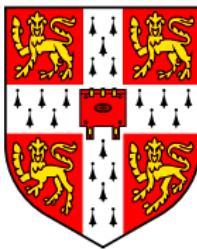


# Collapse and revivals of the photon field in a Landau-Zener process

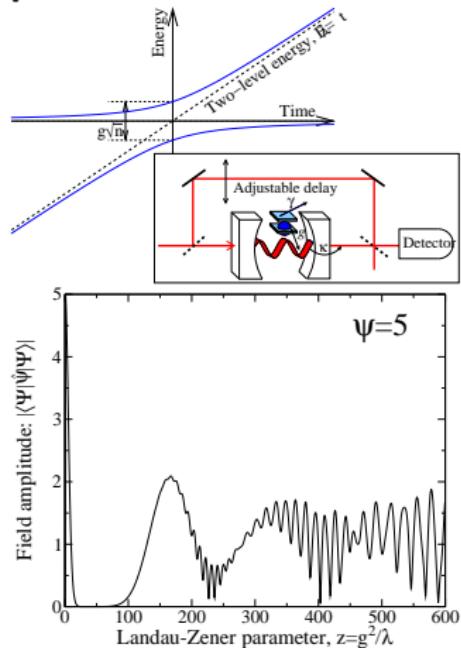
**Jonathan Keeling**

Central European Workshop on Quantum Optics

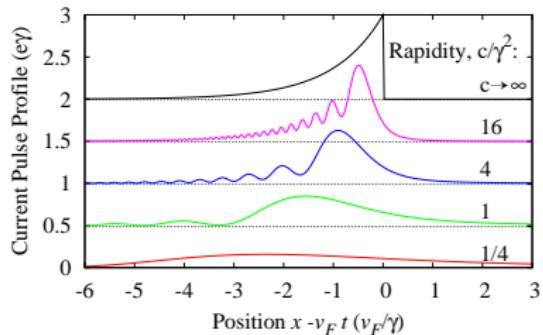
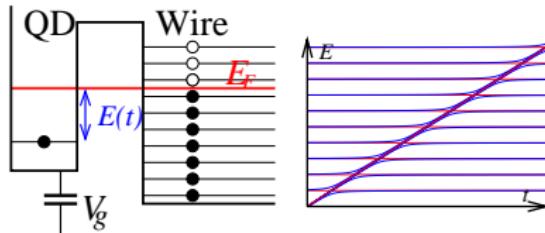


# Generalisations of Landau-Zener problems

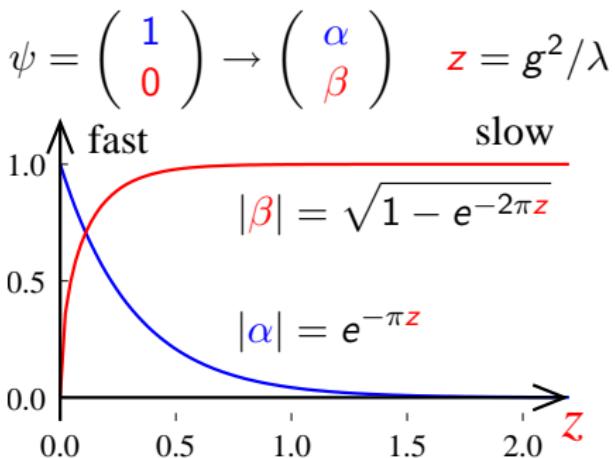
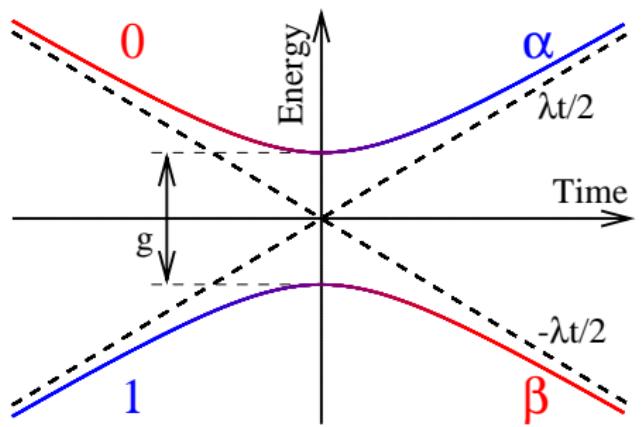
## Two-level atom coupled to photon field



## Quantum dot coupled to 1D wire



# The Landau-Zener problem



$$i\partial_t \psi = \begin{pmatrix} \lambda t/2 & g \\ g & -\lambda t/2 \end{pmatrix} \psi$$

# Coherent field Landau-Zener problem

- Hamiltonian:  $H = \frac{\lambda t}{2} s^z + g (\psi^\dagger s^- + \psi s^+)$ ,

Initial coherent state  $|\psi(0)\rangle = \sum_n \sqrt{n} |n, \downarrow\rangle$

- Each pair  $|n, \downarrow\rangle \rightarrow |n+1, \downarrow\rangle$  undergoes LZ transition

$$H_{n,n+1} = \begin{pmatrix} \lambda t/2 & g\sqrt{n+1} \\ g\sqrt{n+1} & -\lambda t/2 \end{pmatrix}, \quad \begin{pmatrix} 1 \\ 0 \end{pmatrix} \rightarrow \begin{pmatrix} A_{n+1} \\ B_{n+1} \end{pmatrix}$$

Final state:

$$|\Psi(-\infty)\rangle = e^{-iE_F t} \sum_{n=0}^{\infty} \frac{i^n}{\sqrt{n!}} [A_{n+1}(n, \downarrow) + B_{n+1}(n+1, \downarrow)]$$

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→ Each atom in the initial state undergoes LZ transition

$$H_{\text{LZ},1} = \begin{pmatrix} \lambda t/2 & g\sqrt{n+1} \\ g\sqrt{n+1} & -\lambda t/2 \end{pmatrix} \rightarrow \begin{pmatrix} 1 \\ 0 \end{pmatrix} \rightarrow \begin{pmatrix} A_{n+1} \\ B_{n+1} \end{pmatrix}$$

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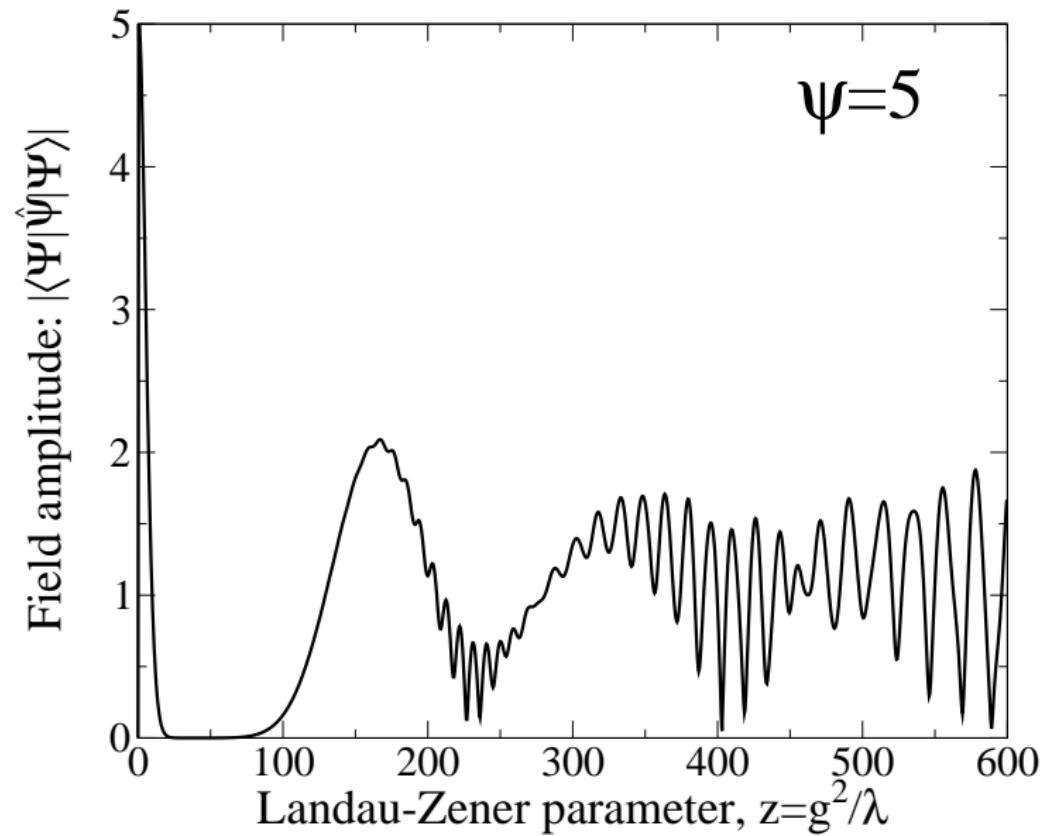
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# Collapse and revivals of field amplitude



# Explaining results: expansion

Adiabatic limit:  $\omega = g^2/\lambda \gg 1$

$$|\Psi(T)\rangle = e^{-|\psi|^2/2} \sum_{n=0}^{\infty} \frac{\psi^n}{\sqrt{n!}} \left[ A_{n+1} |n, \uparrow\rangle + B_{n+1} |n+1, \downarrow\rangle \right]$$

$$\langle \Psi(t) | \Psi \rangle = \psi e^{-|\psi|^2} \sum_n \frac{1}{n!} \rightarrow \sqrt{\frac{e}{\pi}} \frac{1}{\sqrt{n+1}}$$

$$\Delta\phi_n = \omega \left[ (n+1) \ln(n+1) - n \ln n \right]$$

Want  $\Delta\phi_n$  for  $n = |\psi|^2$ .

$$\Delta\phi_{n=|\psi|^2+m} \approx \Delta\phi_{|\psi|^2} + \frac{m}{|\psi|^2} - \frac{m^2}{2|\psi|^4} + \dots$$

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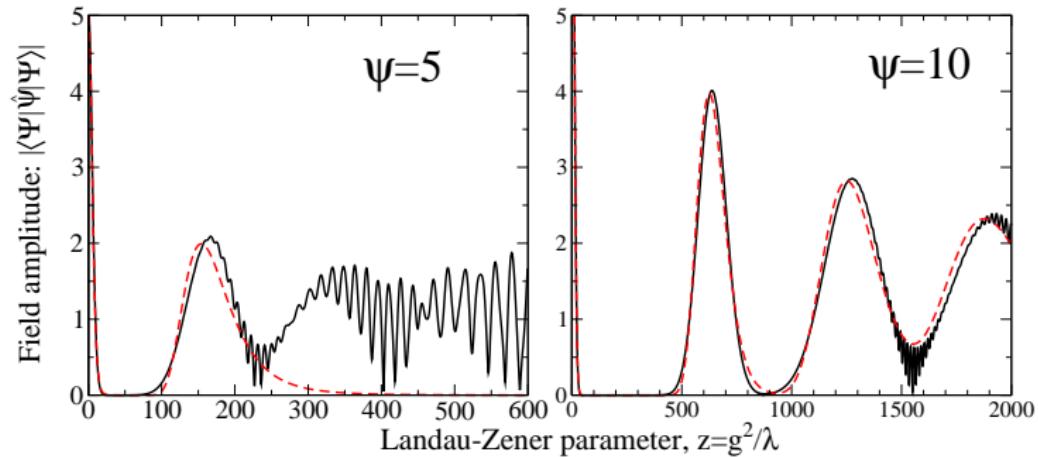
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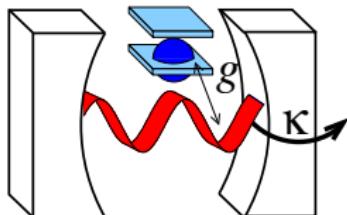
$$\Delta\phi_{n=|\psi|^2+m} \simeq \Delta\phi_{|\psi|^2} + \frac{zm}{|\psi|^2} - \frac{zm^2}{2|\psi|^4} + \dots$$

# Explaining results: comparison

$$\langle \Psi | \psi | \Psi \rangle = \frac{|\psi|}{(1 + z^2/|\psi|^4)^{1/4}} \sum_{N=0}^{N_{\max}} \exp \left[ -\frac{(z - 2\pi N |\psi|^2)^2}{2|\psi|^2(1 + z^2/|\psi|^4)} \right]$$



# Effects of photon decay



$$\partial_t \hat{\rho} = -i [\hat{H}, \hat{\rho}] + L_\kappa[\hat{\rho}]$$

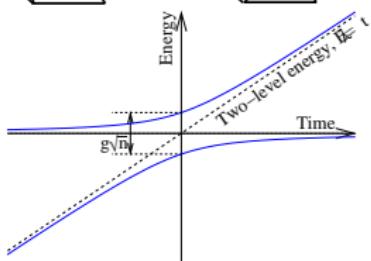
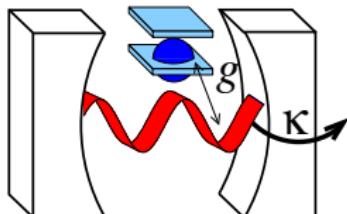
$$L_\kappa[\hat{\rho}] = -\frac{\kappa}{2} (\hat{\psi}^\dagger \hat{\psi} \hat{\rho} + \hat{\rho} \hat{\psi}^\dagger \hat{\psi} - 2 \hat{\psi} \hat{\rho} \hat{\psi}^\dagger)$$

Consider,  $\Lambda_n = \rho_{n,n+1}$

$$\frac{d\Lambda_n}{dt} = i \left[ \frac{d\Delta\phi_{n+1}}{dt} \right] \Lambda_n - \kappa \left[ \left( n - \frac{1}{2} \right) \Lambda_n - \sqrt{n(n+1)} \Lambda_{n+1} \right]$$

When  $|v| \leq g\sqrt{n}/\lambda$ , decay rate  $\kappa v \approx n |d|^2 \gg \gamma$

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Photon loss does not switch branch

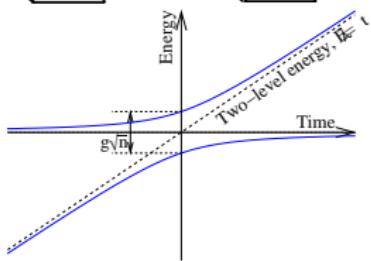
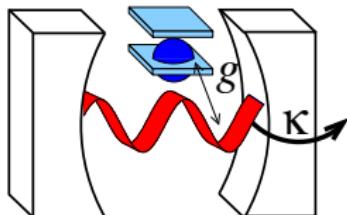
$$P_{\text{trans}} = \frac{|\langle n-1, - | \hat{\psi} | n, + \rangle|^2}{\langle n, + | \hat{\psi}^\dagger \hat{\psi} | n, + \rangle} \leq \frac{27}{256n^2} \ll 1$$

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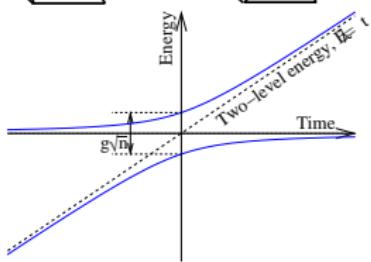
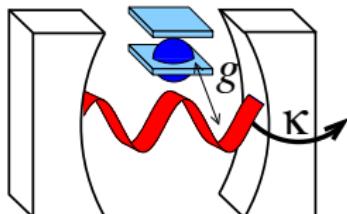
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# Effects of photon decay



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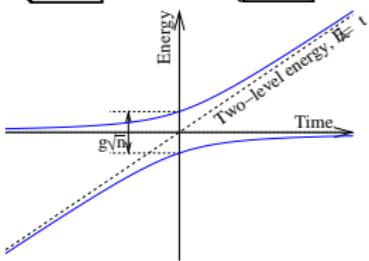
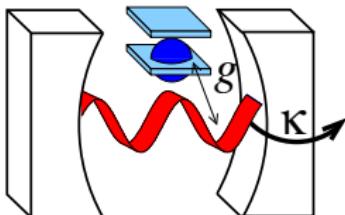
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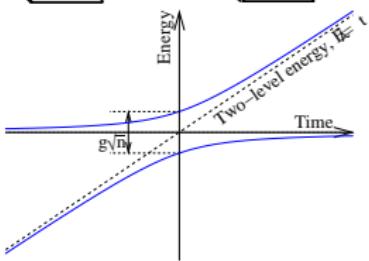
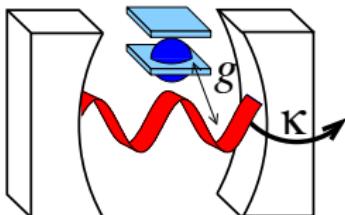
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# Effects of photon decay



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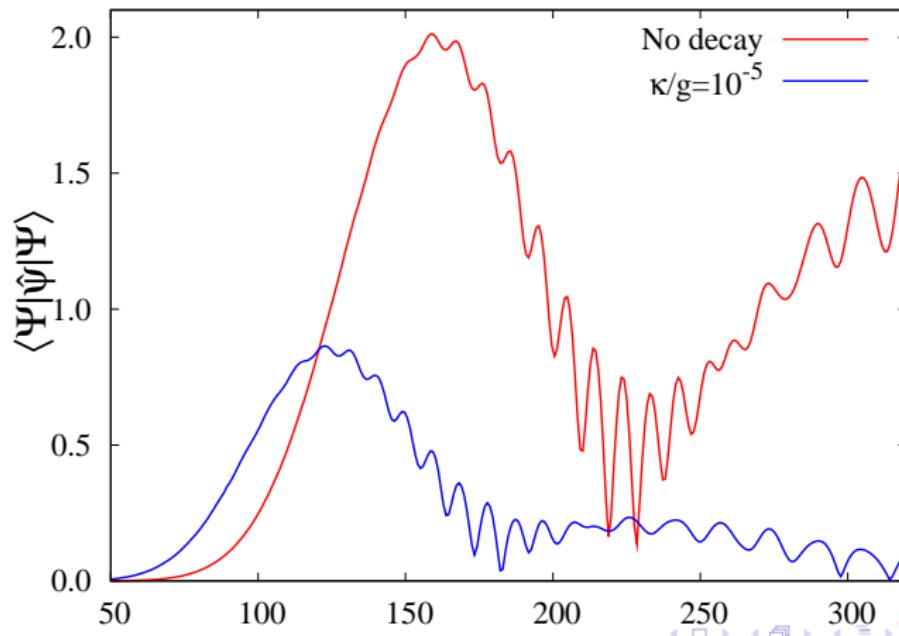
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When  $|t| \leq g\sqrt{n}/\lambda$ , decay rate  $\kappa n \simeq \kappa |\psi|^2 \gg \kappa$ .

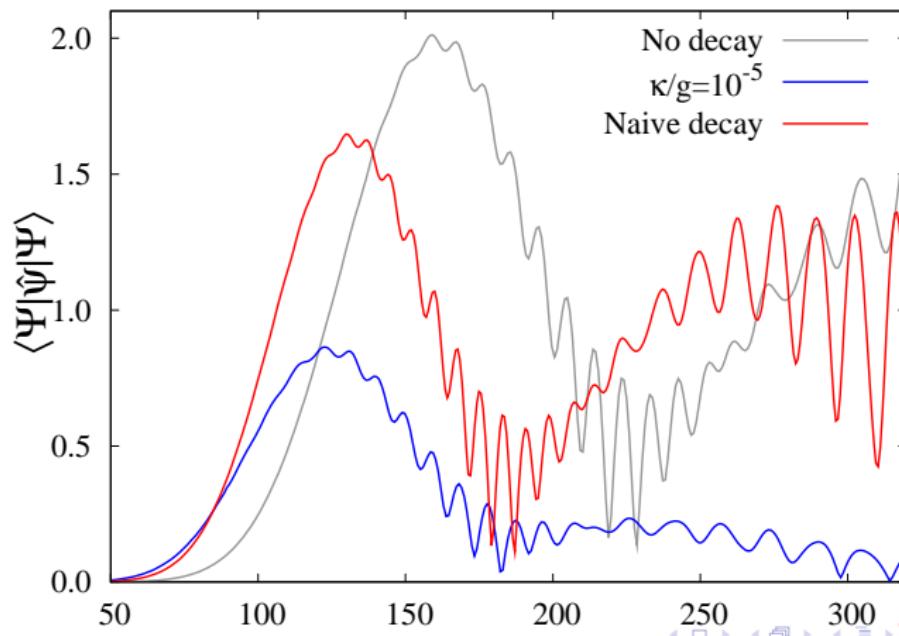
# Comparison of photon decay

$$\langle \Psi | \hat{\psi}(\psi_0, \kappa) | \Psi \rangle$$



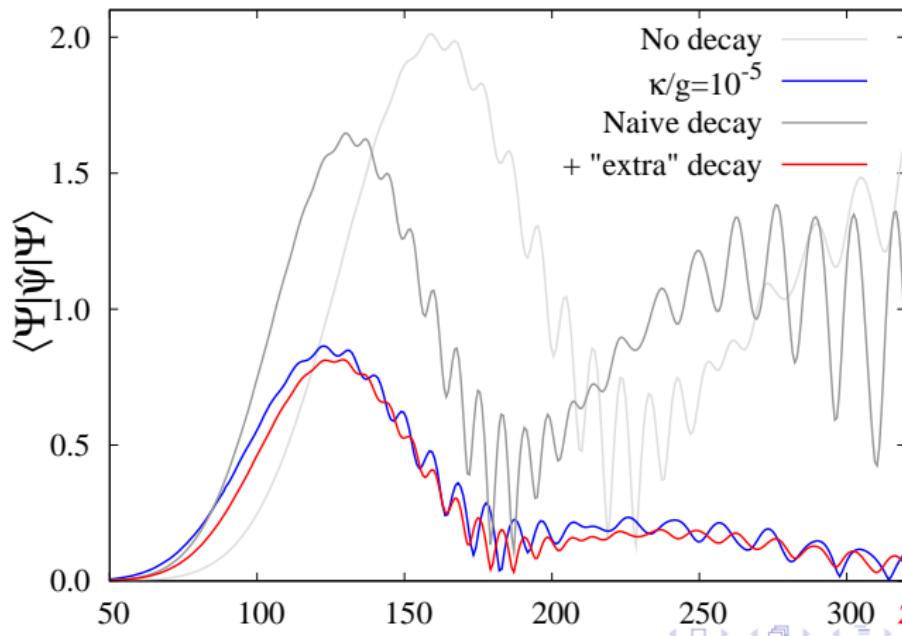
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# Many two-level systems

$$H = \sum_i^N \left( \lambda t s_i^z + g(s_i^+ \psi + s_i^- \psi^\dagger) \right)$$

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For  $N \gg n$ , expand in  $n/N$ :

$$\phi_{\textcolor{red}{n}} = A(N) + B(N) \textcolor{red}{n} + Cz \textcolor{red}{n}^2 + \mathcal{O}\left(\frac{1}{N}\right)$$

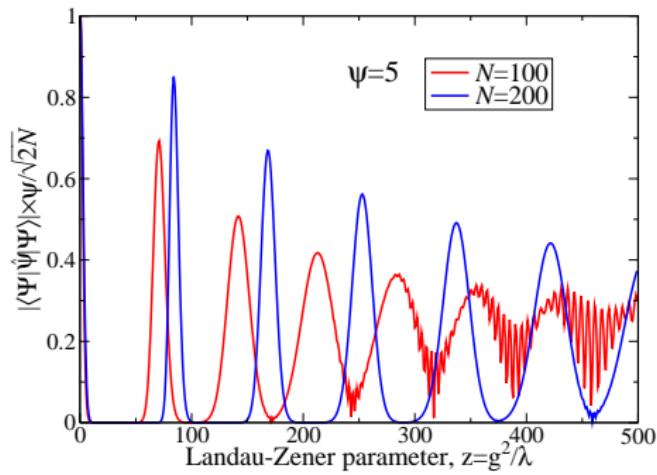
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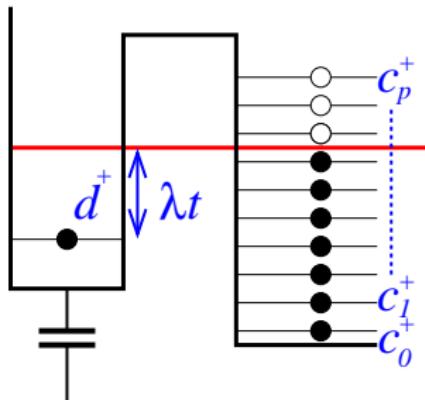
$$\langle \Psi | \psi | \Psi \rangle = \psi e^{-|\psi|^2} \sum_n \frac{|\psi|^{2n}}{n!} \sqrt{\frac{n+1+N}{n+1}} e^{i(\phi_{n+2}-\phi_{n+1})}$$

For  $N \gg n$ , expand in  $n/N$ :

$$\phi_n = A(N) + B(N)n + Cz n^2 + \mathcal{O}\left(\frac{1}{N}\right)$$



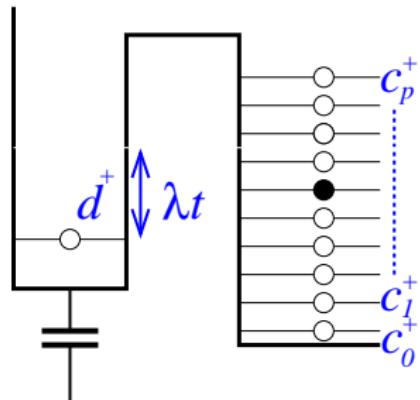
# Quantum dot coupled to 1D wire



$$H = \lambda t d^\dagger d + \sum \varepsilon_p c_p^\dagger c_p + g(d^\dagger c_p + c_p^\dagger d)$$

- Initially: Filled fermi sea  $\prod_{\varepsilon_p < \varepsilon_F} c_p^\dagger |\Omega\rangle$
- Find final state of wire

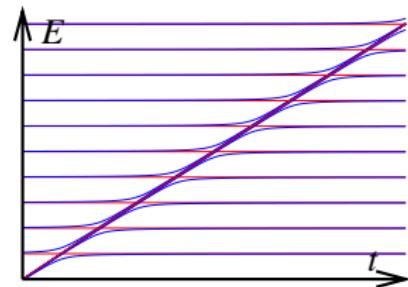
# Quantum dot coupled to 1D wire



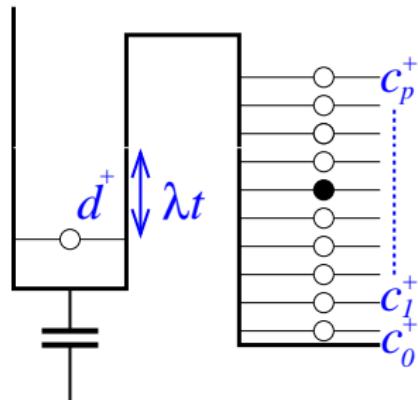
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1 particle: Demkov-Osherov problem



# Quantum dot coupled to 1D wire

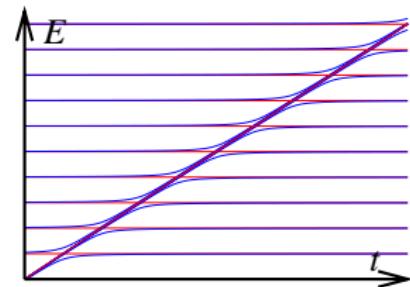


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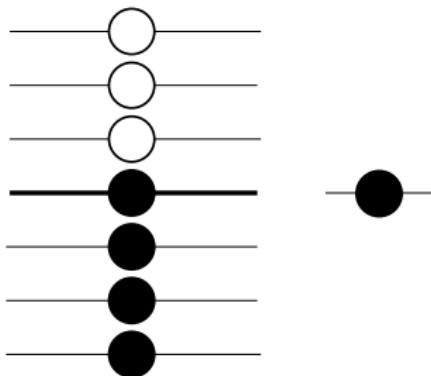
1 particle: Demkov-Osherov problem

$$T(\varepsilon, \varepsilon') \propto \exp \left[ -\frac{g^2 \nu}{\lambda} (\varepsilon - \varepsilon') - \frac{i}{2\lambda} (\varepsilon^2 - \varepsilon'^2) \right]$$



# Many fermion problem

$$T(\varepsilon, \varepsilon') \propto \exp \left[ -\frac{g^2 \nu}{\lambda} (\varepsilon - \varepsilon') - \frac{i}{2\lambda} (\varepsilon^2 - \varepsilon'^2) \right] = C \phi_+(\varepsilon) \langle \phi_-(\varepsilon') \rangle$$



Transitions between fermion states:

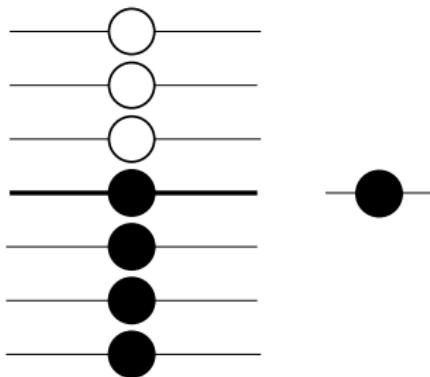
$$\tilde{c}_\varepsilon = \sum_{\varepsilon'} T(\varepsilon, \varepsilon') c_{\varepsilon'}$$

$$\begin{aligned} P_{ab} &= T_{a \rightarrow a} T_{b \rightarrow b} - T_{a \rightarrow b} T_{b \rightarrow a} \\ &= (d|a_a\rangle \langle a_a|b_b\rangle \langle b_b|a_a\rangle) - (d|a_a\rangle \langle a_a|b_b\rangle \langle b_b|a_a\rangle) = 0. \end{aligned}$$

- Max number of particles transferred = rank of  $T = 1$ .
- Particle transferred to state  $|a_a(\cdot)\rangle$

# Many fermion problem

$$T(\varepsilon, \varepsilon') \propto \exp \left[ -\frac{g^2 \nu}{\lambda} (\varepsilon - \varepsilon') - \frac{i}{2\lambda} (\varepsilon^2 - \varepsilon'^2) \right] = C |\phi_+(\varepsilon)\rangle \langle \phi_-(\varepsilon')|$$



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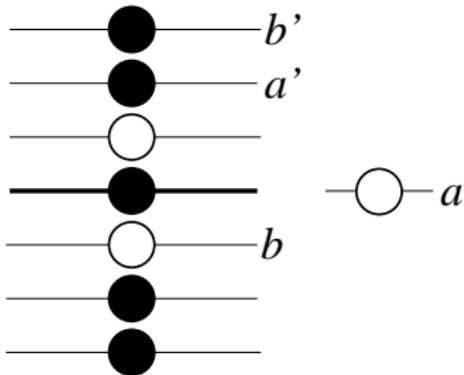
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- Max number of particles transferred = rank of  $T = 1$ .
- Possible transfer to state  $|a_+\rangle \langle a_+|$

# Many fermion problem

$$T(\varepsilon, \varepsilon') \propto \exp \left[ -\frac{g^2 \nu}{\lambda} (\varepsilon - \varepsilon') - \frac{i}{2\lambda} (\varepsilon^2 - \varepsilon'^2) \right] = C |\phi_+(\varepsilon)\rangle \langle \phi_-(\varepsilon')|$$



Transitions between fermion states:

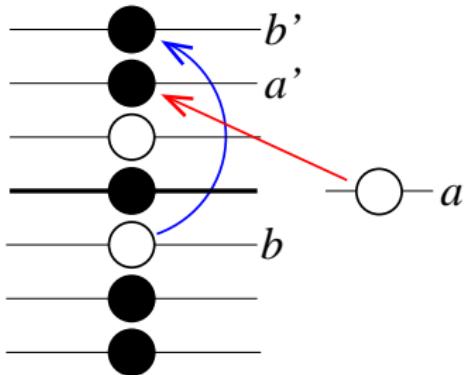
$$\tilde{c}_\varepsilon = \sum_{\varepsilon'} T(\varepsilon, \varepsilon') c_{\varepsilon'}$$

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- Max number of particles transferred = rank of  $T = 1$ .
- Particle transferred to state  $|\phi_+(\varepsilon)\rangle$

# Many fermion problem

$$T(\varepsilon, \varepsilon') \propto \exp \left[ -\frac{g^2 \nu}{\lambda} (\varepsilon - \varepsilon') - \frac{i}{2\lambda} (\varepsilon^2 - \varepsilon'^2) \right] = C |\phi_+(\varepsilon)\rangle \langle \phi_-(\varepsilon')|$$



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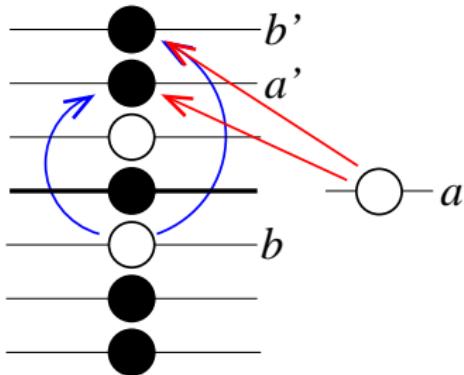
$$\tilde{c}_\varepsilon = \sum_{\varepsilon'} T(\varepsilon, \varepsilon') c_{\varepsilon'}$$

$$\begin{aligned} P_{ab} &= T_{a \rightarrow b} + T_{a \rightarrow b'} - T_{b \rightarrow a} - T_{b \rightarrow a'} \\ &= (\delta_{a,b}) (\phi_+|a\rangle \langle b| \phi_-) (\phi_-|b\rangle) \\ &\quad - (\delta_{a,b'}) (\phi_+|a\rangle \langle b'| \phi_-) (\phi_-|b'\rangle) = 0. \end{aligned}$$

- Max number of particles transferred = rank of  $T = 1$ .
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# Many fermion problem

$$T(\varepsilon, \varepsilon') \propto \exp \left[ -\frac{g^2 \nu}{\lambda} (\varepsilon - \varepsilon') - \frac{i}{2\lambda} (\varepsilon^2 - \varepsilon'^2) \right] = C |\phi_+(\varepsilon)\rangle \langle \phi_-(\varepsilon')|$$



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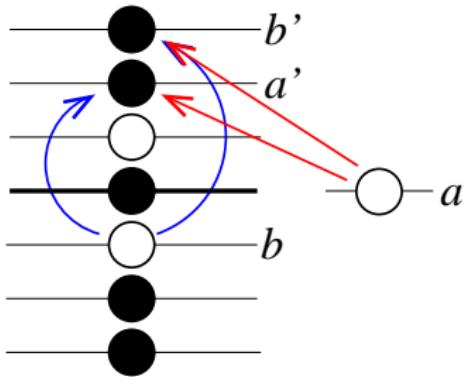
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# Many fermion problem

$$T(\varepsilon, \varepsilon') \propto \exp \left[ -\frac{g^2 \nu}{\lambda} (\varepsilon - \varepsilon') - \frac{i}{2\lambda} (\varepsilon^2 - \varepsilon'^2) \right] = C |\phi_+(\varepsilon)\rangle \langle \phi_-(\varepsilon')|$$



Transitions between fermion states:

$$\tilde{c}_\varepsilon = \sum_{\varepsilon'} T(\varepsilon, \varepsilon') c_{\varepsilon'}$$

$$P_2 = T_{a \rightarrow a'} T_{b \rightarrow b'} - T_{a \rightarrow b'} T_{b \rightarrow a'}$$

$$= (\delta_{a,a'})(\delta_{b,b'}) (\delta_{a',b}) (\delta_{b,a}) = 1$$

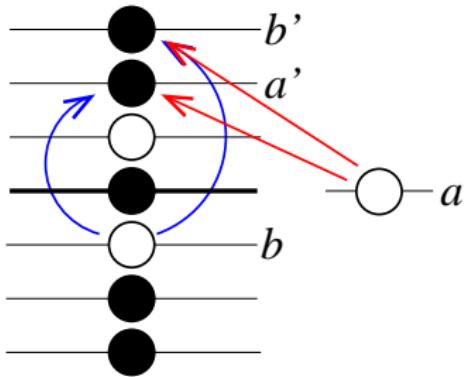
$$- (\delta_{a,a'})(\delta_{b,b'}) (\delta_{a',b}) (\delta_{b,a}) = 0$$

Max number of particles transferred = rank of  $T = 1$ .

Particle transferred to state  $|a_+\rangle$

# Many fermion problem

$$T(\varepsilon, \varepsilon') \propto \exp \left[ -\frac{g^2 \nu}{\lambda} (\varepsilon - \varepsilon') - \frac{i}{2\lambda} (\varepsilon^2 - \varepsilon'^2) \right] = C |\phi_+(\varepsilon)\rangle \langle \phi_-(\varepsilon')|$$



Transitions between fermion states:

$$\tilde{c}_\varepsilon = \sum_{\varepsilon'} T(\varepsilon, \varepsilon') c_{\varepsilon'}$$

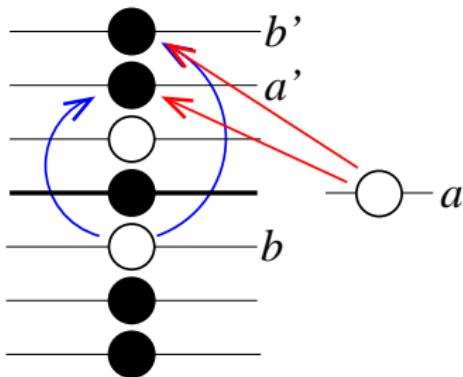
$$\begin{aligned} P_2 &= T_{a \rightarrow a'} T_{b \rightarrow b'} - T_{a \rightarrow b'} T_{b \rightarrow a'} \\ &= \langle a' | \phi_+ \rangle \langle \phi_- | a \rangle \langle b' | \phi_+ \rangle \langle \phi_- | b \rangle \\ &\quad - \langle a' | \phi_+ \rangle \langle \phi_- | b \rangle \langle b' | \phi_+ \rangle \langle \phi_- | a \rangle = 0. \end{aligned}$$

→ Max number of particles transferred = rank of  $T = 1$ .

Particle transferred to state  $|\phi_+(\varepsilon)\rangle$

# Many fermion problem

$$T(\varepsilon, \varepsilon') \propto \exp \left[ -\frac{g^2 \nu}{\lambda} (\varepsilon - \varepsilon') - \frac{i}{2\lambda} (\varepsilon^2 - \varepsilon'^2) \right] = C |\phi_+(\varepsilon)\rangle \langle \phi_-(\varepsilon')|$$



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- Max number of particles transferred = rank of  $T = 1$ .
- Particle transferred to state  $|\phi_+(\varepsilon)\rangle$ .

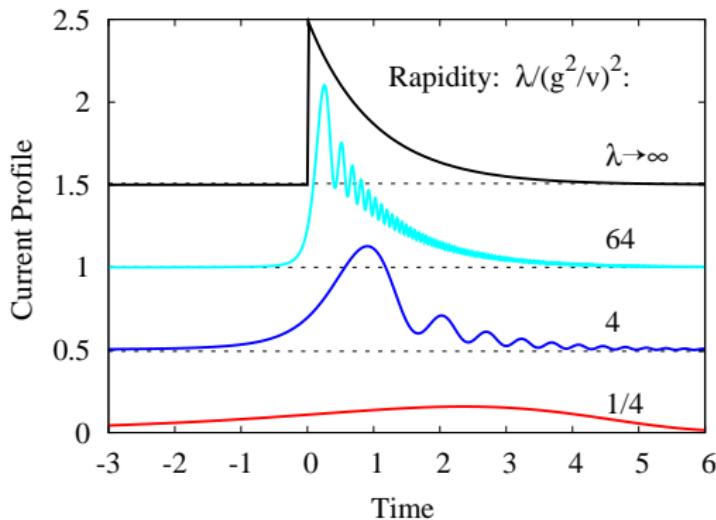
# Current pulse profile

$$\psi(t) = \sqrt{\frac{g^2\nu}{2\pi\lambda}} \int_0^\infty d\varepsilon \exp \left[ -i\varepsilon t - \frac{g^2\nu\varepsilon}{2\lambda} + i\frac{\varepsilon^2}{2\lambda} \right]$$

# Current pulse profile

Driving need not be adiabatic

$$\psi(t) = \sqrt{\frac{g^2 \nu}{2\pi\lambda}} \int_0^\infty d\varepsilon \exp \left[ -i\varepsilon t - \frac{g^2 \nu \varepsilon}{2\lambda} + i\frac{\varepsilon^2}{2\lambda} \right]$$



# Acknowledgements

## Two-level system and photon field

Victor Gurarie.  
Chris Pointon.



## Quantum dot and 1D wire

Leonid Levitov.  
Andriy Shytov.



Engineering and Physical Sciences  
Research Council

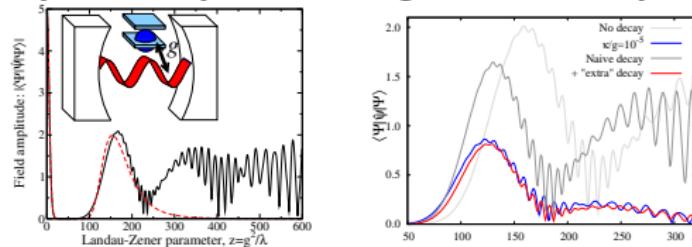
## Funding:



Pembroke College

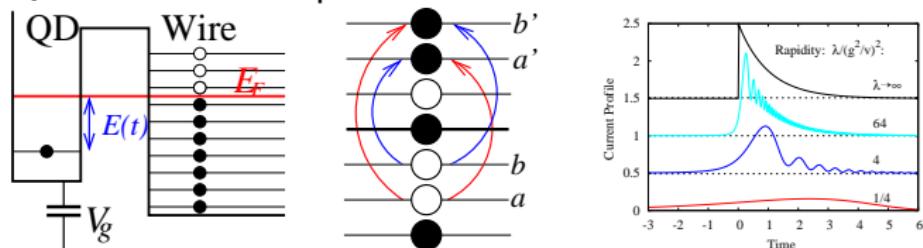
# Conclusions

- Dynamically driven single two-level system:



[Keeling and Gurarie, Phys. Rev. Lett. 101 033001 (2008)]

- Quantum dot coupled to 1D wire



[Keeling, Shytov and Levitov, Phys. Rev. Lett. 101 196404 (2008)]

# Extra material

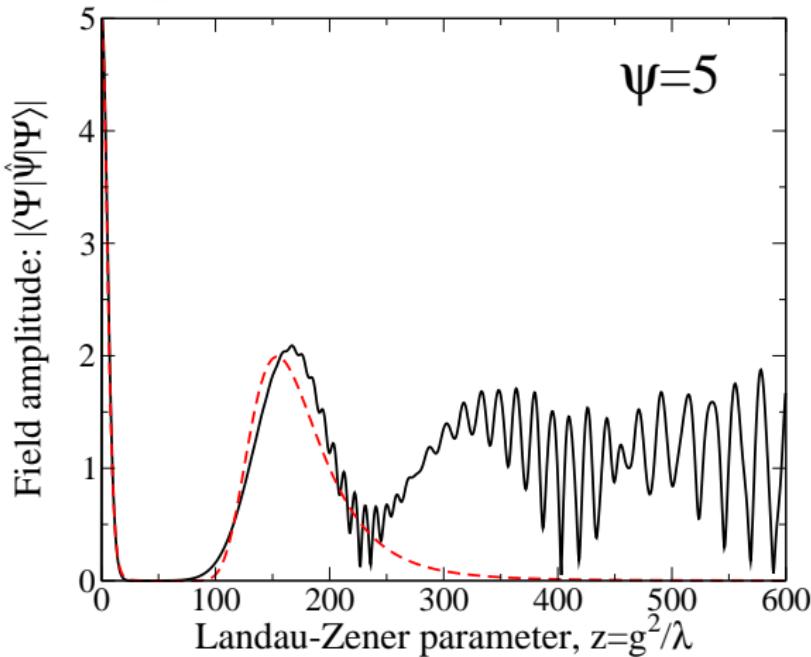
## ③ Coherent Landau-Zener process

- Wigner function
- Graphical description of revivals
- Decay effects
- Systems for LZ problem
- Many spins LZ problem

## ④ Localised fermion coupled to continuum

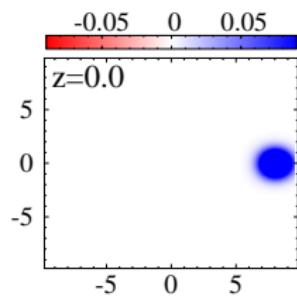
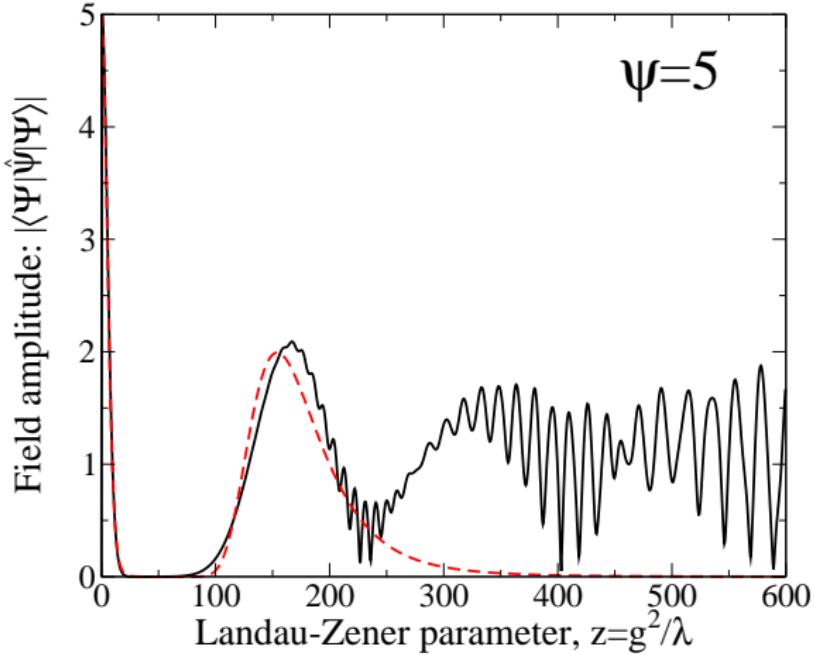
- Finding  $T(\varepsilon, \varepsilon')$
- Periodic driving
- Measuring Noise
- Noisy driving

# Revivals are not coherent states



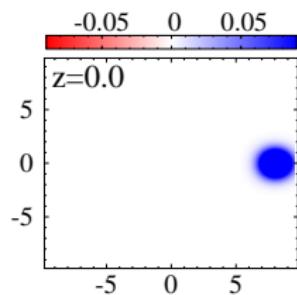
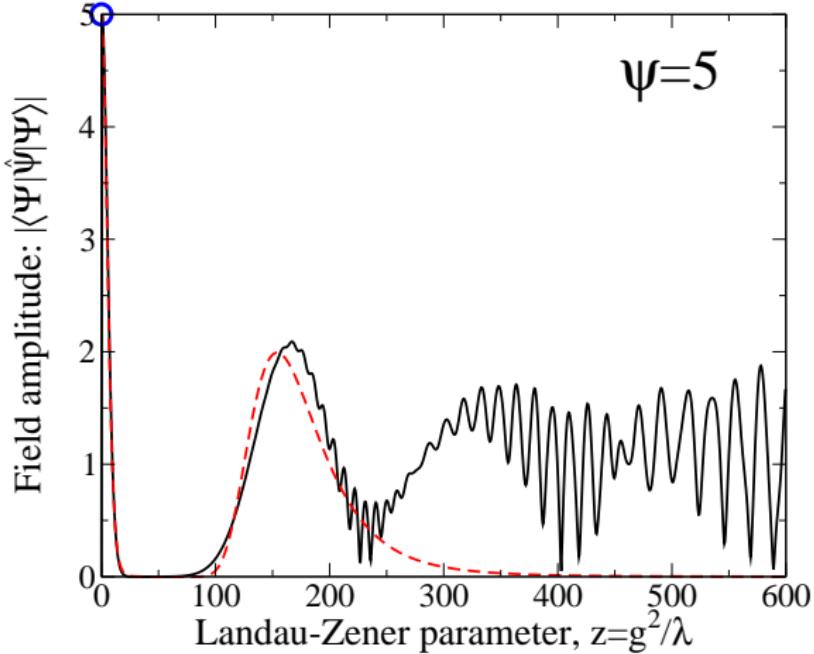
# Revivals are not coherent states

$$W(x, p) = \frac{1}{\pi} \int dy \Psi^*(x + y) \Psi(x - y) e^{2ipy}.$$



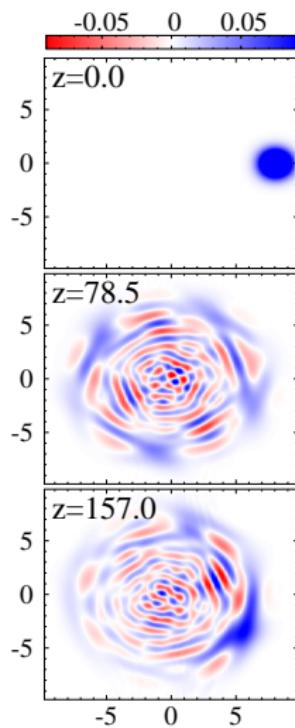
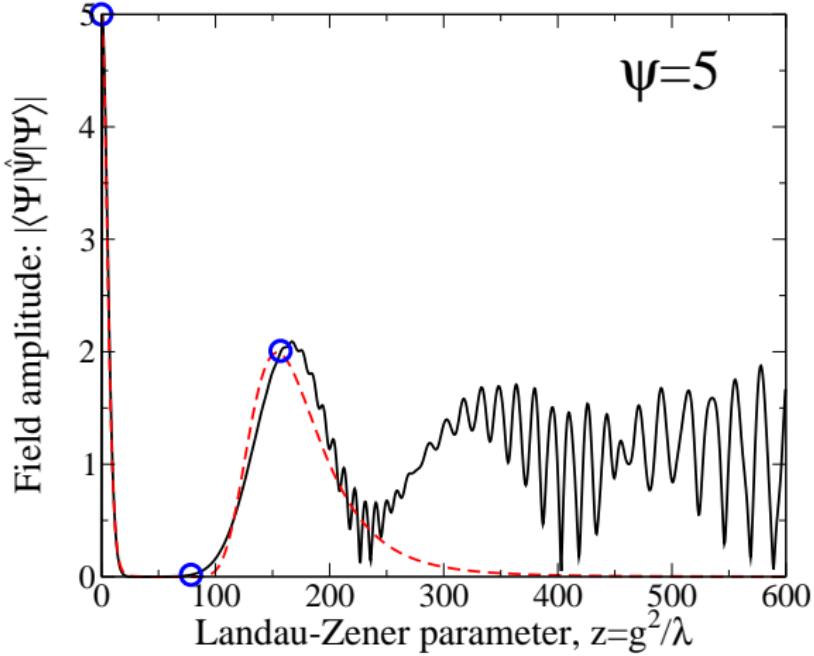
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# Understanding collapse and revival

$$\langle \Psi | \psi | \Psi \rangle = \psi e^{-|\psi|^2} \sum_n \frac{|\psi|^{2n}}{n!} \sqrt{\frac{n+2}{n+1}} e^{i(\phi_{n+2} - \phi_{n+1})}$$

$$\Delta\phi_n = \pi \left[ (n+1)\ln(n+1) - n\ln n \right]$$

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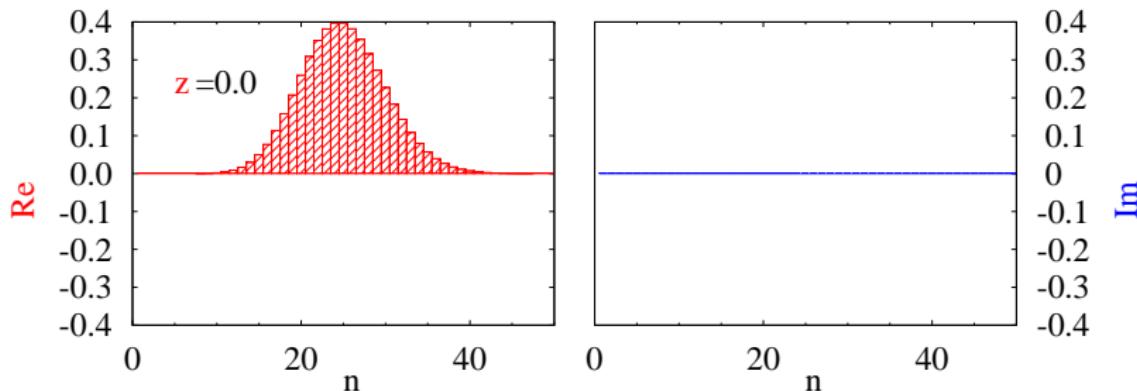
$$\langle \Psi | \psi | \Psi \rangle = \psi e^{-|\psi|^2} \sum_n \frac{|\psi|^{2n}}{n!} e^{i\Delta\phi_{n+1}}$$

$$\Delta\phi_n = \textcolor{red}{z} \left[ (n+1) \ln(n+1) - n \ln n + \ln \left( \frac{\textcolor{red}{z}}{e\lambda T^2} \right) \right]$$

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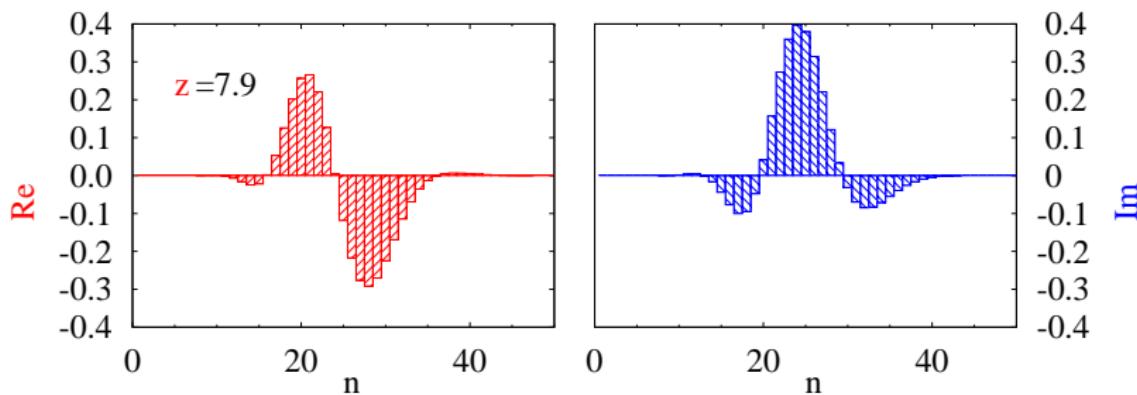
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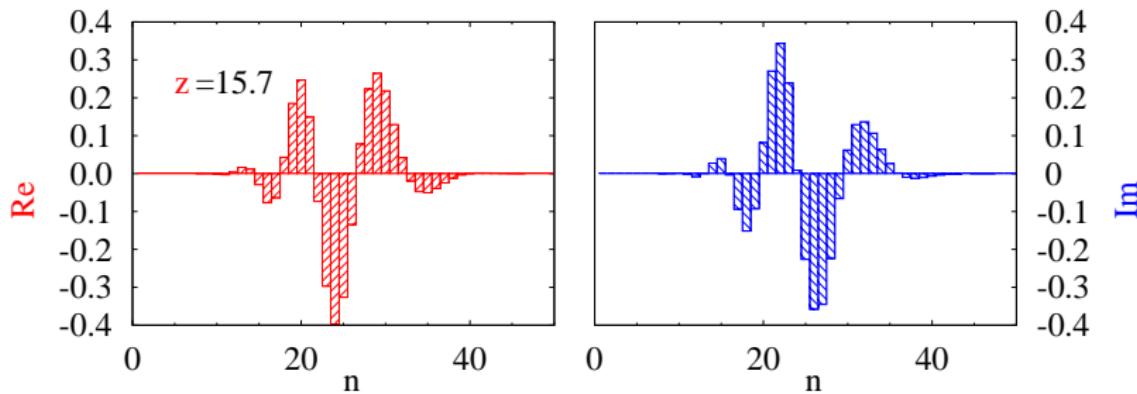
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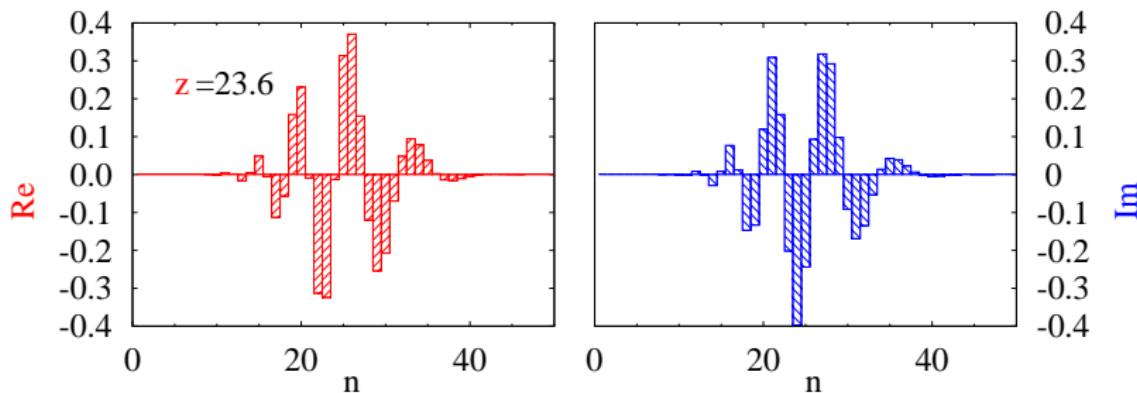
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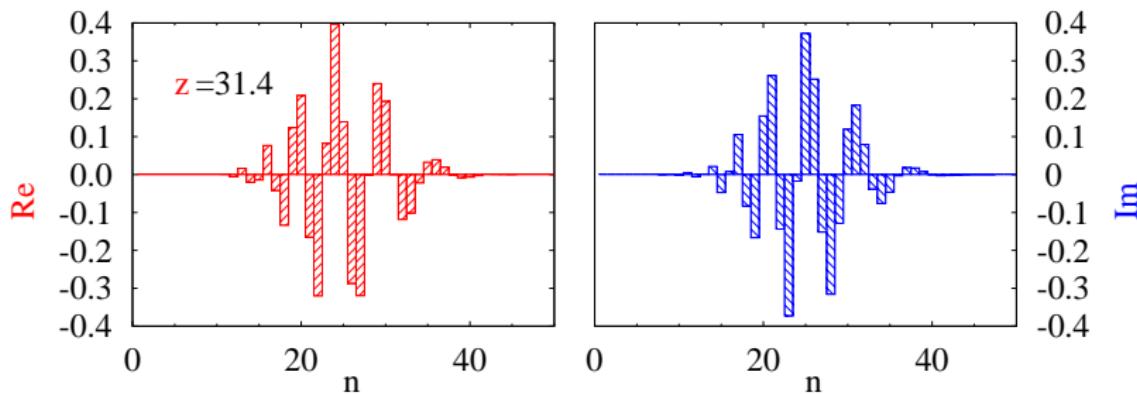
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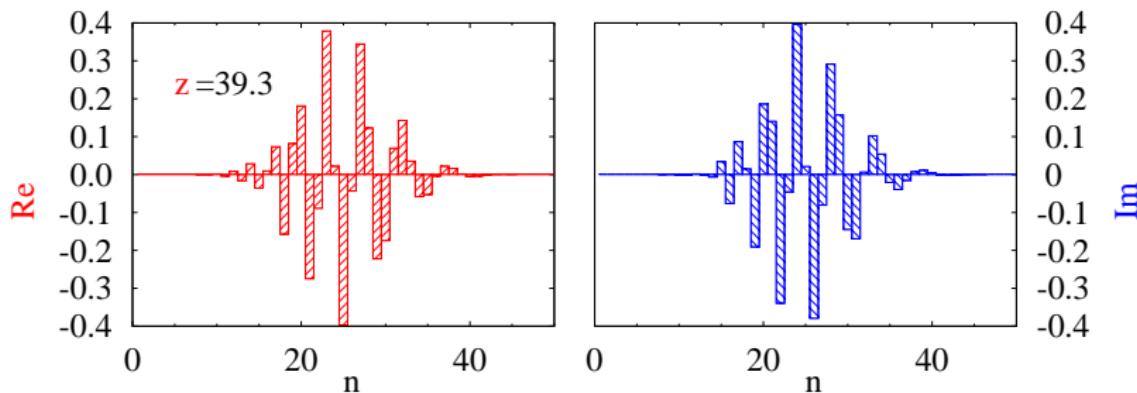
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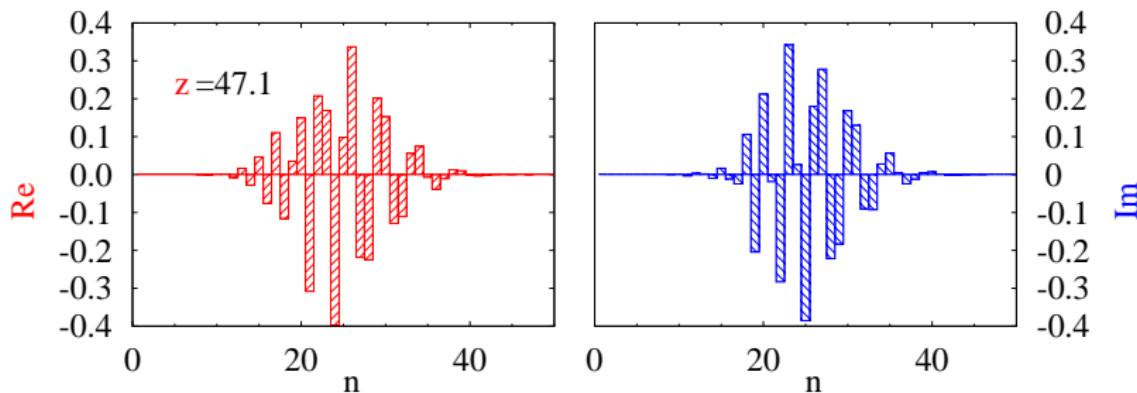
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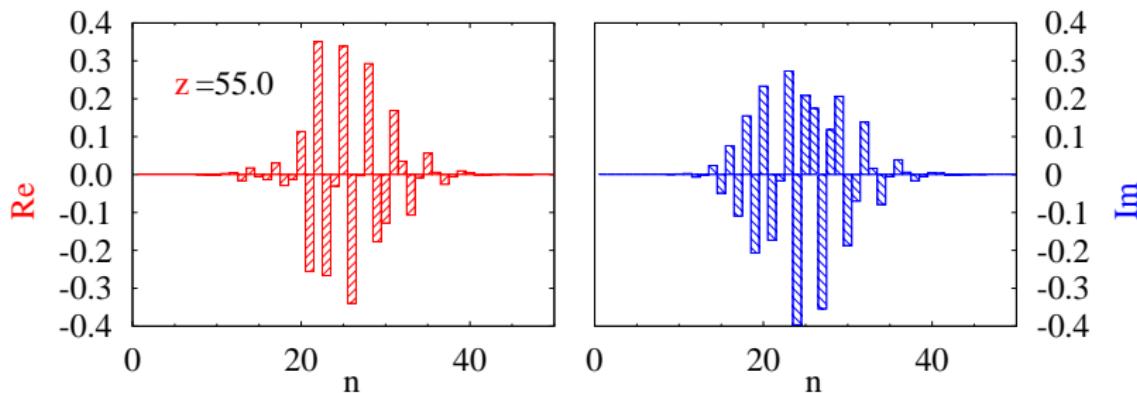
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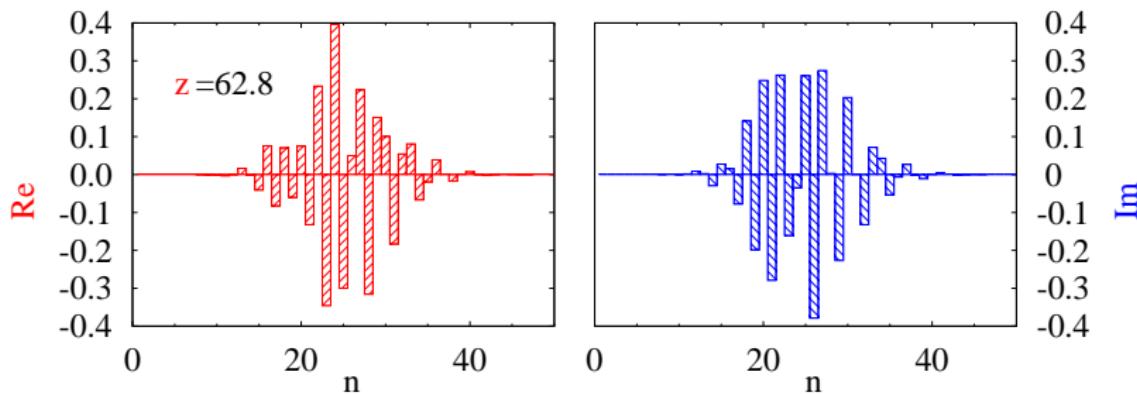
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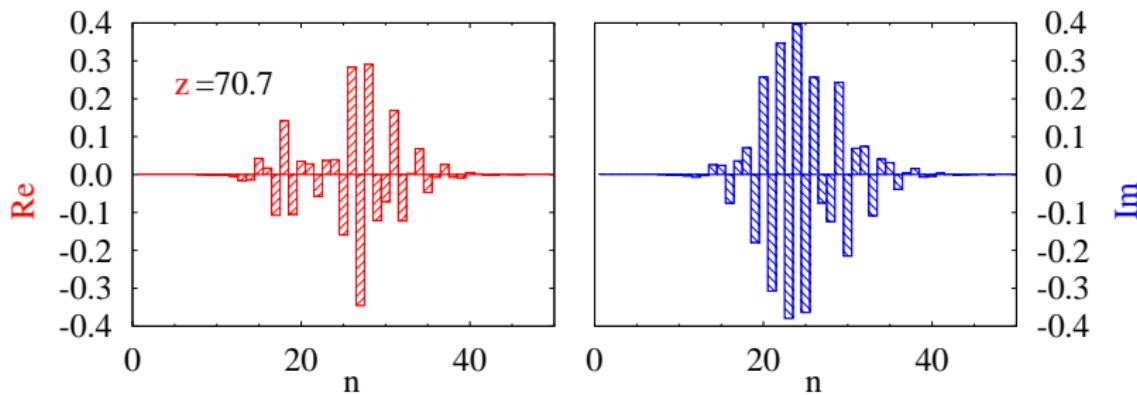
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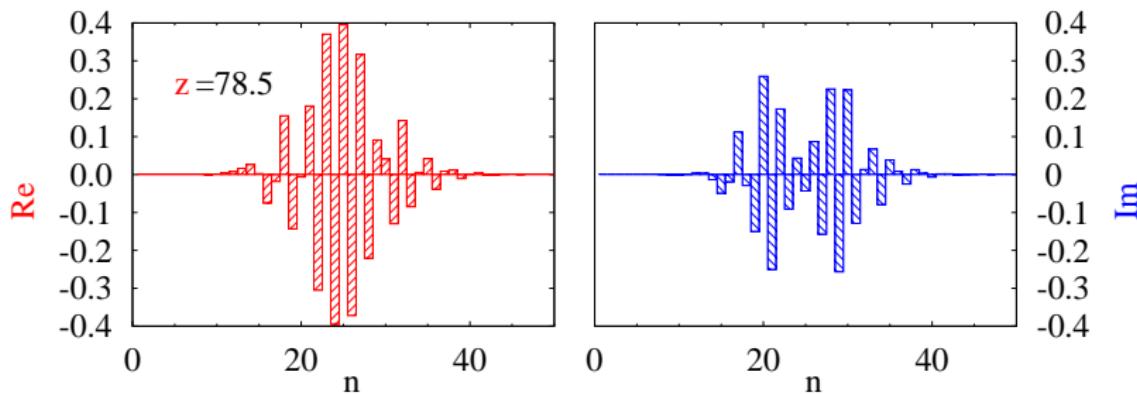
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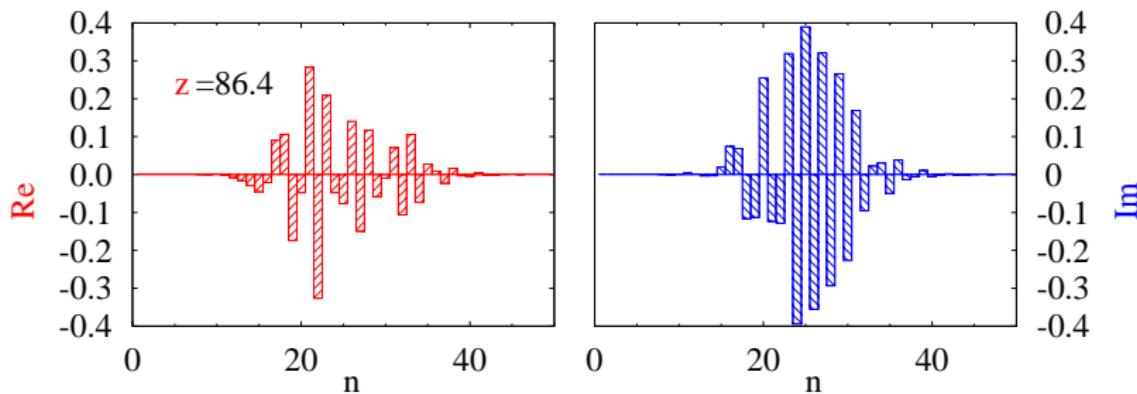
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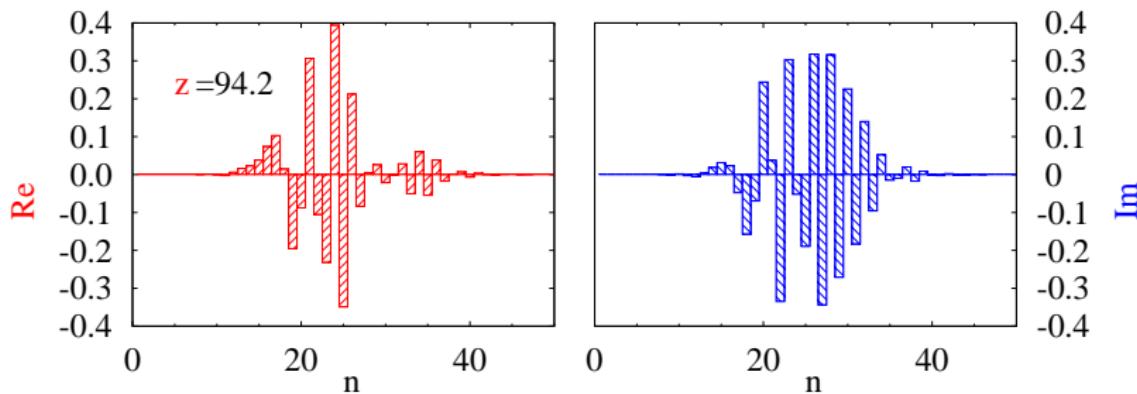
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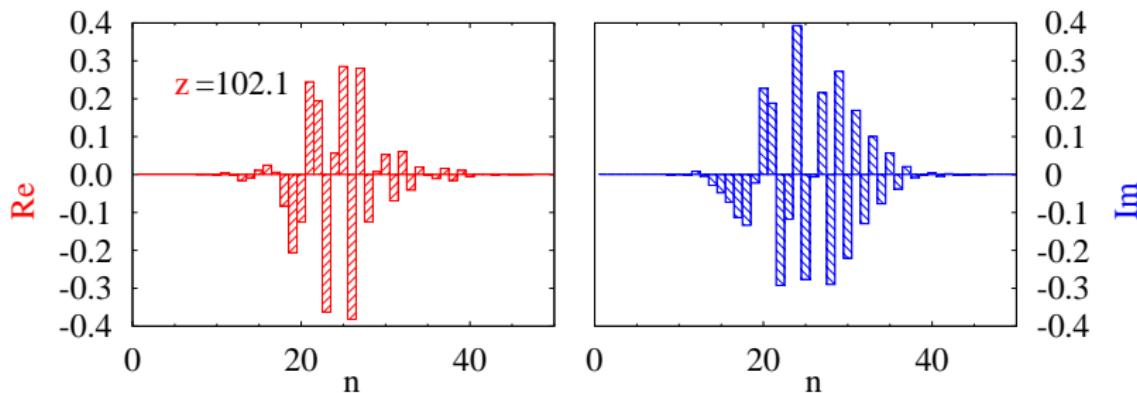
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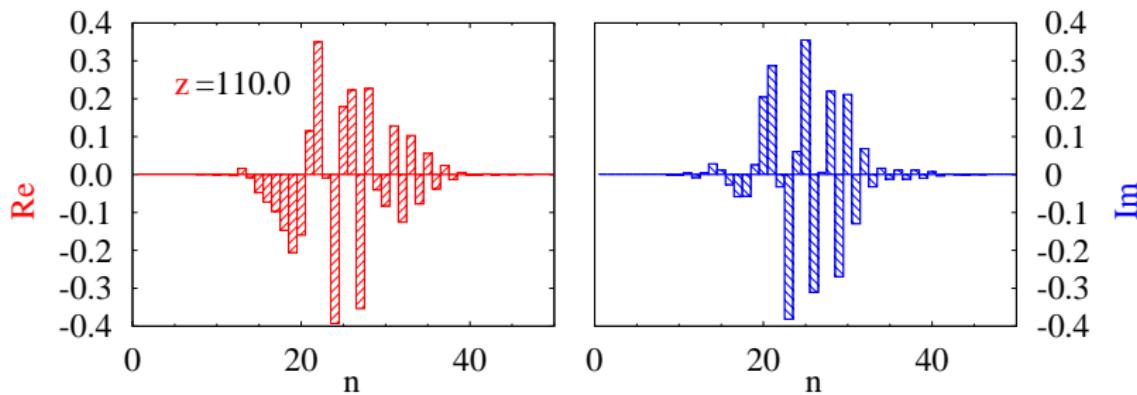
$$\Delta\phi_n = z \left[ (n+1) \ln(n+1) - n \ln n + \ln \left( \frac{z}{e\lambda T^2} \right) \right]$$



# Understanding collapse and revival

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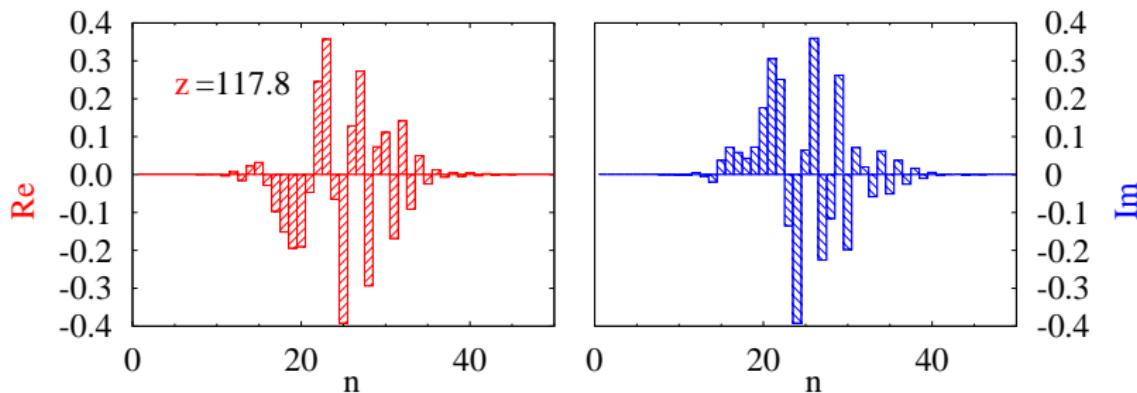
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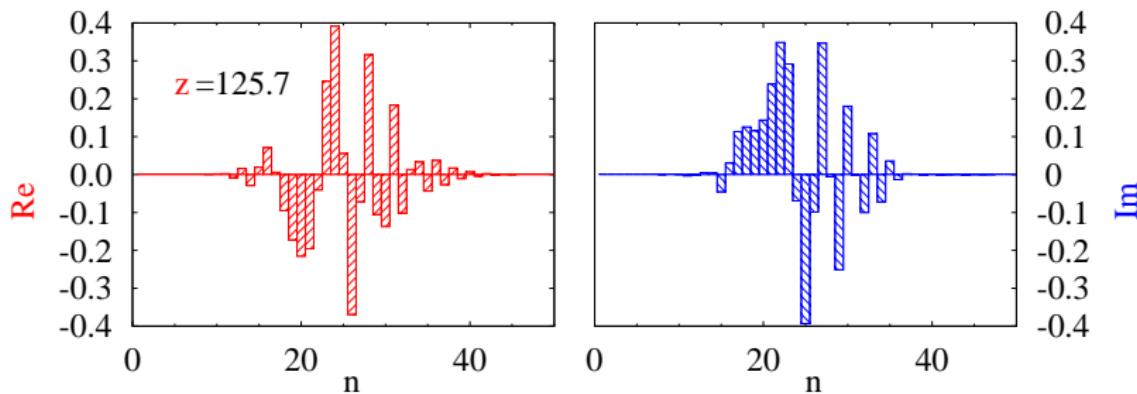
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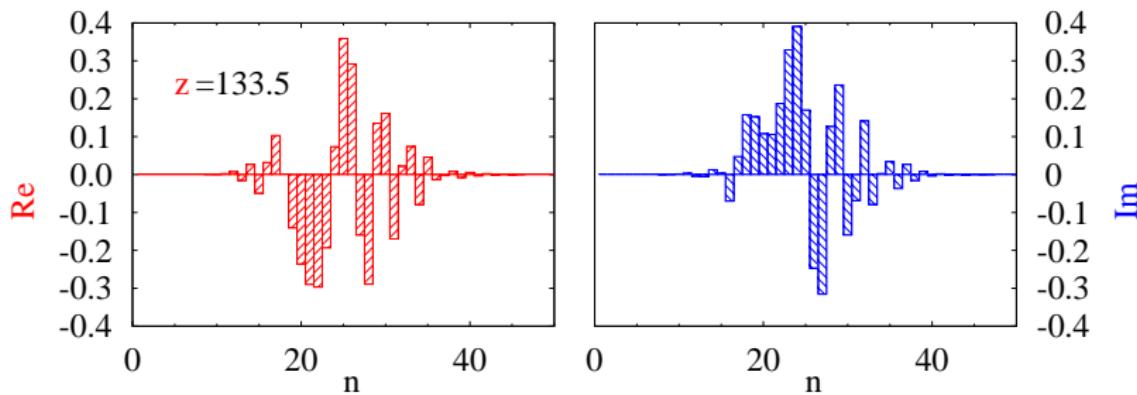
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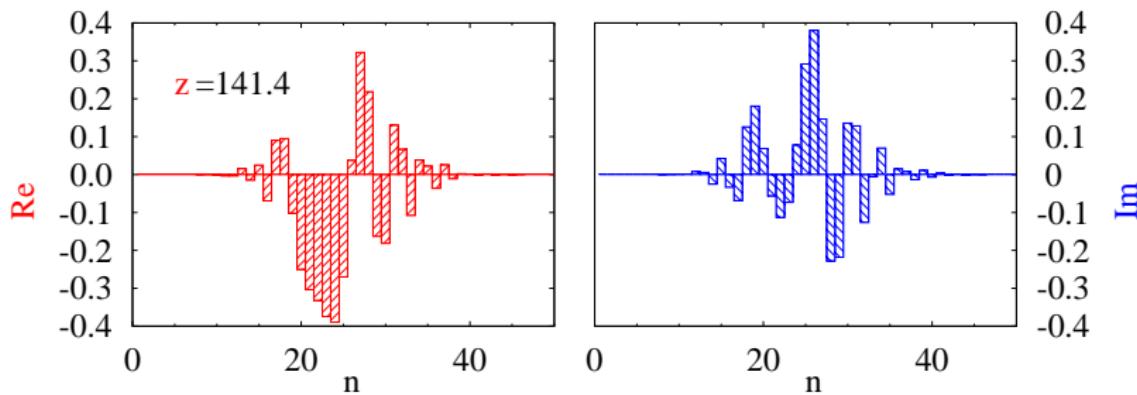
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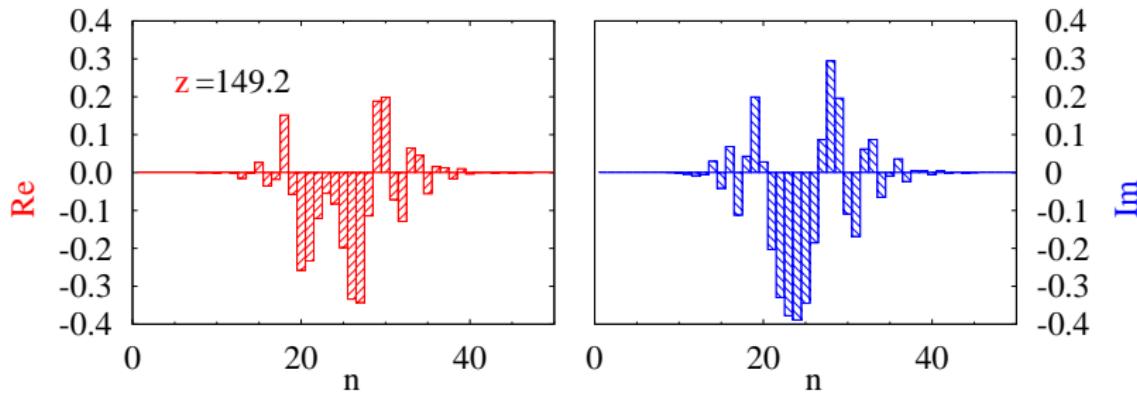
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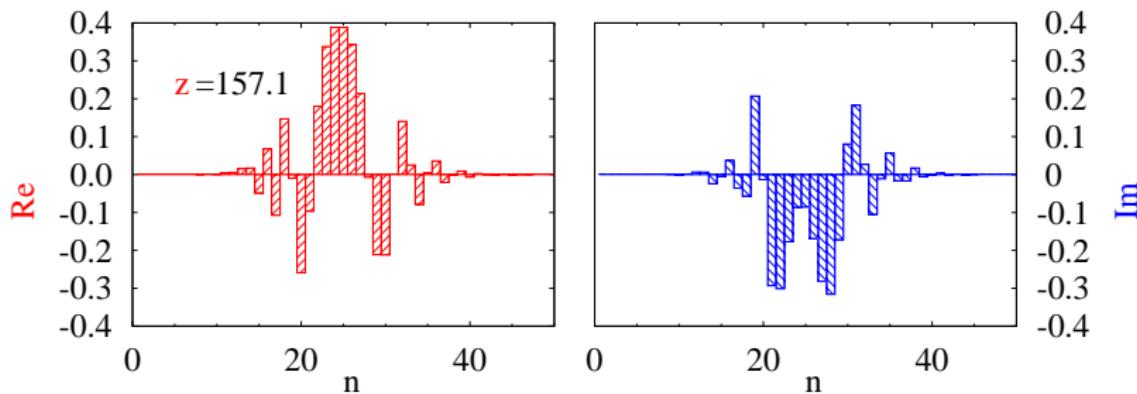
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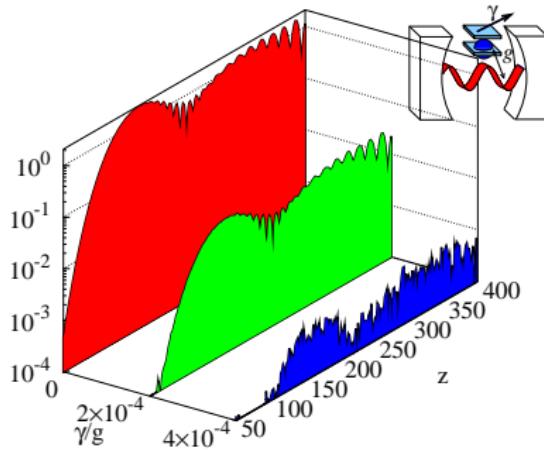
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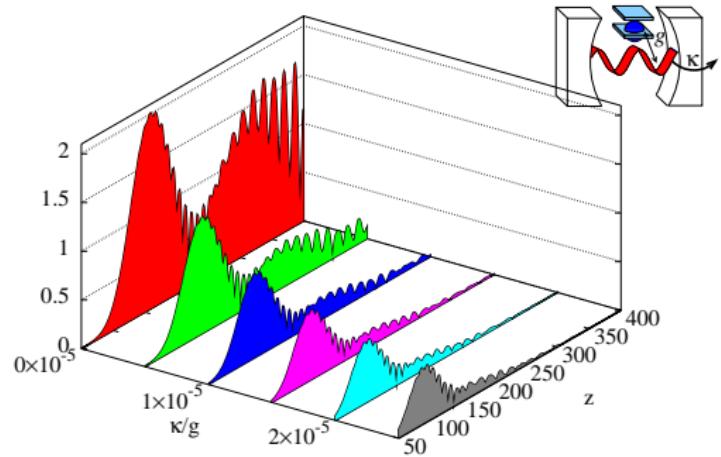
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# Results including decay



$$\gamma_{\max}/g \simeq 2 \times 10^{-4}$$
$$\gamma_{\text{naive}}/g \ll 10 \times 10^{-4}$$

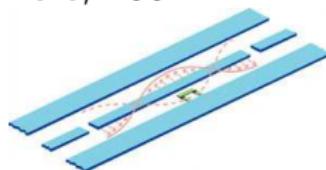


$$\kappa_{\max}/g \simeq 1 \times 10^{-5}$$
$$\kappa_{\text{naive}}/g \ll 100 \times 10^{-5}$$

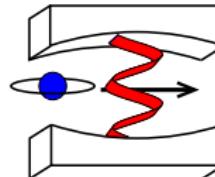
# Possible systems

Requirement:  $\kappa/g \lesssim 1 \times 10^{-5}$ ,  $\gamma/g \lesssim 2 \times 10^{-4}$

System	Source	$\kappa/g$	$\gamma/g$
Quantum dots/Microdisk	CNRS 2005	$2 \times 10^{-1}$	$3 \times 10^{-1}$
Atom/Optical cavity	ETH 2007	$1 \times 10^{-1}$	$3 \times 10^{-1}$
Josephson junction/stripline	Yale, 2004	$2 \times 10^{-2}$	$2 \times 10^{-3}$



Atom/Microwave cavity	ENS, 2004	$7 \times 10^{-3}$	$2 \times 10^{-4}$
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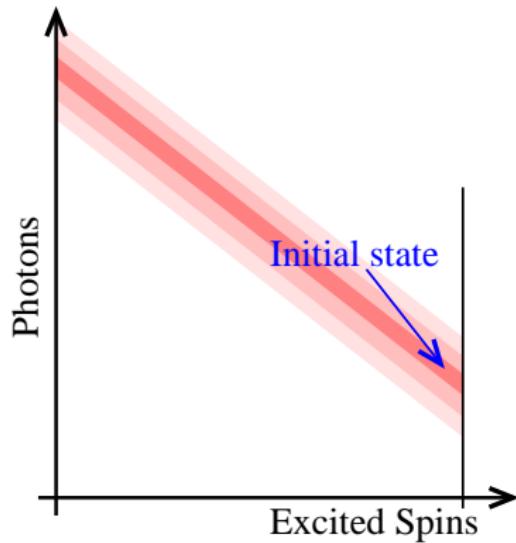
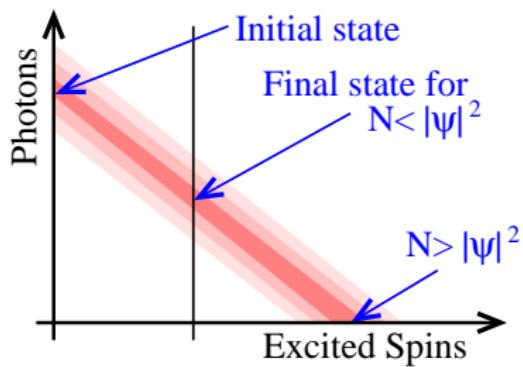


ENS, 2007	$2 \times 10^{-4}$	$2 \times 10^{-4}$
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# Many two-level systems

If many spins:

$$H = \sum_i^N \left( \lambda t s_i^z + g(s_i^+ \psi + s_i^- \psi^\dagger) \right)$$



# Many two-level systems: Converting spins to photons

$$\langle \Psi | \psi | \Psi \rangle = \psi e^{-|\psi|^2} \sum_n \frac{|\psi|^{2n}}{n!} \sqrt{\frac{n+1+N}{n+1}} e^{i(\phi_{n+2}-\phi_{n+1})}$$

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If adiabatic  $\phi$  dominated by logarithm,  $t \rightarrow \pm\infty$

$$H \simeq \begin{pmatrix} \frac{N}{2}\lambda t & g\sqrt{N(n+1)} & 0 & \dots \\ g\sqrt{N(n+1)} & (\frac{N}{2}-1)\lambda t & g\sqrt{2(N-1)(n+2)} & \dots \\ 0 & g\sqrt{2(N-1)(n+2)} & (\frac{N}{2}-2)\lambda t & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

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For  $t \rightarrow -\infty$ :

$$E_- \simeq \frac{N}{2}\lambda t + \frac{g^2 N(n+1)}{\lambda t} - \frac{g^4 N(n+1)}{(\lambda t)^3} [1 + (n+1) - N]$$

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$$\phi_-(n) \simeq \text{Const} + z(n+1) \frac{N}{2} \ln \left[ \frac{z}{\lambda T^2} [1 + (n+1) - N] \right]$$

# Many two-level systems: Converting spins to photons

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Similar expression from  $t \rightarrow +\infty$ . For both:

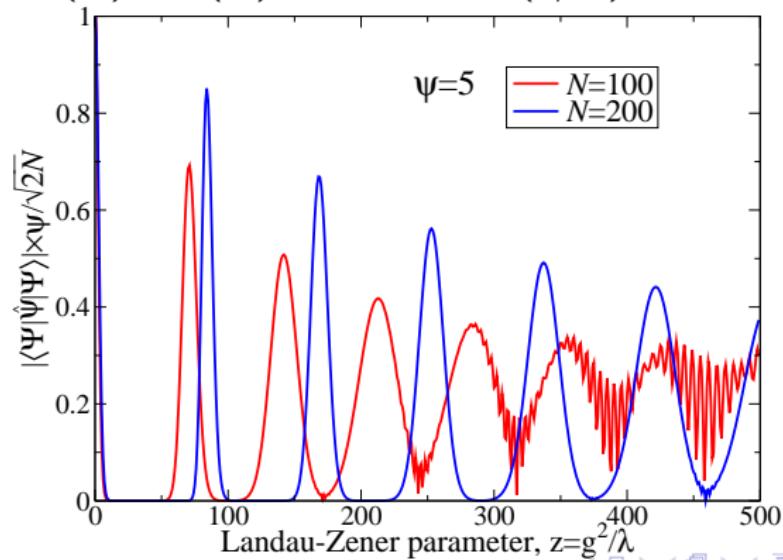
$$\phi = \phi_+ + \phi_- = A(N) + B(N)n + Czn^2 + \mathcal{O}(1/N)$$

# Many two-level systems: Converting spins to photons

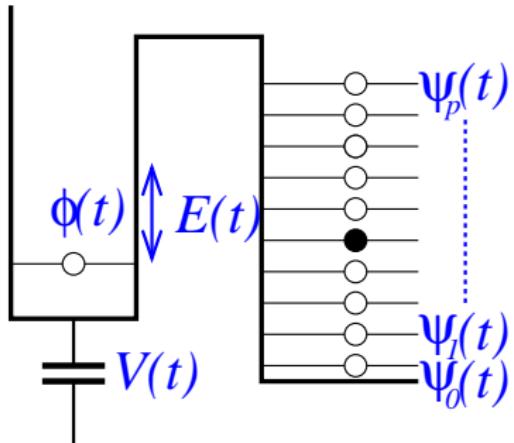
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# Single particle problem



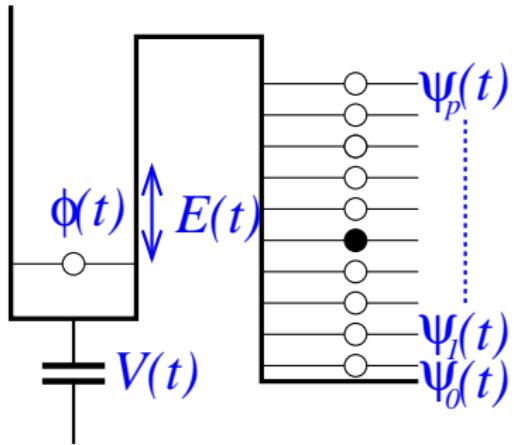
Continuum states:  $v(x, t) = \sum v_n(t) e^{inx}$

Thus, continuum equations:

$$(\partial_t - E(t))v(t) = -\epsilon \int dx v(x, t)\delta(x)$$

$$(\partial_t + i\partial_x)v(x, t) = \epsilon \delta(x)v(t)$$

# Single particle problem



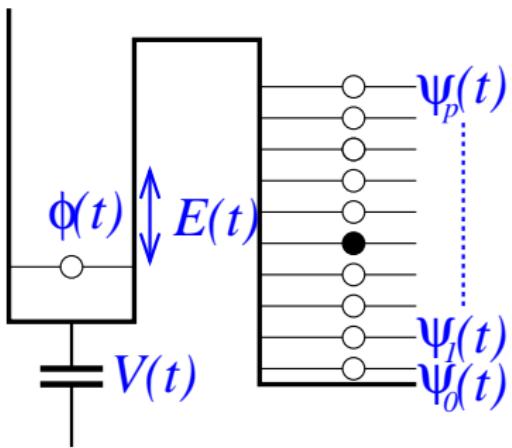
Continuum states:  $\psi(x, t) = \sum \psi_p(t) e^{ipx}$

Thus, continuum equations:

$$[i\partial_t - E(t)]\phi(t) = g \int dx \psi(x, t) \delta(x)$$

$$(i\partial_t + iv\partial_x)\psi(x, t) = g\delta(x)\phi(t)$$

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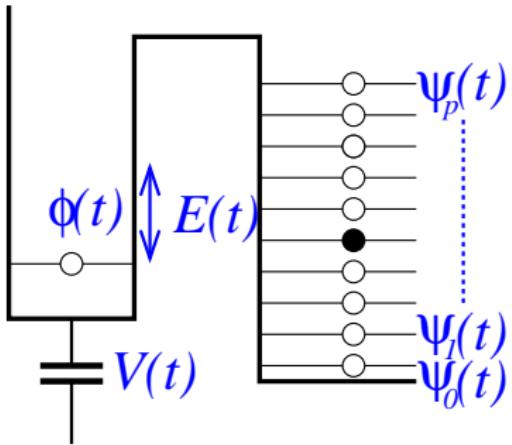
$$(i\partial_t + iv\partial_x)\psi(x, t) = g\delta(x)\phi(t)$$

Transition amplitude:

$$U(\varepsilon, \varepsilon') = 2\pi\delta(\varepsilon - \varepsilon') + T(\varepsilon, \varepsilon')$$

$$T(\varepsilon, \varepsilon') = \iint_{t > t'} dt dt' e^{i(\varepsilon t - \varepsilon' t')} \frac{g^2}{v} \exp \left[ -\frac{g^2}{2v}(t - t') + i \int_{t'}^t E(\tau) d\tau \right]$$

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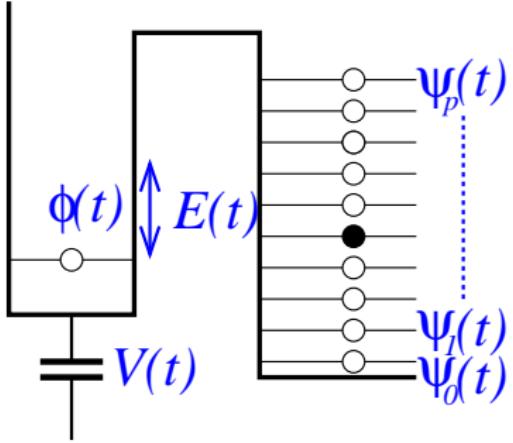
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If:  $E(t) = \lambda t$

# Single particle problem



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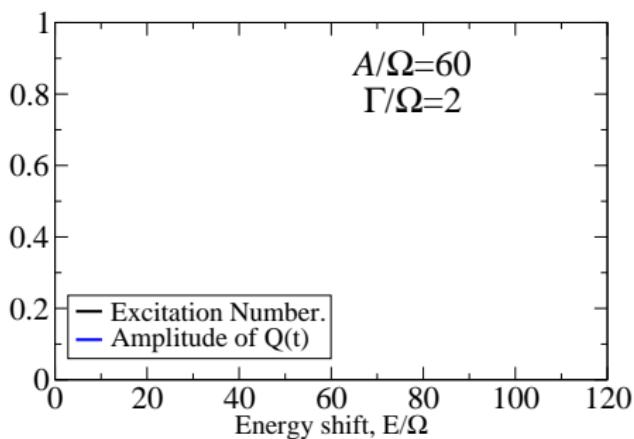
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# Periodic driving: example results

Suppose  $E(t) = E + A\sin(\Omega t)$ : Find  $N^{ex}(t)$  vs  $E$ .

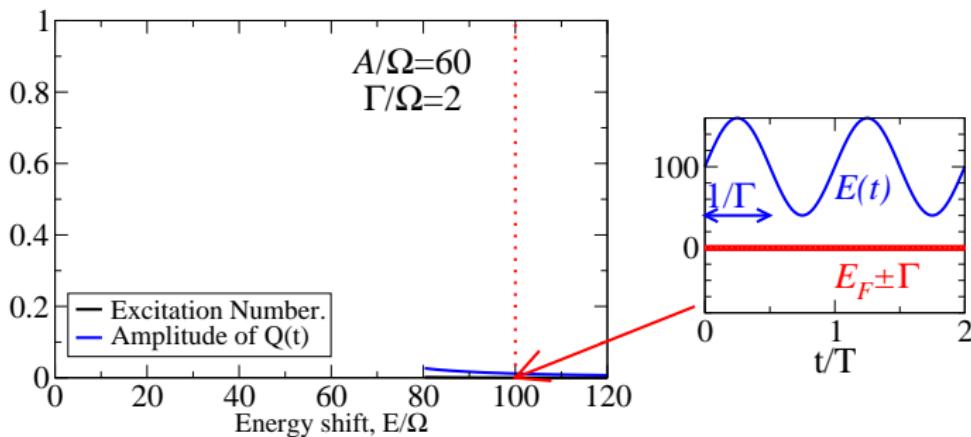
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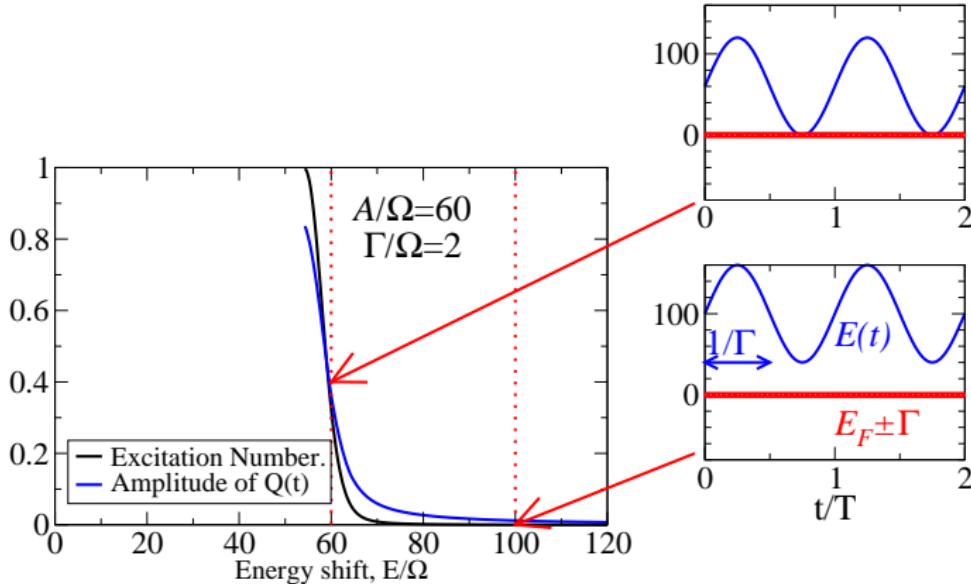
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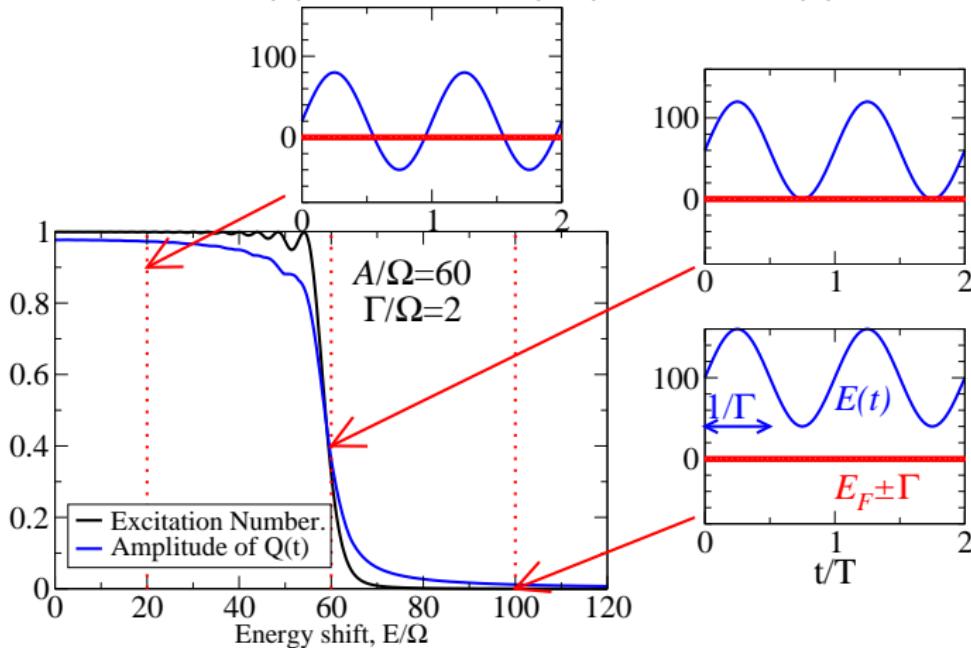
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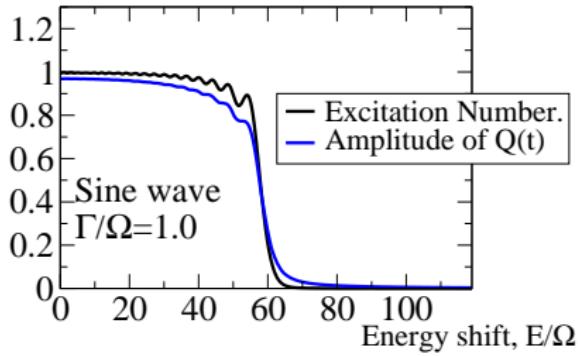
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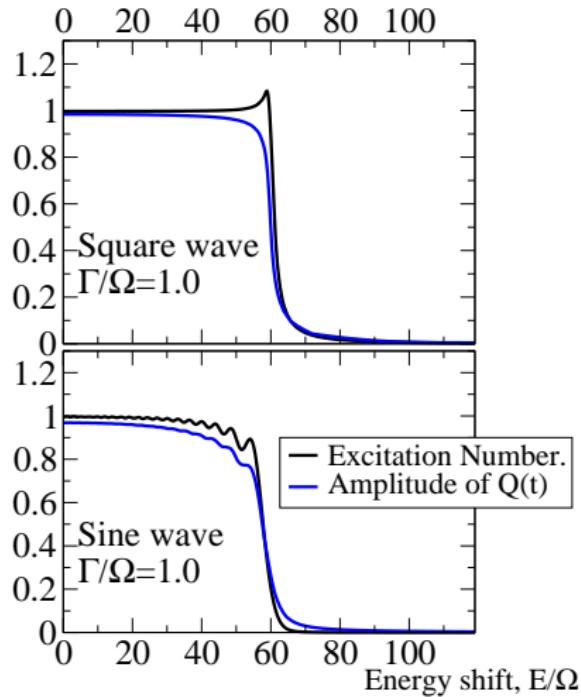


# Comparing $E(t) = E + A \sin(\Omega t)$ and square wave

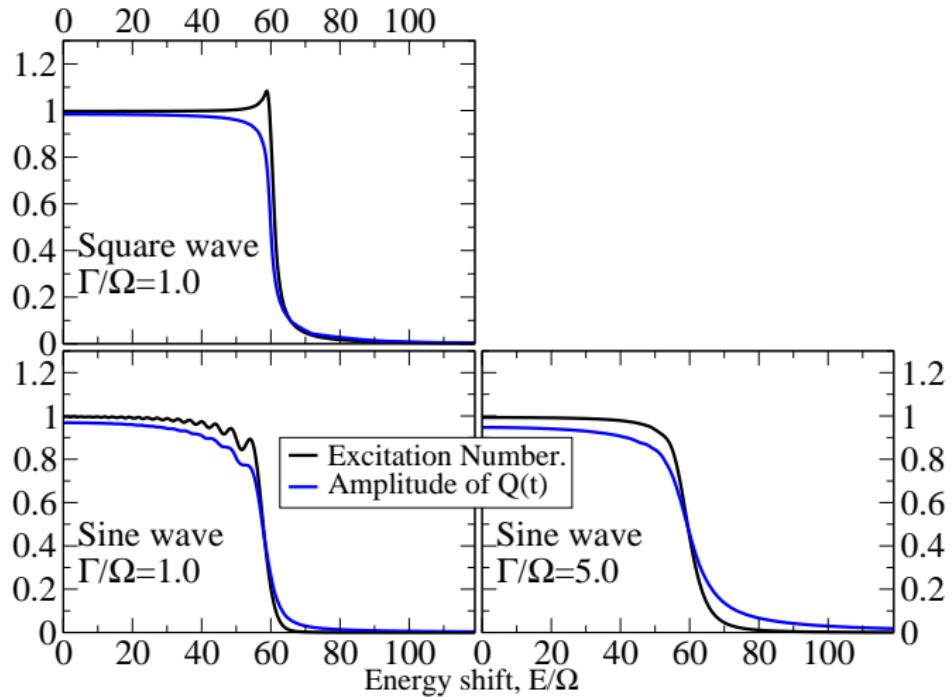
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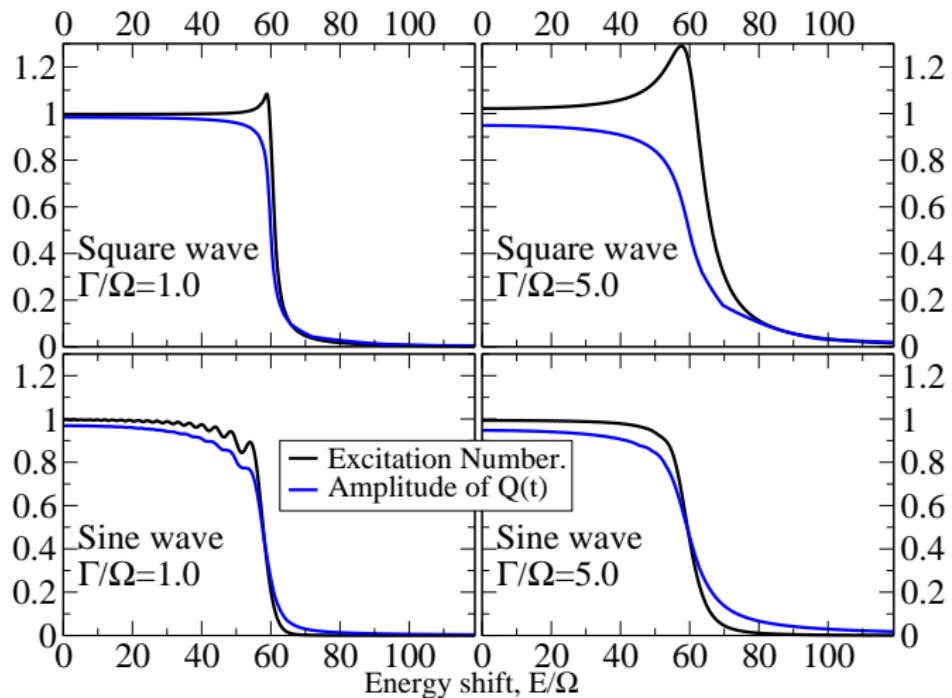
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# General time dependence: measuring noise

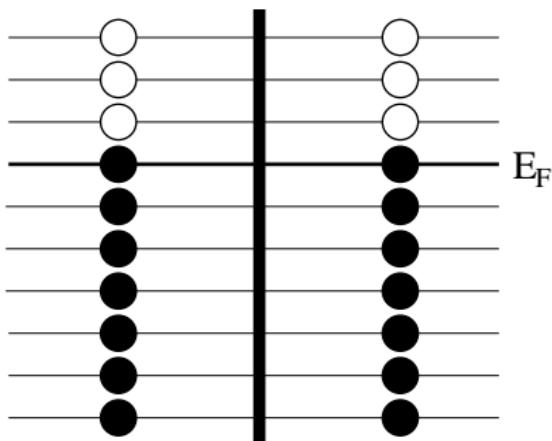
Number of excitations:

$$N^{\text{ex}} = \sum_{\epsilon > \epsilon_F > \epsilon'} |\langle \epsilon | U | \epsilon' \rangle|^2 + \sum_{\epsilon < \epsilon_F < \epsilon'} |\langle \epsilon | U | \epsilon' \rangle|^2$$

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Define:  $N_j = N_j^+ - N_j^-$   
 $g_{j,j} = e(N_j^+ - N_j^-)$   
then

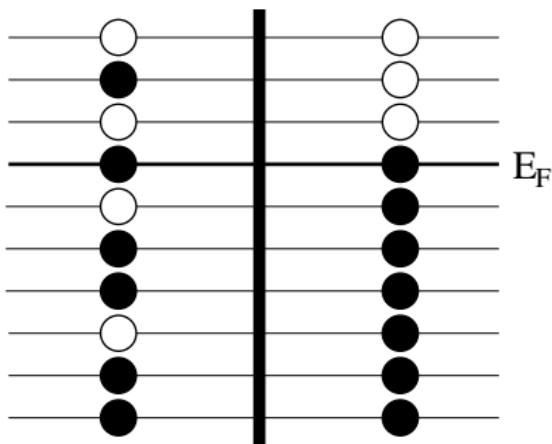
$$\langle g \rangle = T \langle g \rangle$$

$$\langle \Delta g^2 \rangle = T^2 \langle \Delta g^2 \rangle + \delta^2 T (1 - T) \langle N \rangle^2$$

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then

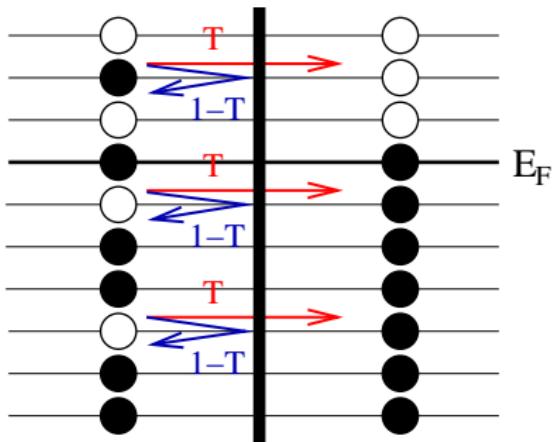
$\langle g \rangle = T(g)$

$(\Delta g^2) = T^2(\Delta g^2) + \beta^2 T(1-T)(N^2)$

# General time dependence: measuring noise

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$$\text{Define: } N_T = N_U - N_D$$
$$g_{TJ} = e(N_U - N_D)$$

then

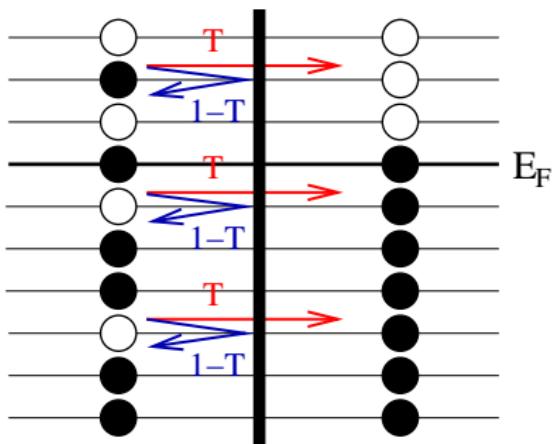
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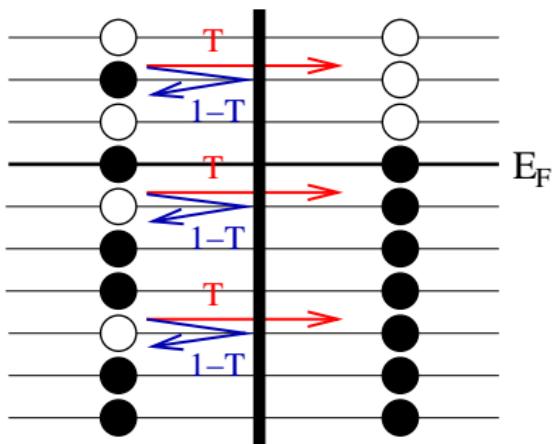


Define:  $N_{r,I}^{\text{ex}} = N_{r,I}^e + N_{r,I}^h$   
 $q_{r,I} = e(N_{r,I}^e - N_{r,I}^h)$

# General time dependence: measuring noise

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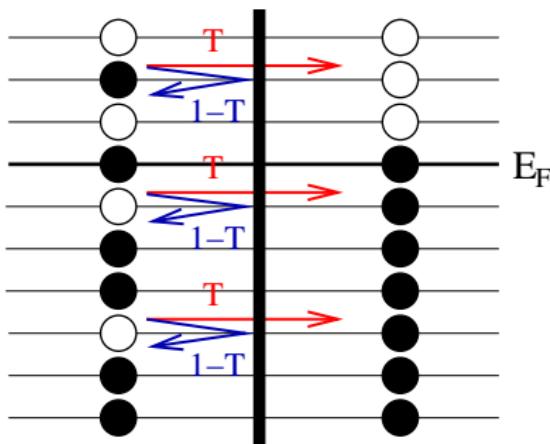
$$\langle q_r \rangle = T \langle q_I \rangle$$

$$(N^e - N^h) + 2T(1-T)(N^h)$$

# General time dependence: measuring noise

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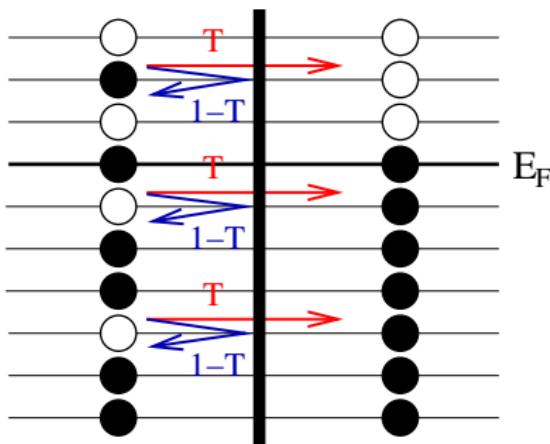
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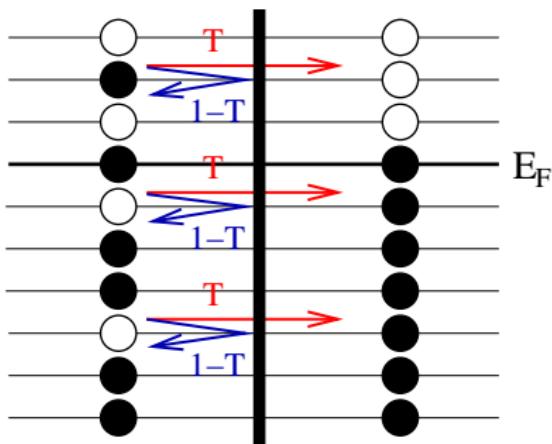
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 $q_{r,I} = e(N_{r,I}^e - N_{r,I}^h)$

Then

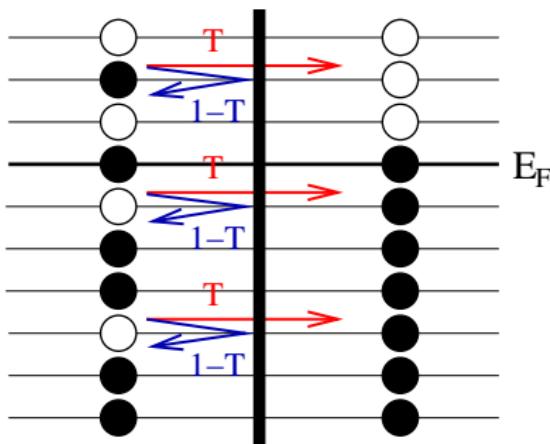
$$\langle q_r \rangle = T \langle q_I \rangle$$

$$\langle \Delta q_r^2 \rangle = T^2 \langle \Delta q_I^2 \rangle$$

# General time dependence: measuring noise

Number of excitations:

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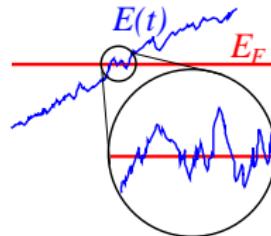
$$\langle q_r \rangle = T \langle q_I \rangle$$

$$\langle \Delta q_r^2 \rangle = T^2 \langle \Delta q_I^2 \rangle + e^2 T (1 - T) \langle N_I^{\text{ex}} \rangle$$

# Noisy driving

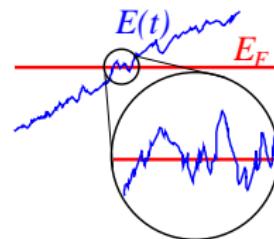
- Suppose  $E(t) = ct + \eta(t)$
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To find  $N^{\infty}$ , need:  $\langle |U(z,z')|^2 \rangle$  thus:



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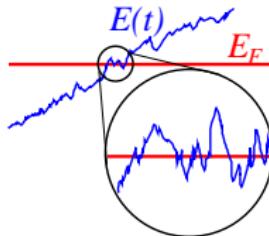


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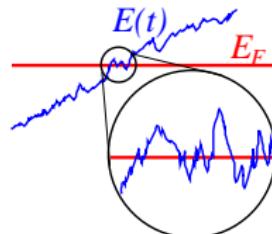
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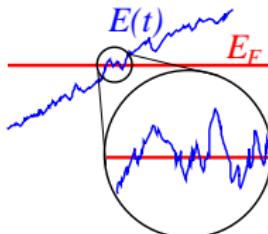
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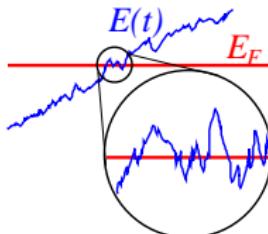
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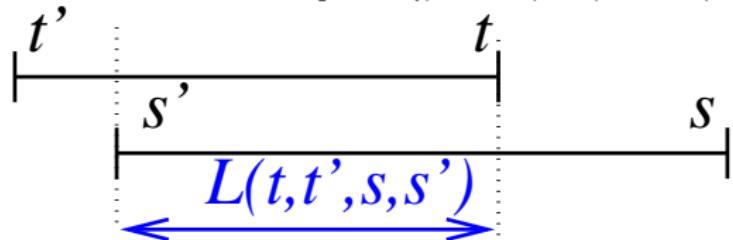
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• Can simplify to  $\Delta = t - s$  and  $\Lambda = t' - t - s - s'$

$$N^{\infty} = \frac{-e^2}{2\pi c} \int_{-\infty}^{\infty} \frac{d\Delta}{(\Delta - i0)^2} \int_0^{\infty} d\Lambda e^{(ie\Delta - \Gamma - \Gamma_2)\Lambda + \Gamma_2 L(\Delta, \Lambda)}$$

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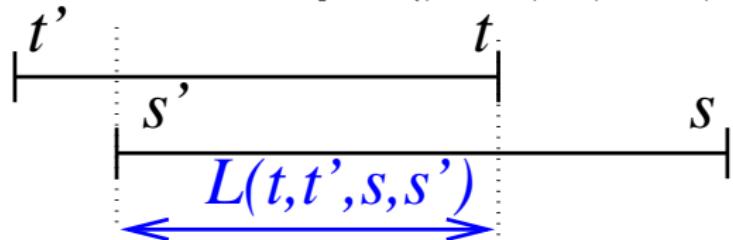


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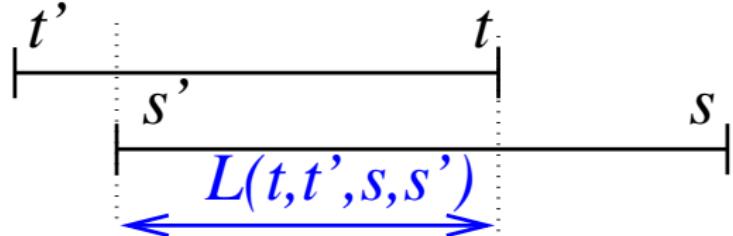


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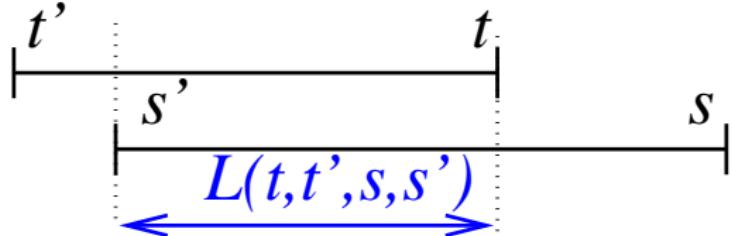


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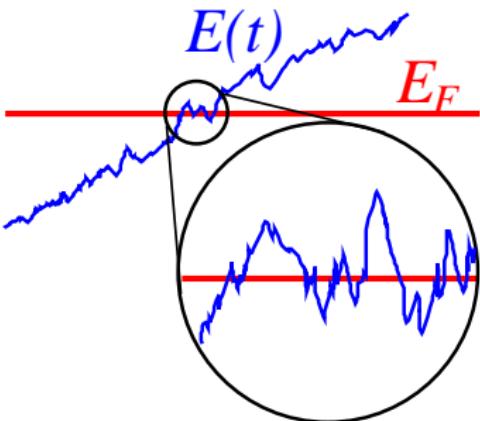
- Integral log divergent: white noise limit
  - infinite no. crossings of Fermi surface
  - Can extract logarithmic contribution

$$N^{\text{ex}} = \begin{cases} 1 & c \gg \Gamma\Gamma_2 \\ \frac{\Gamma^2}{(\Gamma+\Gamma_2)^2} + \frac{2\Gamma\Gamma_2}{\pi c} \ln \frac{\omega_*}{\Gamma+\Gamma_2} & c \ll \Gamma\Gamma_2 \end{cases}$$

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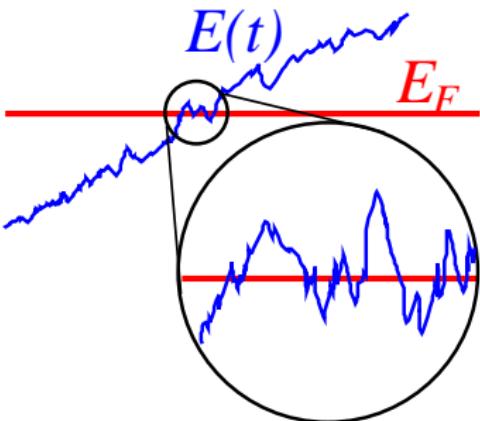
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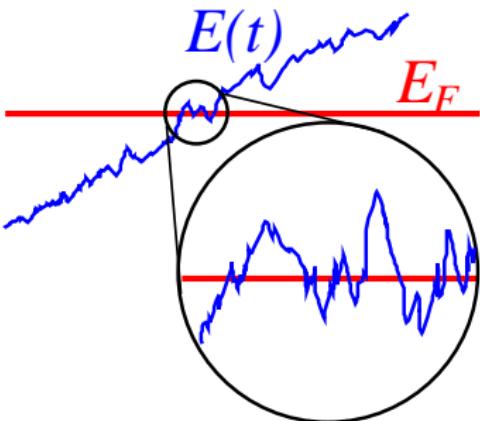
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