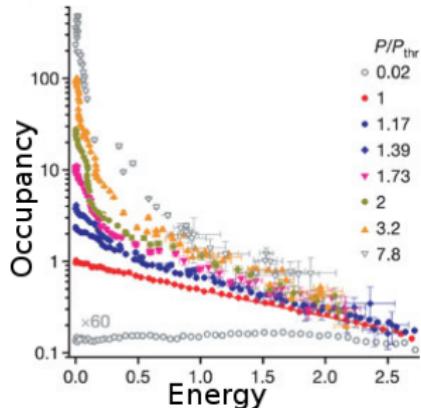


[Kasprzak, et al., Nature, 2006]

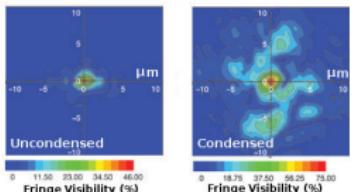
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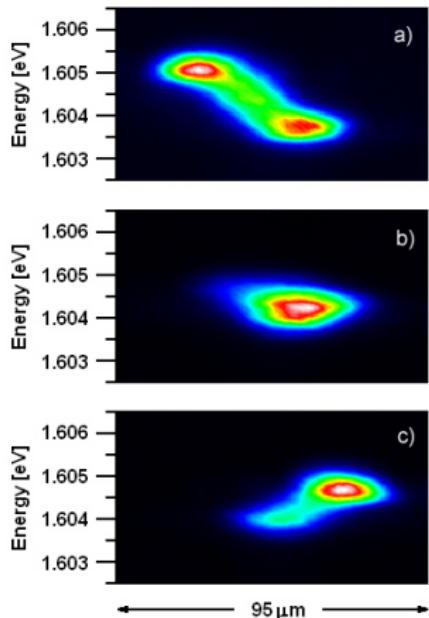
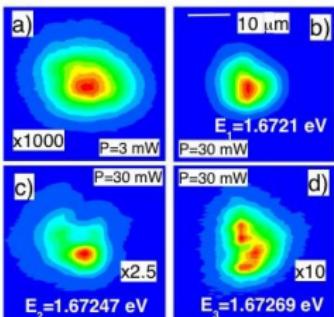
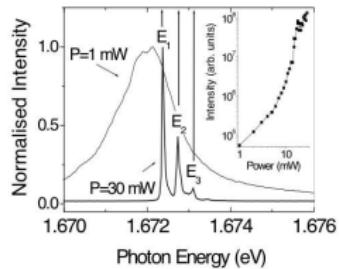


Coherence map:



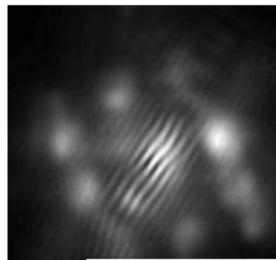
# Other polariton condensation experiments

- Stress traps for polaritons  
[Balili *et al* Science 316 1007 (2007)]
- Temporal coherence and line narrowing  
[Love *et al* Phys. Rev. Lett. 101 067404 (2008)]

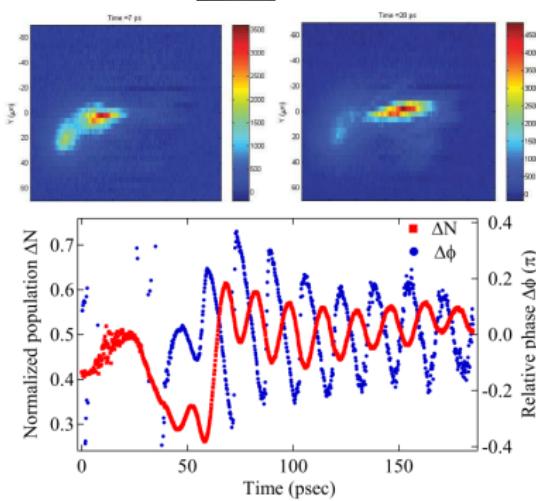


# Other polariton condensation experiments

- Quantised vortices in disorder potential  
[Lagoudakis *et al* Nature Phys. 4, 706 (2008)]

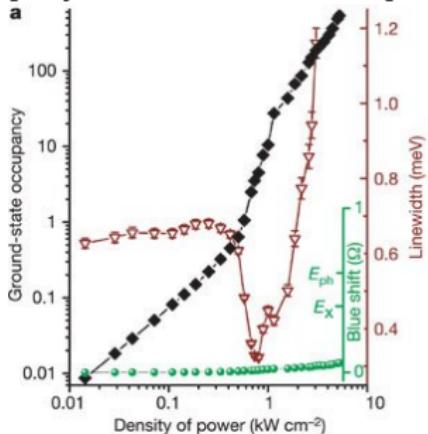


- Soliton propagation  
[Amo *et al* Nature 457 291 (2009)]
- Driven superfluidity  
[Amo *et al* Nature Phys. (2009)]
- Josephson Oscillations  
[Lagoudakis *et al* PRL (2010)]

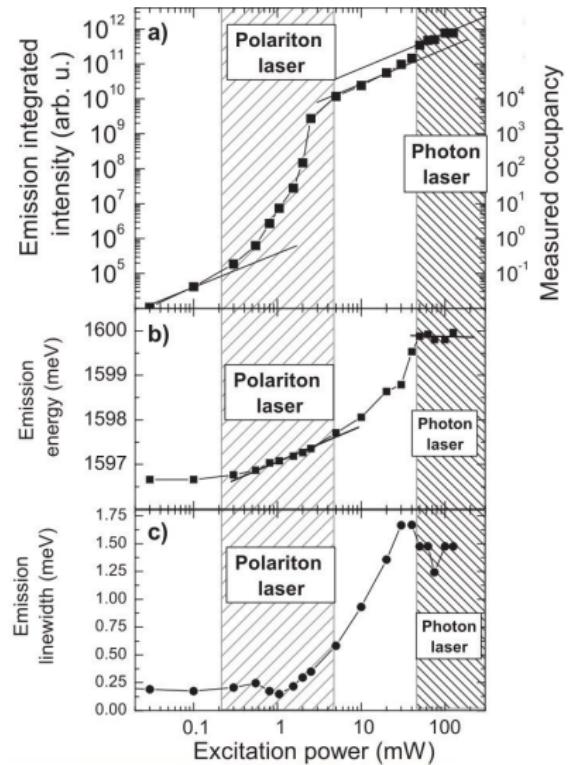


# Polariton experiments: Strong coupling

[Bajoni *et al* PRL 2008]

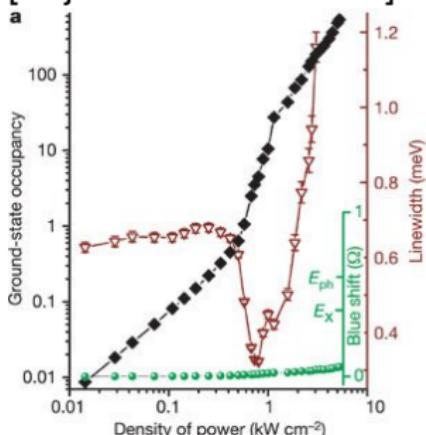


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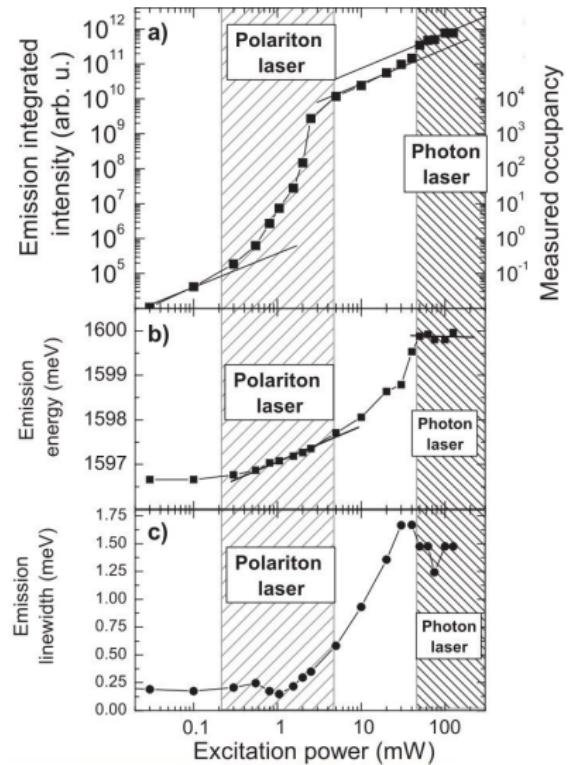
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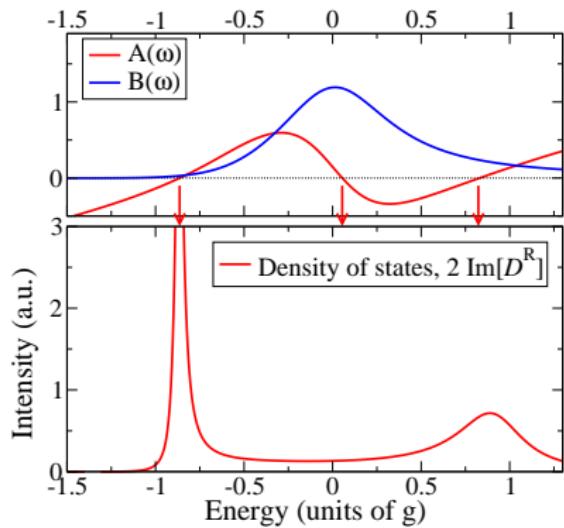
Strong coupling via:

- Small blueshift compared to  $\Omega_R$
- Polaritonic dispersion,  $m > m_{\text{phot}}$
- Separate photon threshold



# Poles of Green's function and stability

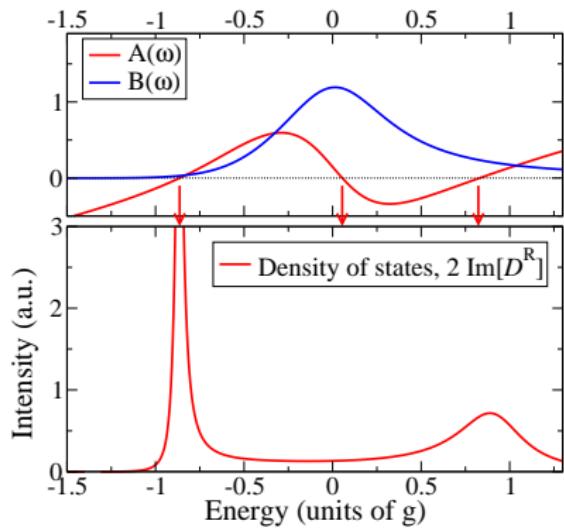
$$\left[ D^R(\omega) \right]^{-1} = (\omega - \omega_k^*) + i\alpha(\omega - \mu_{\text{eff}})$$



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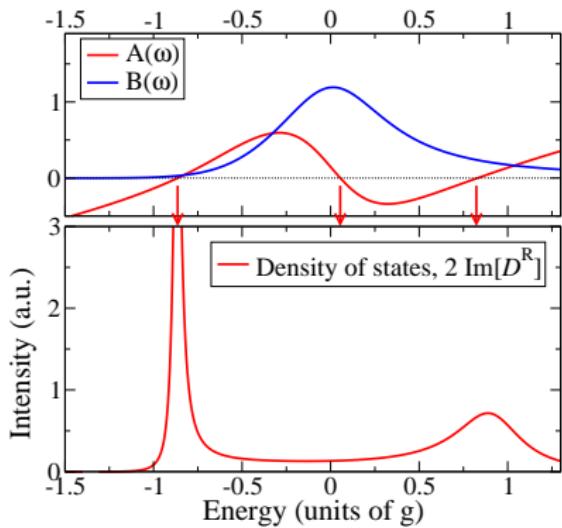
Pole  $\omega = \frac{\omega^* + \alpha^2 \mu_{\text{eff}} + i\alpha(\mu_{\text{eff}} - \omega^*)}{1 + \alpha^2}$



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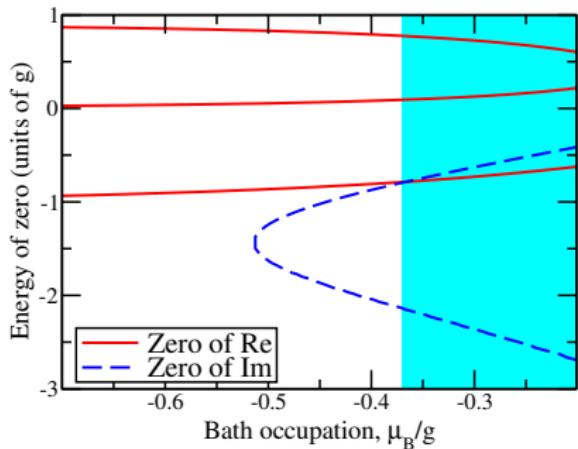
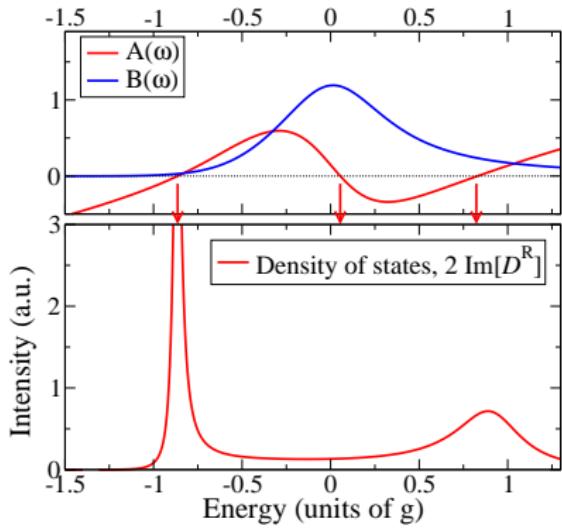
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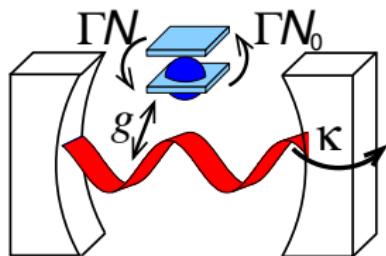
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# Poles and stability for a laser



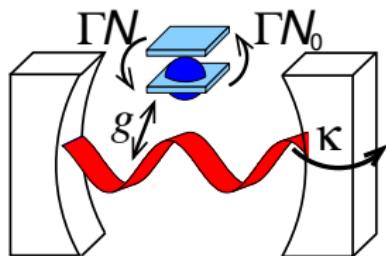
Maxwell-Bloch equations:

$$\partial_t \psi = -i\omega_k \psi - \kappa \psi + g P$$

$$\partial_t P = -2i\epsilon P - 2\gamma P + g\psi N$$

$$\partial_t N = 2\gamma(N_0 - N) - 2g(\psi^* P + P^* \psi)$$

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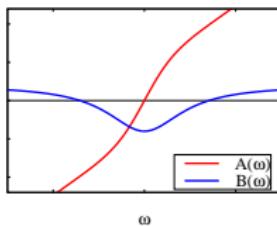
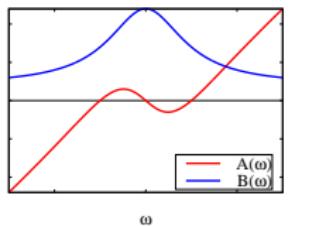
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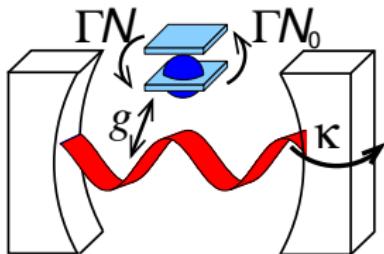
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$$[D^R(\omega)]^{-1} = \omega - \omega_k + i\kappa + \frac{g^2 N_0}{\omega - 2\epsilon + i2\gamma}$$



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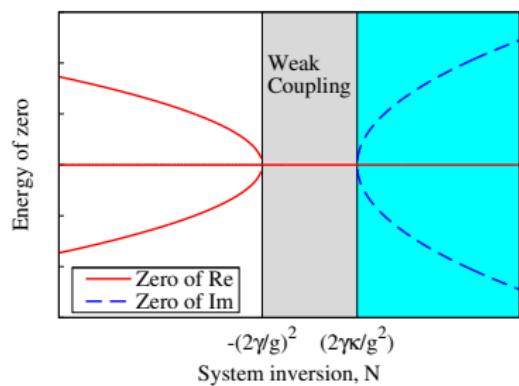
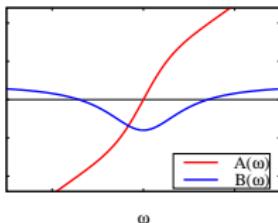
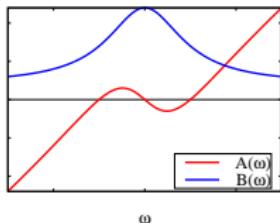
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# Spin in terms of two four-level systems

To include spin, replace 2 level system with 4 levels:  $|0\rangle, |L\rangle, |R\rangle, |LR\rangle$

- Bi-exciton binding  $E_{xx} \leftarrow U_1$
- Mean-field: find polarization given  $\psi_L, \psi_R$ .
- $E_{xx}$  has weak effect on  $T_c$

[Marchetti *et al* PRB, '08]

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$$\hat{h}_\alpha = \begin{pmatrix} 0 & g\psi_L & g\psi_R & 0 \\ g\psi_L^* & \varepsilon_\alpha - \Delta - \mu & 0 & g\psi_L \\ g\psi_R^* & 0 & \varepsilon_\alpha + \Delta - \mu & g\psi_R \\ 0 & g\psi_L^* & g\psi_R^* & 2(\varepsilon_\alpha - \mu) - E_{xx} \end{pmatrix}$$

- Bi-exciton binding  $E_{xx} \leftarrow U_1$
- Mean-field: find polarization given  $\psi_L, \psi_R$ .
- $E_{xx}$  has weak effect on  $T_c$

[Marchetti *et al* PRB, '08]

# Spin in terms of two four-level systems

To include spin, replace 2 level system with 4 levels:  $|0\rangle, |L\rangle, |R\rangle, |LR\rangle$

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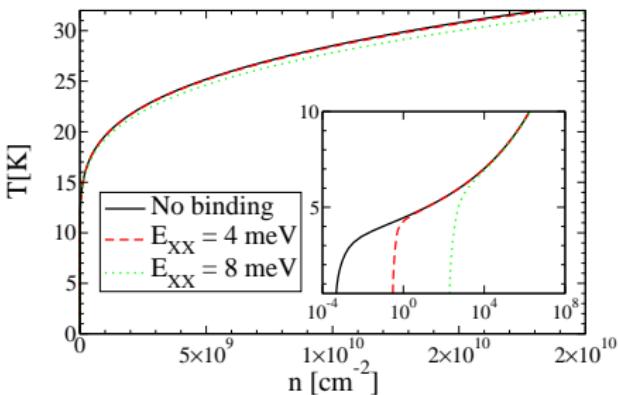
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[Marchetti *et al* PRB, '08]

# Non-equilibrium spinor condensate

$$H = H_0[\psi_L] + H_0[\psi_R] + (U_0 - 2\textcolor{blue}{U}_1)|\psi_L|^2|\psi_R|^2 + \frac{\Delta_{\perp}}{2} \left( |\psi_L|^2 - |\psi_R|^2 \right) + \Delta_{\parallel} (\psi_L^{\dagger}\psi_R + \text{H.c.})$$

Spinor Gross-Pitaevskii equation

$$i\partial_t\psi_L = \left[ -\frac{\nabla^2}{2m} + V(r) + U_0|\psi_L|^2 + i \left( \gamma_{\text{eff}} - \kappa - \Gamma_0|\psi_L|^2 \right) \right] \psi_L$$

- Left-right coupling:  $U_1$
- Cross-spin loss terms  $\Gamma_0$
- Magnetic field:  $\gamma_{\text{eff}}$
- Energy-dependent gain  $\kappa$
- [Moulton et al. PRL 103]

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[Wouters *et al* PRB '10]

# Non-equilibrium spinor system: two-mode model

Two-mode case (neglect spatial variation) [Wouters PRB '08]

$$i\partial_t \psi_L = \left[ U_0 |\psi_L|^2 + (U_0 - 2\textcolor{blue}{U}_1) |\psi_R|^2 + \frac{\Delta_{\perp}}{2} + i(\gamma_{\text{net}} - \Gamma_0 |\psi_L|^2) \right] \psi_L + \Delta_{\parallel} \psi_R$$

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$$\psi_L = \sqrt{R+z} e^{i\phi + \textcolor{blue}{i}\theta/2},$$

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Simple case  $\Gamma_1 = \eta = 0$

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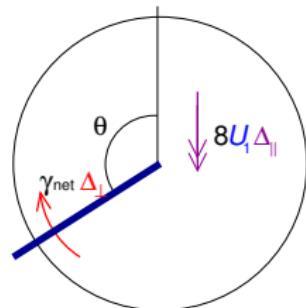
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Simple case  $\Gamma_1 = \eta = 0$

Josephson regime:  $\Delta_{\parallel} \ll U_1 R, z \ll R$ .

Damped, driven pendulum

$$\ddot{\theta} + 2\gamma_{\text{net}} \dot{\theta} = 8U_1 \Delta_{\parallel} \frac{\gamma_{\text{net}}}{\Gamma_0} \sin(\theta) - 2\gamma_{\text{net}} \Delta_{\perp}$$

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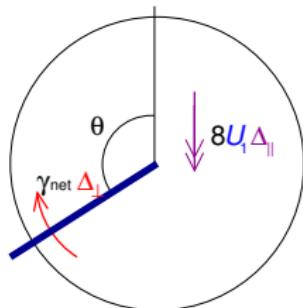
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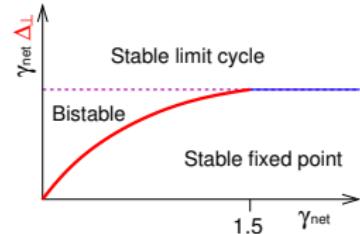
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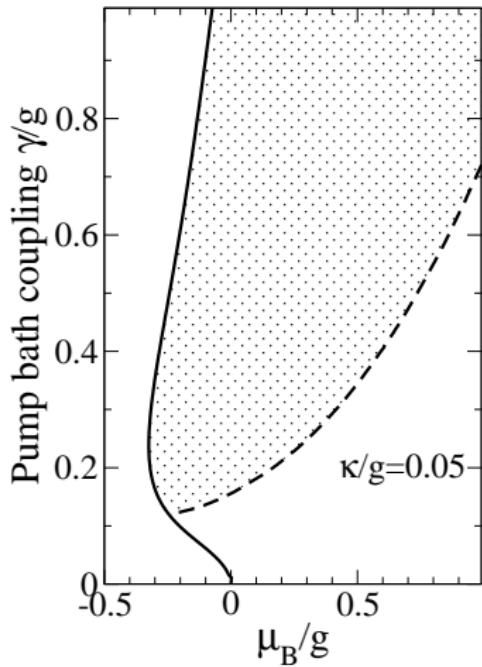
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Cartoon:



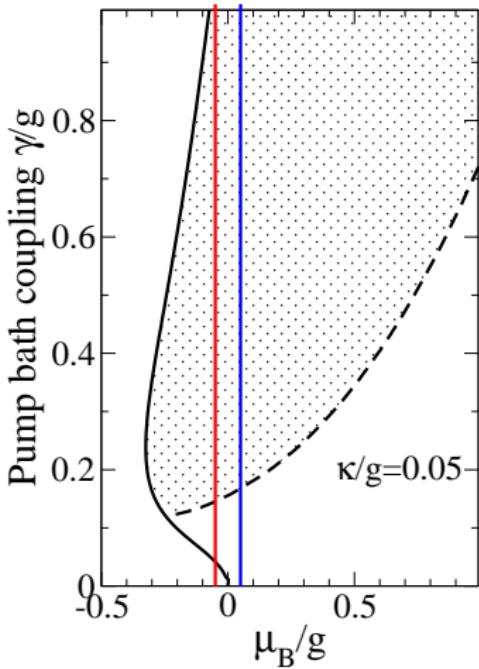
# Zero temperature phase diagram

$$\tilde{\omega}_0 - i\kappa = g^2 \gamma \sum_{\alpha} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(\tilde{\epsilon}_{\alpha} + i\gamma)}{[(\nu - E_{\alpha})^2 + \gamma^2][(v + E_{\alpha})^2 + \gamma^2]}.$$



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