

# Nonequilibrium quantum condensates: from microscopic theory to macroscopic phenomenology

**J. M. J. Keeling**

N. G. Berloff, P. B. Littlewood, F. M. Marchetti, M. H. Szymanska.

Quantum Aggregates, April 2009



# Acknowledgements

## People:



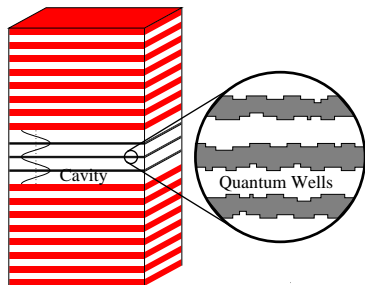
## Funding:

**EPSRC** Engineering and Physical Sciences  
Research Council

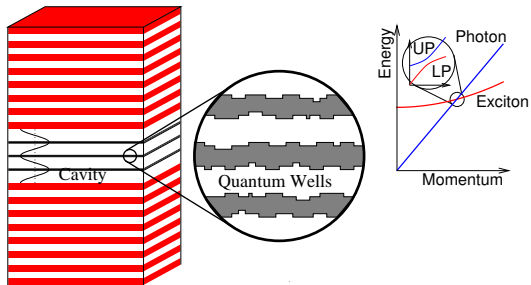


Pembroke College

# Microcavity Polaritons



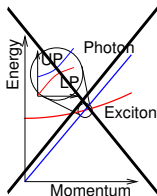
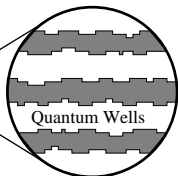
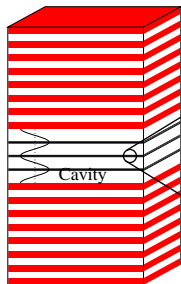
# Microcavity Polaritons



[Pekar, JETP(1958)]

[Hopfield, Phys. Rev.(1958)]

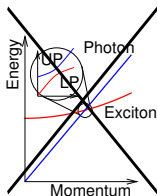
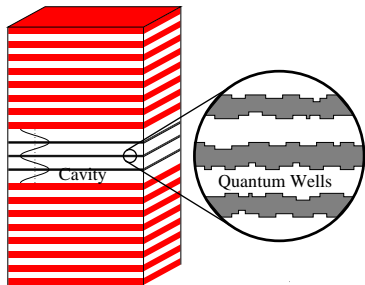
# Microcavity Polaritons



[Pekar, JETP(1958)]

[Hopfield, Phys. Rev.(1958)]

# Microcavity Polaritons



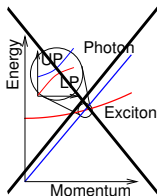
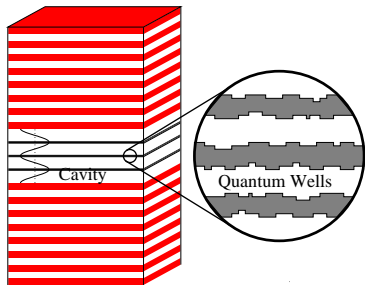
[Pekar, JETP(1958)]

[Hopfield, Phys. Rev.(1958)]

Cavity photons:

$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2/2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$

# Microcavity Polaritons



[Pekar, JETP(1958)]

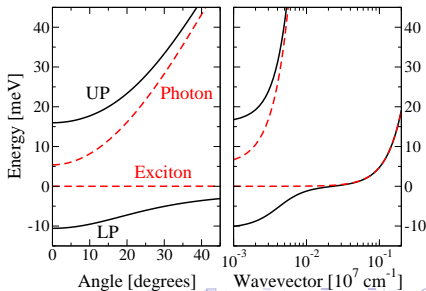
[Hopfield, Phys. Rev.(1958)]

Cavity photons:

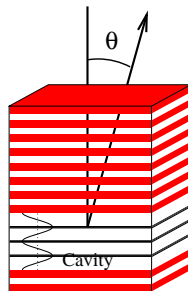
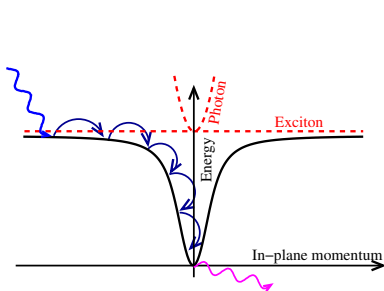
$$\omega_k = \sqrt{\omega_0^2 + c^2 k^2}$$

$$\simeq \omega_0 + \frac{k^2}{2m^*}$$

$$m^* \sim 10^{-4} m_e$$

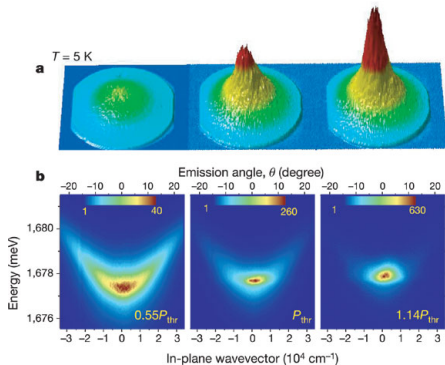


# Non-equilibrium system

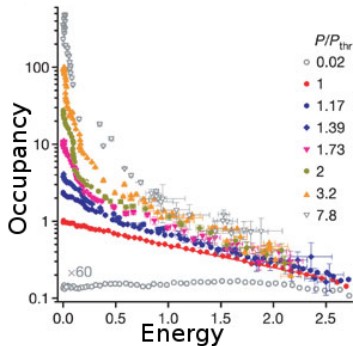




# Polariton experiments: Momentum/Energy distribution

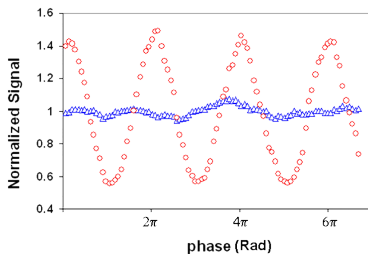
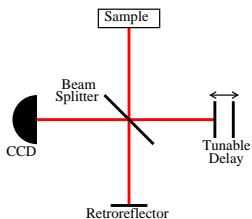


[Kasprzak, et al., Nature, 2006]

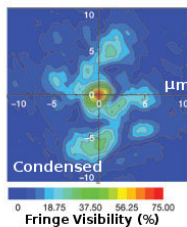
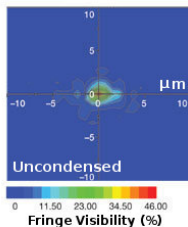
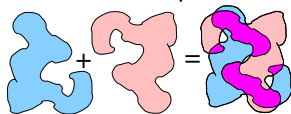


# Polariton experiments: Coherence

Basic idea:



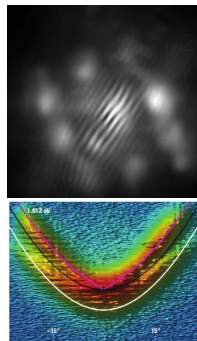
Coherence map:



[Kasprzak, et al., Nature, 2006]

# Other polariton condensation experiments

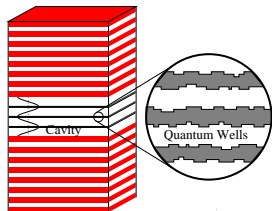
- Old measurements of  $\langle N(t)N(t + \tau) \rangle$   
[Deng *et al* PNAS 100 15318 (2003)]
- Stress traps for polaritons  
[Balili *et al* Science 316 1007 (2007)]
- Temporal coherence and line narrowing  
[Love *et al* Phys. Rev. Lett. 101 067404 (2008)]
- Quantised vortices in disorder potential  
[Lagoudakis *et al* Nature Phys. 4, 706 (2008)]
- Changes to excitation spectrum  
[Utsunomiya *et al* Nature Phys. 4 700 (2008)]
- Soliton propagation  
[Amo *et al* Nature 457 291 (2009)]



# Overview

- 1 Introduction to microcavity polaritons
- 2 Model and review of equilibrium results
  - Disorder-localised exciton model
  - Equilibrium mean field theory
- 3 Non-equilibrium model and mean-field theory
  - Meaning of mean-field condition
- 4 Macroscopic phenomenology
  - Gross Pitaevskii equation in an harmonic trap
- 5 Fluctuations and correlations
  - Fluctuations about mean-field theory

# Excitons in a disorderd Quantum well



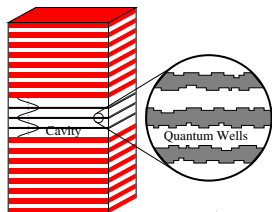
Exciton states in disorder:

$$\left[ -\frac{\nabla_{\mathbf{R}}^2}{2m_{\chi}} + V(\mathbf{R}) \right] \Phi_{\alpha}(\mathbf{R}) = \varepsilon_{\alpha} \Phi_{\alpha}(\mathbf{R})$$

$V(\vec{R})$  smoothed by exciton Bohr radius

[PRL 96 066405 (2006); PRB 76 115326 (2007)]

# Excitons in a disorderd Quantum well



Exciton states in disorder:

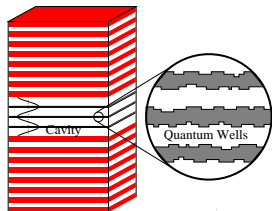
$$\left[ -\frac{\nabla_{\mathbf{R}}^2}{2m_{\chi}} + V(\mathbf{R}) \right] \Phi_{\alpha}(\mathbf{R}) = \varepsilon_{\alpha} \Phi_{\alpha}(\mathbf{R})$$

$V(\vec{R})$  smoothed by exciton Bohr radius

Want: Energies  $\varepsilon_{\alpha}$  Oscillator strengths:  $g_{\alpha,\mathbf{p}} \propto \psi_{1s}(0)\Phi_{\alpha,\mathbf{p}}$

[PRL 96 066405 (2006); PRB 76 115326 (2007)]

# Excitons in a disorderd Quantum well

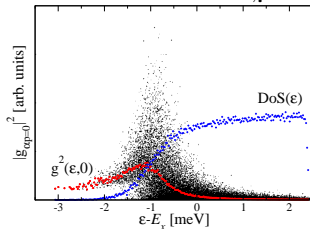


Exciton states in disorder:

$$\left[ -\frac{\nabla_{\mathbf{R}}^2}{2m_{\chi}} + V(\mathbf{R}) \right] \Phi_{\alpha}(\mathbf{R}) = \varepsilon_{\alpha} \Phi_{\alpha}(\mathbf{R})$$

$V(\vec{R})$  smoothed by exciton Bohr radius

Want: Energies  $\varepsilon_{\alpha}$  Oscillator strengths:  $g_{\alpha,p} \propto \psi_{1s}(0)\Phi_{\alpha,p}$

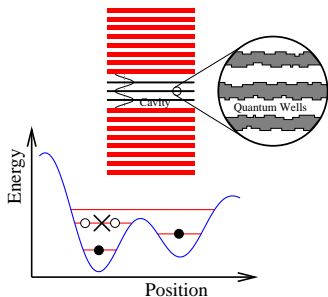


[PRL 96 066405 (2006); PRB 76 115326 (2007)]

# Polariton system model

## Polariton model

- Disorder-localised excitons
- Treat disorder sites as two-level (exciton/no-exciton)
- Propagating (2D) photons
- Exciton-photon coupling  $g$ .

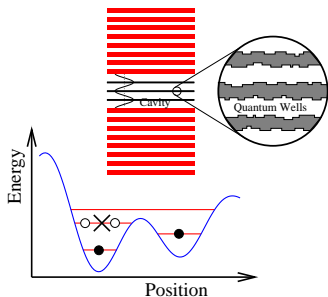




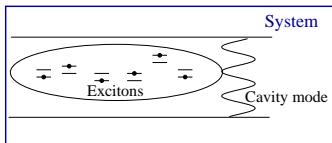
# Polariton system model

## Polariton model

- Disorder-localised excitons
- Treat disorder sites as two-level (exciton/no-exciton)
- Propagating (2D) photons
- Exciton-photon coupling  $g$ .

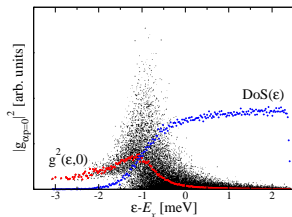


$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} \left[ \epsilon_{\alpha} \left( b_{\alpha}^{\dagger} b_{\alpha} - a_{\alpha}^{\dagger} a_{\alpha} \right) + \frac{1}{\sqrt{\text{Area}}} g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} b_{\alpha}^{\dagger} a_{\alpha} + \text{H.c.} \right]$$



# Equilibrium: Mean-field theory

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} \left[ \epsilon_{\alpha} \left( b_{\alpha}^{\dagger} b_{\alpha} - a_{\alpha}^{\dagger} a_{\alpha} \right) + \frac{g_{\alpha, \mathbf{k}}}{\sqrt{A}} \psi_{\mathbf{k}} b_{\alpha}^{\dagger} a_{\alpha} + \text{H.c.} \right]$$



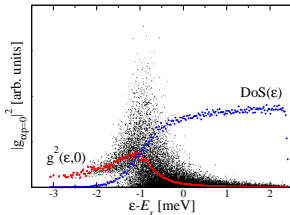
# Equilibrium: Mean-field theory

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} \left[ \epsilon_{\alpha} \left( b_{\alpha}^{\dagger} b_{\alpha} - a_{\alpha}^{\dagger} a_{\alpha} \right) + \frac{g_{\alpha, \mathbf{k}}}{\sqrt{A}} \psi_{\mathbf{k}} b_{\alpha}^{\dagger} a_{\alpha} + \text{H.c.} \right]$$

Mean-field theory:

Self-consistent polarisation and field

$$\left[ i\partial_t + \omega_0 - \frac{\nabla^2}{2m} \right] \psi = \frac{1}{\sqrt{A}} \sum_{\alpha} g_{\alpha} P_{\alpha}$$



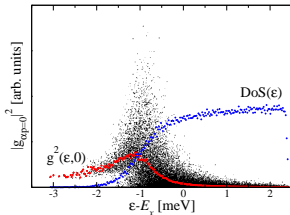
# Equilibrium: Mean-field theory

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} \left[ \epsilon_{\alpha} \left( b_{\alpha}^{\dagger} b_{\alpha} - a_{\alpha}^{\dagger} a_{\alpha} \right) + \frac{g_{\alpha, \mathbf{k}}}{\sqrt{A}} \psi_{\mathbf{k}} b_{\alpha}^{\dagger} a_{\alpha} + \text{H.c.} \right]$$

Mean-field theory:

Self-consistent polarisation and field

$$\left[ i\partial_t + \omega_0 - \frac{\nabla^2}{2m} \right] \psi = \frac{1}{\sqrt{A}} \sum_{\alpha} g_{\alpha} \langle a_{\alpha}^{\dagger} b_{\alpha} \rangle$$



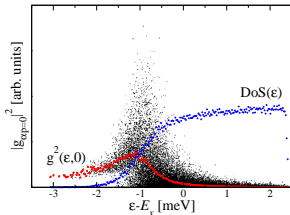
# Equilibrium: Mean-field theory

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \sum_{\alpha} \left[ \epsilon_{\alpha} \left( b_{\alpha}^\dagger b_{\alpha} - a_{\alpha}^\dagger a_{\alpha} \right) + \frac{g_{\alpha, \mathbf{k}}}{\sqrt{A}} \psi_{\mathbf{k}} b_{\alpha}^\dagger a_{\alpha} + \text{H.c.} \right]$$

Mean-field theory:

Self-consistent polarisation and field

$$\left[ i\partial_t + \omega_0 - \frac{\nabla^2}{2m} \right] \psi = \frac{1}{\sqrt{A}} \sum_{\alpha} g_{\alpha} \frac{g_{\alpha} \psi}{2E_{\alpha}} \tanh(\beta E_{\alpha})$$
$$E_{\alpha}^2 = \left( \frac{\tilde{\epsilon}_{\alpha}}{2} \right)^2 + g_{\alpha}^2 \psi^2, \quad \tilde{\epsilon}_{\alpha} = \epsilon_{\alpha} - \mu$$



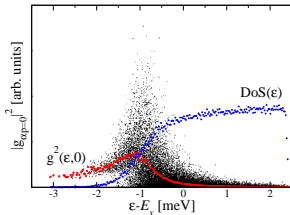
# Equilibrium: Mean-field theory

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \sum_{\alpha} \left[ \epsilon_{\alpha} \left( b_{\alpha}^\dagger b_{\alpha} - a_{\alpha}^\dagger a_{\alpha} \right) + \frac{g_{\alpha, \mathbf{k}}}{\sqrt{A}} \psi_{\mathbf{k}} b_{\alpha}^\dagger a_{\alpha} + \text{H.c.} \right]$$

Mean-field theory:

Self-consistent polarisation and field

$$\left[ -\mu + \omega_0 - \frac{\nabla^2}{2m} \right] \psi = \frac{1}{\sqrt{A}} \sum_{\alpha} g_{\alpha} \frac{g_{\alpha} \psi}{2E_{\alpha}} \tanh(\beta E_{\alpha})$$
$$E_{\alpha}^2 = \left( \frac{\tilde{\epsilon}_{\alpha}}{2} \right)^2 + g_{\alpha}^2 \psi^2, \quad \tilde{\epsilon}_{\alpha} = \epsilon_{\alpha} - \mu$$



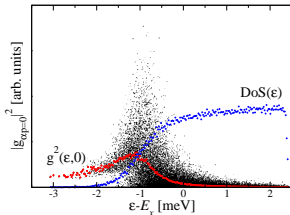
# Equilibrium: Mean-field theory

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} \left[ \epsilon_{\alpha} \left( b_{\alpha}^{\dagger} b_{\alpha} - a_{\alpha}^{\dagger} a_{\alpha} \right) + \frac{g_{\alpha, \mathbf{k}}}{\sqrt{A}} \psi_{\mathbf{k}} b_{\alpha}^{\dagger} a_{\alpha} + \text{H.c.} \right]$$

Mean-field theory:

Self-consistent polarisation and field

$$\left[ -\mu + \omega_0 - \frac{\nabla^2}{2m} \right] \psi = \frac{1}{\sqrt{A}} \sum_{\alpha} g_{\alpha} \frac{g_{\alpha} \psi}{2E_{\alpha}} \tanh(\beta E_{\alpha})$$
$$E_{\alpha}^2 = \left( \frac{\tilde{\epsilon}_{\alpha}}{2} \right)^2 + g_{\alpha}^2 \psi^2, \quad \tilde{\epsilon}_{\alpha} = \epsilon_{\alpha} - \mu$$



Density

$$\rho = |\psi|^2 + \frac{1}{A} \sum_{\alpha} \left[ \frac{1}{2} - \frac{\tilde{\epsilon}_{\alpha}}{4E_{\alpha}} \tanh(\beta E_{\alpha}) \right]$$

# Equilibrium: Mean-field theory

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \sum_{\alpha} \left[ \epsilon_{\alpha} \left( b_{\alpha}^\dagger b_{\alpha} - a_{\alpha}^\dagger a_{\alpha} \right) + \frac{g_{\alpha, \mathbf{k}}}{\sqrt{A}} \psi_{\mathbf{k}} b_{\alpha}^\dagger a_{\alpha} + \text{H.c.} \right]$$

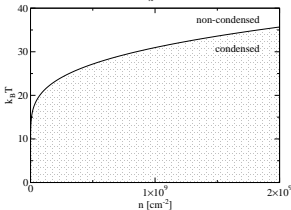
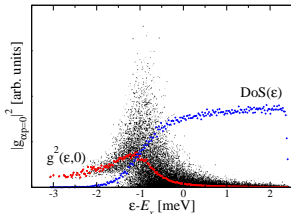
Mean-field theory:

Self-consistent polarisation and field

$$\left[ -\mu + \omega_0 - \frac{\nabla^2}{2m} \right] \psi = \frac{1}{\sqrt{A}} \sum_{\alpha} g_{\alpha} \frac{g_{\alpha} \psi}{2E_{\alpha}} \tanh(\beta E_{\alpha})$$
$$E_{\alpha}^2 = \left( \frac{\tilde{\epsilon}_{\alpha}}{2} \right)^2 + g_{\alpha}^2 \psi^2, \quad \tilde{\epsilon}_{\alpha} = \epsilon_{\alpha} - \mu$$

Density

$$\rho = |\psi|^2 + \frac{1}{A} \sum_{\alpha} \left[ \frac{1}{2} - \frac{\tilde{\epsilon}_{\alpha}}{4E_{\alpha}} \tanh(\beta E_{\alpha}) \right]$$





# Equilibrium: Mean-field theory

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^{\dagger} \psi_{\mathbf{k}} + \sum_{\alpha} \left[ \epsilon_{\alpha} \left( b_{\alpha}^{\dagger} b_{\alpha} - a_{\alpha}^{\dagger} a_{\alpha} \right) + \frac{g_{\alpha, \mathbf{k}}}{\sqrt{A}} \psi_{\mathbf{k}} b_{\alpha}^{\dagger} a_{\alpha} + \text{H.c.} \right]$$

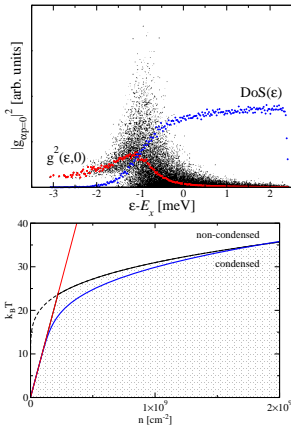
Mean-field theory:

Self-consistent polarisation and field

$$\left[ -\mu + \omega_0 - \frac{\nabla^2}{2m} \right] \psi = \frac{1}{\sqrt{A}} \sum_{\alpha} g_{\alpha} \frac{g_{\alpha} \psi}{2E_{\alpha}} \tanh(\beta E_{\alpha})$$
$$E_{\alpha}^2 = \left( \frac{\tilde{\epsilon}_{\alpha}}{2} \right)^2 + g_{\alpha}^2 \psi^2, \quad \tilde{\epsilon}_{\alpha} = \epsilon_{\alpha} - \mu$$

Density

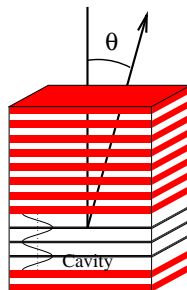
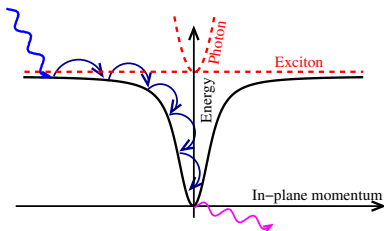
$$\rho = |\psi|^2 + \frac{1}{A} \sum_{\alpha} \left[ \frac{1}{2} - \frac{\tilde{\epsilon}_{\alpha}}{4E_{\alpha}} \tanh(\beta E_{\alpha}) \right]$$



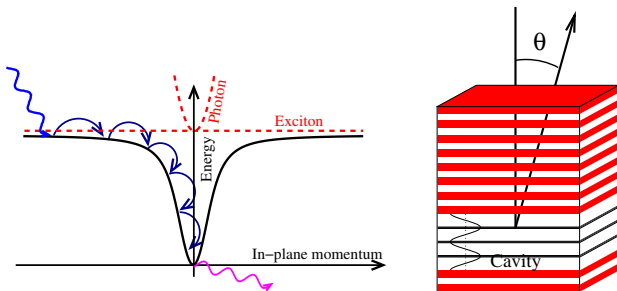
# Overview

- 1 Introduction to microcavity polaritons
- 2 Model and review of equilibrium results
  - Disorder-localised exciton model
  - Equilibrium mean field theory
- 3 Non-equilibrium model and mean-field theory**
  - Meaning of mean-field condition**
- 4 Macroscopic phenomenology
  - Gross Pitaevskii equation in an harmonic trap
- 5 Fluctuations and correlations
  - Fluctuations about mean-field theory

# Non-equilibrium system



# Non-equilibrium system

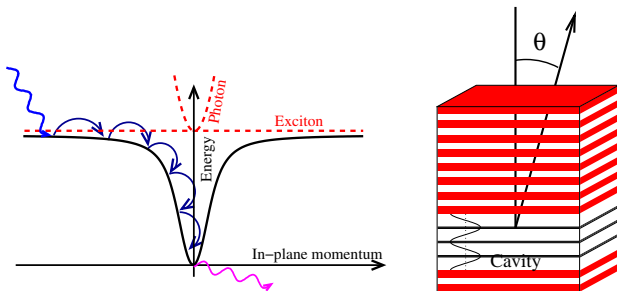


	Lifetime	Thermalisation
Atoms	10s	10ms
Excitons <sup>a</sup>	50ns	0.2ns
<b>Polaritons</b>	<b>5ps</b>	<b>0.5ps</b>
Magnons <sup>b</sup>	1 $\mu$ s(??)	100ns(?)

<sup>a</sup>Coupled quantum wells. [Hammack et al PRB 76 193308 (2007)]

<sup>b</sup>Yttrium Iron Garnett. [Demokritov et al, Nature 443 430 (2007)]

# Non-equilibrium system

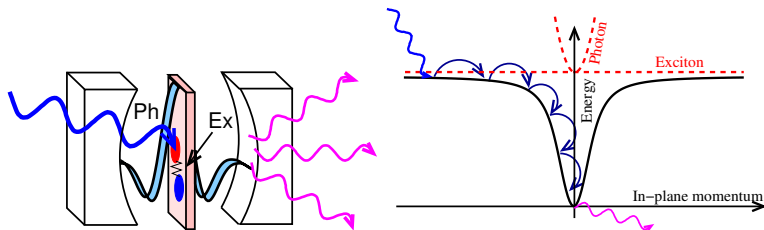


	Lifetime	Thermalisation	Linewidth	Temperature	
Atoms	10s	10ms	$2.5 \times 10^{-13}$ meV	$10^{-8}$ K	$10^{-9}$ meV
Excitons <sup>a</sup>	50ns	0.2ns	$5 \times 10^{-5}$ meV	1K	0.1meV
<b>Polaritons</b>	<b>5ps</b>	<b>0.5ps</b>	<b>0.5meV</b>	<b>20K</b>	<b>2meV</b>
Magnons <sup>b</sup>	$1\mu\text{s}(??)$	100ns(?)	$2.5 \times 10^{-6}$ meV	300K	30meV

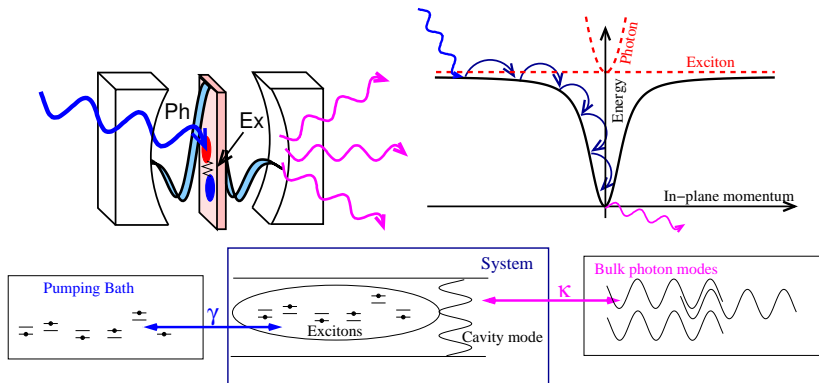
<sup>a</sup>Coupled quantum wells. [Hammack et al PRB 76 193308 (2007)]

<sup>b</sup>Yttrium Iron Garnett. [Demokritov et al, Nature 443 430 (2007)]

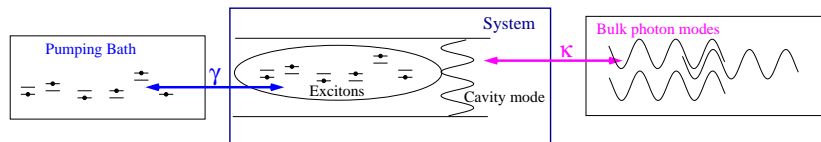
# Non-equilibrium: flux and baths



# Non-equilibrium: flux and baths



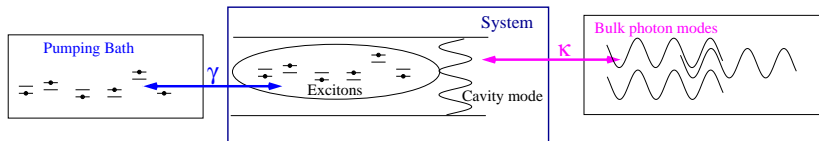
# Non-equilibrium model: baths



$$H = H_{\text{sys}} + H_{\text{sys,bath}} + H_{\text{bath}}$$



# Non-equilibrium model: baths

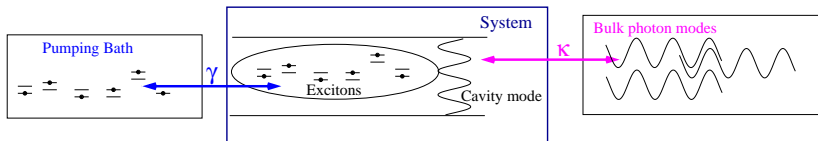


$$H = H_{\text{sys}} + H_{\text{sys,bath}} + H_{\text{bath}}$$

Schematically: pump  $\gamma$ , decay  $\kappa$

$$H_{\text{sys,bath}} \simeq \sum_{\mathbf{p}, \vec{k}} \sqrt{\kappa} \psi_{\mathbf{k}} \Psi_{\mathbf{p}}^{\dagger} + \sum_{\alpha, \beta} \sqrt{\gamma} \left( a_{\alpha}^{\dagger} A_{\beta} + b_{\alpha}^{\dagger} B_{\beta} \right) + \text{H.c.}$$

# Non-equilibrium model: baths



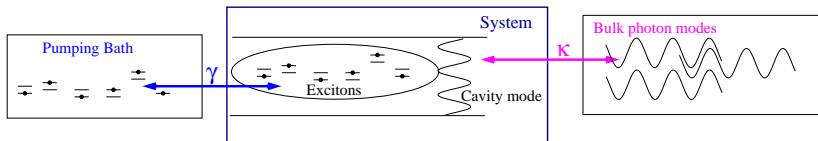
$$H = H_{\text{sys}} + H_{\text{sys,bath}} + H_{\text{bath}}$$

Schematically: pump  $\gamma$ , decay  $\kappa$

$$H_{\text{sys,bath}} \simeq \sum_{\mathbf{p}, \vec{k}} \sqrt{\kappa} \psi_{\mathbf{k}} \Psi_{\mathbf{p}}^{\dagger} + \sum_{\alpha, \beta} \sqrt{\gamma} \left( a_{\alpha}^{\dagger} A_{\beta} + b_{\alpha}^{\dagger} B_{\beta} \right) + \text{H.c.}$$

Bath correlations,  $\langle \Psi^{\dagger} \Psi \rangle$ ,  $\langle A^{\dagger} A \rangle$ ,  $\langle B^{\dagger} B \rangle$  fixed:

# Non-equilibrium model: baths

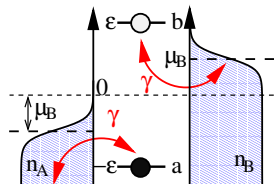


$$H = H_{\text{sys}} + H_{\text{sys,bath}} + H_{\text{bath}}$$

Schematically: pump  $\gamma$ , decay  $\kappa$

$$H_{\text{sys,bath}} \simeq \sum_{\mathbf{p}, \vec{k}} \sqrt{\kappa} \psi_{\mathbf{k}} \Psi_{\mathbf{p}}^{\dagger} + \sum_{\alpha, \beta} \sqrt{\gamma} \left( a_{\alpha}^{\dagger} A_{\beta} + b_{\alpha}^{\dagger} B_{\beta} \right) + \text{H.c.}$$

Bath correlations,  $\langle \Psi^{\dagger} \Psi \rangle$ ,  $\langle A^{\dagger} A \rangle$ ,  $\langle B^{\dagger} B \rangle$  fixed:  
 $\Psi$  bath is empty. Pumping bath thermal,  $\mu_B, T$ :



# Non-equilibrium theory; mean-field

Look for mean-field solution,  $\psi(\mathbf{r}, t) = \psi_0 e^{-i\mu_s t}$ . Gap equation:

$$(\mu_s - \omega_0 + i\kappa) \psi_0 = P = \chi(\psi_0, \mu_s) \psi_0$$

# Non-equilibrium theory; mean-field

Look for mean-field solution,  $\psi(\mathbf{r}, t) = \psi_0 e^{-i\mu_s t}$ . Gap equation:

$$(\mu_s - \omega_0 + i\kappa) \psi_0 = P = \chi(\psi_0, \mu_s) \psi_0$$

Susceptibility:

$$\chi(\psi_0, \mu_s) = -g^2 \gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

# Non-equilibrium theory; mean-field

Look for mean-field solution,  $\psi(\mathbf{r}, t) = \psi_0 e^{-i\mu_s t}$ . Gap equation:

$$(\mu_s - \omega_0 + i\kappa) \psi_0 = P = \chi(\psi_0, \mu_s) \psi_0$$

Susceptibility:  $E_\alpha^2 = \epsilon_\alpha^2 + g^2 |\psi_0|^2$

$$\chi(\psi_0, \mu_s) = -g^2 \gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

# Non-equilibrium theory; mean-field

Look for mean-field solution,  $\psi(\mathbf{r}, t) = \psi_0 e^{-i\mu_s t}$ . Gap equation:

$$(\mu_s - \omega_0 + i\kappa) \psi_0 = P = \chi(\psi_0, \mu_s) \psi_0$$

Susceptibility:  $E_\alpha^2 = \epsilon_\alpha^2 + g^2 |\psi_0|^2$ ,  $F_{a,b}(\nu) = F[\nu \mp (\frac{1}{2}\mu_s - \mu_B)]$

$$\chi(\psi_0, \mu_s) = -g^2 \gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

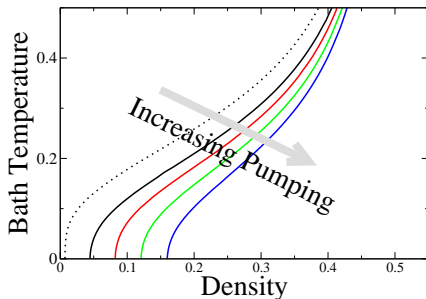
# Non-equilibrium theory; mean-field

Look for mean-field solution,  $\psi(\mathbf{r}, t) = \psi_0 e^{-i\mu_s t}$ . Gap equation:

$$(\mu_s - \omega_0 + i\kappa) \psi_0 = P = \chi(\psi_0, \mu_s) \psi_0$$

Susceptibility:  $E_\alpha^2 = \epsilon_\alpha^2 + g^2 |\psi_0|^2$ ,  $F_{a,b}(\nu) = F[\nu \mp (\frac{1}{2}\mu_s - \mu_B)]$

$$\chi(\psi_0, \mu_s) = -g^2 \gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$





# Limits of gap equation

Gap equation:

$$\mu_s - \omega_0 + i\kappa = -g^2\gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

# Limits of gap equation

Gap equation:

$$\mu_s - \omega_0 + i\kappa = -g^2\gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

- Laser Limit Imaginary part: Gain vs Loss.

# Limits of gap equation

Gap equation:

$$\mu_s - \omega_0 + i\kappa = -g^2\gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

- Laser Limit Imaginary part: Gain vs Loss. If  $T \gg \gamma$

$$\kappa = -g^2\gamma \sum_{\text{excitons}} \frac{[F_b(E_\alpha) - F_a(E_\alpha)]}{4E_\alpha^2 + 4\gamma^2}$$

# Limits of gap equation

Gap equation:

$$\mu_s - \omega_0 + i\kappa = -g^2\gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

- Laser Limit Imaginary part: Gain vs Loss. If  $T \gg \gamma$

$$\kappa = -g^2\gamma \sum_{\text{excitons}} \frac{[F_b(E_\alpha) - F_a(E_\alpha)]}{4E_\alpha^2 + 4\gamma^2} = \frac{g^2}{2\gamma} \times \text{Inversion}$$

# Limits of gap equation

Gap equation:

$$\mu_s - \omega_0 + i\kappa = -g^2\gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

- Laser Limit Imaginary part: Gain vs Loss. If  $T \gg \gamma$

$$\kappa = -g^2\gamma \sum_{\text{excitons}} \frac{[F_b(E_\alpha) - F_a(E_\alpha)]}{4E_\alpha^2 + 4\gamma^2} = \frac{g^2}{2\gamma} \times \text{Inversion}$$

- Equilibrium limit: finite T set by pumping, need  $\kappa \ll \gamma$ .

# Limits of gap equation

Gap equation:

$$\mu_S - \omega_0 + i\kappa = -g^2\gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_S)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

- Laser Limit Imaginary part: Gain vs Loss. If  $T \gg \gamma$

$$\kappa = -g^2\gamma \sum_{\text{excitons}} \frac{[F_b(E_\alpha) - F_a(E_\alpha)]}{4E_\alpha^2 + 4\gamma^2} = \frac{g^2}{2\gamma} \times \text{Inversion}$$

- Equilibrium limit: finite T set by pumping, need  $\kappa \ll \gamma$ .  
Require:  $F_a = F_b$  so  $\mu_S = 2\mu_B$

# Limits of gap equation

Gap equation:

$$\mu_S - \omega_0 + i\kappa = -g^2\gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_S)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

- Laser Limit Imaginary part: Gain vs Loss. If  $T \gg \gamma$

$$\kappa = -g^2\gamma \sum_{\text{excitons}} \frac{[F_b(E_\alpha) - F_a(E_\alpha)]}{4E_\alpha^2 + 4\gamma^2} = \frac{g^2}{2\gamma} \times \text{Inversion}$$

- Equilibrium limit: finite  $T$  set by pumping, need  $\kappa \ll \gamma$ .  
Require:  $F_a = F_b$  so  $\mu_S = 2\mu_B$

# Limits of gap equation

Gap equation:

$$\mu_s - \omega_0 + i\kappa = -g^2\gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(\nu + E_\alpha)^2 + \gamma^2]}$$

- Laser Limit Imaginary part: Gain vs Loss. If  $T \gg \gamma$

$$\kappa = -g^2\gamma \sum_{\text{excitons}} \frac{[F_b(E_\alpha) - F_a(E_\alpha)]}{4E_\alpha^2 + 4\gamma^2} = \frac{g^2}{2\gamma} \times \text{Inversion}$$

- Equilibrium limit: finite T set by pumping, need  $\kappa \ll \gamma$ .  
Require:  $F_a = F_b$  so  $\mu_s = 2\mu_B$

$$\omega_0 - \mu_s = \frac{g^2}{2E} \tanh\left(\frac{\beta E}{2}\right)$$



# Limits of gap equation

Gap equation:

$$(\mu_s - \omega_0 + i\kappa)\psi = \chi(\psi, \mu_s)\psi$$

- Local density limit:

# Limits of gap equation

Gap equation:

$$(\mu_s - \omega_0 + i\kappa)\psi = \chi(\psi, \mu_s)\psi$$

- Local density limit: Gross-Pitaevskii equation

$$\left( i\hbar\partial_t + i\kappa - \left[ V(r) - \frac{\hbar^2\nabla^2}{2m} \right] \right) \psi(r) = \chi(\psi(r, t))\psi(r, t)$$

Nonlinear, complex susceptibility  $\chi(\psi(r, t))$

# Limits of gap equation

Gap equation:

$$(\mu_s - \omega_0 + i\kappa)\psi = \chi(\psi, \mu_s)\psi$$

- Local density limit: Gross-Pitaevskii equation

$$\left( i\hbar\partial_t + i\kappa - \left[ V(r) - \frac{\hbar^2\nabla^2}{2m} \right] \right) \psi(r) = \chi(\psi(r, t))\psi(r, t)$$

Nonlinear, complex susceptibility  $\chi[E(\psi(r, t))]$ ,  $E_\alpha^2 = \epsilon_\alpha^2 + g^2|\psi_0|^2$

$$i\hbar\partial_t\psi|_{\text{nl}} = U|\psi|^2\psi$$

# Limits of gap equation

Gap equation:

$$(\mu_s - \omega_0 + i\kappa)\psi = \chi(\psi, \mu_s)\psi$$

- Local density limit: Gross-Pitaevskii equation

$$\left( i\hbar\partial_t + i\kappa - \left[ V(r) - \frac{\hbar^2\nabla^2}{2m} \right] \right) \psi(r) = \chi(\psi(r, t))\psi(r, t)$$

Nonlinear, complex susceptibility  $\chi[E(\psi(r, t))]$ ,  $E_\alpha^2 = \epsilon_\alpha^2 + g^2|\psi_0|^2$

$$i\hbar\partial_t\psi|_{\text{nl}} = U|\psi|^2\psi$$

$$i\hbar\partial_t\psi|_{\text{loss}} = -i\kappa\psi \quad i\hbar\partial_t\psi|_{\text{gain}} = i\gamma_{\text{eff}}(\mu_B)\psi$$

# Limits of gap equation

Gap equation:

$$(\mu_s - \omega_0 + i\kappa)\psi = \chi(\psi, \mu_s)\psi$$

- Local density limit: Gross-Pitaevskii equation

$$\left( i\hbar\partial_t + i\kappa - \left[ V(r) - \frac{\hbar^2\nabla^2}{2m} \right] \right) \psi(r) = \chi(\psi(r, t))\psi(r, t)$$

Nonlinear, complex susceptibility  $\chi[E(\psi(r, t))]$ ,  $E_\alpha^2 = \epsilon_\alpha^2 + g^2|\psi_0|^2$

$$i\hbar\partial_t\psi|_{\text{nl}} = U|\psi|^2\psi$$

$$i\hbar\partial_t\psi|_{\text{loss}} = -i\kappa\psi \quad i\hbar\partial_t\psi|_{\text{gain}} = i\gamma_{\text{eff}}(\mu_B)\psi - i\Gamma|\psi|^2\psi$$

# Limits of gap equation

Gap equation:

$$(\mu_s - \omega_0 + i\kappa) \psi = \chi(\psi, \mu_s) \psi$$

- Local density limit: Gross-Pitaevskii equation

$$\left( i\hbar\partial_t + i\kappa - \left[ V(r) - \frac{\hbar^2\nabla^2}{2m} \right] \right) \psi(r) = \chi(\psi(r, t))\psi(r, t)$$

Nonlinear, complex susceptibility  $\chi[E(\psi(r, t))]$ ,  $E_\alpha^2 = \epsilon_\alpha^2 + g^2|\psi_0|^2$

$$i\hbar\partial_t\psi|_{\text{nl}} = U|\psi|^2\psi$$

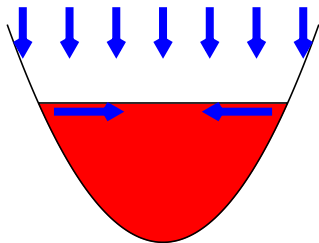
$$i\hbar\partial_t\psi|_{\text{loss}} = -i\kappa\psi$$

$$i\hbar\partial_t\psi|_{\text{gain}} = i\gamma_{\text{eff}}(\mu_B)\psi - i\Gamma|\psi|^2\psi$$

$$i\partial_t\psi = \left[ -\frac{\hbar^2\nabla^2}{2m} + V(r) + U|\psi|^2 + i(\gamma_{\text{eff}}(\mu_B) - \kappa - \Gamma|\psi|^2) \right] \psi$$

# Gross-Pitaevskii equation: Harmonic trap

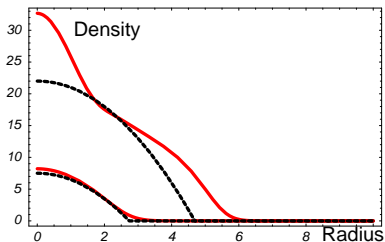
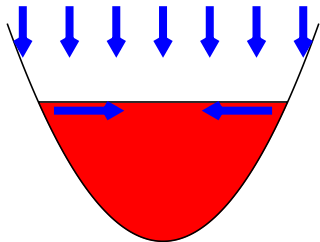
$$i\hbar\partial_t\psi = \left[ -\frac{\hbar^2\nabla^2}{2m} + \frac{m\omega^2}{2}r^2 + U|\psi|^2 + i(\gamma_{\text{eff}} - \kappa - \Gamma|\psi|^2) \right] \psi$$



[Keeling & Berloff, PRL, '08]

# Gross-Pitaevskii equation: Harmonic trap

$$i\hbar\partial_t\psi = \left[ -\frac{\hbar^2\nabla^2}{2m} + \frac{m\omega^2}{2}r^2 + U|\psi|^2 + i(\gamma_{\text{eff}} - \kappa - \Gamma|\psi|^2) \right] \psi$$

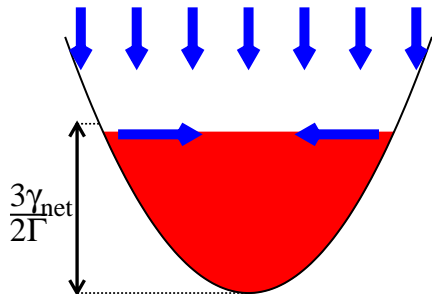


[Keeling & Berloff, PRL, '08]



# Stability of Thomas-Fermi solution

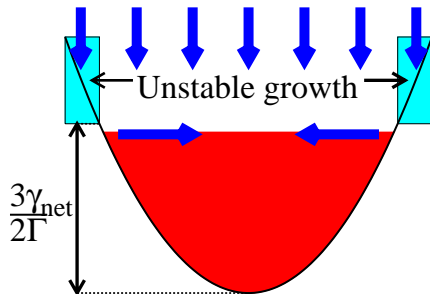
$$\frac{1}{2}\partial_t\rho + \nabla \cdot (\rho\mathbf{v}) = \frac{1}{\hbar}(\gamma_{\text{net}} - \Gamma\rho)\rho$$



# Stability of Thomas-Fermi solution

High  $m$  modes:  $\delta\rho_{n,m} \simeq e^{im\theta} r^m \dots$

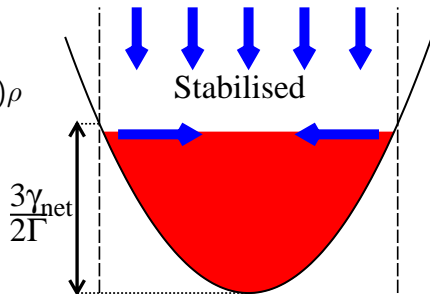
$$\frac{1}{2}\partial_t\rho + \nabla \cdot (\rho\mathbf{v}) = \frac{1}{\hbar}(\gamma_{\text{net}} - \Gamma\rho)\rho$$



# Stability of Thomas-Fermi solution

High  $m$  modes:  $\delta\rho_{n,m} \simeq e^{im\theta} r^m \dots$

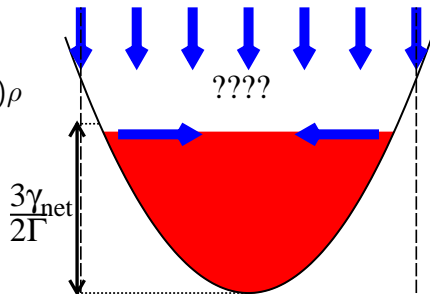
$$\frac{1}{2}\partial_t\rho + \nabla \cdot (\rho\mathbf{v}) = \frac{1}{\hbar}(\gamma_{\text{net}}\Theta(R-r) - \Gamma\rho)\rho$$



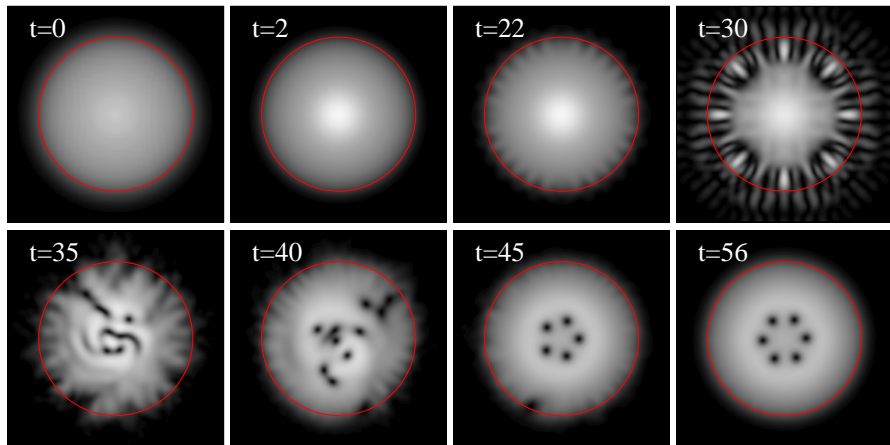
# Stability of Thomas-Fermi solution

High  $m$  modes:  $\delta\rho_{n,m} \simeq e^{im\theta} r^m \dots$

$$\frac{1}{2}\partial_t\rho + \nabla \cdot (\rho\mathbf{v}) = \frac{1}{\hbar}(\gamma_{\text{net}}\Theta(R-r) - \Gamma\rho)\rho$$



# Time evolution:



[Keeling & Berloff, PRL, '08]

# Why vortices

Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{\mu}{U} = \frac{\gamma_{\text{net}}}{\Gamma} \Theta(R-r)$$

[Keeling & Berloff, PRL, '08]

# Why vortices

Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

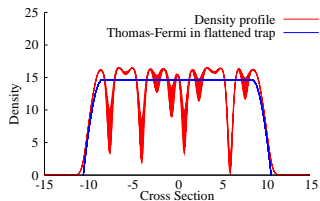
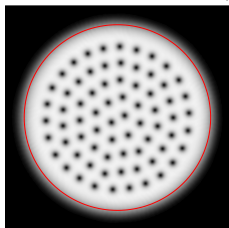
$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{\mu}{U} = \frac{\gamma_{\text{net}}}{\Gamma} \Theta(R - r)$$

[Keeling & Berloff, PRL, '08]

# Why vortices

Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{\mu}{U} = \frac{\gamma_{\text{net}}}{\Gamma} \Theta(R - r)$$



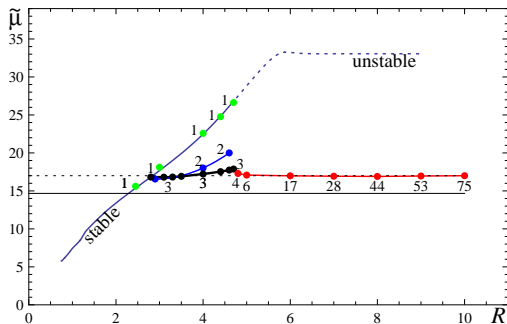
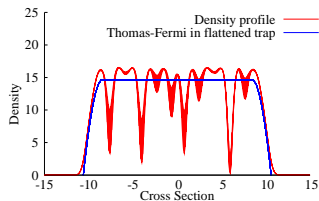
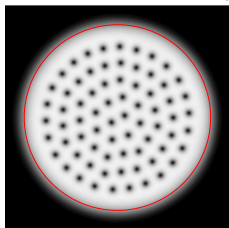
[Keeling & Berloff, PRL, '08]



# Why vortices

Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{\mu}{U} = \frac{\gamma_{\text{net}}}{\Gamma} \Theta(R - r)$$



[Keeling & Berloff, PRL, '08]

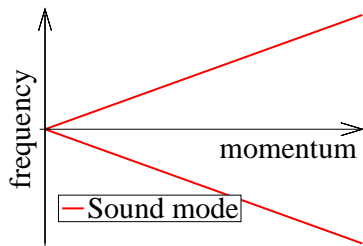
# Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = \omega^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = c|\mathbf{k}|$$



[Szymańska et al., PRL '06; PRB '07]

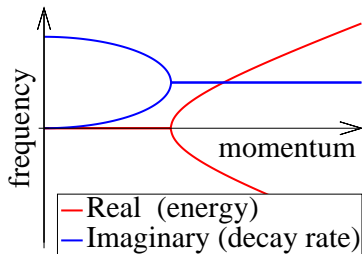
# Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = (\omega + ix)^2 + x^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = -ix \pm \sqrt{c^2 \mathbf{k}^2 - x^2}$$



[Szymańska et al., PRL '06; PRB '07]

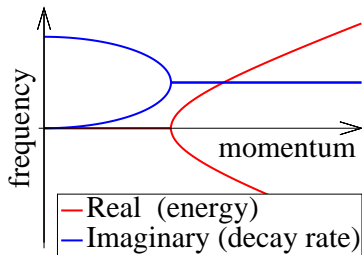
# Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = (\omega + ix)^2 + x^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = -ix \pm \sqrt{c^2 \mathbf{k}^2 - x^2}$$



$$\text{Correlations (in 2D): } \langle \psi^\dagger(\mathbf{r}, t) \psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$$

[Szymańska et al., PRL '06; PRB '07]

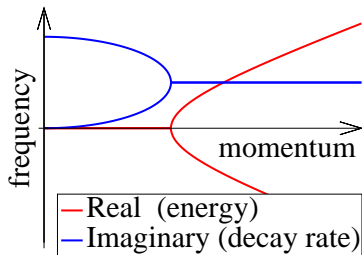
# Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = (\omega + ix)^2 + x^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = -ix \pm \sqrt{c^2 \mathbf{k}^2 - x^2}$$



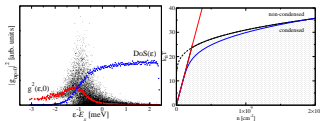
Correlations (in 2D):  $\langle \psi^\dagger(\mathbf{r}, t) \psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$

$$\langle \psi^\dagger(\mathbf{r}, t) \psi(0, 0) \rangle \simeq |\psi_0|^2 \exp \left[ -\eta \begin{cases} \ln(r/\xi) & r \rightarrow \infty, t \simeq 0 \\ \frac{1}{2} \ln(c^2 t/x\xi^2) & r \simeq, t \rightarrow \infty \end{cases} \right]$$

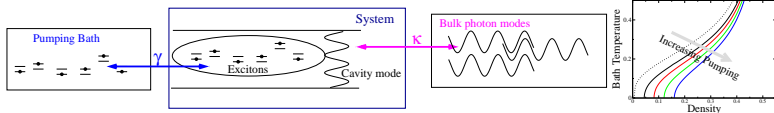
[Szymańska et al., PRL '06; PRB '07]

# Conclusions

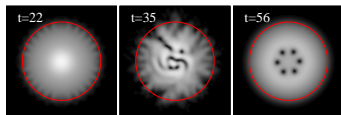
- Localised two-level system model
- Mean-field and fluctuations



- Effects of pumping on mean-field theory



- Modification to Thomas-Fermi profile
- Spontaneous rotating vortex lattice



- 6 Equilibrium spectrum
- 7 Fluctuation corrections
- 8 Vortices
- 9 Superfluidity
- 10 Zero temperature Keldysh boundaries
- 11 Non-equilibrium Fluctuations
  - Finite size effects: single vs many modes

# Equilibrium: Fluctuations about mean-field

Fluctuations  $\psi \rightarrow \psi + \delta\psi$ ; correction to action

$$S = S_0 + \sum_{\nu, \mathbf{p}, \mathbf{q}} \begin{pmatrix} \delta\psi_{\mathbf{p}}^* \\ \delta\psi_{-\mathbf{p}} \end{pmatrix}^T \mathcal{G}_{\mathbf{p}\mathbf{q}}^{-1}(\nu) \begin{pmatrix} \delta\psi_{\mathbf{q}} \\ \delta\psi_{-\mathbf{q}}^* \end{pmatrix}$$



# Equilibrium: Fluctuations about mean-field

Fluctuations  $\psi \rightarrow \psi + \delta\psi$ ; correction to action

$$S = S_0 + \sum_{\nu, \mathbf{p}, \mathbf{q}} \begin{pmatrix} \delta\psi_{\mathbf{p}}^* \\ \delta\psi_{-\mathbf{p}} \end{pmatrix}^T \mathcal{G}_{\mathbf{p}\mathbf{q}}^{-1}(\nu) \begin{pmatrix} \delta\psi_{\mathbf{q}} \\ \delta\psi_{-\mathbf{q}}^* \end{pmatrix}$$

$$\mathcal{G}_{\mathbf{p}\mathbf{q}}^{-1}(\nu) = (\omega_{\mathbf{p}} + i\nu) \mathbb{1} \delta_{\mathbf{p}\mathbf{q}} + \sum_{\alpha} g_{\alpha\mathbf{p}}^* g_{\alpha\mathbf{q}} \begin{pmatrix} \chi_{\alpha}^{(1)}(\nu) & \chi_{\alpha}^{(2)}(\nu) \\ \chi_{\alpha}^{(2)}(\nu) & \chi_{\alpha}^{(1)*}(\nu) \end{pmatrix}$$

# Equilibrium: Fluctuations about mean-field

Fluctuations  $\psi \rightarrow \psi + \delta\psi$ ; correction to action

$$S = S_0 + \sum_{\nu, \mathbf{p}, \mathbf{q}} \begin{pmatrix} \delta\psi_{\mathbf{p}}^* \\ \delta\psi_{-\mathbf{p}} \end{pmatrix}^T \mathcal{G}_{\mathbf{p}\mathbf{q}}^{-1}(\nu) \begin{pmatrix} \delta\psi_{\mathbf{q}} \\ \delta\psi_{-\mathbf{q}}^* \end{pmatrix}$$

$$\mathcal{G}_{\mathbf{p}\mathbf{q}}^{-1}(\nu) = (\omega_{\mathbf{p}} + i\nu) \mathbb{1} \delta_{\mathbf{p}\mathbf{q}} + \sum_{\alpha} g_{\alpha\mathbf{p}}^* g_{\alpha\mathbf{q}} \begin{pmatrix} \chi_{\alpha}^{(1)}(\nu) & \chi_{\alpha}^{(2)}(\nu) \\ \chi_{\alpha}^{(2)}(\nu) & \chi_{\alpha}^{(1)*}(\nu) \end{pmatrix}$$

- Optical response:

Treat  $\mathbf{p} \neq \mathbf{q}$  perturbatively [D. M. Whittaker PRL 80 4791]

- ▶ Spectral weight  $W(\nu, \mathbf{p}) = 2\Im [\mathcal{G}_{\mathbf{p}\mathbf{p}}^{11}(i\nu)]$

# Equilibrium: Fluctuations about mean-field

Fluctuations  $\psi \rightarrow \psi + \delta\psi$ ; correction to action

$$S = S_0 + \sum_{\nu, \mathbf{p}, \mathbf{q}} \begin{pmatrix} \delta\psi_{\mathbf{p}}^* \\ \delta\psi_{-\mathbf{p}} \end{pmatrix}^T \mathcal{G}_{\mathbf{p}\mathbf{q}}^{-1}(\nu) \begin{pmatrix} \delta\psi_{\mathbf{q}} \\ \delta\psi_{-\mathbf{q}}^* \end{pmatrix}$$

$$\mathcal{G}_{\mathbf{p}\mathbf{q}}^{-1}(\nu) = (\omega_{\mathbf{p}} + i\nu) \mathbb{1} \delta_{\mathbf{p}\mathbf{q}} + \sum_{\alpha} g_{\alpha\mathbf{p}}^* g_{\alpha\mathbf{q}} \begin{pmatrix} \chi_{\alpha}^{(1)}(\nu) & \chi_{\alpha}^{(2)}(\nu) \\ \chi_{\alpha}^{(2)}(\nu) & \chi_{\alpha}^{(1)*}(\nu) \end{pmatrix}$$

- Optical response:

Treat  $\mathbf{p} \neq \mathbf{q}$  perturbatively [D. M. Whittaker PRL 80 4791]

▶ Spectral weight	$W(\nu, \mathbf{p})$	=	$2\Im [\mathcal{G}_{\mathbf{p}\mathbf{p}}^{11}(i\nu)]$
▶ Emission	$P_{\text{emit}}(\nu, \mathbf{p})$	=	$n_B(\nu) W(\nu, \mathbf{p})$
▶ Absorption	$P_{\text{absorb}}(\nu, \mathbf{p})$	=	$(1 + n_B(\nu)) W(\nu, \mathbf{p})$

# Equilibrium: Fluctuations about mean-field

Fluctuations  $\psi \rightarrow \psi + \delta\psi$ ; correction to action

$$S = S_0 + \sum_{\nu, \mathbf{p}, \mathbf{q}} \begin{pmatrix} \delta\psi_{\mathbf{p}}^* \\ \delta\psi_{-\mathbf{p}} \end{pmatrix}^T \mathcal{G}_{\mathbf{p}\mathbf{q}}^{-1}(\nu) \begin{pmatrix} \delta\psi_{\mathbf{q}} \\ \delta\psi_{-\mathbf{q}}^* \end{pmatrix}$$

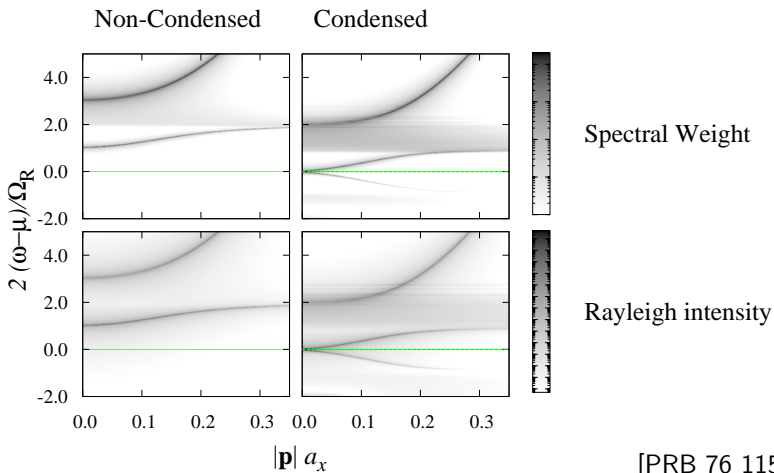
$$\mathcal{G}_{\mathbf{p}\mathbf{q}}^{-1}(\nu) = (\omega_{\mathbf{p}} + i\nu) \mathbb{1} \delta_{\mathbf{p}\mathbf{q}} + \sum_{\alpha} g_{\alpha\mathbf{p}}^* g_{\alpha\mathbf{q}} \begin{pmatrix} \chi_{\alpha}^{(1)}(\nu) & \chi_{\alpha}^{(2)}(\nu) \\ \chi_{\alpha}^{(2)}(\nu) & \chi_{\alpha}^{(1)*}(\nu) \end{pmatrix}$$

- Optical response:

Treat  $\mathbf{p} \neq \mathbf{q}$  perturbatively [D. M. Whittaker PRL 80 4791]

- |                       |                                       |   |  |
|-----------------------|---------------------------------------|---|--|
| ▶ Spectral weight     | $W(\nu, \mathbf{p})$                  | = | $2\Im [\mathcal{G}_{\mathbf{p}\mathbf{p}}^{11}(i\nu)]$ |
| ▶ Emission            | $P_{\text{emit}}(\nu, \mathbf{p})$    | = | $n_B(\nu) W(\nu, \mathbf{p})$                          |
| ▶ Absorption          | $P_{\text{absorb}}(\nu, \mathbf{p})$  | = | $(1 + n_B(\nu)) W(\nu, \mathbf{p})$                    |
| ▶ Rayleigh scattering | $I_{\mathbf{p} \neq \mathbf{q}}(\nu)$ | = | $ \mathcal{G}_{\mathbf{p}\mathbf{q}}^{11}(i\nu) ^2$    |

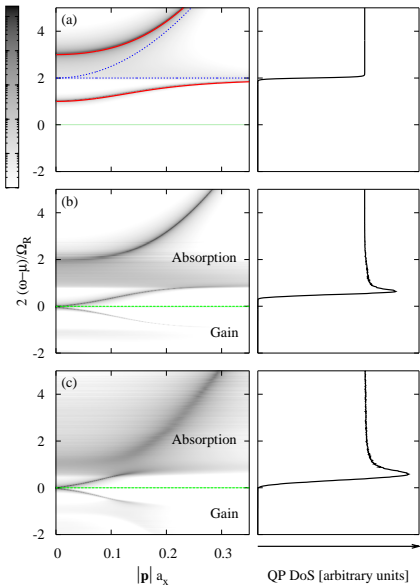
# Fluctuations and optical spectra



[PRB 76 115326 (2007)]

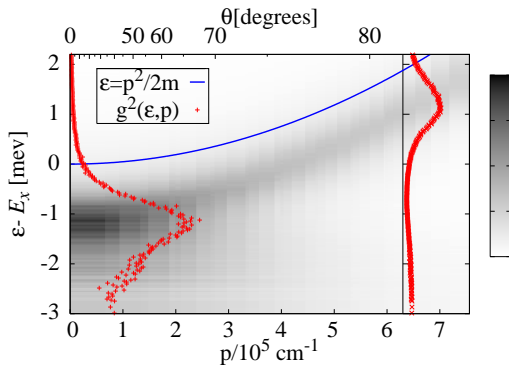
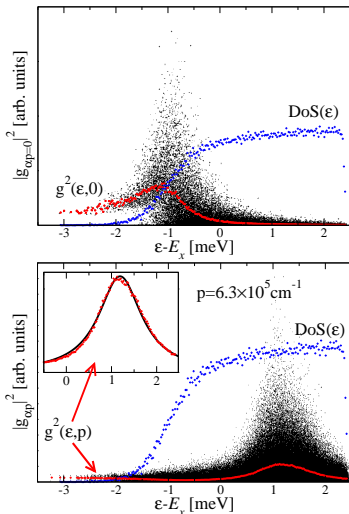
- Phase sensitive Rayleigh  $\rightarrow$  “negative energy” Bogoliubov modes.

# Exciton disorder and polariton density of states



- “Dark” exciton states left at exciton energy.
- Dark states not truly dark, but weak coupling.
- No gap in condensate due to weak coupling tail.

# Disorder localised states



# Fluctuation corrections to phase boundary

Fluctuation corrections to density

$$\rho \rightarrow \rho_0 + \sum_q \langle \delta\psi^\dagger \delta\psi \rangle$$

In 2D system: modified critical condition:

$$\rho_s = \rho - \rho_{\text{normal}} = \# \frac{2m^2}{\hbar} k_B T$$



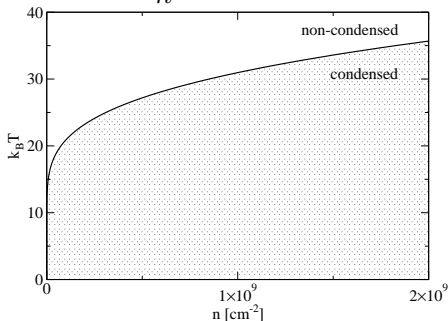
# Fluctuation corrections to phase boundary

Fluctuation corrections to density

$$\rho \rightarrow \rho_0 + \sum_q \langle \delta\psi^\dagger \delta\psi \rangle$$

In 2D system: modified critical condition:

$$\rho_s = \rho - \rho_{\text{normal}} = \# \frac{2m^2}{\hbar} k_B T$$



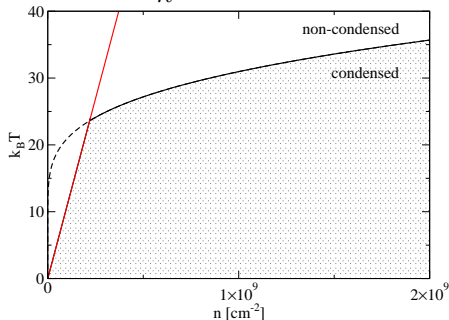
# Fluctuation corrections to phase boundary

Fluctuation corrections to density

$$\rho \rightarrow \rho_0 + \sum_q \langle \delta\psi^\dagger \delta\psi \rangle$$

In 2D system: modified critical condition:

$$\rho_s = \rho - \rho_{\text{normal}} = \# \frac{2m^2}{\hbar} k_B T$$



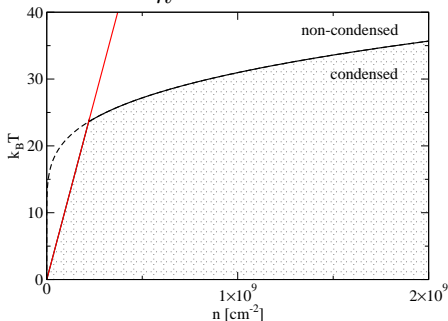
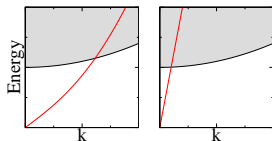
# Fluctuation corrections to phase boundary

Fluctuation corrections to density

$$\rho \rightarrow \rho_0 + \sum_q \langle \delta\psi^\dagger \delta\psi \rangle$$

In 2D system: modified critical condition:

$$\rho_s = \rho - \rho_{\text{normal}} = \# \frac{2m^2}{\hbar} k_B T$$



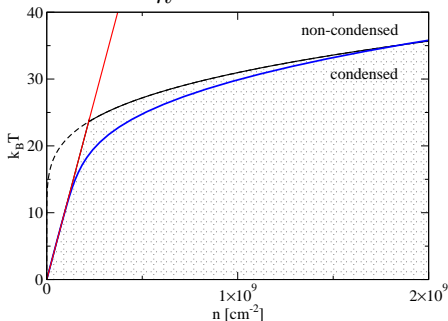
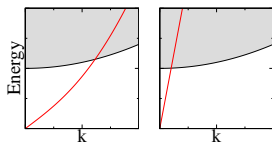
# Fluctuation corrections to phase boundary

Fluctuation corrections to density

$$\rho \rightarrow \rho_0 + \sum_q \langle \delta\psi^\dagger \delta\psi \rangle$$

In 2D system: modified critical condition:

$$\rho_s = \rho - \rho_{\text{normal}} = \# \frac{2m^2}{\hbar} k_B T$$



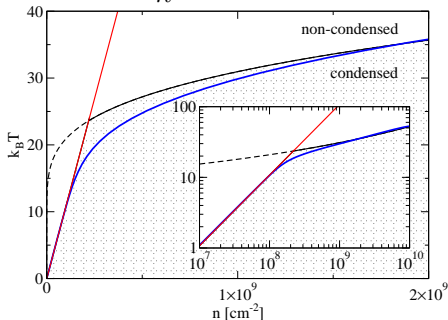
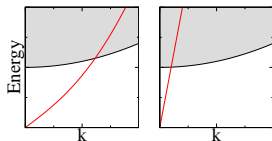
# Fluctuation corrections to phase boundary

Fluctuation corrections to density

$$\rho \rightarrow \rho_0 + \sum_q \langle \delta\psi^\dagger \delta\psi \rangle$$

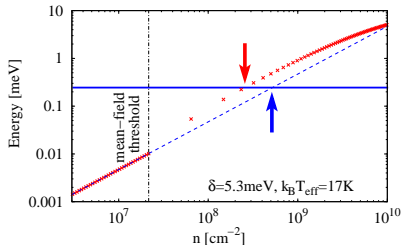
In 2D system: modified critical condition:

$$\rho_s = \rho - \rho_{\text{normal}} = \# \frac{2m^2}{\hbar} k_B T$$



# Blueshift and experimental phase boundary

Blueshift:



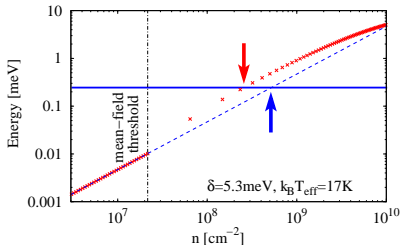
Clean limit:

$$\delta E_{\text{LP}} \simeq \mathcal{R}y_X a_X^2 n + \Omega_R a_X^2 n$$

Here:  $\Omega_R a_X^2 \rightarrow \Omega_R \xi^2$   
[PRB 77 235313]

# Blueshift and experimental phase boundary

Blueshift:

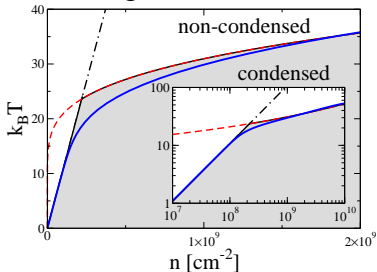


Clean limit:

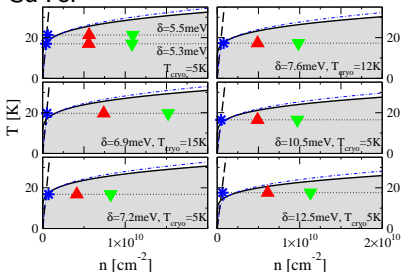
$$\delta E_{LP} \simeq \mathcal{R}_{YX} a_X^2 n + \Omega_R a_X^2 n$$

Here:  $\Omega_R a_X^2 \rightarrow \Omega_R \xi^2$   
[PRB 77 235313]

Phase diagram:



CdTe:



# Instability of Thomas-Fermi: details

$$\psi = \sqrt{\rho} e^{i\phi} \quad \mathbf{v} = \frac{\hbar}{m} \nabla \phi$$

$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}} - \Gamma \rho) \rho$$

$$\partial_t \mathbf{v} + \nabla \left( U \rho + \frac{m\omega^2}{2} r^2 + \frac{m}{2} |\mathbf{v}|^2 \right) = 0$$



# Instability of Thomas-Fermi: details

$$\psi = \sqrt{\rho} e^{i\phi} \quad \mathbf{v} = \frac{\hbar}{m} \nabla \phi$$

$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}} - \Gamma \rho) \rho$$

$$\partial_t \mathbf{v} + \nabla \left( U \rho + \frac{m\omega^2}{2} r^2 + \frac{m}{2} |\mathbf{v}|^2 \right) = 0$$

Consider  $\rho \rightarrow \rho + \delta\rho$ ,  $\mathbf{v} \rightarrow \mathbf{v} + \delta\mathbf{v}$ .

# Instability of Thomas-Fermi: details

$$\psi = \sqrt{\rho} e^{i\phi} \quad \mathbf{v} = \frac{\hbar}{m} \nabla \phi$$

$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}} - \Gamma \rho) \rho$$

$$\partial_t \mathbf{v} + \nabla \left( U \rho + \frac{m\omega^2}{2} r^2 + \frac{m}{2} |\mathbf{v}|^2 \right) = 0$$

If  $\gamma_{\text{net}}, \Gamma \rightarrow 0$ , can find normal modes in 2D trap:

$$\delta \rho_{n,m}(r, \theta, t) = e^{im\theta} h_{n,m}(r) e^{i\omega_{n,m} t}$$

$$\omega_{n,m} = \omega 2 \sqrt{m(1+2n) + 2n(n+1)}$$

Consider  $\rho \rightarrow \rho + \delta\rho, \mathbf{v} \rightarrow \mathbf{v} + \delta\mathbf{v}$ .

# Instability of Thomas-Fermi: details

$$\psi = \sqrt{\rho} e^{i\phi} \quad \mathbf{v} = \frac{\hbar}{m} \nabla \phi$$

$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}} - \Gamma \rho) \rho$$

$$\partial_t \mathbf{v} + \nabla \left( U \rho + \frac{m\omega^2}{2} r^2 + \frac{m}{2} |\mathbf{v}|^2 \right) = 0$$

If  $\gamma_{\text{net}}, \Gamma \rightarrow 0$ , can find normal modes in 2D trap:

$$\delta \rho_{n,m}(r, \theta, t) = e^{im\theta} h_{n,m}(r) e^{i\omega_{n,m} t}$$

$$\omega_{n,m} = \omega \sqrt{2m(1+2n) + 2n(n+1)}$$

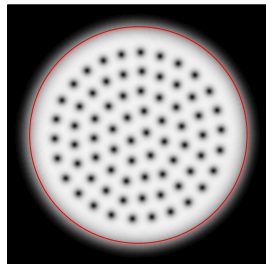
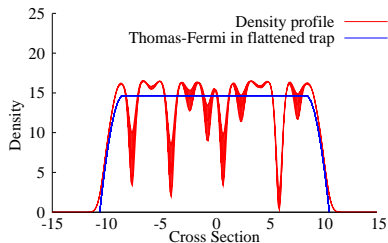
Consider  $\rho \rightarrow \rho + \delta\rho$ ,  $\mathbf{v} \rightarrow \mathbf{v} + \delta\mathbf{v}$ .

Add weak pumping/decay:

$$\omega_{n,n} \rightarrow \omega_{n,m} + i\gamma_{\text{net}} \left[ \frac{m(1+2n) + 2n(n+1) - m^2}{2m(1+2n) + 4n(n+1) + m^2} \right]$$

Instability

# Why vortices

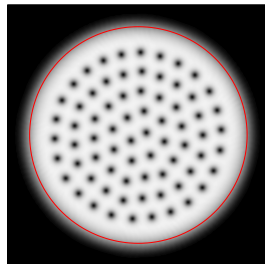
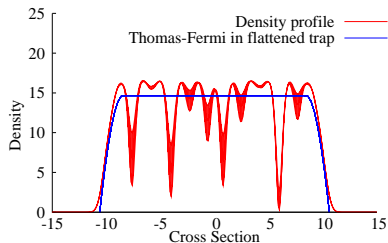


$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{max}} \Theta(R-r) - \Gamma \rho) \rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U \rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{\gamma_{\text{max}}}{\Gamma} \Theta(R-r) = \frac{\mu}{U}$$

# Why vortices



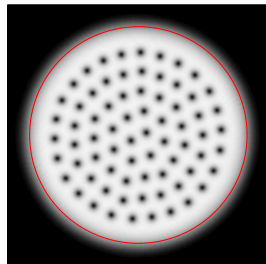
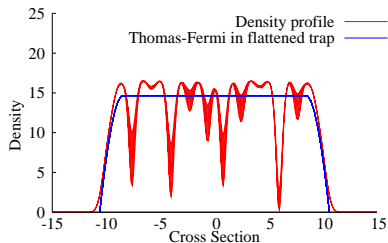
Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{max}} \Theta(R-r) - \Gamma \rho) \rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U\rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{\gamma_{\text{max}}}{\Gamma} \Theta(R-r) = \frac{\mu}{U}$$

# Why vortices



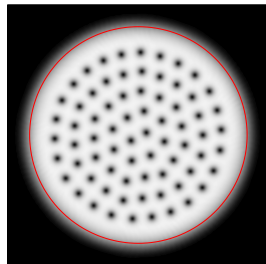
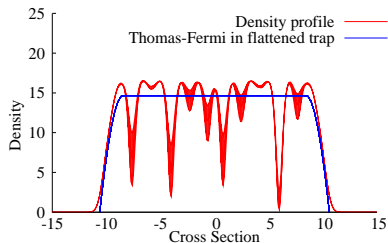
Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{net}}\Theta(R - r) - \Gamma\rho)\rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U\rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{m\omega}{\Gamma} \Theta(R - r) = \frac{\mu}{U}$$

# Why vortices



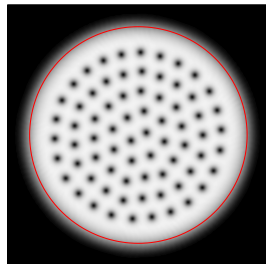
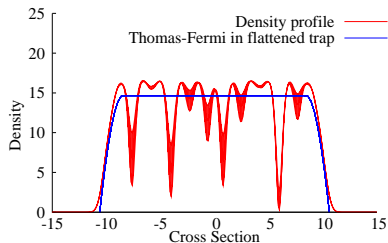
Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{net}}\Theta(R - r) - \Gamma\rho)\rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U\rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{m\omega}{\Gamma} \Theta(R - r) = \frac{\mu}{U}$$

# Why vortices



Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

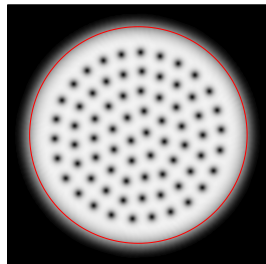
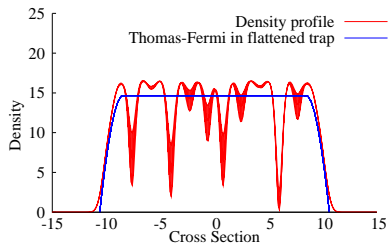
$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{net}}\Theta(R - r) - \Gamma\rho)\rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U\rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{m\omega}{\Gamma} \Theta(R - r) = \frac{\rho_0}{\Gamma}$$



# Why vortices



Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

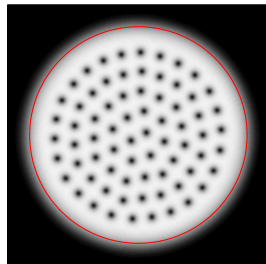
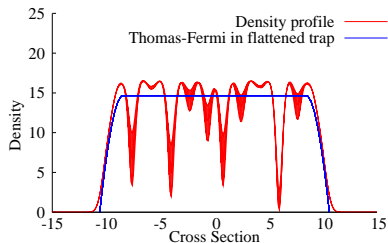
$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{net}}\Theta(R - r) - \Gamma\rho)\rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U\rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega,$$

$$\rho = \frac{\gamma_{\text{net}}}{U} \Theta(R - r) = \frac{\rho_0}{U}$$

# Why vortices



Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

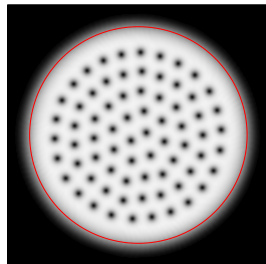
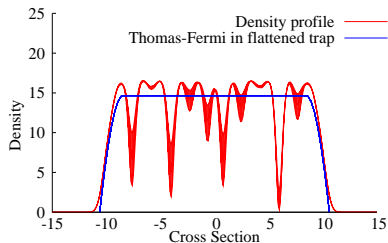
$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{net}}\Theta(R - r) - \Gamma\rho)\rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U\rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega,$$

$$\rho = \frac{\gamma_{\text{net}}}{U} \Theta(R - r) = \frac{\rho_0}{U}$$

# Why vortices



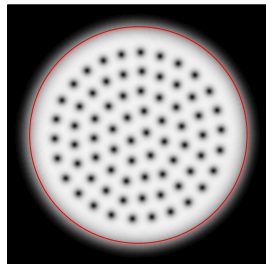
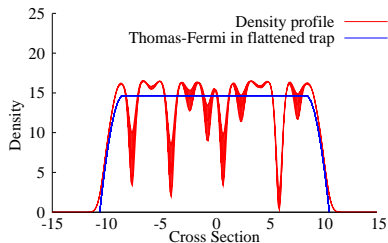
Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{net}}\Theta(R - r) - \Gamma\rho)\rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U\rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{\gamma_{\text{net}}}{\Gamma} \Theta(R - r)$$

# Why vortices



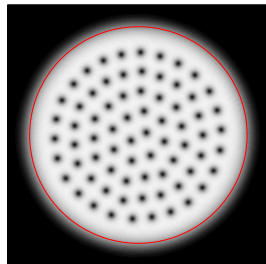
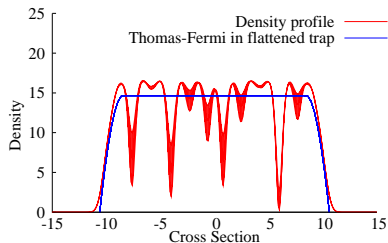
Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{net}}\Theta(R - r) - \Gamma\rho)\rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U\rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{\gamma_{\text{net}}}{\Gamma} \Theta(R - r)$$

# Why vortices



Rotating solution:  $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

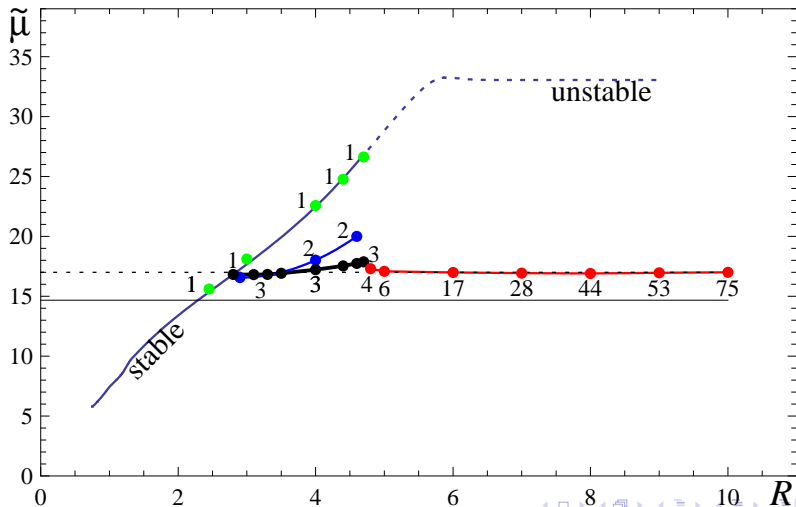
$$\nabla \cdot [\rho(\mathbf{v} - \Omega \times \mathbf{r})] = (\gamma_{\text{net}}\Theta(R - r) - \Gamma\rho)\rho,$$

$$\mu = \frac{m}{2} |\mathbf{v} - \Omega \times \mathbf{r}|^2 + \frac{m}{2} r^2 (\omega^2 - \Omega^2) + U\rho - \frac{\hbar^2 \nabla^2 \sqrt{\rho}}{2m\sqrt{\rho}}$$

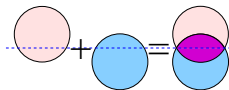
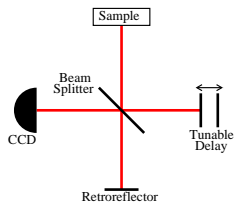
$$\mathbf{v} = \Omega \times \mathbf{r}, \quad \Omega = \omega, \quad \rho = \frac{\gamma_{\text{net}}}{\Gamma} \Theta(R - r) = \frac{\mu}{U}$$

# Why vortices: chemical potential vs size

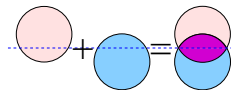
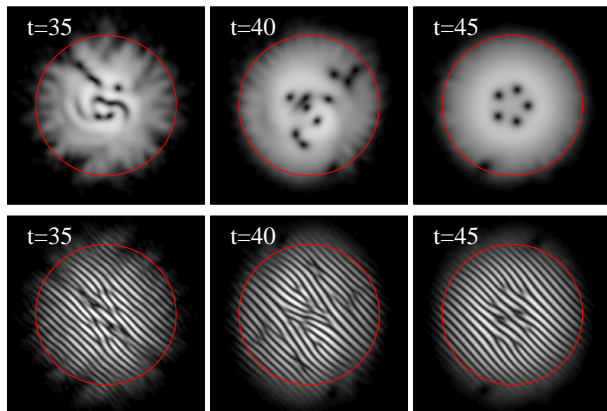
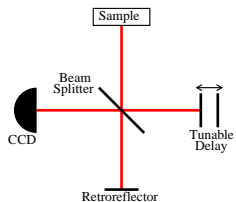
$$\text{Thomas-Fermi : } \mu \propto R^2 \quad \text{Vortex : } \mu = \frac{U\gamma_{\text{net}}}{\Gamma}$$



# Observing vortices: fringe pattern



# Observing vortices: fringe pattern





# Superfluidity

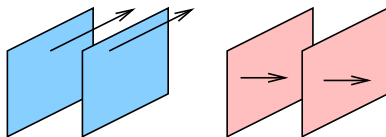
Superfluidity:

$$\mathbf{J} = \rho \mathbf{v} = \Psi^\dagger i \hbar \nabla \Psi = |\Psi|^2 \nabla \phi$$

# Superfluidity

Superfluidity:

$$\mathbf{J} = \rho \mathbf{v} = \Psi^\dagger i \hbar \nabla \Psi = |\Psi|^2 \nabla \phi$$

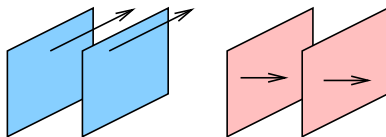


$$\begin{aligned} \chi_{ij}(\omega = 0, \mathbf{q} \rightarrow 0) &= \langle J_i(\mathbf{q}) J_j(\mathbf{q}) \rangle \\ &= \chi_T(q) \left( \delta_{ij} - \frac{q_i q_j}{q^2} \right) + \chi_L(q) \frac{q_i q_j}{q^2} \end{aligned}$$

# Superfluidity

Superfluidity:

$$\mathbf{J} = \rho \mathbf{v} = \Psi^\dagger i \hbar \nabla \Psi = |\Psi|^2 \nabla \phi$$



$$\begin{aligned} \chi_{ij}(\omega = 0, \mathbf{q} \rightarrow 0) &= \langle J_i(\mathbf{q}) J_j(\mathbf{q}) \rangle \\ &= \chi_T(q) \left( \delta_{ij} - \frac{q_i q_j}{q^2} \right) + \chi_L(q) \frac{q_i q_j}{q^2} \end{aligned}$$

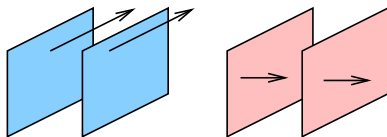
Superfluid part,

$$\rho_s \propto \lim_{q \rightarrow 0} (\chi_L - \chi_T).$$

# Superfluidity

Superfluidity:

$$\mathbf{J} = \rho \mathbf{v} = \Psi^\dagger i \hbar \nabla \Psi = |\Psi|^2 \nabla \phi$$



$$\chi_{ij}(\omega = 0, \mathbf{q} \rightarrow 0) = \langle J_i(\mathbf{q}) J_j(\mathbf{q}) \rangle$$

$$= \chi_T(q) \left( \delta_{ij} - \frac{q_i q_j}{q^2} \right) + \chi_L(q) \frac{q_i q_j}{q^2}$$

Superfluid part,

$$\rho_s \propto \lim_{q \rightarrow 0} (\chi_L - \chi_T).$$

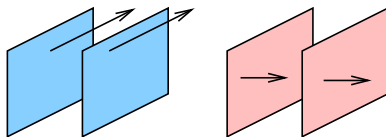
$$J_i(\mathbf{q}) = \langle \psi^\dagger \gamma_i(\mathbf{q}) \psi \rangle$$

$$\Delta \chi_{ij}(q) = \begin{array}{c} \gamma_i(\mathbf{q}, 0) \psi_0 \qquad \gamma_j(\mathbf{q}, 0) \psi_0 \\ \text{---} \bullet \text{---} \blacktriangleright \text{---} \bullet \text{---} \text{wavy} + \dots \\ \mathcal{G}(\omega = 0, \mathbf{q}) \end{array}$$

# Superfluidity

Superfluidity:

$$\mathbf{J} = \rho \mathbf{v} = \Psi^\dagger i \hbar \nabla \Psi = |\Psi|^2 \nabla \phi$$



$$\chi_{ij}(\omega = 0, \mathbf{q} \rightarrow 0) = \langle J_i(\mathbf{q}) J_j(\mathbf{q}) \rangle$$

$$= \chi_T(q) \left( \delta_{ij} - \frac{q_i q_j}{q^2} \right) + \chi_L(q) \frac{q_i q_j}{q^2}$$

Superfluid part,

$$\rho_s \propto \lim_{q \rightarrow 0} (\chi_L - \chi_T).$$

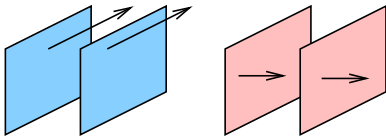
$$J_i(\mathbf{q}) = \langle \psi^\dagger \gamma_i(\mathbf{q}) \psi \rangle$$

$$\begin{aligned} \Delta \chi_{ij}(q) &= \text{---} \overset{\gamma_i(\mathbf{q}, 0) \psi_0}{\bullet} \text{---} \xrightarrow{\mathcal{G}(\omega = 0, \mathbf{q})} \text{---} \overset{\gamma_j(\mathbf{q}, 0) \psi_0}{\bullet} \text{---} + \dots \\ &= \gamma_i(\mathbf{q}) \mathcal{G}_R(\omega = 0, \vec{q}) \gamma_j(\mathbf{q}) + \dots \end{aligned}$$

# Superfluidity

Superfluidity:

$$\mathbf{J} = \rho \mathbf{v} = \Psi^\dagger i \hbar \nabla \Psi = |\Psi|^2 \nabla \phi$$



$$\chi_{ij}(\omega = 0, \mathbf{q} \rightarrow 0) = \langle J_i(\mathbf{q}) J_j(\mathbf{q}) \rangle$$

$$= \chi_T(q) \left( \delta_{ij} - \frac{q_i q_j}{q^2} \right) + \chi_L(q) \frac{q_i q_j}{q^2}$$

Superfluid part,

$$\rho_s \propto \lim_{q \rightarrow 0} (\chi_L - \chi_T).$$

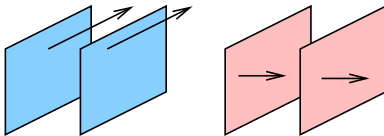
$$J_i(\mathbf{q}) = \langle \psi^\dagger \gamma_i(\mathbf{q}) \psi \rangle$$

$$\begin{aligned} \Delta \chi_{ij}(q) &= \text{---} \overset{\gamma_i(\mathbf{q}, 0) \psi_0}{\bullet} \text{---} \xrightarrow{\mathcal{G}(\omega = 0, \mathbf{q})} \text{---} \overset{\gamma_j(\mathbf{q}, 0) \psi_0}{\bullet} \text{---} + \dots \\ &= \gamma_i(\mathbf{q}) \mathcal{G}_R(\omega = 0, \vec{q}) \gamma_j(\mathbf{q}) + \dots \\ &\propto \frac{q_i (\dots) q_j}{(\omega + ix)^2 + x^2 - c^2 \mathbf{q}^2} + \dots \end{aligned}$$

# Superfluidity

Superfluidity:

$$\mathbf{J} = \rho \mathbf{v} = \Psi^\dagger i \hbar \nabla \Psi = |\Psi|^2 \nabla \phi$$



$$\chi_{ij}(\omega = 0, \mathbf{q} \rightarrow 0) = \langle J_i(\mathbf{q}) J_j(\mathbf{q}) \rangle$$

$$= \chi_T(q) \left( \delta_{ij} - \frac{q_i q_j}{q^2} \right) + \chi_L(q) \frac{q_i q_j}{q^2}$$

Superfluid part,

$$\rho_s \propto \lim_{q \rightarrow 0} (\chi_L - \chi_T).$$

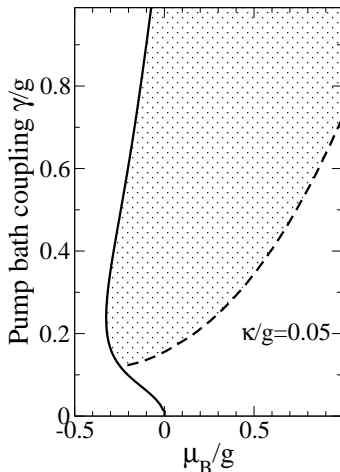
$$J_i(\mathbf{q}) = \langle \psi^\dagger \gamma_i(\mathbf{q}) \psi \rangle$$

Static  $\rho_S$  survives

$$\begin{aligned} \Delta \chi_{ij}(q) &= \text{---} \overset{\gamma_i(\mathbf{q}, 0) \psi_0}{\bullet} \text{---} \xrightarrow{\mathcal{G}(\omega = 0, \mathbf{q})} \text{---} \overset{\gamma_j(\mathbf{q}, 0) \psi_0}{\bullet} \text{---} + \dots \\ &= \gamma_i(\mathbf{q}) \mathcal{G}_R(\omega = 0, \vec{q}) \gamma_j(\mathbf{q}) + \dots \\ &\propto \frac{q_i (\dots) q_j}{(\omega + ix)^2 + x^2 - c^2 \mathbf{q}^2} + \dots \end{aligned}$$

# Zero temperature phase diagram

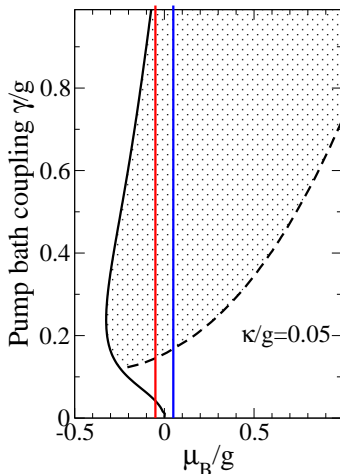
$$\tilde{\omega}_0 - i\kappa = g^2 \gamma \sum_{\alpha} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(\tilde{\epsilon}_{\alpha} + i\gamma)}{[(\nu - E_{\alpha})^2 + \gamma^2][(\nu + E_{\alpha})^2 + \gamma^2]}.$$





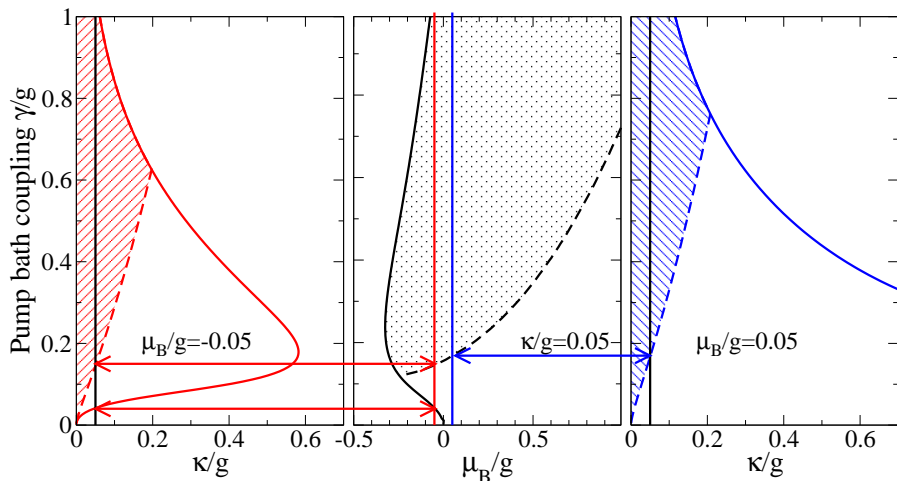
# Zero temperature phase diagram

$$\tilde{\omega}_0 - i\kappa = g^2 \gamma \sum_{\alpha} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(\tilde{\epsilon}_{\alpha} + i\gamma)}{[(\nu - E_{\alpha})^2 + \gamma^2][(\nu + E_{\alpha})^2 + \gamma^2]}.$$



# Zero temperature phase diagram

$$\tilde{\omega}_0 - i\kappa = g^2 \gamma \sum_{\alpha} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(\tilde{\epsilon}_{\alpha} + i\gamma)}{[(\nu - E_{\alpha})^2 + \gamma^2][(\nu + E_{\alpha})^2 + \gamma^2]}.$$



# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

Keldysh approach:

$$\mathcal{G}_{R,A} = \mp i\theta[\pm(t-t')] \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$$

$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle$$

# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

Keldysh approach:

$$\mathcal{G}_R - \mathcal{G}_A = -i \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$$

$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle$$

# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

Keldysh approach:

$$\mathcal{G}_R - \mathcal{G}_A = -i \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$$

$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle = (2n(\omega) + 1)(\mathcal{G}_R - \mathcal{G}_A)$$

# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

Keldysh approach:

$$\mathcal{G}_R - \mathcal{G}_A = -i \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$$

$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle = (2n(\omega) + 1)(\mathcal{G}_R - \mathcal{G}_A)$$

# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

Keldysh approach:

$$\mathcal{G}_R - \mathcal{G}_A = -i \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$$

$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle = (2n(\omega) + 1)(\mathcal{G}_R - \mathcal{G}_A)$$

$$\mathcal{G}_K = -\frac{1}{\mathcal{G}_R^{-1}} \mathcal{G}_K^{-1} \frac{1}{[\mathcal{G}_R^{-1}]^\dagger}$$

$$\mathcal{G}_R - \mathcal{G}_A = \frac{1}{\mathcal{G}_R^{-1}} - \frac{1}{[\mathcal{G}_R^{-1}]^\dagger}$$

# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

Keldysh approach:

$$\mathcal{G}_R - \mathcal{G}_A = -i \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$$

$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle = (2n(\omega) + 1)(\mathcal{G}_R - \mathcal{G}_A)$$

$$\mathcal{G}_K = \frac{-\mathcal{G}_K^{-1}}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

$$\mathcal{G}_R - \mathcal{G}_A = \frac{2\text{Im}[\mathcal{G}_R^{-1}]}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$



# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

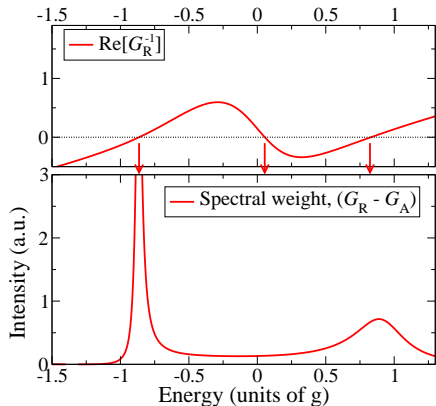
Keldysh approach:  $\mathcal{G}_R - \mathcal{G}_A = -i \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$

$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle = (2n(\omega) + 1)(\mathcal{G}_R - \mathcal{G}_A)$$

$$\mathcal{G}_K = \frac{-\mathcal{G}_K^{-1}}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

$$\mathcal{G}_R - \mathcal{G}_A = \frac{2\text{Im}[\mathcal{G}_R^{-1}]}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

$$\mathcal{G}_R^{-1}(\omega, k) = (\omega - \omega_k^*)$$



# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

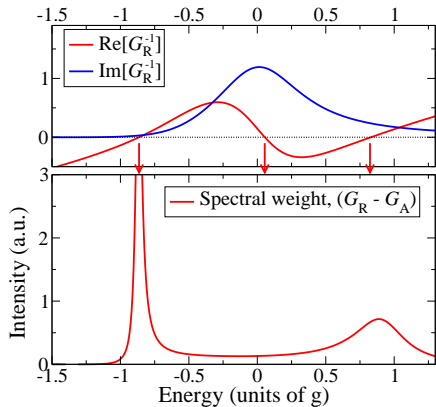
Keldysh approach:  $\mathcal{G}_R - \mathcal{G}_A = -i \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$

$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle = (2n(\omega) + 1)(\mathcal{G}_R - \mathcal{G}_A)$$

$$\mathcal{G}_K = \frac{-\mathcal{G}_K^{-1}}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

$$\mathcal{G}_R - \mathcal{G}_A = \frac{2\text{Im}[\mathcal{G}_R^{-1}]}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

$$\mathcal{G}_R^{-1}(\omega, k) = (\omega - \omega_k^*) + i\alpha$$



# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

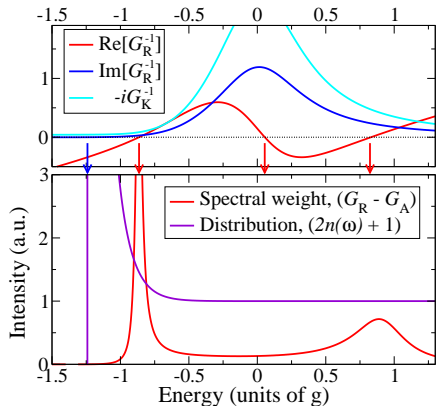
Keldysh approach:  $\mathcal{G}_R - \mathcal{G}_A = -i \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$

$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle = (2n(\omega) + 1)(\mathcal{G}_R - \mathcal{G}_A)$$

$$\mathcal{G}_K = \frac{-\mathcal{G}_K^{-1}}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

$$\mathcal{G}_R - \mathcal{G}_A = \frac{2\text{Im}[\mathcal{G}_R^{-1}]}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

$$\mathcal{G}_R^{-1}(\omega, k) = (\omega - \omega_k^*) + i\alpha(\omega - \mu_{\text{eff}})$$



# Fluctuations $\rightarrow$ Stability, Luminescence, Absorption

Keldysh approach:  $\mathcal{G}_R - \mathcal{G}_A = -i \left\langle \left[ \psi^\dagger, \psi \right]_- \right\rangle$

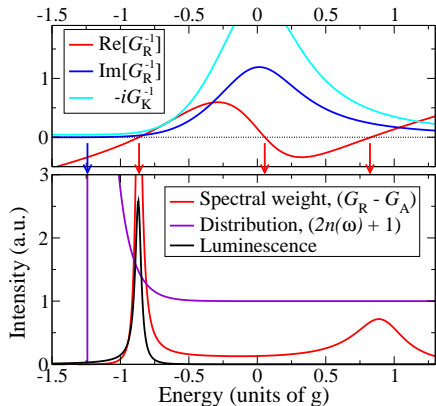
$$\mathcal{G}_K = -i \left\langle \left[ \psi^\dagger, \psi \right]_+ \right\rangle = (2n(\omega) + 1)(\mathcal{G}_R - \mathcal{G}_A)$$

$$\mathcal{L} = \langle \psi^\dagger \psi \rangle = \frac{i}{2} [\mathcal{G}_K + (\mathcal{G}_R - \mathcal{G}_A)]$$

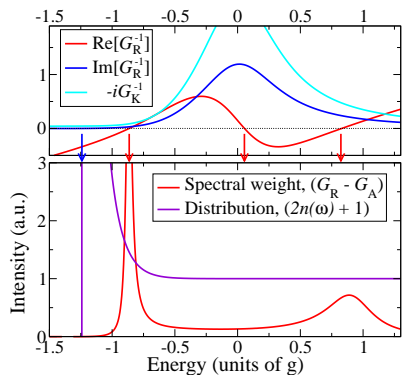
$$\mathcal{G}_K = \frac{-\mathcal{G}_K^{-1}}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

$$\mathcal{G}_R - \mathcal{G}_A = \frac{2\text{Im}[\mathcal{G}_R^{-1}]}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

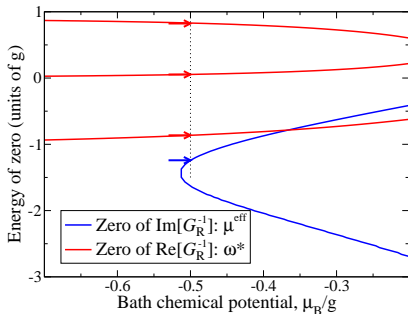
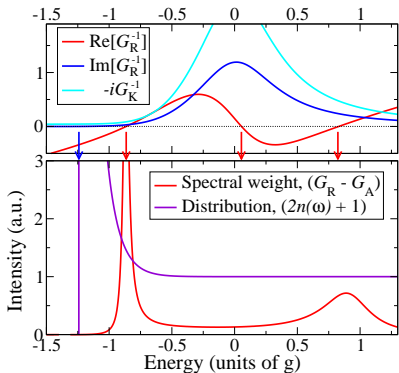
$$\mathcal{G}_R^{-1}(\omega, k) = (\omega - \omega_k^*) + i\alpha(\omega - \mu_{\text{eff}})$$



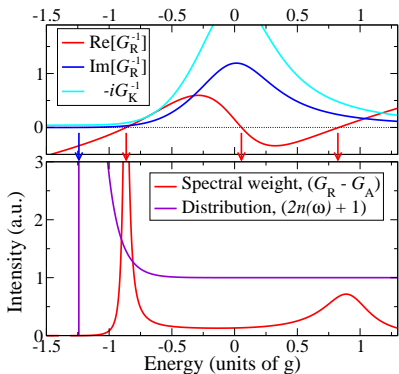
# Linewidth, inverse Green's function and gap equation



# Linewidth, inverse Green's function and gap equation

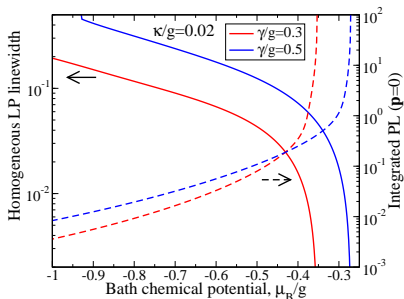
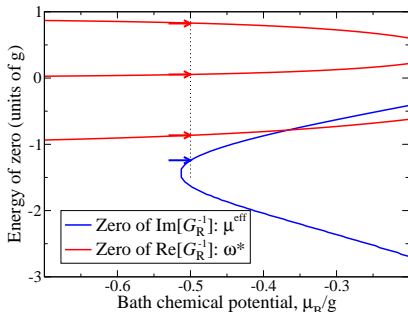


# Linewidth, inverse Green's function and gap equation



At transition, Gap Equation implies:

$$G_R^{-1}(\omega = \mu_S, k = 0) = 0$$



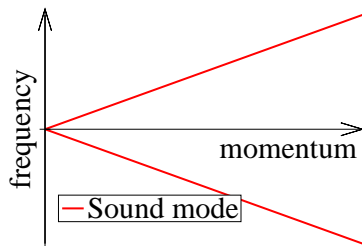
# Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = \omega^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = c|\mathbf{k}|$$



[Szymańska et al., PRL '06; PRB '07]



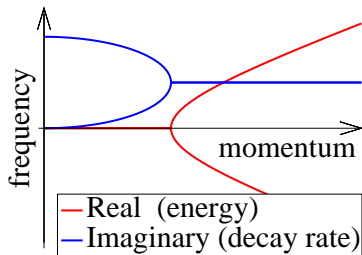
# Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = (\omega + ix)^2 + x^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = -ix \pm \sqrt{c^2 \mathbf{k}^2 - x^2}$$



[Szymańska et al., PRL '06; PRB '07]

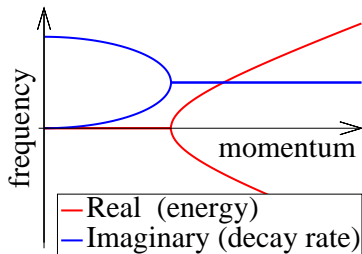
# Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = (\omega + ix)^2 + x^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = -ix \pm \sqrt{c^2 \mathbf{k}^2 - x^2}$$



$$\text{Correlations (in 2D): } \langle \psi^\dagger(\mathbf{r}, t) \psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$$

[Szymańska et al., PRL '06; PRB '07]

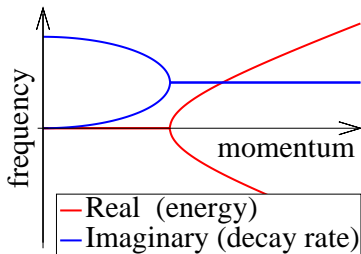
# Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = (\omega + ix)^2 + x^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = -ix \pm \sqrt{c^2 \mathbf{k}^2 - x^2}$$



Correlations (in 2D):  $\langle \psi^\dagger(\mathbf{r}, t) \psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$

$$\langle \psi^\dagger(\mathbf{r}, t) \psi(0, 0) \rangle \simeq |\psi_0|^2 \exp \left[ -\eta \begin{cases} \ln(r/\xi) & r \rightarrow \infty, t \simeq 0 \\ \frac{1}{2} \ln(c^2 t/x\xi^2) & r \simeq, t \rightarrow \infty \end{cases} \right]$$

[Szymańska et al., PRL '06; PRB '07]

## Finite size effects: Single mode vs many mode

$$\langle \psi^\dagger(r, t) \psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$$

# Finite size effects: Single mode vs many mode

$$\langle \psi^\dagger(r, t) \psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$$

$\mathcal{D}_{\phi\phi}(t, r, r')$  from sum of phase modes. Study  $ct \gg r$  limit:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

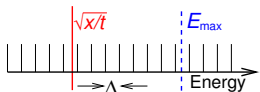
# Finite size effects: Single mode vs many mode

$$\langle \psi^\dagger(r, t) \psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$$

$\mathcal{D}_{\phi\phi}(t, r, r')$  from sum of phase modes. Study  $ct \gg r$  limit:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{\max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

$$\Delta \ll \sqrt{x/t} \ll E_{\max}$$



$$\mathcal{D}_{\phi\phi} \sim 1 + \ln(E_{\max} \sqrt{t/x})$$

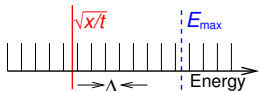
# Finite size effects: Single mode vs many mode

$$\langle \psi^\dagger(r, t) \psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$$

$\mathcal{D}_{\phi\phi}(t, r, r')$  from sum of phase modes. Study  $ct \gg r$  limit:

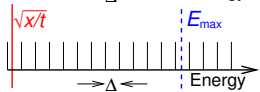
$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{\max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

$$\Delta \ll \sqrt{x/t} \ll E_{\max}$$



$$\mathcal{D}_{\phi\phi} \sim 1 + \ln(E_{\max} \sqrt{t/x})$$

$$\sqrt{x/t} \ll \Delta \ll E_{\max}$$



$$\mathcal{D}_{\phi\phi} \sim \left(\frac{\pi C}{2x}\right) \left(\frac{t}{2x}\right)$$

# Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]

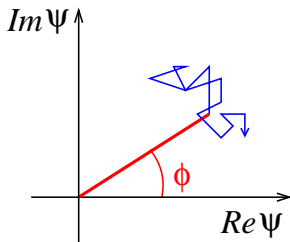


# Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{\max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]



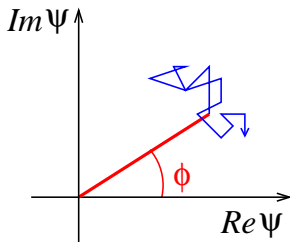
$$\partial_t \phi = U \delta N$$

# Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]



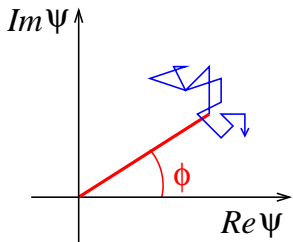
$$\begin{aligned} \partial_t \phi &= U \delta N \\ \partial_t \delta N &= -\Gamma \delta N + F(t), \quad \langle F(t)F(t') \rangle = C \delta(t - t') \end{aligned}$$

# Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{\max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]



$$\partial_t \phi = U \delta N$$

$$\partial_t \delta N = -\Gamma \delta N + F(t), \quad \langle F(t) F(t') \rangle = C \delta(t - t')$$

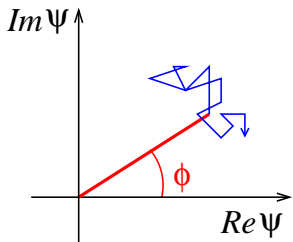
$$\mathcal{D}_{\phi\phi}(t) = \int \frac{d\omega}{2\pi} (1 - e^{i\omega t}) \langle \phi_\omega \phi_{-\omega} \rangle$$

# Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{\max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]



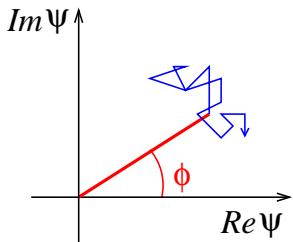
$$\begin{aligned}\partial_t \phi &= U \delta N \\ \partial_t \delta N &= -\Gamma \delta N + F(t), \quad \langle F(t)F(t') \rangle = C \delta(t - t') \\ \mathcal{D}_{\phi\phi}(t) &= \int \frac{d\omega}{2\pi} (1 - e^{i\omega t}) \frac{CU^2}{|\omega(\omega + i\Gamma)|^2}\end{aligned}$$

# Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]



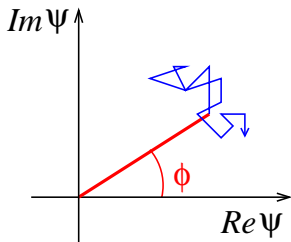
$$\begin{aligned}\partial_t \phi &= U \delta N \\ \partial_t \delta N &= -\Gamma \delta N + F(t), \quad \langle F(t)F(t') \rangle = C \delta(t - t') \\ \mathcal{D}_{\phi\phi}(t) &= \int \frac{d\omega}{2\pi} (1 - e^{i\omega t}) \frac{CU^2}{|\omega(\omega + i\Gamma)|^2} \\ &= \frac{CU^2}{2\Gamma^3} \left[ \Gamma t - 1 + e^{-\Gamma t} \right]\end{aligned}$$

# Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]



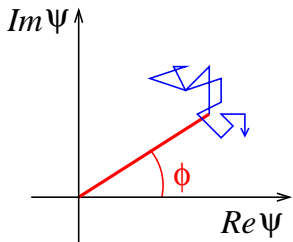
$$\begin{aligned}\partial_t \phi &= U \delta N \\ \partial_t \delta N &= -\Gamma \delta N + F(t), \quad \langle F(t)F(t') \rangle = C \delta(t - t') \\ \mathcal{D}_{\phi\phi}(t) &= \int \frac{d\omega}{2\pi} (1 - e^{i\omega t}) \frac{CU^2}{|\omega(\omega + i\Gamma)|^2} \\ &= \frac{CU^2}{2\Gamma^3} \begin{cases} \Gamma^2 t^2 / 2 & t \ll 1/\Gamma \\ \Gamma t & t \gg 1/\Gamma \end{cases}\end{aligned}$$

# Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]



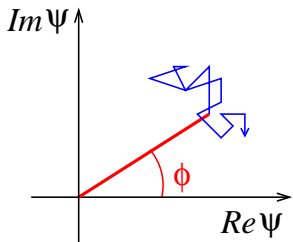
$$\begin{aligned}\partial_t \phi &= U \delta N \\ \partial_t \delta N &= -\Gamma \delta N + F(t), \quad \langle F(t)F(t') \rangle = C \delta(t - t') \\ \mathcal{D}_{\phi\phi}(t) &= \int \frac{d\omega}{2\pi} (1 - e^{i\omega t}) \frac{CU^2}{|\omega(\omega + i\Gamma)|^2} \\ &= \frac{CU^2}{2\Gamma^3} \begin{cases} \Gamma^2 t^2 / 2 & t \ll 1/\Gamma \\ \Gamma t & t \gg 1/\Gamma \end{cases}\end{aligned}$$

# Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2|^2}$$

Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]



$$\begin{aligned}\partial_t \phi &= U \delta N \\ \partial_t \delta N &= -\Gamma \delta N + F(t), \quad \langle F(t)F(t') \rangle = C \delta(t - t') \\ \mathcal{D}_{\phi\phi}(t) &= \int \frac{d\omega}{2\pi} (1 - e^{i\omega t}) \frac{CU^2}{|\omega(\omega + i\Gamma)|^2} \\ &= \frac{CU^2}{2\Gamma^3} \begin{cases} \Gamma^2 t^2 / 2 & t \ll 1/\Gamma \\ \Gamma t & t \gg 1/\Gamma \end{cases}\end{aligned}$$