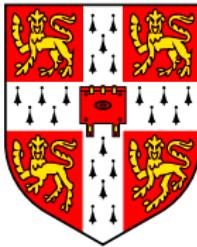


Nonequilibrium quantum condensates: from microscopic theory to macroscopic phenomenology

J. M. J. Keeling

N. G. Berloff, P. B. Littlewood, F. M. Marchetti, M. H. Szymanska.

Sheffield LDSD Seminar. February 2009



Acknowledgements

People:



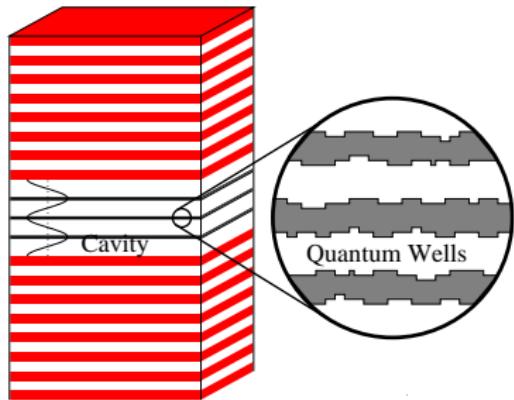
Funding:

EPSRC Engineering and Physical Sciences Research Council

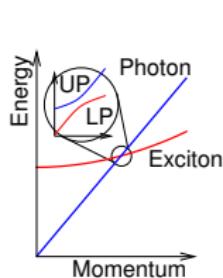
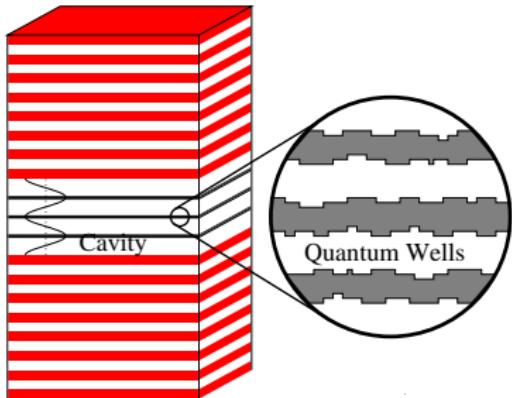


Pembroke College

Microcavity Polaritons

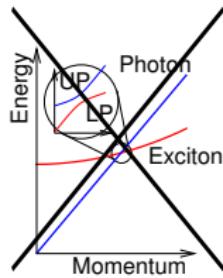
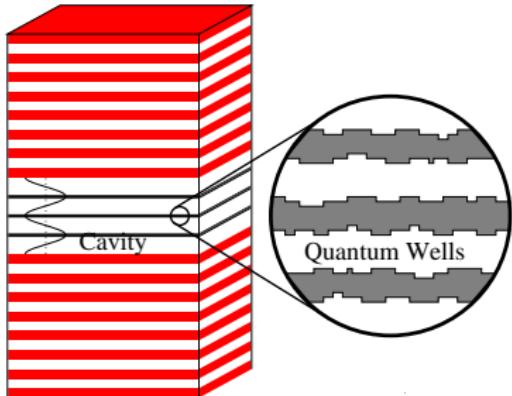


Microcavity Polaritons



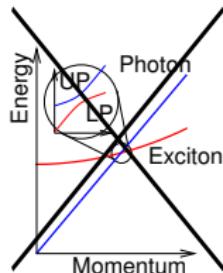
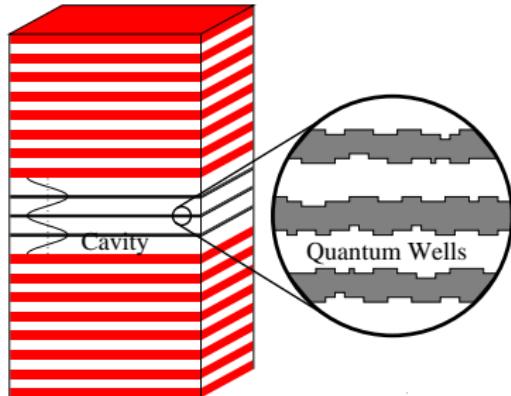
[Pekar, JETP(1958)]
[Hopfield, Phys. Rev.(1958)]

Microcavity Polaritons



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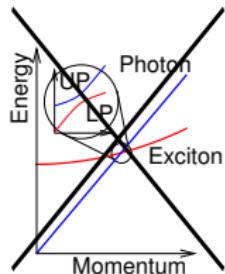
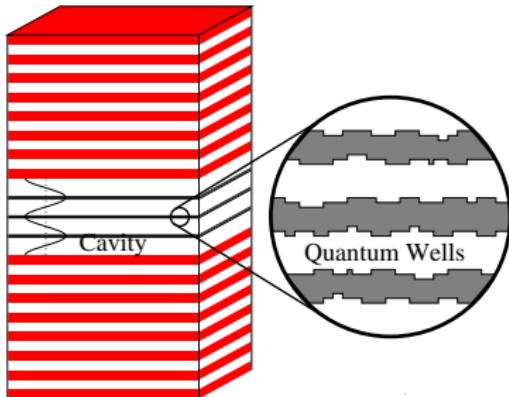
Cavity photons:

$$\omega_k = \sqrt{\omega_0^2 + c^2 k^2}$$

$$\simeq \omega_0 + k^2 / 2m^*$$

$$m^* \sim 10^{-4} m_e$$

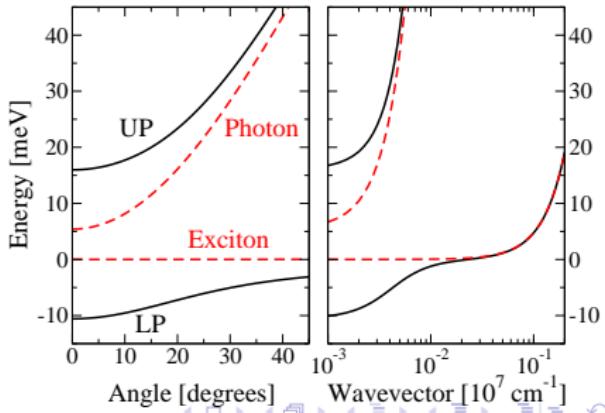
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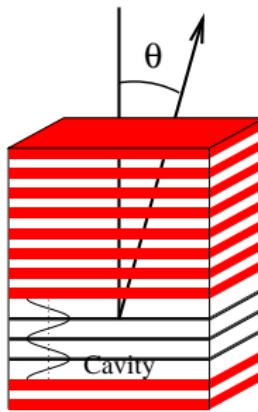
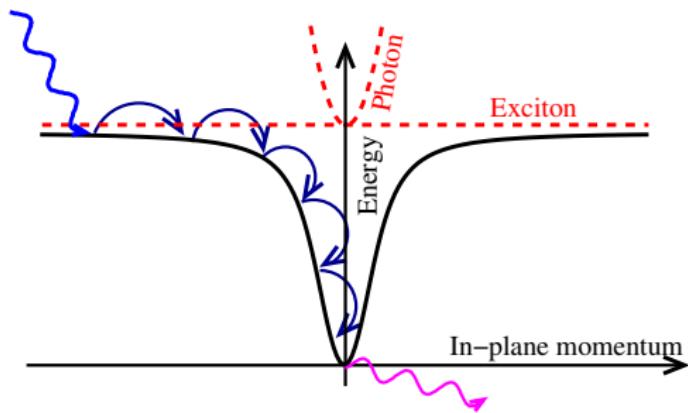
[Pekar, JETP(1958)]
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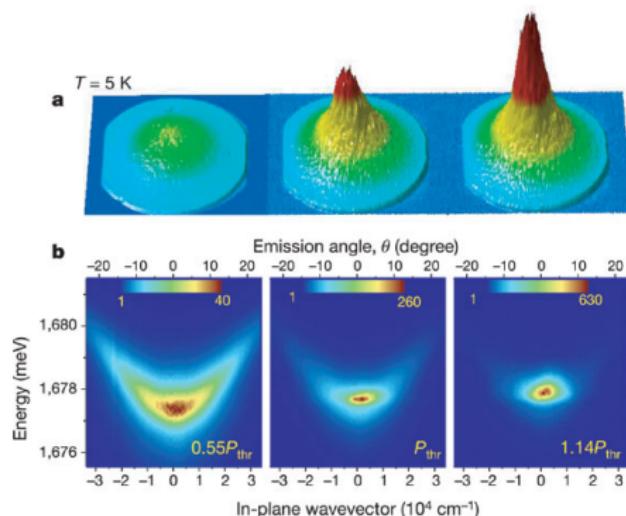
$$\begin{aligned}\omega_k &= \sqrt{\omega_0^2 + c^2 k^2} \\ &\simeq \omega_0 + k^2 / 2m^* \\ m^* &\sim 10^{-4} m_e\end{aligned}$$



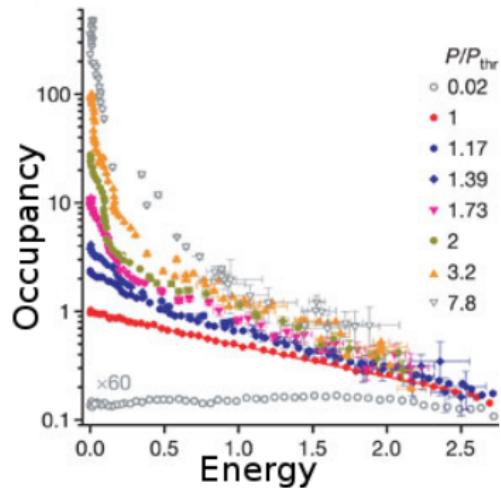
Nonequilibrium system



Polariton experiments: Momentum/Energy distribution

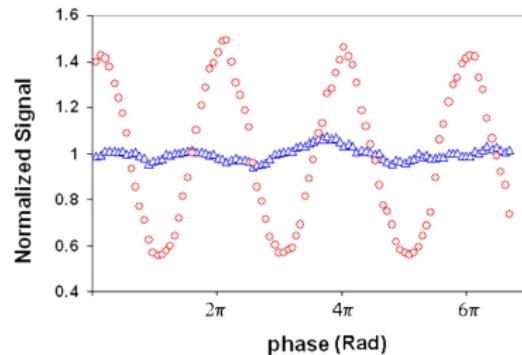
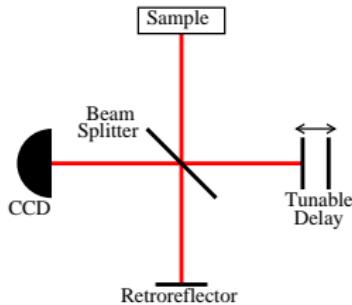


[Kasprzak, et al., Nature, 2006]



Polariton experiments: Coherence

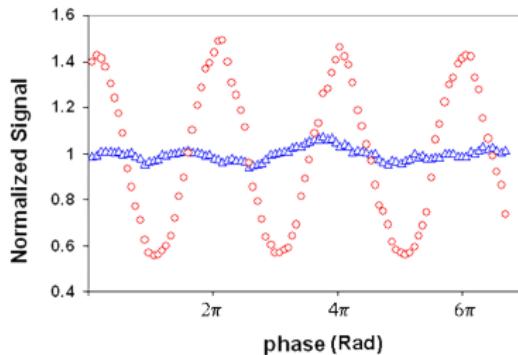
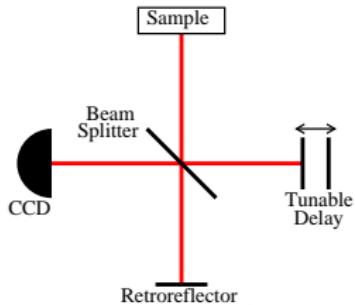
Basic idea:



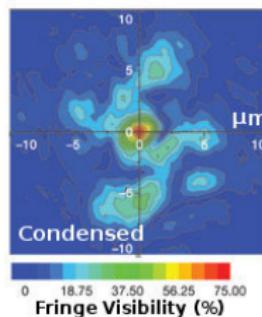
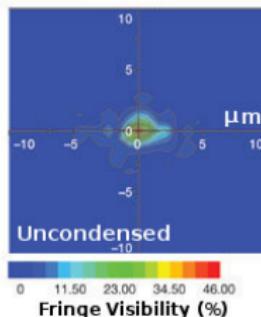
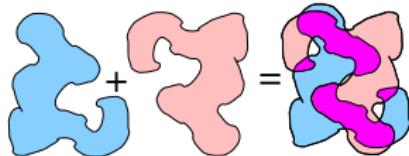
[Kasprzak, et al., Nature, 2006]

Polariton experiments: Coherence

Basic idea:



Coherence map:



[Kasprzak, et al., Nature, 2006]

Other polariton condensation experiments

- Stress traps for polaritons
[Balili *et al* Science 316 1007 (2007)]

Stress-induced polariton trapping
[Love *et al* Phys. Rev. Lett. 101 067404 (2008)]

- Quantised vortices in disorder potential
[Lagoudakis *et al* Nature Phys. 4, 705 (2008)]
- Changes to excitation spectrum
[Utsunomiya *et al* Nature Phys. 4, 700 (2008)]

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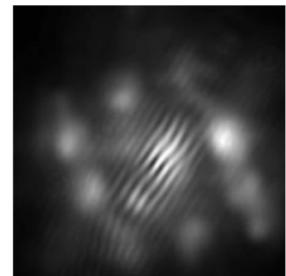
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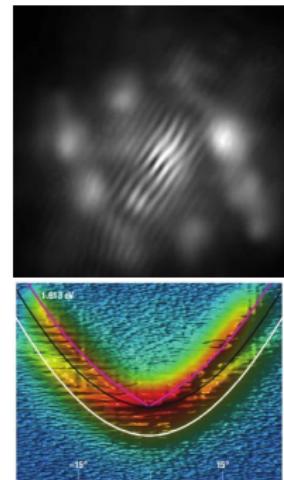


Quantized vortices in disorder potential

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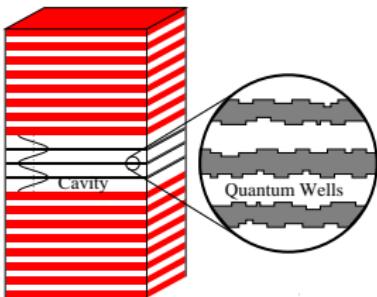
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Overview

- ① Introduction to microcavity polaritons
- ② Model and review of equilibrium results
 - Disorder-localised exciton model
 - Basic equilibrium results
- ③ Non-equilibrium model and mean-field theory
 - Meaning of mean-field condition
- ④ Macroscopic phenomenology
 - Gross Pitaevskii equation in an harmonic trap
- ⑤ Fluctuations and correlations
 - Fluctuations about mean-field theory
 - Finite size effects: single vs many modes

Excitons in a disorderd Quantum well



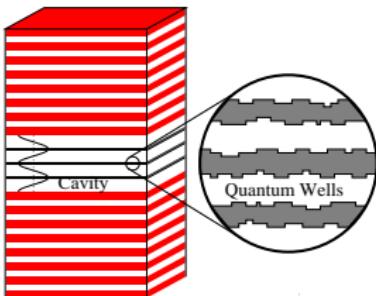
Exciton states in disorder:

$$\left[-\frac{\nabla_{\mathbf{R}}^2}{2m_x} + V(\mathbf{R}) \right] \Phi_{\alpha}(\mathbf{R}) = \varepsilon_{\alpha} \Phi_{\alpha}(\mathbf{R})$$

$V(\vec{R})$ smoothed by exciton Bohr radius

[PRL 96 066405 (2006); PRB 76 115326 (2007)]

Excitons in a disorderd Quantum well



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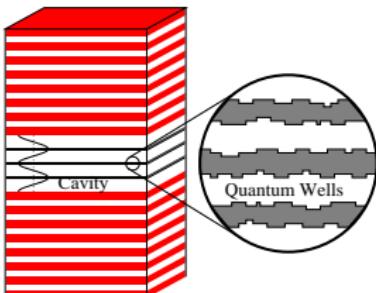
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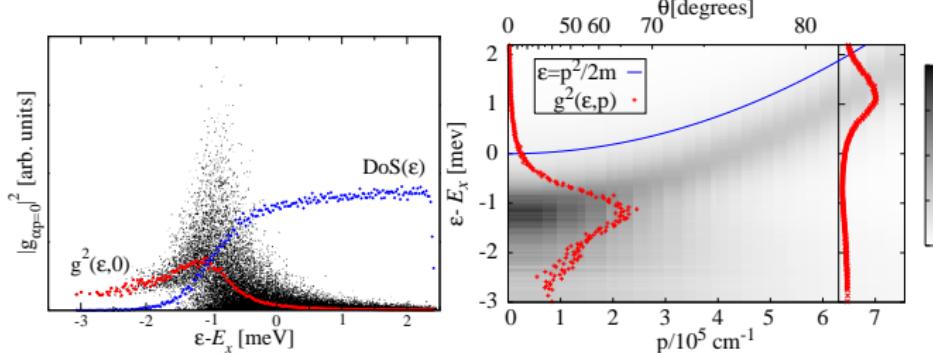


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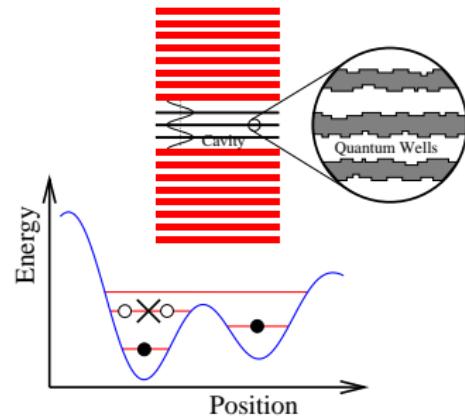


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Polariton system model

Polariton model

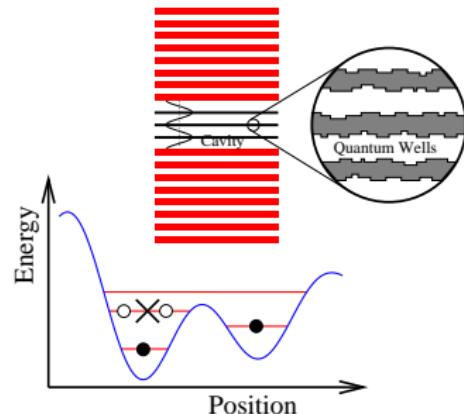
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- Propagating (2D) photons
- Exciton–photon coupling g .



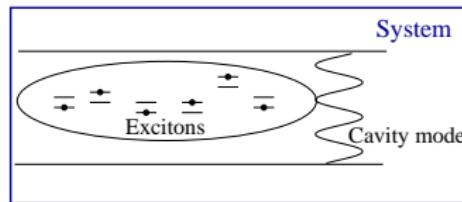
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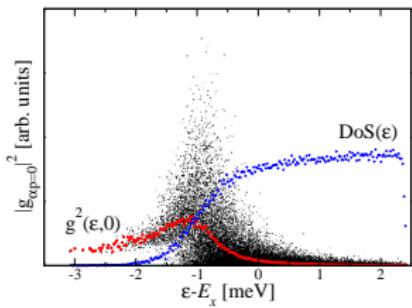


$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \sum_{\alpha} \left[\epsilon_{\alpha} (b_{\alpha}^\dagger b_{\alpha} - a_{\alpha}^\dagger a_{\alpha}) + \frac{1}{\sqrt{\text{Area}}} g_{\alpha, \mathbf{k}} \psi_{\mathbf{k}} b_{\alpha}^\dagger a_{\alpha} + \text{H.c.} \right]$$



Equilibrium: Mean-field theory

$$H_{\text{sys}} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} \psi_{\mathbf{k}}^\dagger \psi_{\mathbf{k}} + \sum_{\alpha} \left[\epsilon_{\alpha} \left(b_{\alpha}^\dagger b_{\alpha} - a_{\alpha}^\dagger a_{\alpha} \right) + \frac{g_{\alpha,\mathbf{k}}}{\sqrt{A}} \psi_{\mathbf{k}} b_{\alpha}^\dagger a_{\alpha} + \text{H.c.} \right]$$



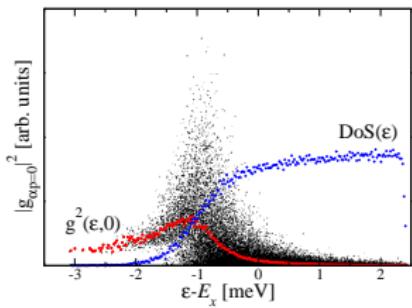
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Mean-field theory:

Self-consistent polarisation and field

$$\left[i\partial_t + \omega_0 - \frac{\nabla^2}{2m} \right] \psi = \frac{1}{\sqrt{A}} \sum_{\alpha} g_{\alpha} P_{\alpha}$$



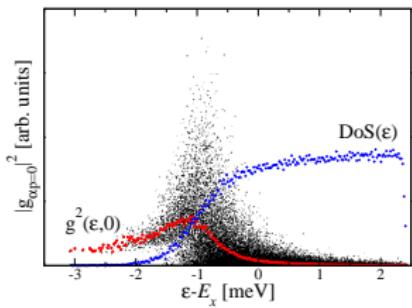
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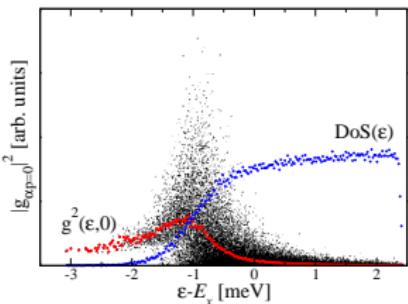
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$$E_{\alpha}^2 = \left(\frac{\tilde{\epsilon}_{\alpha}}{2} \right)^2 + g_{\alpha}^2 \psi^2, \quad \tilde{\epsilon}_{\alpha} = \epsilon_{\alpha} - \mu$$



Equilibrium: Mean-field theory

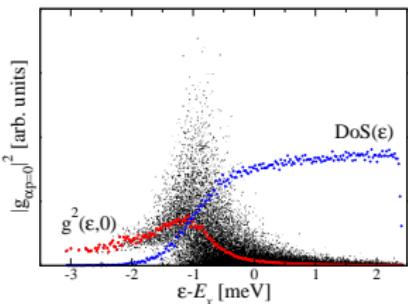
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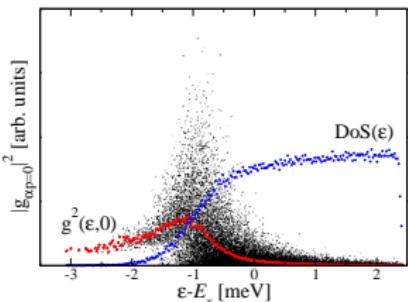
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Density

$$\rho = |\psi|^2 + \frac{1}{A} \sum_{\alpha} \left[\frac{1}{2} - \frac{\tilde{\epsilon}_{\alpha}}{4E_{\alpha}} \tanh(\beta E_{\alpha}) \right]$$

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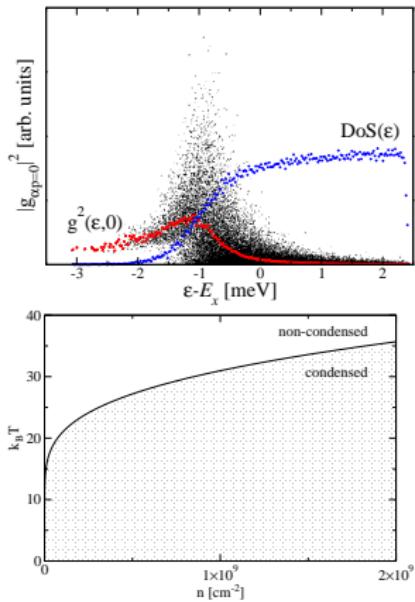
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Equilibrium: Fluctuations about mean-field

Fluctuations $\psi \rightarrow \psi + \delta\psi$; correction to action

$$S = S_0 + \sum_{\nu, \mathbf{p}, \mathbf{q}} \begin{pmatrix} \delta\psi_{\mathbf{p}}^* \\ \delta\psi_{-\mathbf{p}} \end{pmatrix}^T \mathcal{G}_{\mathbf{pq}}^{-1}(\nu) \begin{pmatrix} \delta\psi_{\mathbf{q}} \\ \delta\psi_{-\mathbf{q}}^* \end{pmatrix}$$

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- Optical response:

Treat $\mathbf{p} \neq \mathbf{q}$ perturbatively [D. M. Whittaker PRL 80 4791]

- ▶ Spectral weight $W(\nu, \mathbf{p}) = 2\Im [\mathcal{G}_{\mathbf{pp}}^{11}(i\nu)]$

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- Emission $P_{\text{emit}}(\nu, \mathbf{p}) = n_B(\nu)W(\nu, \mathbf{p})$
- Absorption $P_{\text{absorb}}(\nu, \mathbf{p}) = (1 + n_B(\nu))W(\nu, \mathbf{p})$

Equilibrium: Fluctuations about mean-field

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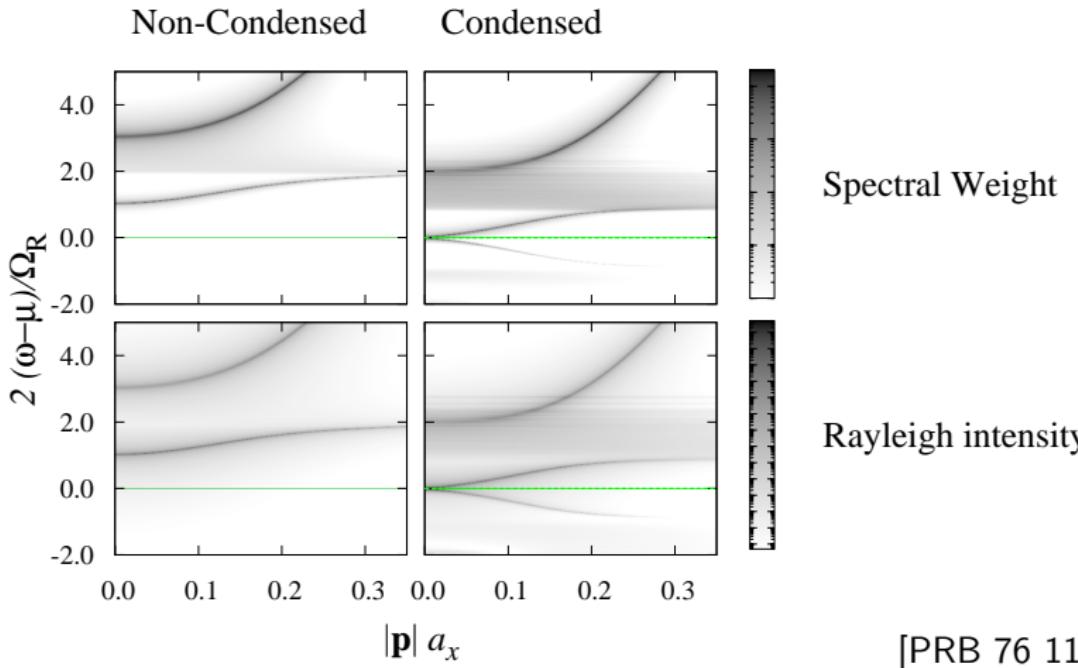
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Treat $\mathbf{p} \neq \mathbf{q}$ perturbatively [D. M. Whittaker PRL 80 4791]

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- ▶ Emission $P_{\text{emit}}(\nu, \mathbf{p})$ = $n_B(\nu)W(\nu, \mathbf{p})$
- ▶ Absorption $P_{\text{absorb}}(\nu, \mathbf{p})$ = $(1 + n_B(\nu))W(\nu, \mathbf{p})$
- ▶ Rayleigh scattering $I_{\mathbf{p} \neq \mathbf{q}}(\nu)$ = $|\mathcal{G}_{\mathbf{pq}}^{11}(i\nu)|^2$

Fluctuations and optical spectra



- Phase sensitive Rayleigh \rightarrow “negative energy” Bogoliubov modes.

Fluctuation corrections to phase boundary

Fluctuation corrections to density

$$\rho \rightarrow \rho_0 + \sum_q \langle \delta\psi^\dagger \delta\psi \rangle$$

In 2D system: modified critical condition:

$$\rho_s = \rho - \rho_{\text{normal}} = \# \frac{2m^2}{\hbar} k_B T$$

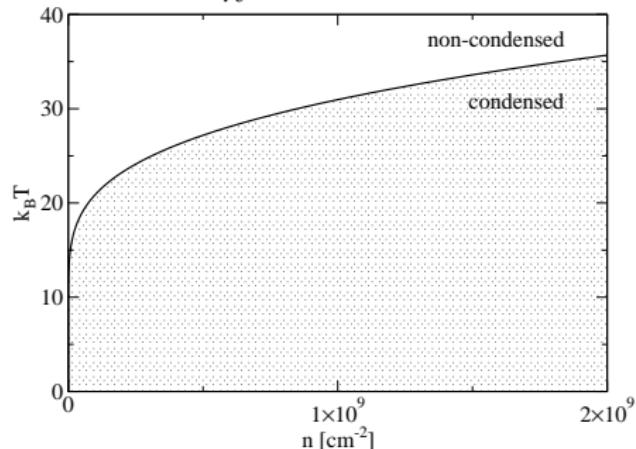
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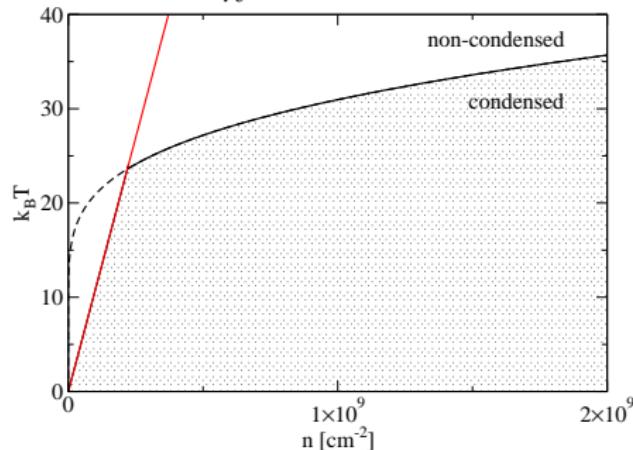
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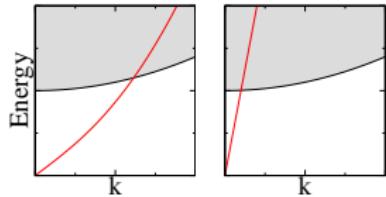


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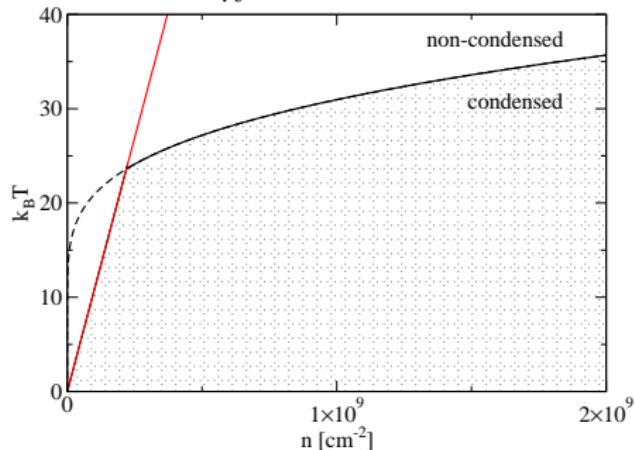
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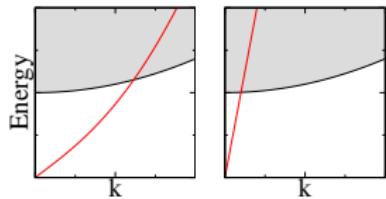


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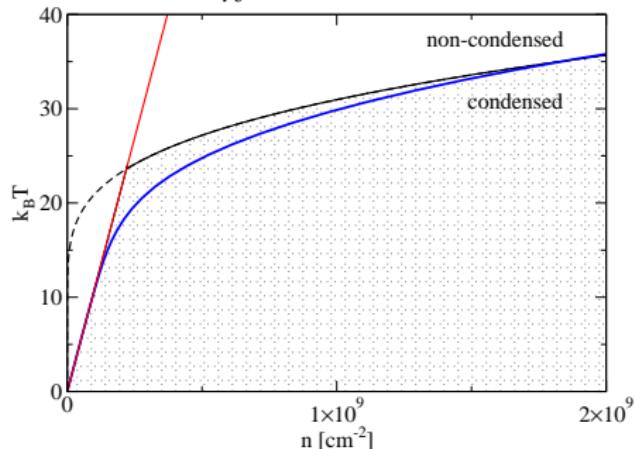
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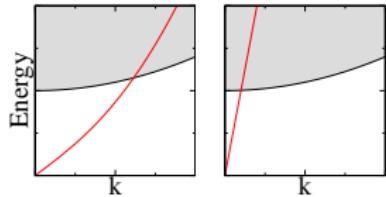


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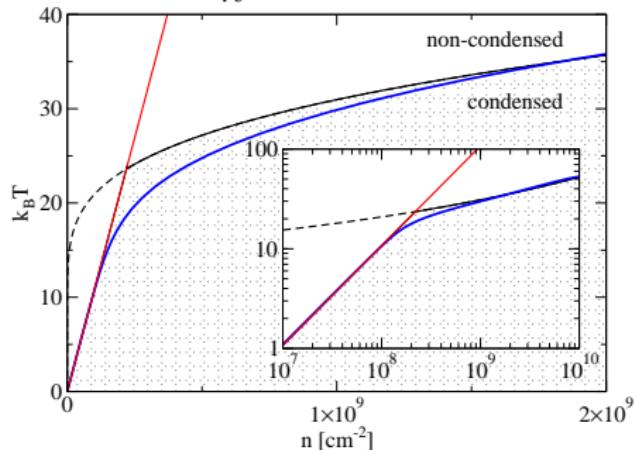
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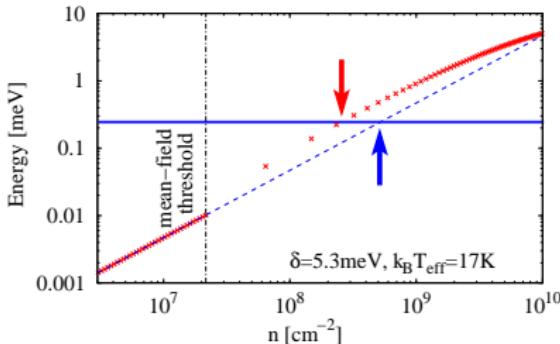


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Blueshift and experimental phase boundary

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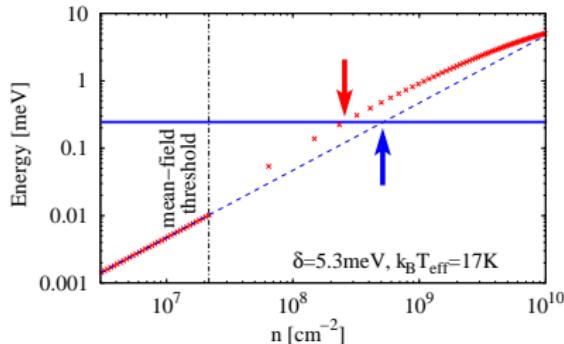
Clean limit:

$$\delta E_{\text{LP}} \simeq \mathcal{R} y_X a_X^2 n + \Omega_R a_X^2 n$$

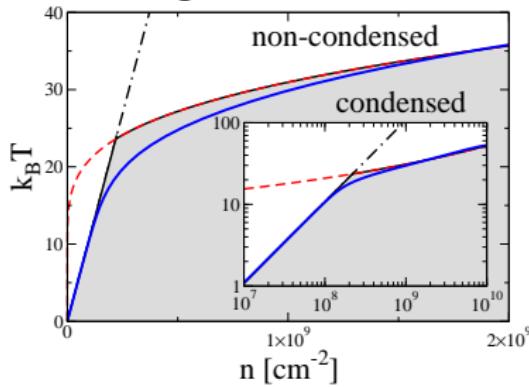
Here: $\Omega_R a_X^2 \rightarrow \Omega_R \xi^2$
[PRB 77 235313]

Blueshift and experimental phase boundary

Blueshift:



Phase diagram:

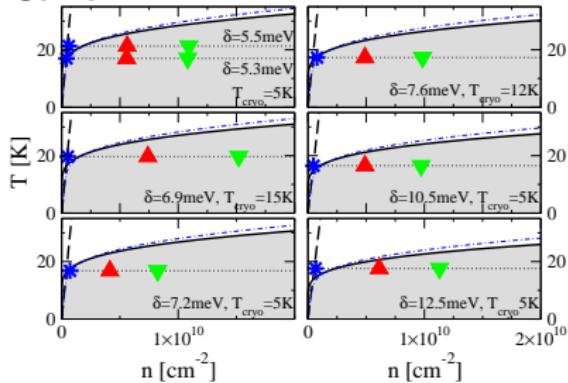


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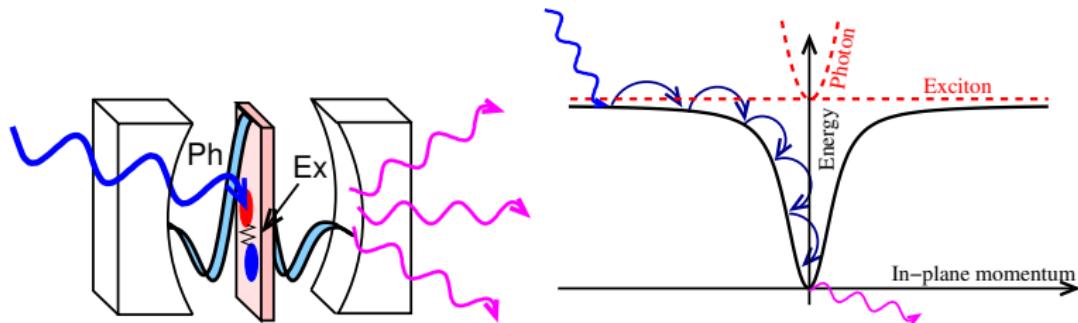
CdTe:



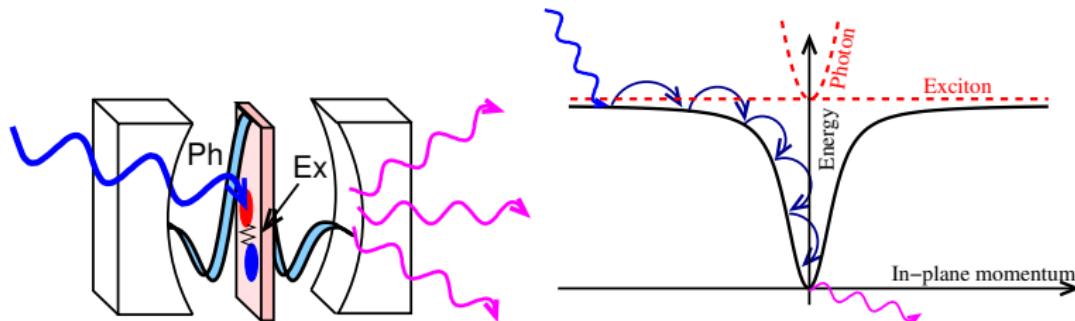
Overview

- 1 Introduction to microcavity polaritons
- 2 Model and review of equilibrium results
 - Disorder-localised exciton model
 - Basic equilibrium results
- 3 Non-equilibrium model and mean-field theory
 - Meaning of mean-field condition
- 4 Macroscopic phenomenology
 - Gross Pitaevskii equation in an harmonic trap
- 5 Fluctuations and correlations
 - Fluctuations about mean-field theory
 - Finite size effects: single vs many modes

Non-equilibrium: Timescales



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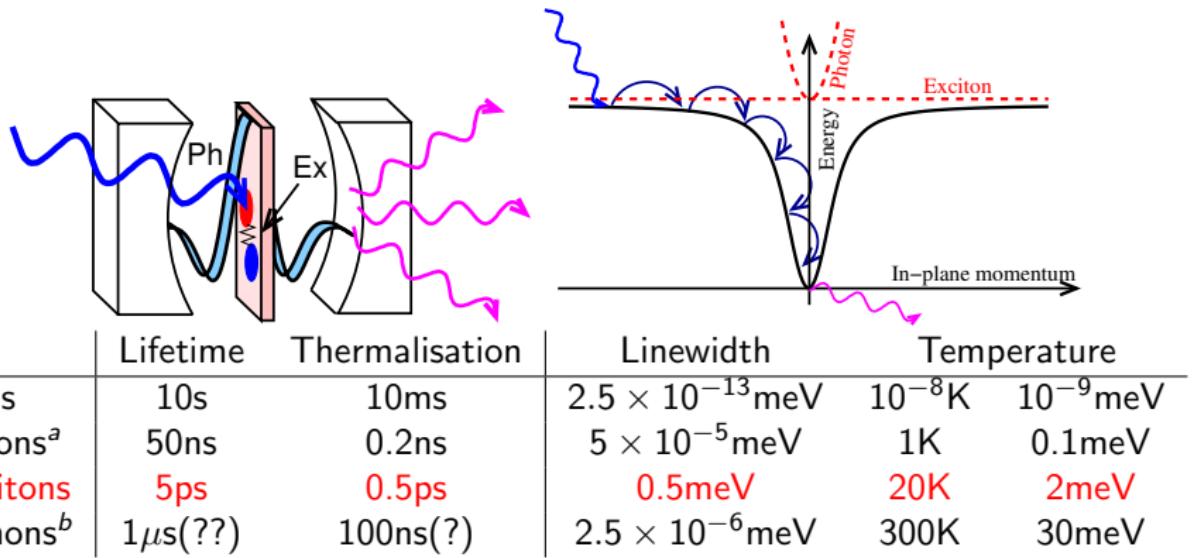


	Lifetime	Thermalisation
Atoms	10s	10ms
Excitons ^a	50ns	0.2ns
Polaritons	5ps	0.5ps
Magnons ^b	1μs(???)	100ns(?)

^aCoupled quantum wells. [Hammack et al PRB 76 193308 (2007)]

^bYttrium Iron Garnett. [Demokritov et al, Nature 443 430 (2007)]

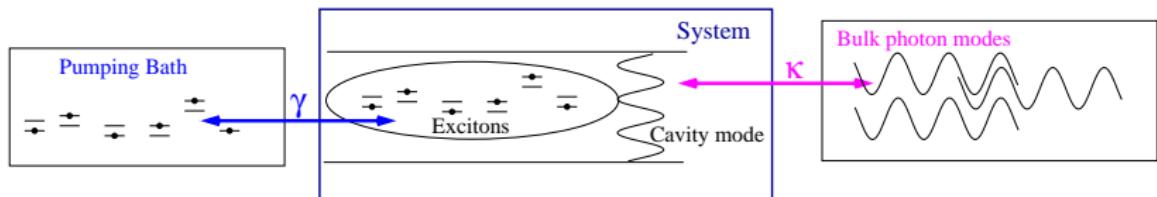
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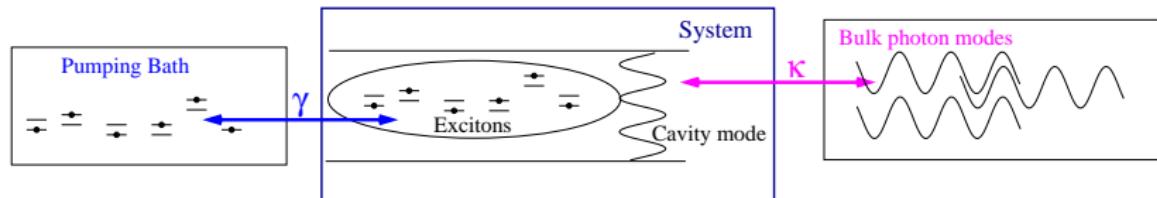
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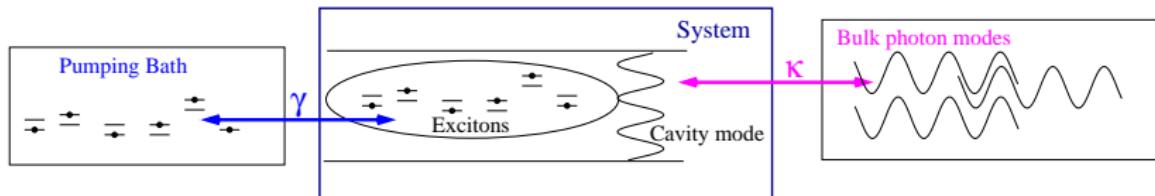


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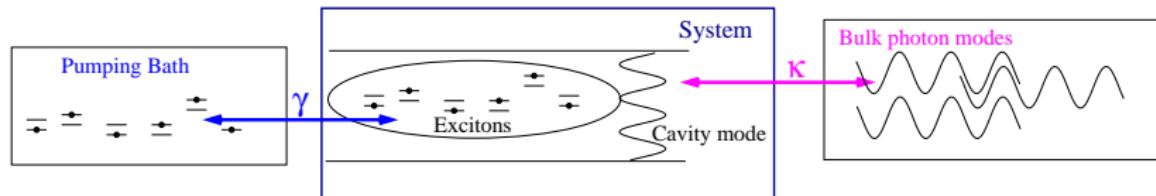
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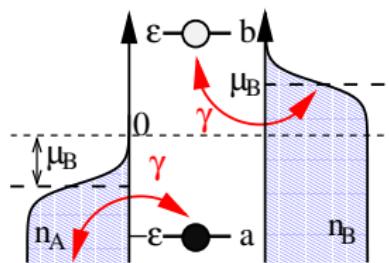


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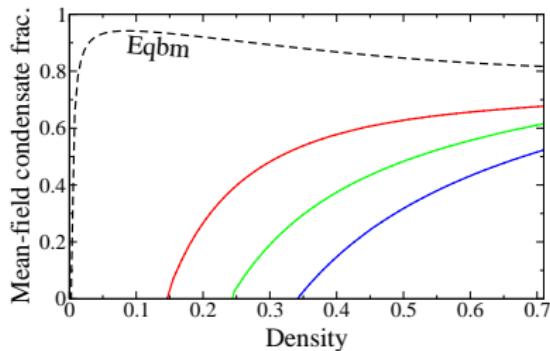
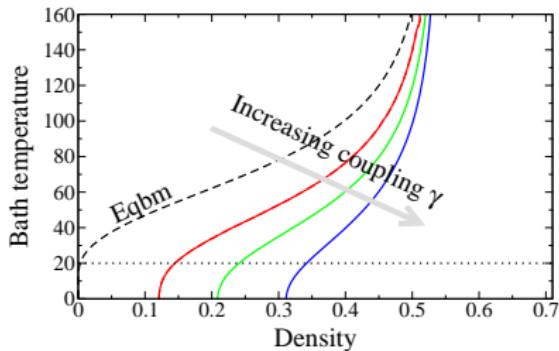
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Nonlinear, complex susceptibility $\chi(\psi(r, t))$

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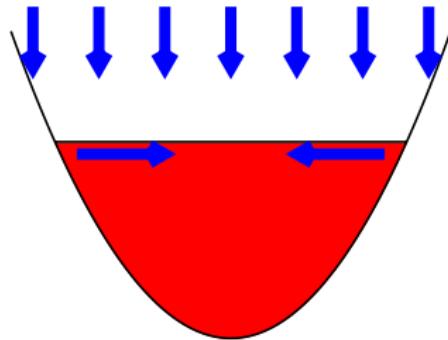
$$i\hbar\partial_t \psi|_{\text{nonlin}} = U|\psi|^2 \psi$$

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Gross-Pitaevskii equation: Harmonic trap

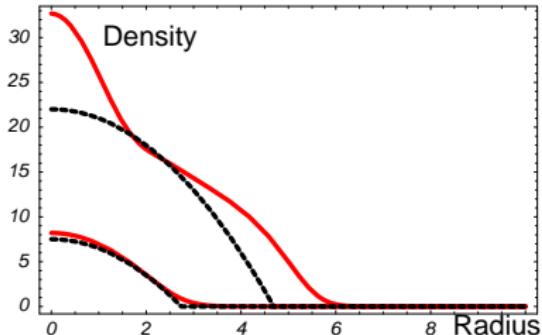
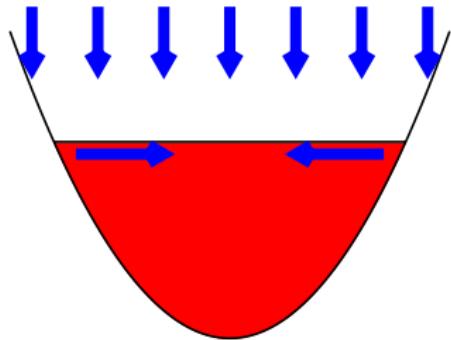
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[Keeling & Berloff, PRL, '08]

Gross-Pitaevskii equation: Harmonic trap

$$i\hbar\partial_t\psi = \left[-\frac{\hbar^2\nabla^2}{2m} + \frac{m\omega^2}{2}r^2 + U|\psi|^2 + i(\gamma_{\text{eff}} - \kappa - \Gamma|\psi|^2) \right] \psi$$



[Keeling & Berloff, PRL, '08]

Instability of Thomas-Fermi: details

$$\psi = \sqrt{\rho} e^{i\phi} \quad \mathbf{v} = \frac{\hbar}{m} \nabla \phi$$

$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}} - \Gamma \rho) \rho$$

$$\partial_t \mathbf{v} + \nabla (U \rho + \frac{m\omega^2}{2} r^2 + \frac{m}{2} |\mathbf{v}|^2) = 0$$

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Consider $\rho \rightarrow \rho + \delta\rho$, $\mathbf{v} \rightarrow \mathbf{v} + \delta\mathbf{v}$.

Instability of Thomas-Fermi: details

$$\psi = \sqrt{\rho} e^{i\phi} \quad \mathbf{v} = \frac{\hbar}{m} \nabla \phi$$

If $\gamma_{\text{net}}, \Gamma \rightarrow 0$, can find normal modes in 2D trap:

$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = (\gamma_{\text{net}} - \Gamma \rho) \rho$$

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$$\delta \rho_{n,m}(r, \theta, t) = e^{im\theta} h_{n,m}(r) e^{i\omega_{n,m} t}$$

$$\omega_{n,m} = \omega 2 \sqrt{m(1+2n) + 2n(n+1)}$$

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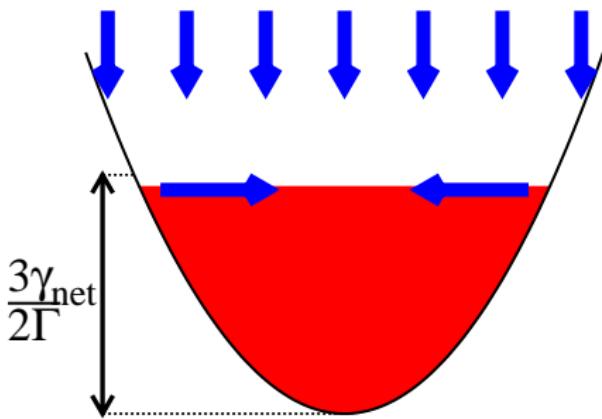
Add weak pumping/decay:

$$\omega_{n,n} \rightarrow \omega_{n,m} + i\gamma_{\text{net}} \left[\frac{m(1+2n) + 2n(n+1) - m^2}{2m(1+2n) + 4n(n+1) + m^2} \right]$$

Instability

Stability of Thomas-Fermi solution

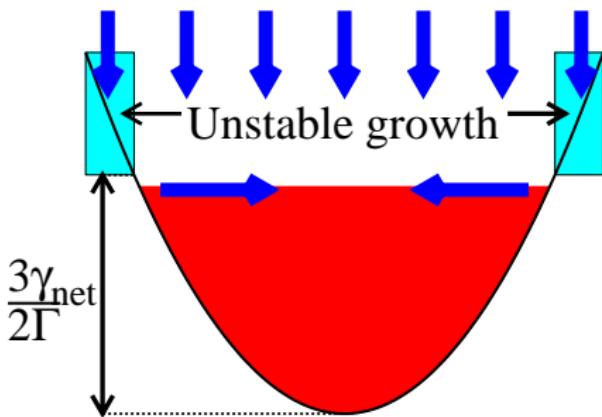
$$\frac{1}{2} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = \frac{1}{\hbar} (\gamma_{\text{net}} - \Gamma \rho) \rho$$



Stability of Thomas-Fermi solution

High m modes: $\delta\rho_{n,m} \simeq e^{im\theta} r^m \dots$

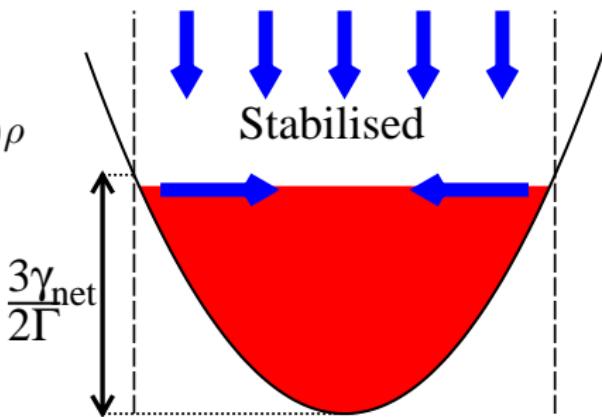
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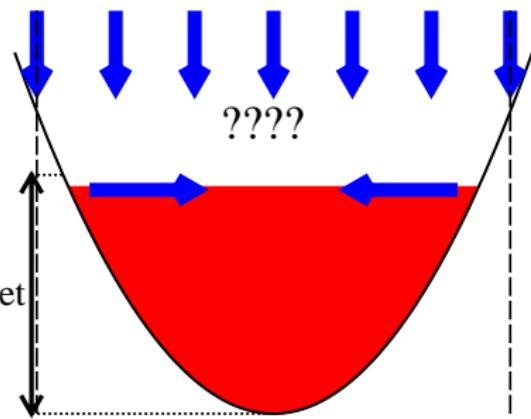
$$\frac{1}{2}\partial_t\rho + \nabla \cdot (\rho \mathbf{v}) = \frac{1}{\hbar}(\gamma_{\text{net}}\Theta(R-r) - \Gamma\rho)\rho$$



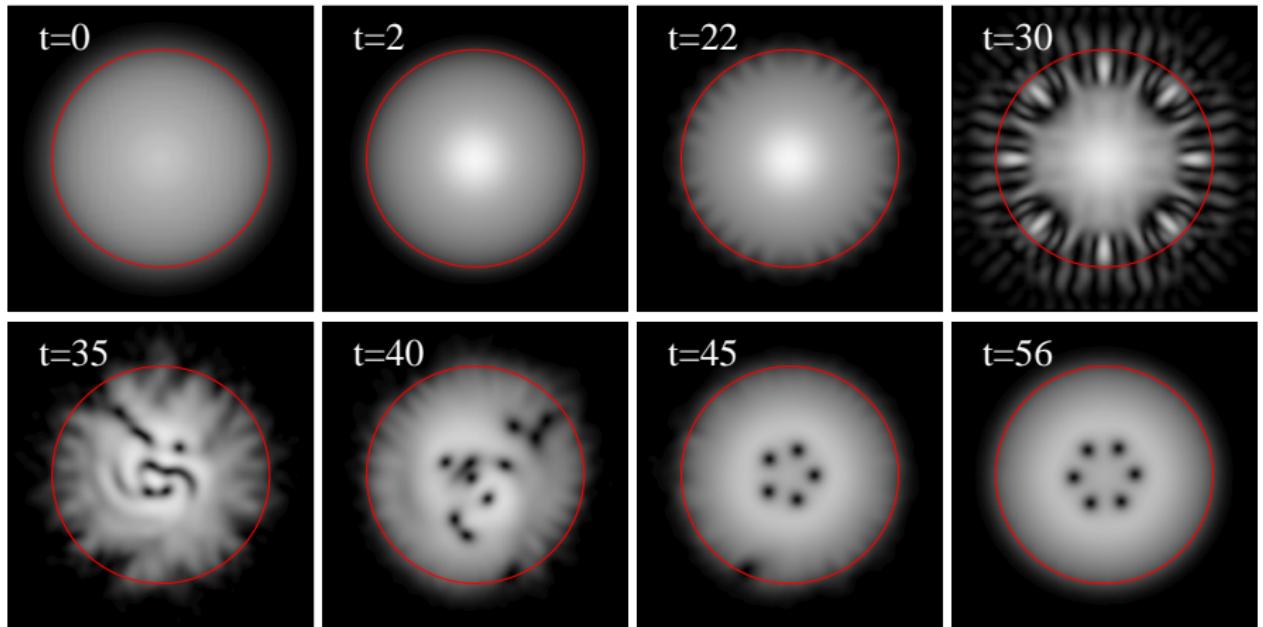
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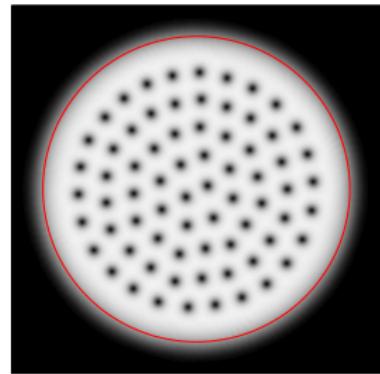
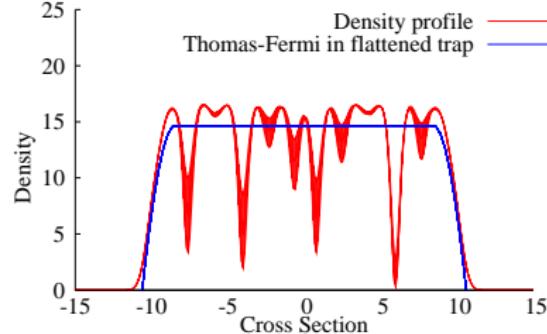


Time evolution:



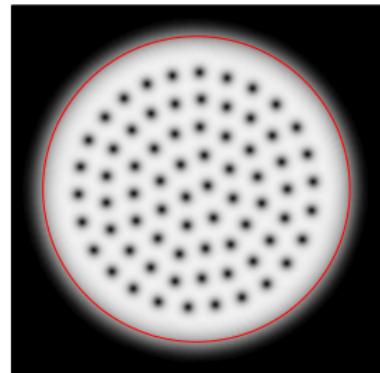
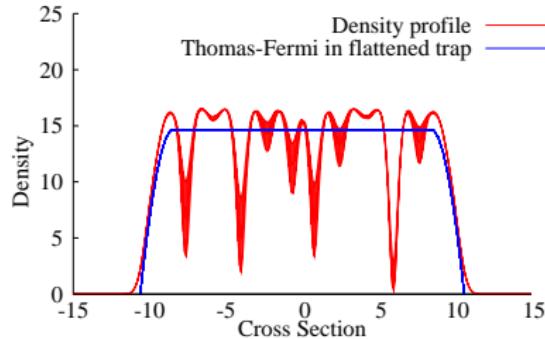
[Keeling & Berloff, PRL, '08]

Why vortices



$$\nabla \cdot [p(v - \Omega \times r)] = (\gamma_{\text{ho}} \Theta(R - r) - \Gamma_p) p$$
$$p = \frac{\hbar}{2} |v - \Omega \times r|^2 + \frac{\hbar^2}{2} r^2 (\omega^2 - \Omega^2) + U_p - \frac{\hbar^2 \nabla^2 \sqrt{p}}{2m_e \beta}$$
$$v = \Omega \times r, \quad \Omega = \omega, \quad \beta = \frac{2m_e}{\hbar^2} \gamma_{\text{ho}} (\Omega - r) = \frac{R}{\hbar}$$

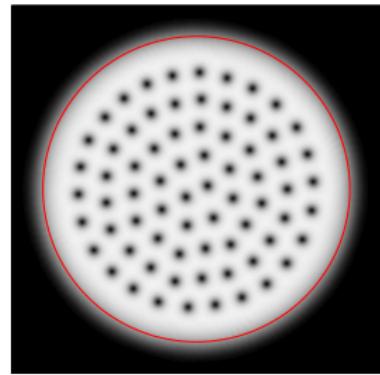
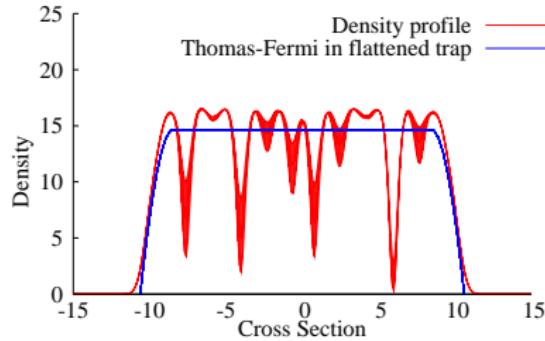
Why vortices



Rotating solution: $i\partial_t\psi = (\mu - 2\Omega L_z)\psi$

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Why vortices

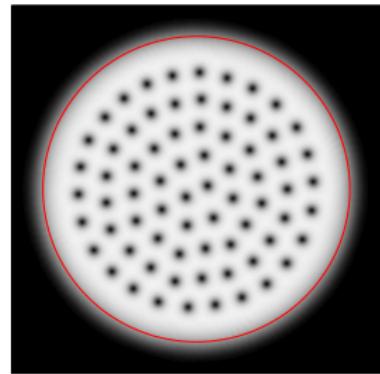
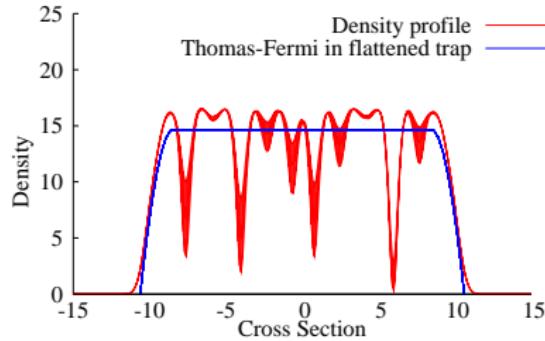


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Why vortices



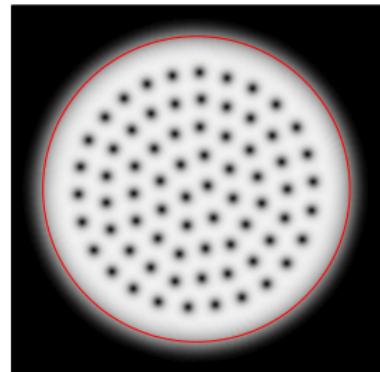
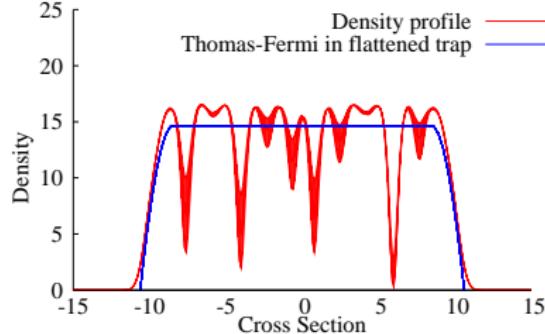
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Why vortices



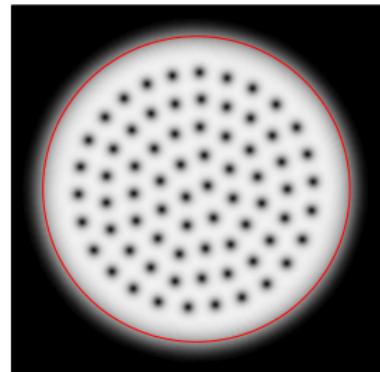
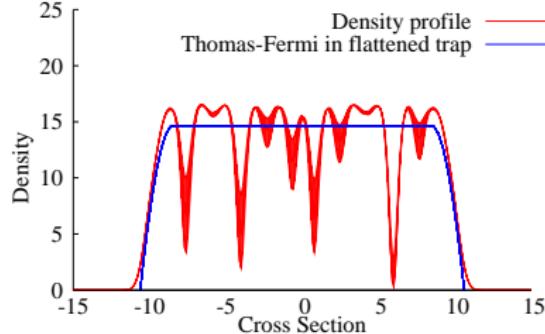
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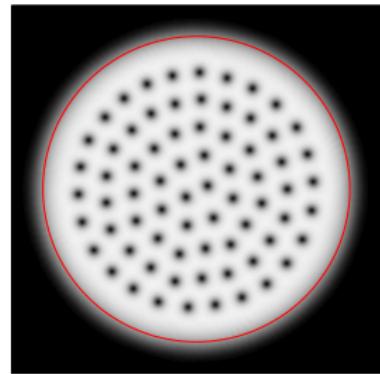
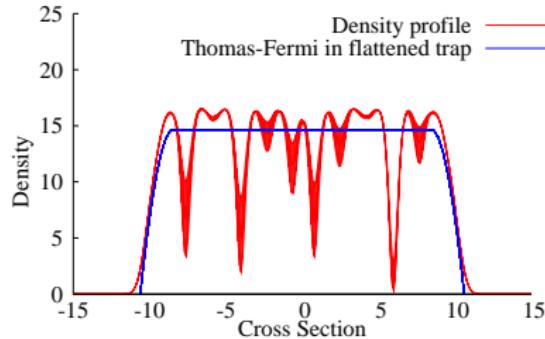


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Why vortices

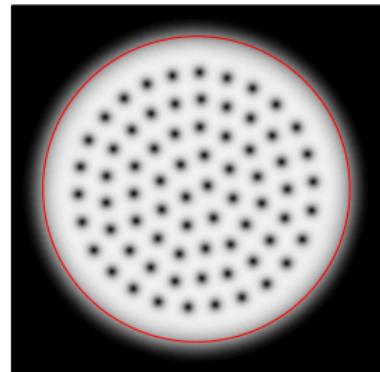
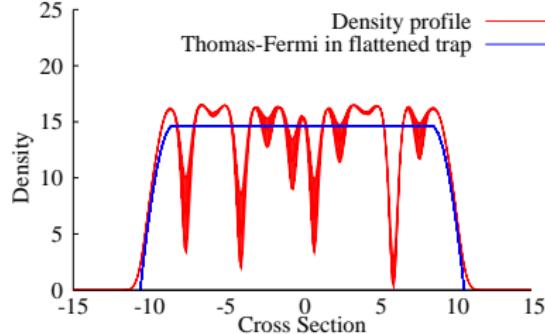


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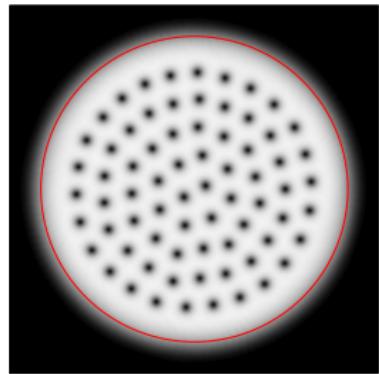
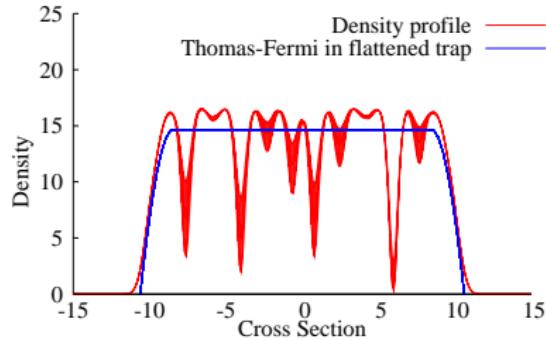
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Why vortices



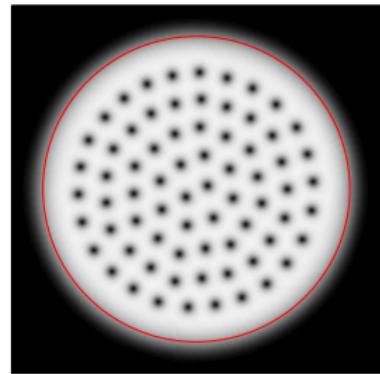
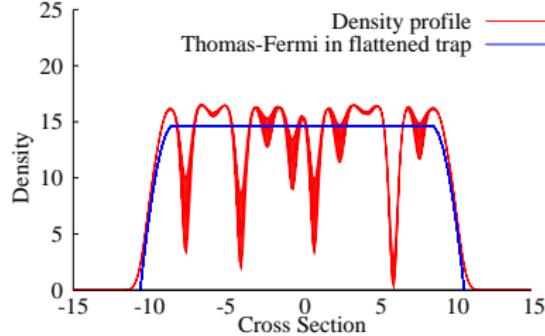
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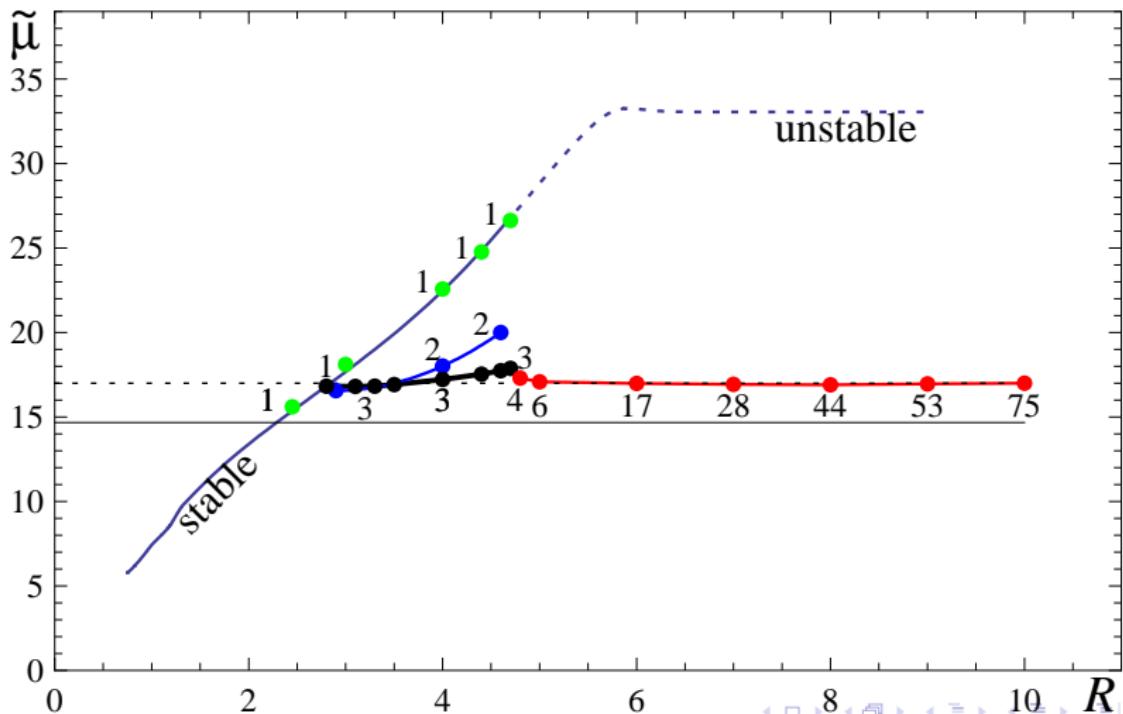
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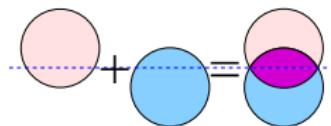
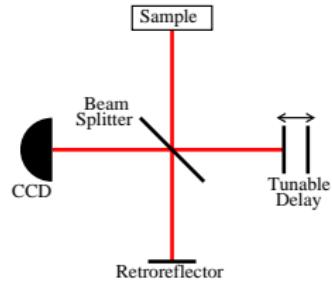
$$\mathbf{v} = \boldsymbol{\Omega} \times \mathbf{r}, \quad \boldsymbol{\Omega} = \omega, \quad \rho = \frac{\gamma_{\text{net}}}{\Gamma}\Theta(R - r) = \frac{\mu}{U}$$

Why vortices: chemical potential vs size

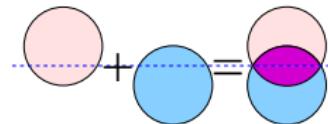
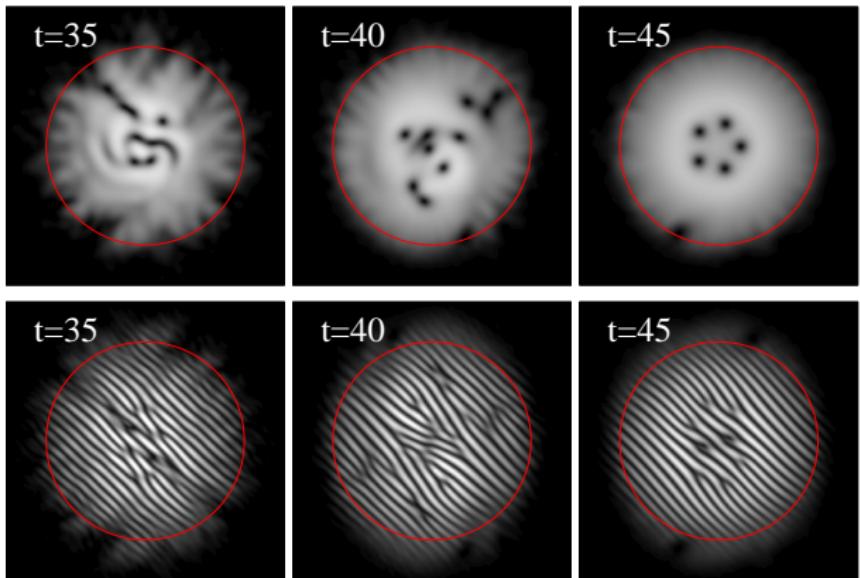
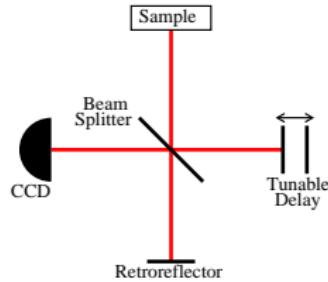
$$\text{Thomas-Fermi : } \mu \propto R^2 \quad \text{Vortex : } \mu = \frac{U\gamma_{\text{net}}}{\Gamma}$$



Observing vortices: fringe pattern



Observing vortices: fringe pattern



Overview

- 1 Introduction to microcavity polaritons
- 2 Model and review of equilibrium results
 - Disorder-localised exciton model
 - Basic equilibrium results
- 3 Non-equilibrium model and mean-field theory
 - Meaning of mean-field condition
- 4 Macroscopic phenomenology
 - Gross Pitaevskii equation in an harmonic trap
- 5 Fluctuations and correlations
 - Fluctuations about mean-field theory
 - Finite size effects: single vs many modes

Fluctuations → Stability, Luminescence, Absorption

Keldysh approach:

$$\mathcal{G}_{R,A} = \mp i\theta[\pm(t - t')] \left\langle [\psi^\dagger, \psi]_- \right\rangle$$

$$\mathcal{G}_K = -i \left\langle [\psi^\dagger, \psi]_+ \right\rangle$$

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$$\mathcal{G}_R - \mathcal{G}_A = \frac{1}{\mathcal{G}_R^{-1}} - \frac{1}{[\mathcal{G}_R^{-1}]^\dagger}$$

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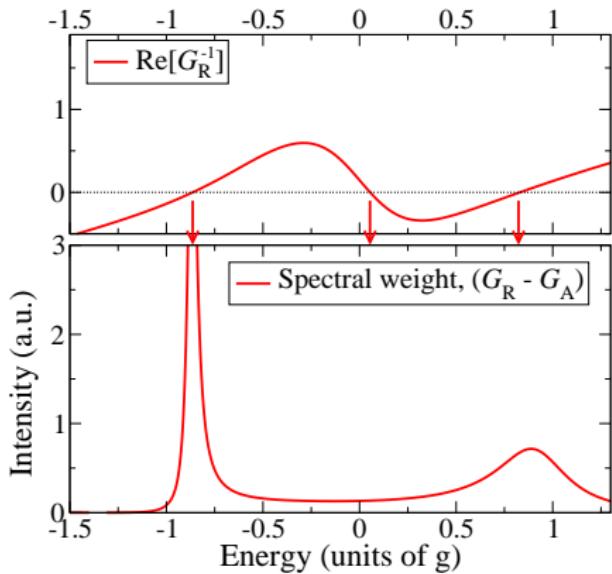
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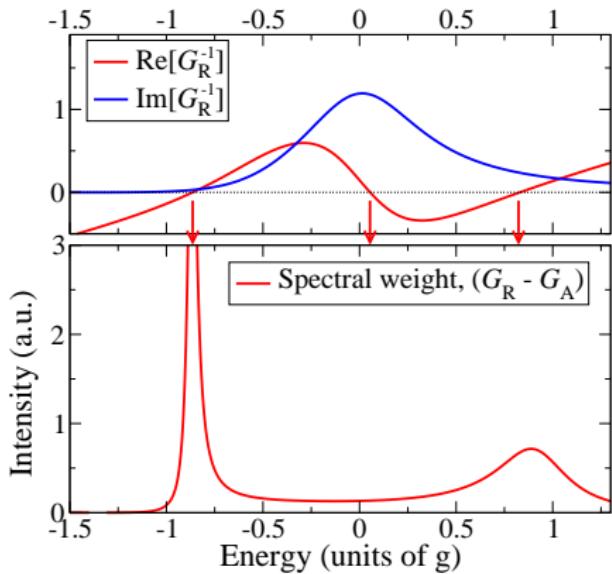
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Fluctuations → Stability, Luminescence, Absorption

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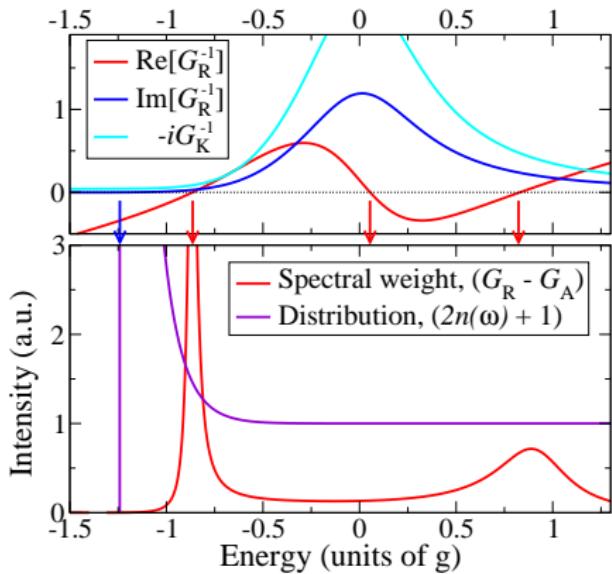
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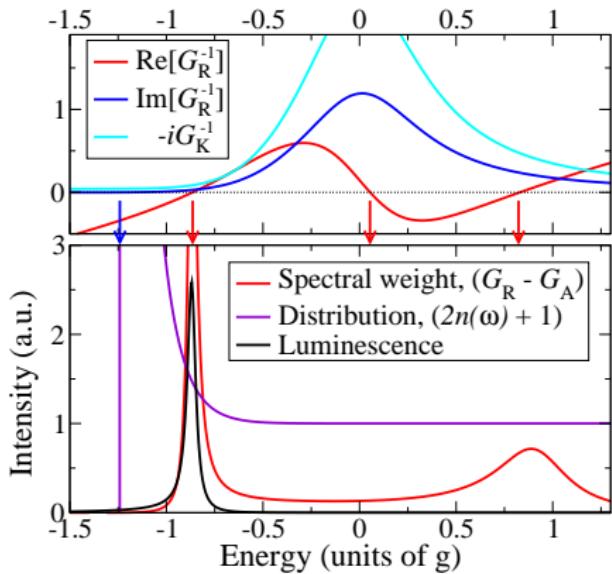
$$\mathcal{G}_K = -i \left\langle [\psi^\dagger, \psi]_+ \right\rangle = (2n(\omega) + 1)(\mathcal{G}_R - \mathcal{G}_A)$$

$$\mathcal{L} = \left\langle \psi^\dagger \psi \right\rangle = \frac{i}{2} [\mathcal{G}_K + (\mathcal{G}_R - \mathcal{G}_A)]$$

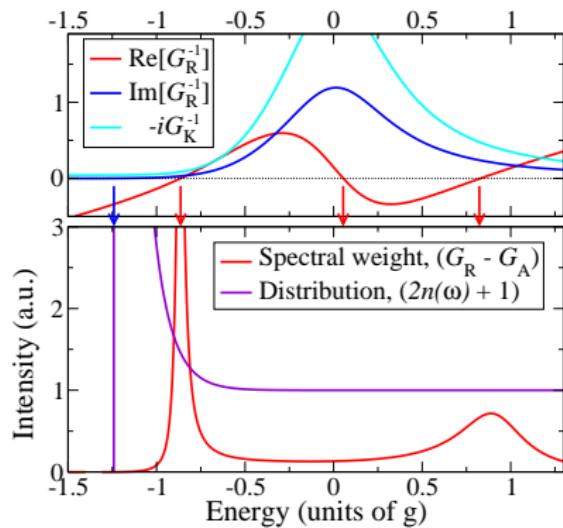
$$\mathcal{G}_K = \frac{-\mathcal{G}_K^{-1}}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

$$\mathcal{G}_R - \mathcal{G}_A = \frac{2\text{Im}[\mathcal{G}_R^{-1}]}{\text{Im}[\mathcal{G}_R^{-1}]^2 + \text{Re}[\mathcal{G}_R^{-1}]^2}$$

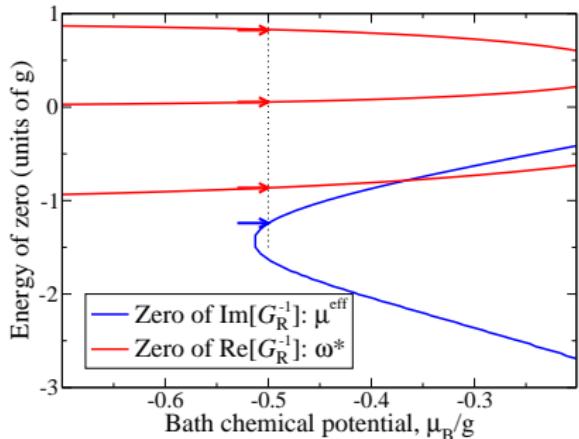
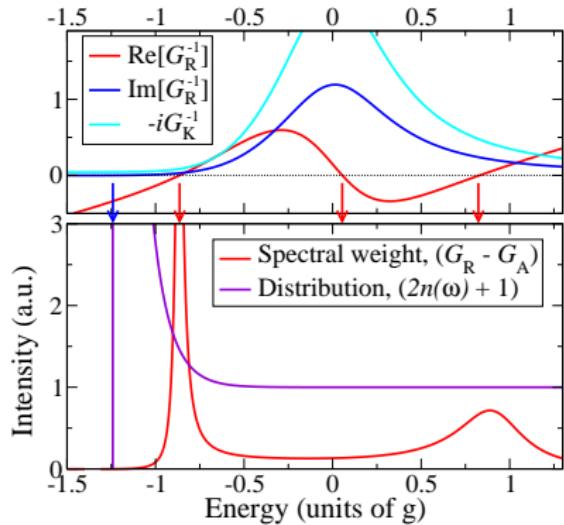
$$\mathcal{G}_R^{-1}(\omega, k) = (\omega - \omega_k^*) + i\alpha(\omega - \mu_{\text{eff}})$$



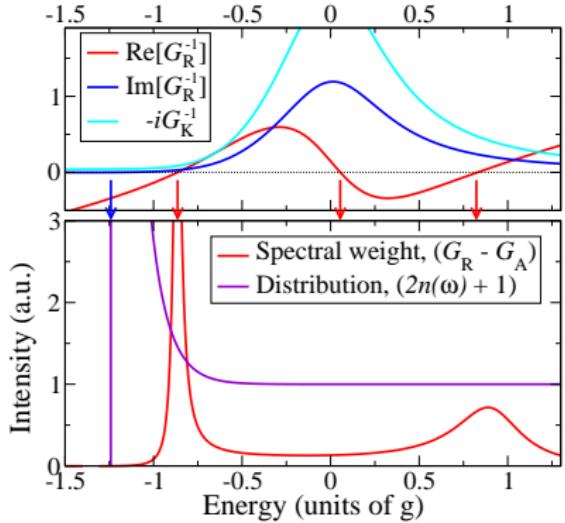
Linewidth, inverse Green's function and gap equation



Linewidth, inverse Green's function and gap equation

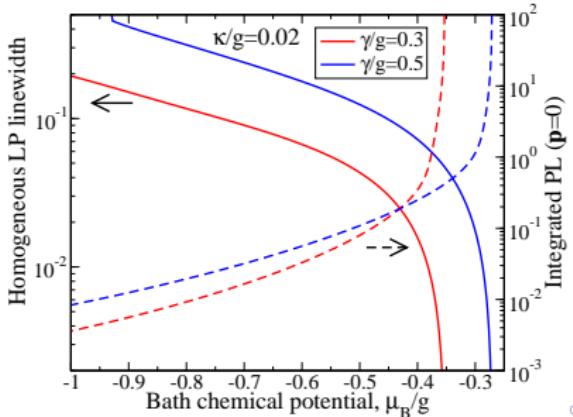
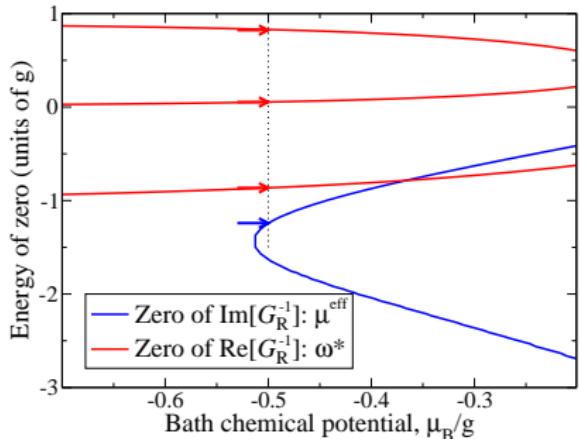


Linewidth, inverse Green's function and gap equation



At transition, Gap Equation implies:

$$\mathcal{G}_R^{-1}(\omega = \mu_S, k = 0) = 0$$



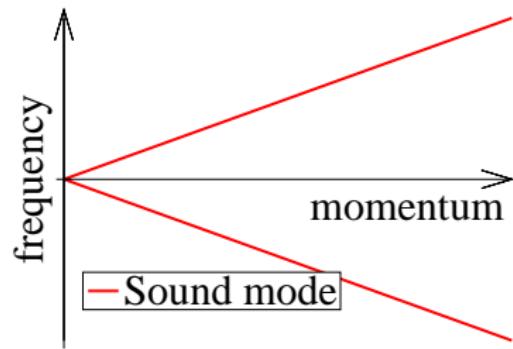
Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = \omega^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = c|\mathbf{k}|$$



[Szymańska et al., PRL '06; PRB '07]

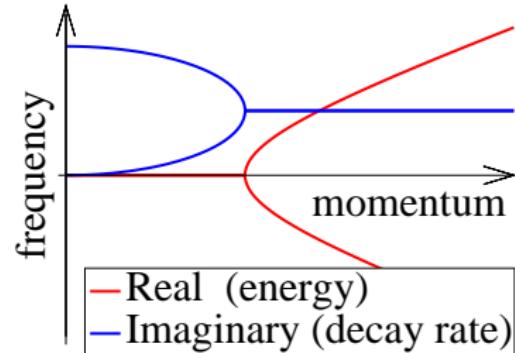
Fluctuations above transition

When condensed

$$\text{Det} [\mathcal{G}_R^{-1}(\omega, \mathbf{k})] = (\omega + ix)^2 + x^2 - c^2 \mathbf{k}^2$$

Poles:

$$\omega^* = -ix \pm \sqrt{c^2 \mathbf{k}^2 - x^2}$$



[Szymańska et al., PRL '06; PRB '07]

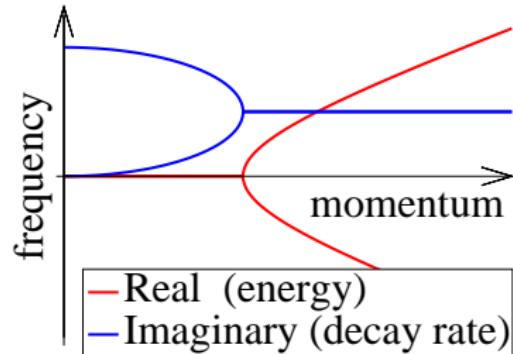
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Correlations (in 2D): $\langle \psi^\dagger(\mathbf{r}, t)\psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$

[Szymańska et al., PRL '06; PRB '07]

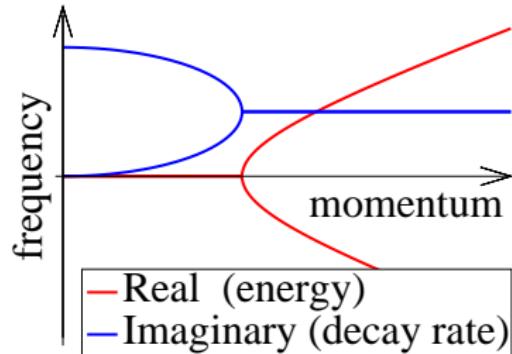
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Correlations (in 2D): $\langle \psi^\dagger(\mathbf{r}, t)\psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$

$$\langle \psi^\dagger(\mathbf{r}, t)\psi(0, 0) \rangle \simeq |\psi_0|^2 \exp \left[-\eta \begin{cases} \ln(r/\xi) & r \rightarrow \infty, t \simeq 0 \\ \frac{1}{2} \ln(c^2 t / x \xi^2) & r \simeq, t \rightarrow \infty \end{cases} \right]$$

[Szymańska et al., PRL '06; PRB '07]

Finite size effects: Single mode vs many mode

$$\langle \psi^\dagger(r, t)\psi(0, r') \rangle \simeq |\psi_0|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$$

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$\mathcal{D}_{\phi\phi}(t, r, r')$ from sum of phase modes. Study $ct \gg r$ limit:

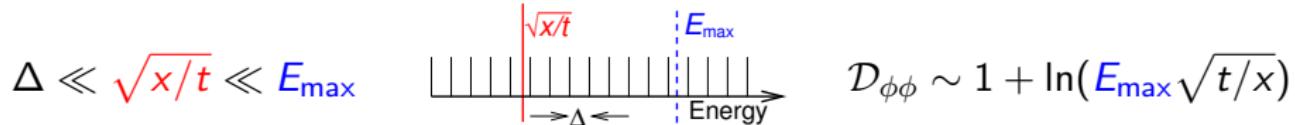
$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

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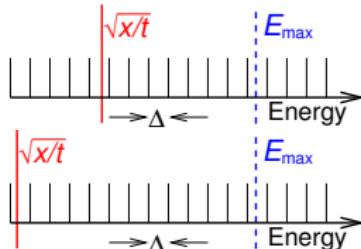
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$$\Delta \ll \sqrt{x/t} \ll E_{\max}$$



$$\mathcal{D}_{\phi\phi} \sim 1 + \ln(E_{\max} \sqrt{t/x})$$

$$\sqrt{x/t} \ll \Delta \ll E_{\max}$$

$$\mathcal{D}_{\phi\phi} \sim \left(\frac{\pi C}{2x}\right) \left(\frac{t}{2x}\right)$$

Relating finite-size spectrum to self phase modulation

Finite-size spectrum:

$$\mathcal{D}_{\phi\phi}(t, r, r) \propto \sum_n^{n_{max}} \int \frac{d\omega}{2\pi} \frac{|\varphi_n(r)|^2 (1 - e^{i\omega t})}{|(\omega + ix)^2 + x^2 - (n\Delta)^2|^2}$$

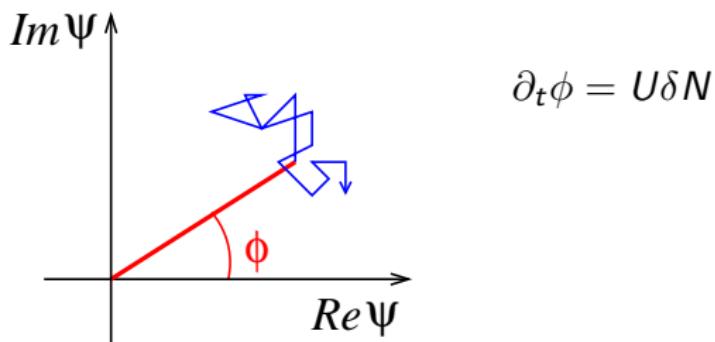
Connection to Kubo Formula [Eastham and Whittaker, arXiv:0811.4333]

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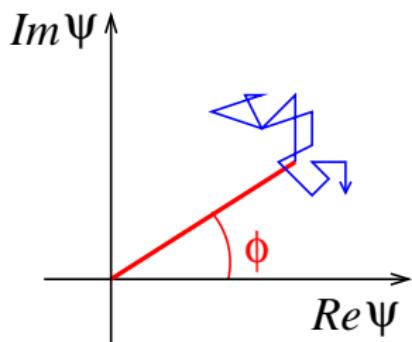


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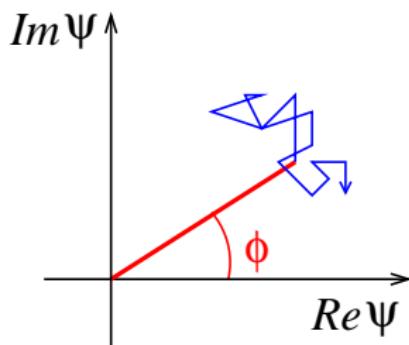
$$\begin{aligned}\partial_t \phi &= U \delta N \\ \partial_t \delta N &= -\Gamma \delta N + F(t), \quad \langle F(t) F(t') \rangle = C \delta(t - t')\end{aligned}$$

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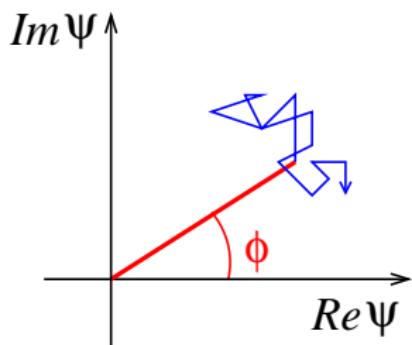
$$\begin{aligned}\partial_t \phi &= U \delta N \\ \partial_t \delta N &= -\Gamma \delta N + F(t), \quad \langle F(t) F(t') \rangle = C \delta(t - t') \\ \mathcal{D}_{\phi\phi}(t) &= \int \frac{d\omega}{2\pi} (1 - e^{i\omega t}) \langle \phi_\omega \phi_{-\omega} \rangle\end{aligned}$$

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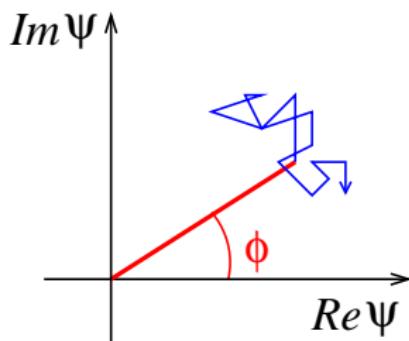
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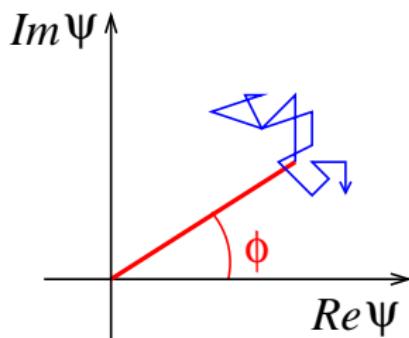
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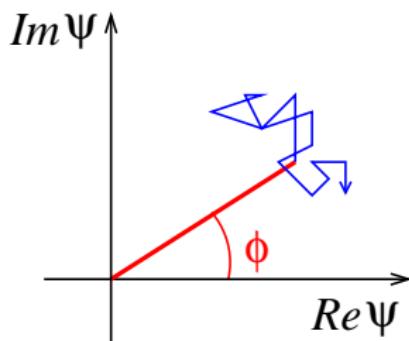
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Relating finite-size spectrum to self phase modulation

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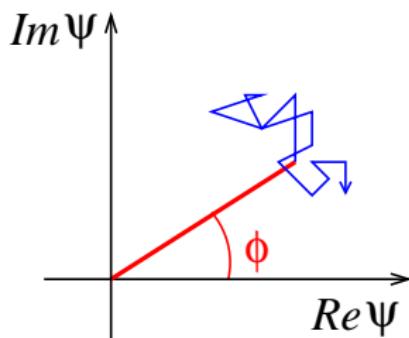
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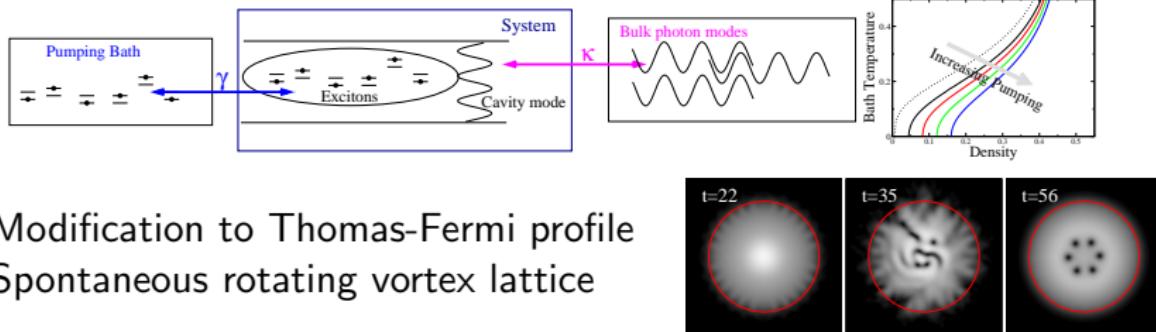
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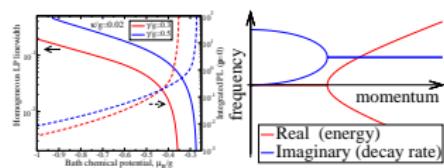
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Conclusions

- Localised two-level system model
- Mean-field and fluctuations
- Effects of pumping on mean-field theory



- Modification to Thomas-Fermi profile
- Spontaneous rotating vortex lattice
- Change to spectrum and correlations
- Phase modes and finite size



Acknowledgements

People:



Funding:

EPSRC Engineering and Physical Sciences Research Council



Pembroke College

Extra slides

6 Superfluidity

7 Zero temperature Keldysh boundaries

Superfluidity

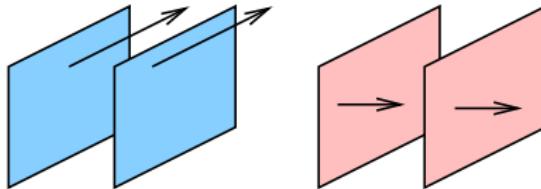
Superfluidity:

$$\mathbf{J} = \rho \mathbf{v} = \Psi^\dagger i\hbar \nabla \Psi = |\Psi|^2 \nabla \phi$$

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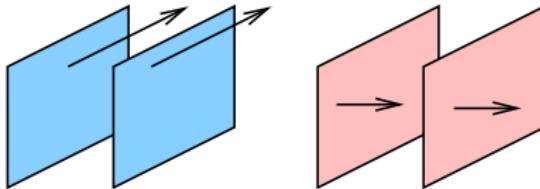


$$\begin{aligned}\chi_{ij}(\omega = 0, \mathbf{q} \rightarrow 0) &= \langle J_i(\mathbf{q}) J_j(\mathbf{q}) \rangle \\ &= \chi_T(q) \left(\delta_{ij} - \frac{q_i q_j}{q^2} \right) + \chi_L(q) \frac{q_i q_j}{q^2}\end{aligned}$$

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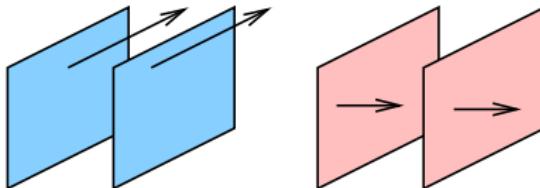
Superfluid part,

$$\rho_s \propto \lim_{q \rightarrow 0} (\chi_L - \chi_T).$$

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Superfluid part,
 $\rho_s \propto \lim_{q \rightarrow 0} (\chi_L - \chi_T)$.

$$J_i(q) = \langle \psi^\dagger \gamma_i(q) \psi \rangle$$

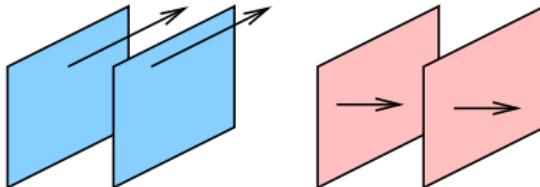
$$\Delta \chi_{ij}(q) = \text{---} \bullet \xrightarrow[\mathcal{G}(\omega = 0, \mathbf{q})]{} \bullet \text{---} + \dots$$

$\gamma_i(\mathbf{q}, 0) \psi_0$ $\gamma_j(\mathbf{q}, 0) \psi_0$

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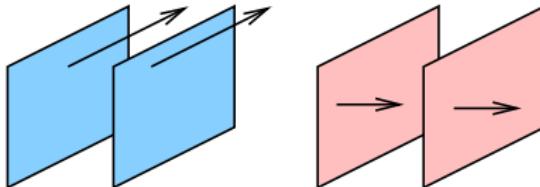
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$$\begin{aligned}\Delta \chi_{ij}(q) &= \text{---} \overset{\gamma_i(\mathbf{q}, 0)\psi_0}{\curvearrowleft} \overset{\gamma_j(\mathbf{q}, 0)\psi_0}{\curvearrowright} \underset{\mathcal{G}(\omega = 0, \mathbf{q})}{\text{---}} + \dots \\ &= \gamma_i(q) \mathcal{G}_R(\omega = 0, \vec{q}) \gamma_j(q) + \dots\end{aligned}$$

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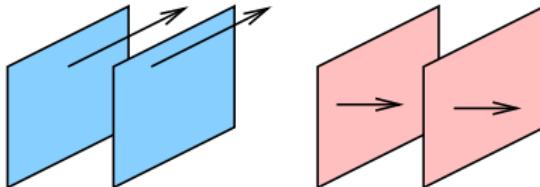
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 $J_i(q) = \langle \psi^\dagger \gamma_i(q) \psi \rangle$

$$\begin{aligned}\Delta \chi_{ij}(q) &= \text{Diagram showing two wavy lines connected by a horizontal arrow labeled } \mathcal{G}(\omega = 0, \mathbf{q}) \text{ with vertices } \gamma_i(\mathbf{q}, 0)\psi_0 \text{ and } \gamma_j(\mathbf{q}, 0)\psi_0 \\ &= \gamma_i(q) \mathcal{G}_R(\omega = 0, \vec{q}) \gamma_j(q) + \dots \\ &\propto \frac{q_i (\dots) q_j}{(\omega + ix)^2 + x^2 - c^2 \mathbf{q}^2} + \dots\end{aligned}$$

Superfluidity

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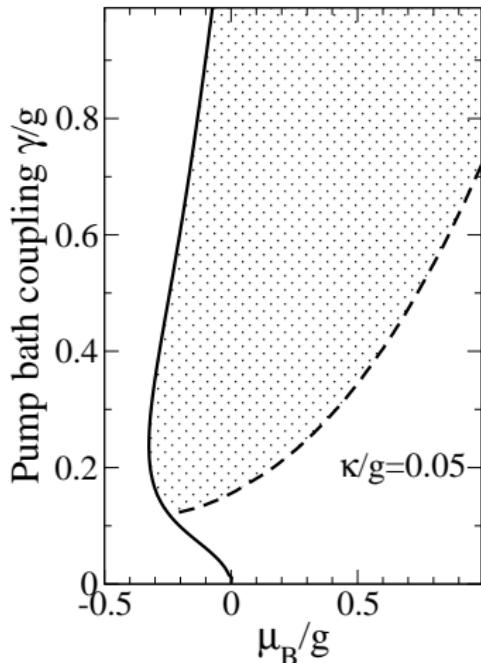
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Static ρ_S survives

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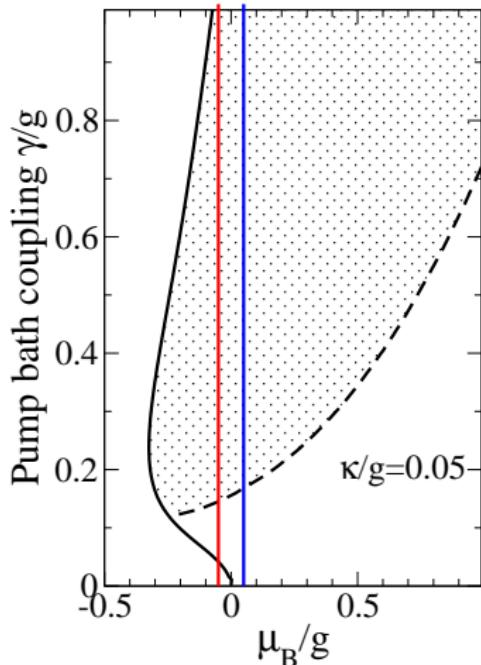
Zero temperature phase diagram

$$\tilde{\omega}_0 - i\kappa = g^2 \gamma \sum_{\alpha} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(\tilde{\epsilon}_{\alpha} + i\gamma)}{[(\nu - E_{\alpha})^2 + \gamma^2][(v + E_{\alpha})^2 + \gamma^2]}.$$



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